

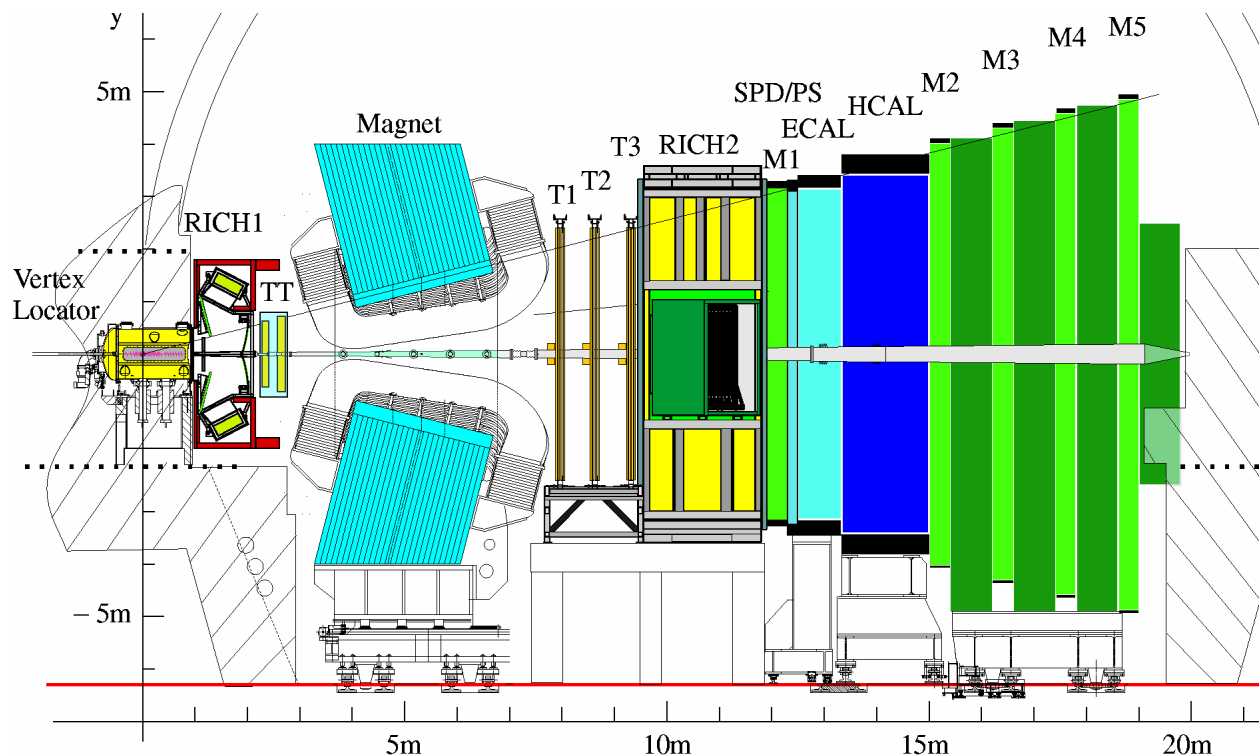
LHCb semileptonic highlights from 2010 and prospects for 2011

Robert W. Lambert

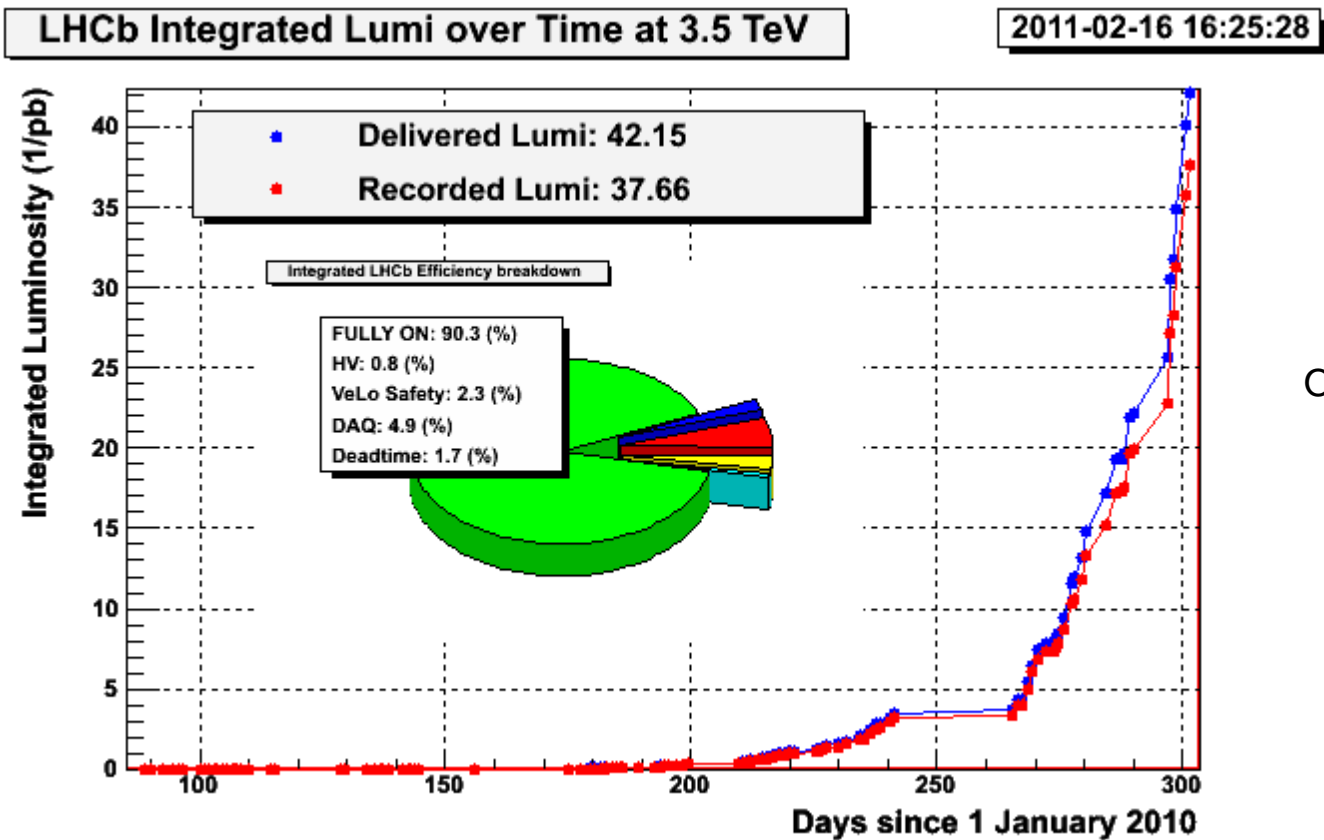
on behalf of the LHCb Collaboration

- LHC and LHCb
- Semileptonics in general
- Motivation and measurement of $B_s^0 \rightarrow D_{s1,2}^{*\pm} \mu^{\mp} \nu X^0$
- Flavour specific asymmetry

- LHCb is a dedicated, precision, *b*-physics experiment
- High statistics: we're in the forward region, and at LHC

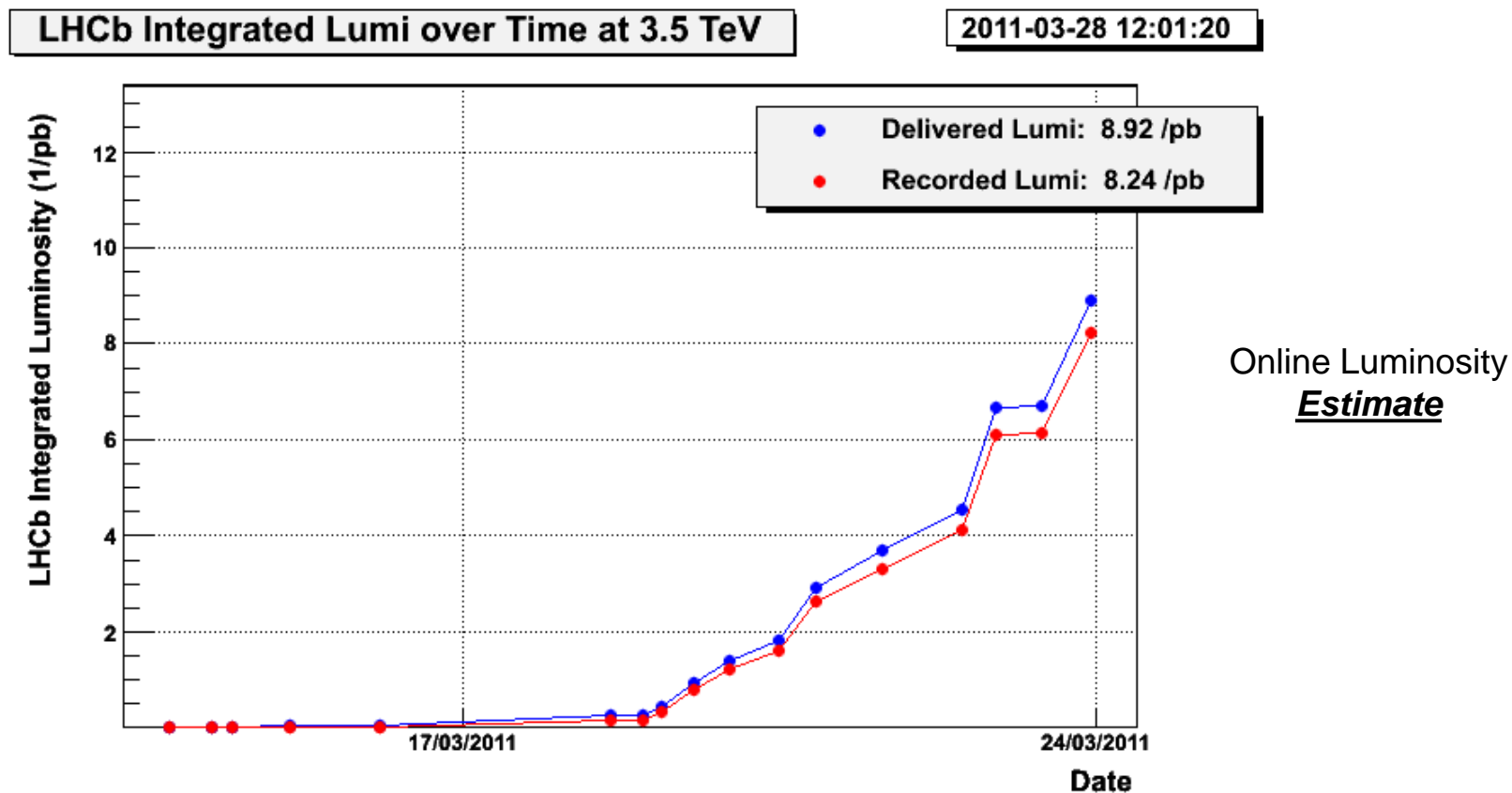


- The LHC delivered $\sim 40 \text{ pb}^{-1}$ in 2010



Online Luminosity
Estimate

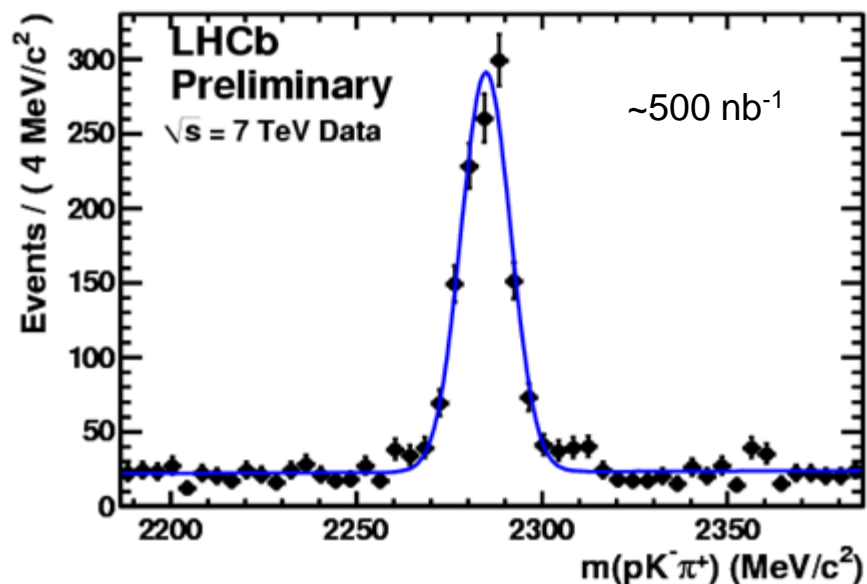
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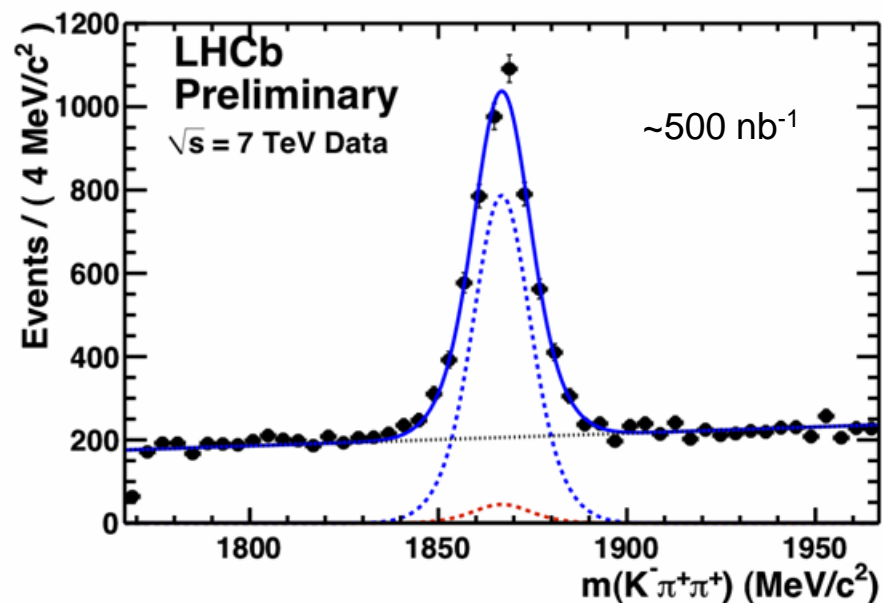
- 25% more already in 2011!

- LHCb reconstructs charmed and charmless B-decays
 - Backgrounds from open charm and other light mesons

$$\Lambda_b \rightarrow \Lambda_c \mu \nu X$$



$$X_b \rightarrow D^+ \mu \nu X$$



- ... just a sample of the many reconstructed modes

Production and decay

- Spectroscopy
- Hadronization fractions
- Form factors
- V_{cb} , V_{ub}

B-meson mixing

- Lifetime
- Mass difference
- Width difference
- Flavour-specific asymmetry

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2010, 2011

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2010, 2011

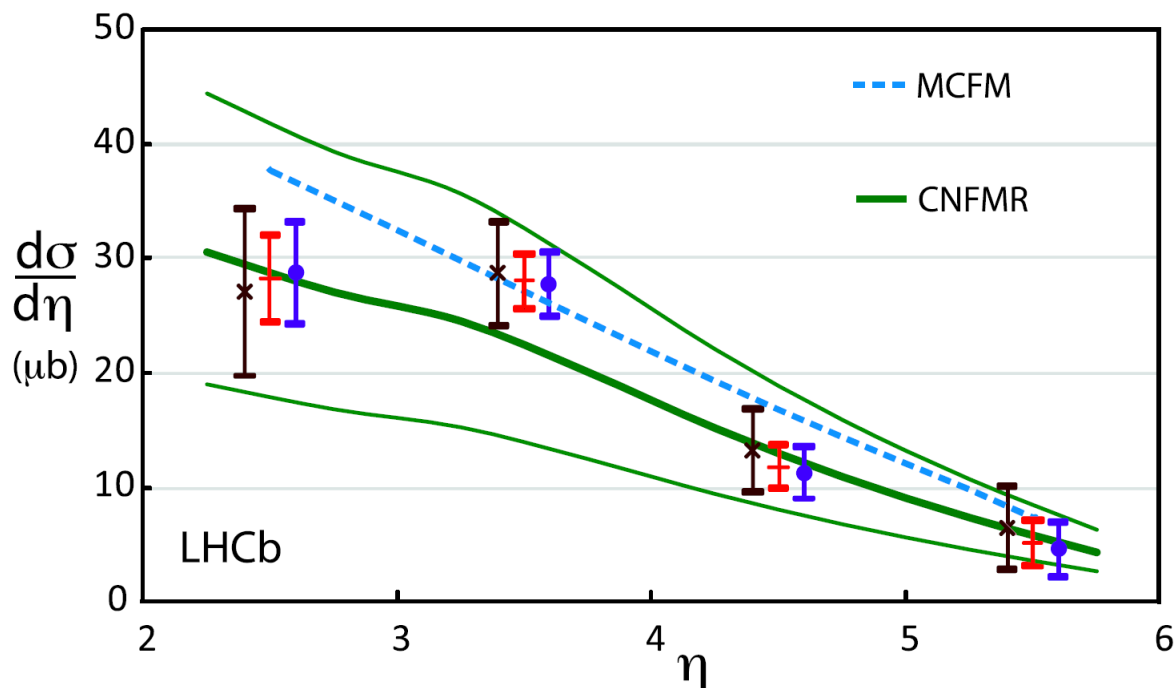
Highlights to cover here:

- First LHC *b*-physics paper
- New D^{**} decays of *B*-mesons
- Prospects for flavour-specific asymmetry

See Niels Tuning
from Tuesday,
and
Stefania Vecchi
from Monday

Phys. Lett. B 694 pp. 209
arxiv:1009.2731

- First LHCb b -physics paper, $\mathcal{L}=(15.0\pm 1.5)\text{nb}^{-1}$
- $\sigma_{b\bar{b}}$ from semileptonic decays of the form $X_b \rightarrow D^0 \mu \nu X$

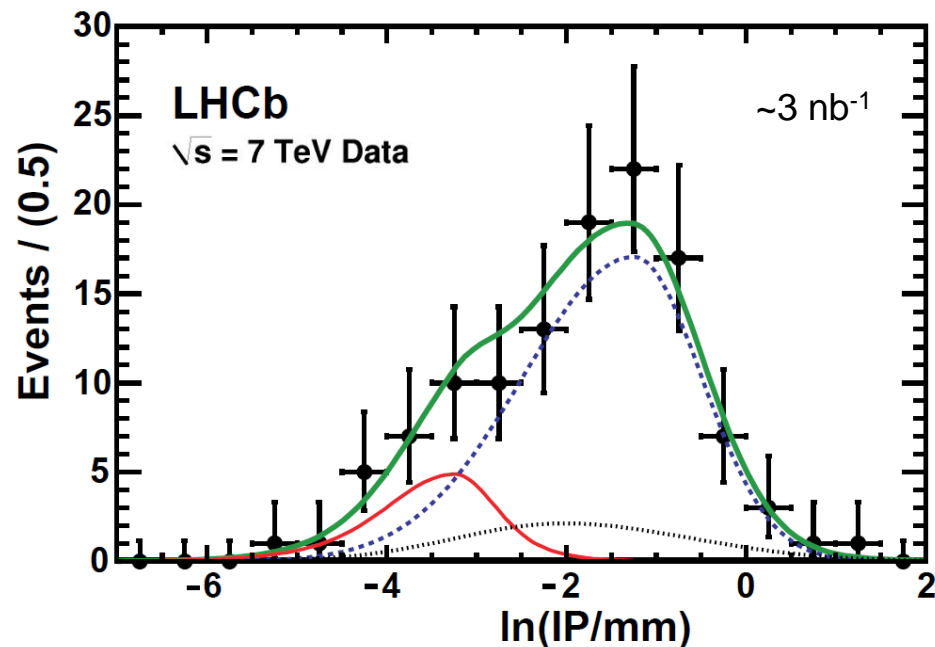
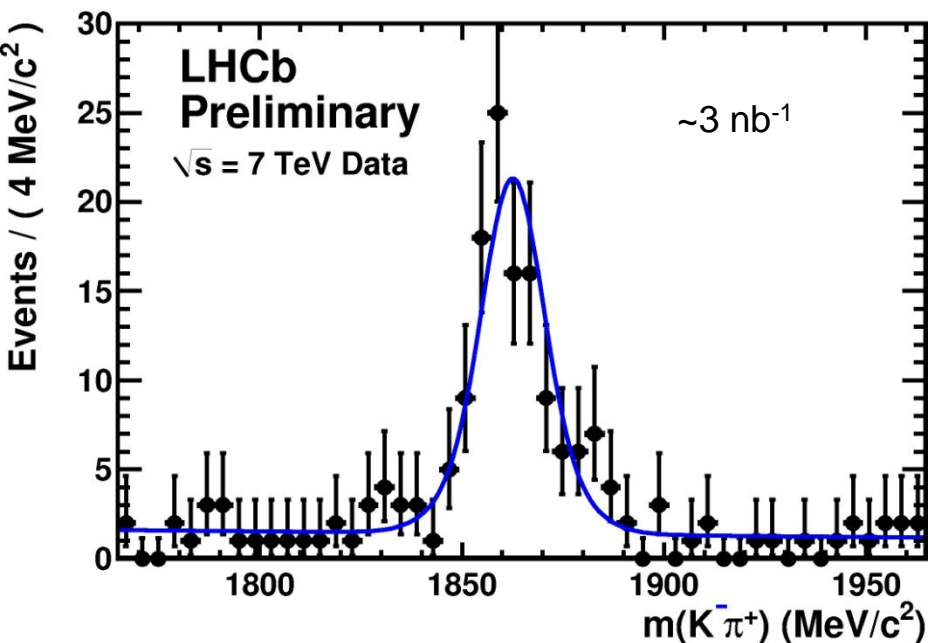


See Niels Tuning
from Tuesday

- Scale to full rapidity: $\sigma(pp \rightarrow b\bar{b} X) = (284 \pm 20 \pm 49) \mu\text{b}$

Phys. Lett. B 694 pp. 209
arxiv:1009.2731

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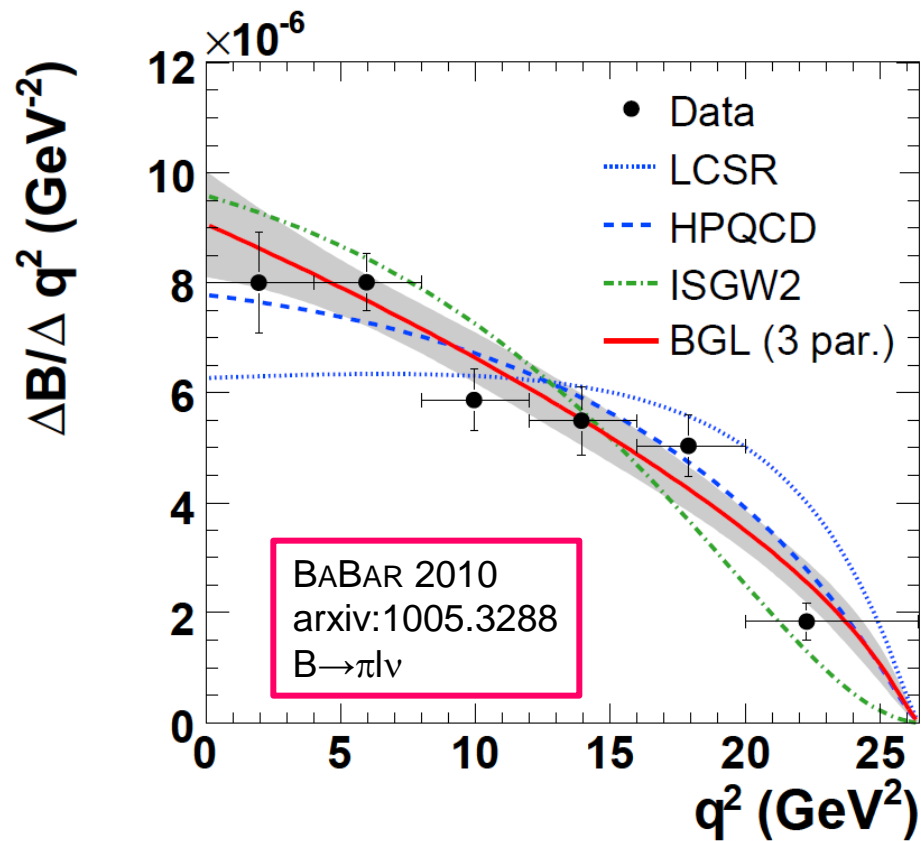
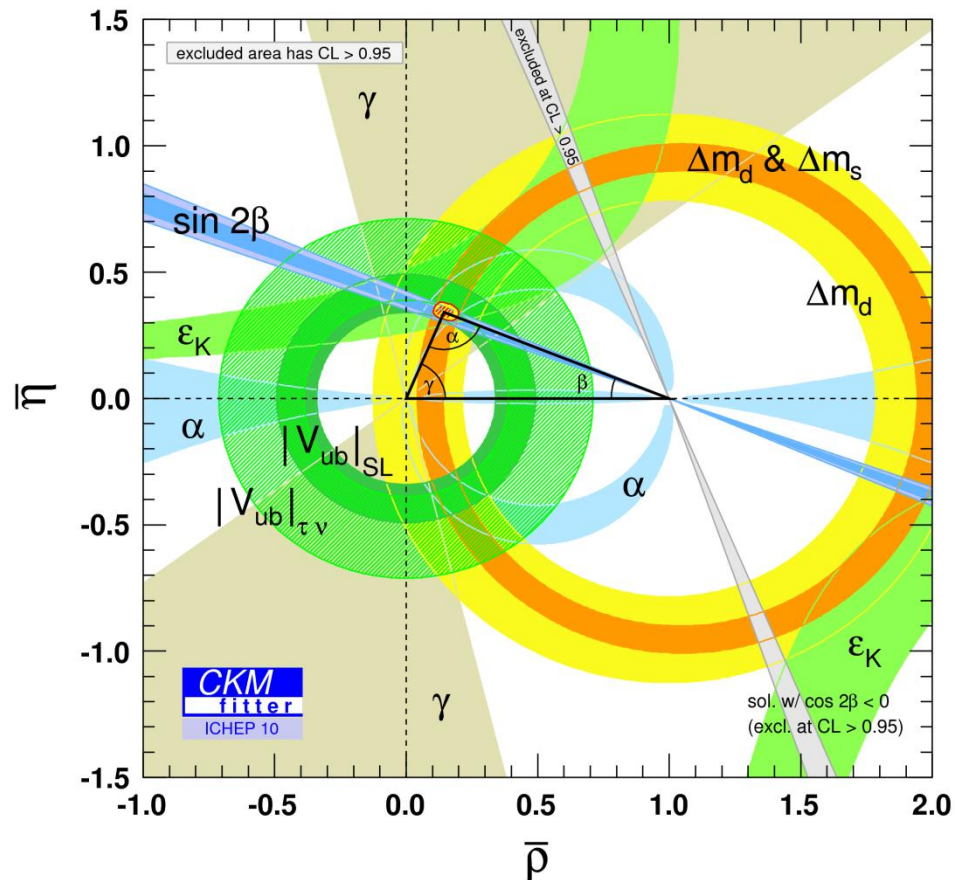


- $\ln(\text{IP})$ distribution used to discriminate prompt contribution

Progress towards $V_{(u/c)b}$

Measurement of $B_s^0 \rightarrow D_{s1,2}^{*\pm} \mu^{\mp} \nu X^0$

- Constrains UT and/or measures form factors

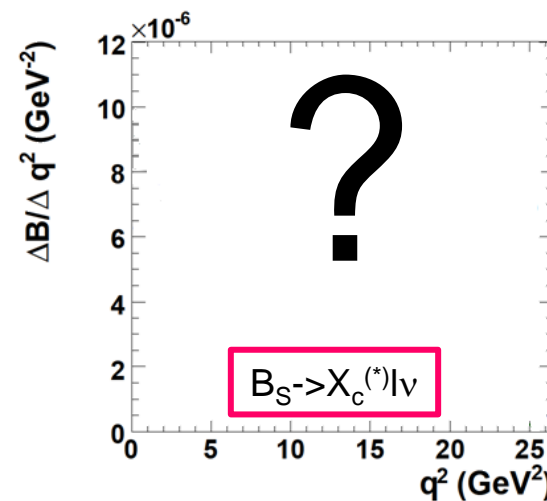
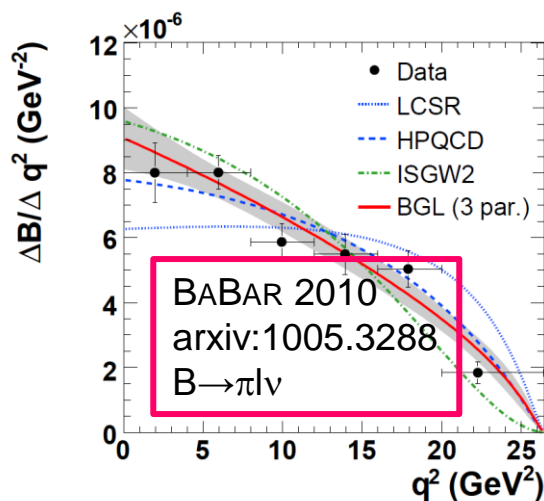


- LHCb also probes $V_{(u/c)b}$ in the B_s system

➤ To determine V_{ub} , we must first understand V_{cb}

X In the B_s system, very little is known phenomenologically

X Form factors unknown, expected to differ in q^2 shape



✓ Theory interest and input encouraged

➤ To determine V_{ub} , we must first understand V_{cb}

X In the B_s system, very little is known experimentally

X Excited D resonances unknown

B^0 DECAY MODES	Fraction (Γ_i/Γ)	Scale factor/ Confidence level	ρ (MeV/c)
$D\ell^+\nu_\ell$ anything	(9.3 \pm 0.9) %	—	—
$D^-\ell^+\nu_\ell$	[nnn] (2.17 \pm 0.12) %	2309	—
$D^-\tau^+\nu_\tau$	(1.1 \pm 0.4) %	1909	—
$D^*(2010)^-\ell^+\nu_\ell$	[nnn] (5.01 \pm 0.12) %	2257	—
$D^*(2010)^-\tau^+\nu_\tau$	(1.5 \pm 0.5) %	1837	—
$\bar{D}^0\pi^-\ell^+\nu_\ell$	(4.3 \pm 0.6) $\times 10^{-3}$	S=1.3	2308
$D_0^{*+}(2400)^-\ell^+\nu_\ell \times$ $B(D_0^{*+} \rightarrow \bar{D}^0\pi^-)$	(3.0 \pm 1.2) $\times 10^{-3}$	S=1.8	—
$D_2^{*+}(2460)^-\ell^+\nu_\ell \times$ $B(D_2^{*+} \rightarrow \bar{D}^0\pi^-)$	(2.2 \pm 0.6) $\times 10^{-3}$	—	2067
$\bar{D}^{(*)}\pi\pi^+\ell^+\nu_\ell$ (n \geq 1)	(2.3 \pm 0.5) %	—	—
$\bar{D}^{*0}\pi^-\ell^+\nu_\ell$	(4.9 \pm 0.8) $\times 10^{-3}$	2256	—
$D_1(2420)^-\ell^+\nu_\ell \times B(D_1^- \rightarrow \bar{D}^{*0}\pi^-)$	(2.80 \pm 0.28) $\times 10^{-3}$	—	—
$D_1'(2430)^-\ell^+\nu_\ell \times B(D_1'^- \rightarrow \bar{D}^{*0}\pi^-)$	(3.1 \pm 0.9) $\times 10^{-3}$	—	—
$D_2^{*+}(2460)^-\ell^+\nu_\ell \times$ $B(D_2^{*+} \rightarrow \bar{D}^{*0}\pi^-)$	(1.2 \pm 0.5) $\times 10^{-3}$	S=2.7	2067
$\rho^-\ell^+\nu_\ell$	[nnn] (2.47 \pm 0.33) $\times 10^{-4}$	—	2583
$\pi^-\ell^+\nu_\ell$	[nnn] (1.34 \pm 0.08) $\times 10^{-4}$	—	2638



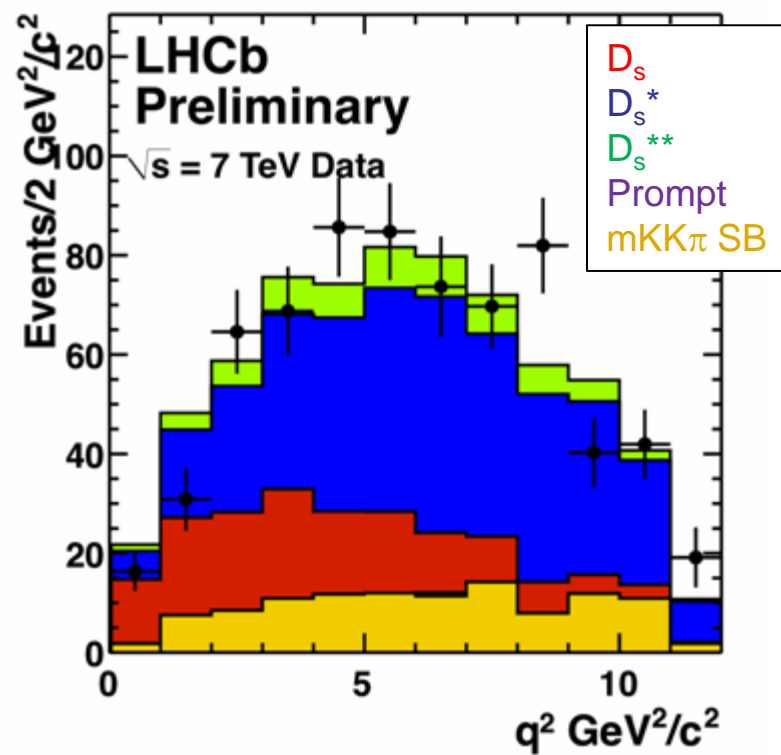
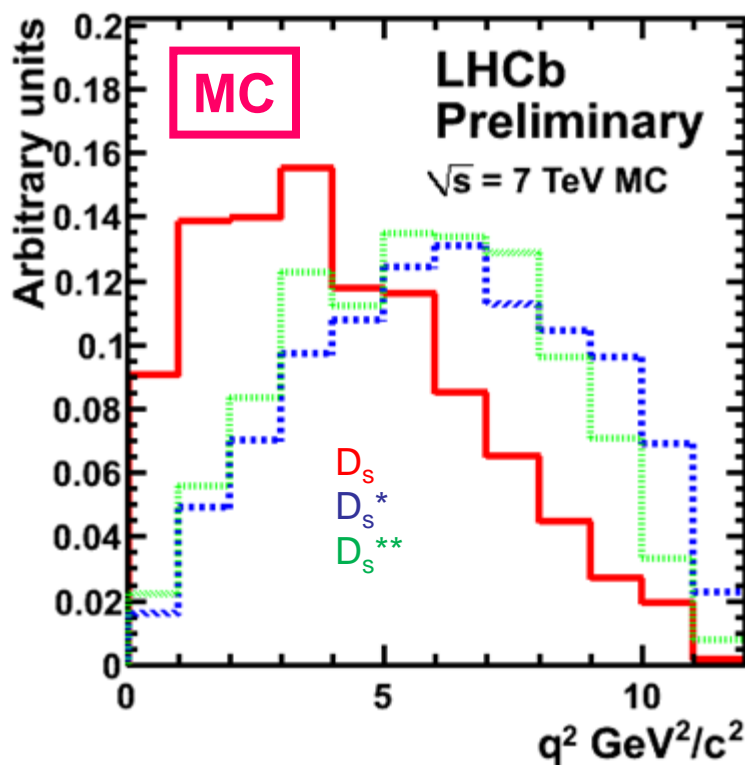
(my copy of
the PDG 2010)



B_s^0 DECAY MODES	Fraction (Γ_i/Γ)	Confidence level	ρ (MeV/c)
D_s^- anything	(93 \pm 25) %	—	—
$D_s^-\ell^+\nu_\ell$ anything	[8888] (7.9 \pm 2.4) %	—	—
$D_{s1}(2536)^-\mu^+\nu_\mu \times$ $B(D_{s1}^- \rightarrow D^{*-}K_S^0)$	(2.4 \pm 0.7) $\times 10^{-3}$	—	—
$D_s^-\pi^+$	(3.2 \pm 0.5) $\times 10^{-3}$	—	2320
$D_s^-\pi^+\pi^+\pi^-$	(8.4 \pm 3.3) $\times 10^{-3}$	—	2301
$D_s^+K^\pm$	(3.0 \pm 0.7) $\times 10^{-4}$	—	2292
$D_s^+D_s^-$	(1.04 \pm 0.35) %	—	1823
$D_s^{*+}D_s^-$	< 12.1 %	90%	1742
$D_s^{*+}D_s^{*-}$	< 25.7 %	90%	1655

✓ There are many more resonances to be found!

- The form factors in q^2 differ for the different excited states
 - Neither branching ratios nor form factors are well known!



- Need to experimentally constrain ratio of $D^{**} : D^*$

Phys. Rev. D 32,
pp. 189–231 (1985)

➤ Spectroscopy of D_s states is very rich

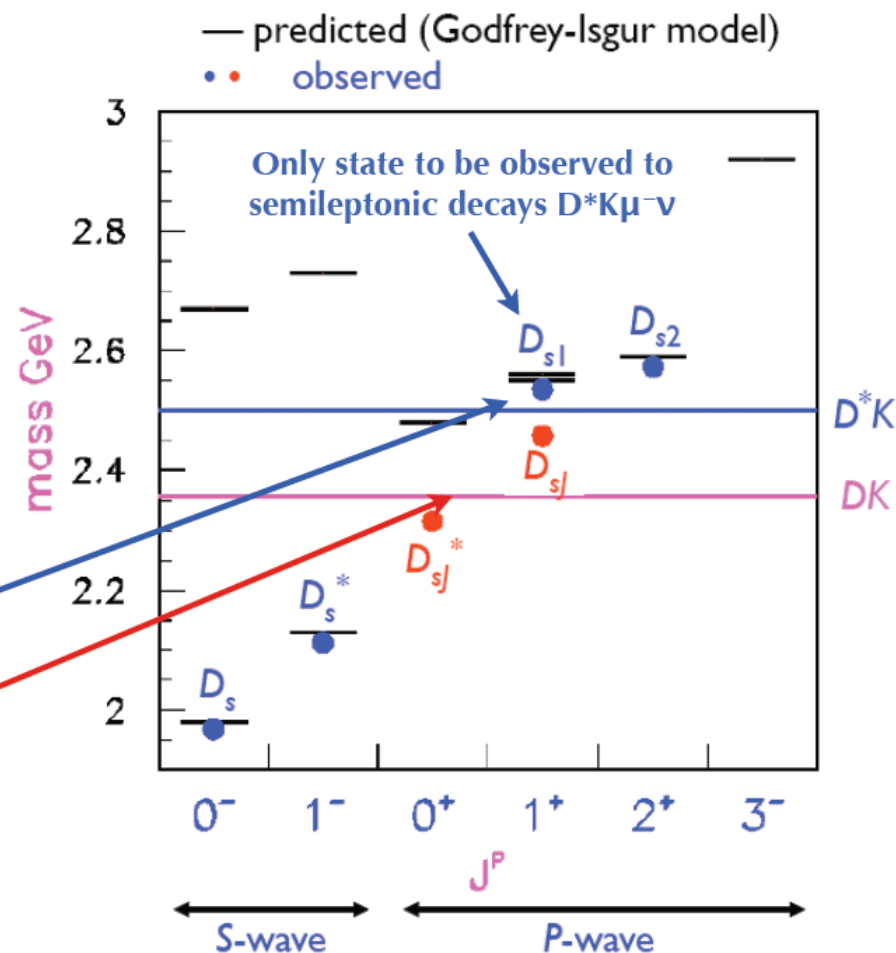
✓ D_{s1}, D_{s2} decay to $D^{(*)}K$

∴ $D^{(*)}K$ resonance search

➤ Can measure relative BR

$D_{s1}', D_{s2}^* \rightarrow D^{(*)} K$, small $D_s^{(*)} + n \pi$ contribution

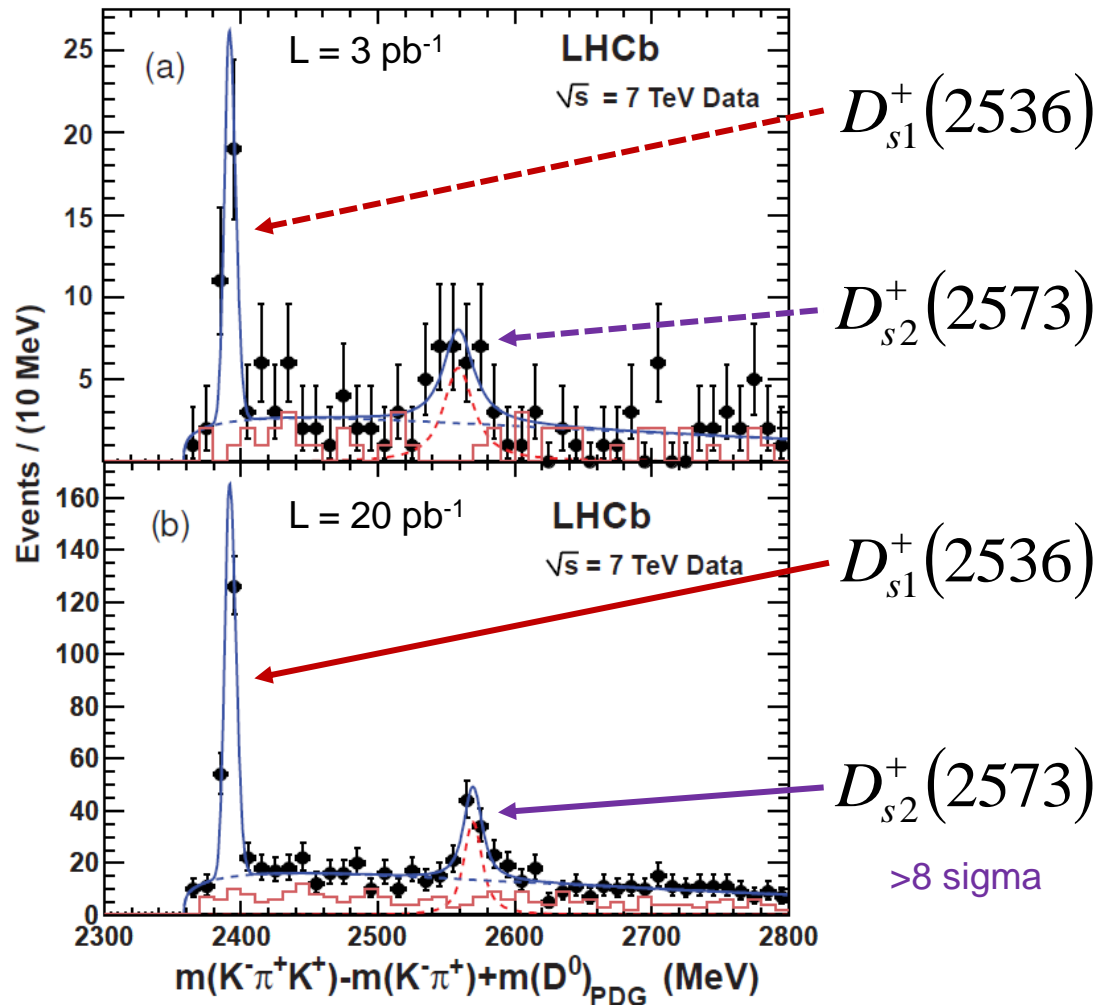
$D_{sj}^{(*)} \rightarrow D_s^{(*)} + n \pi^0/\gamma$



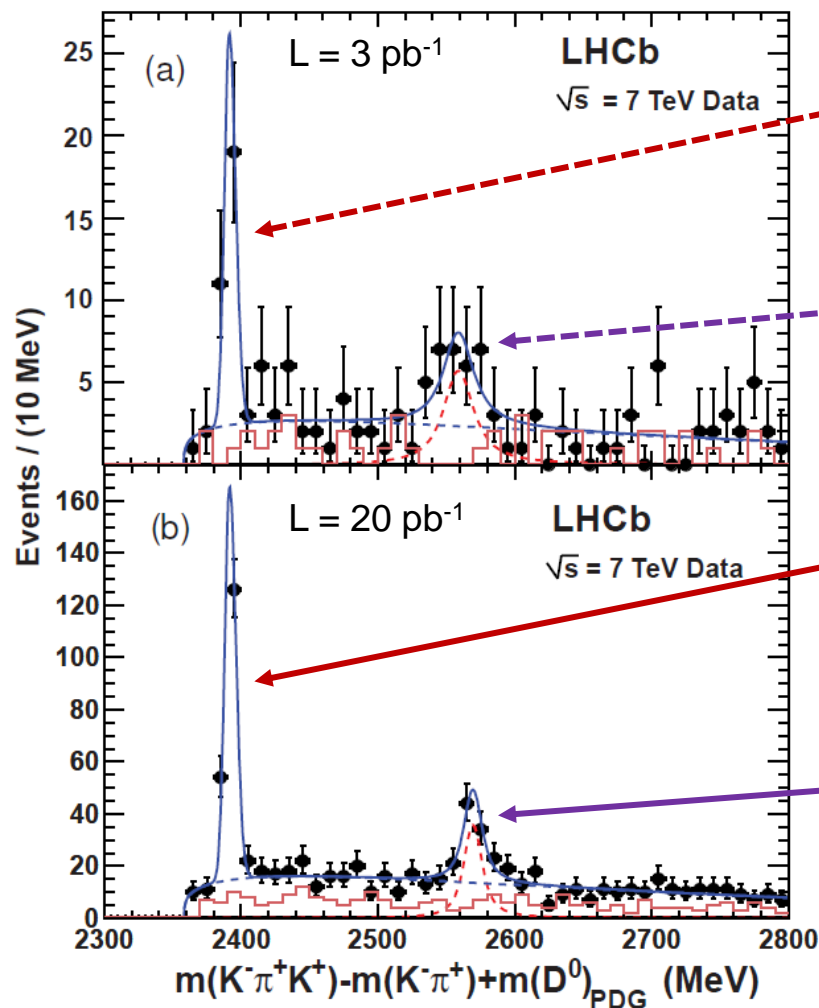
Phys. Lett. B 698 pp.14-20,2011
arXiv:1102.0348

- Quickly changing 2010 conditions
- Need to subdivide into two datasets
 - 3 pb⁻¹, small pileup, inclusive single-muon trigger
 - 20 pb⁻¹, larger pileup, any trigger (hadronic, muonic, random)
- 1. Observe D⁰K modes in both samples
- 2. Measure significance, widths and masses in 20 pb⁻¹
- 3. Determine ratio of D⁰K yields in 20 pb⁻¹
- 4. Normalize to total SL-width within the 3 pb⁻¹
 - (due to the inclusive triggering)

Phys. Lett. B 698 pp.14-20,2011
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$D_{s1}^+(2536)$

Axial vector ... need an extra π^0/γ
 Lower mass from missing π^0/γ
 MC prediction 2392 MeVc⁻²

$D_{s2}^+(2573)$

No need for extra π^0/γ
 Almost at PDG mass (~2570)

$D_{s1}^+(2536)$

Events	155	\pm	15
Mass [MeVc ⁻²]	2391.6	\pm	0.5
Width [MeVc ⁻²]	4.0	\pm	0.4

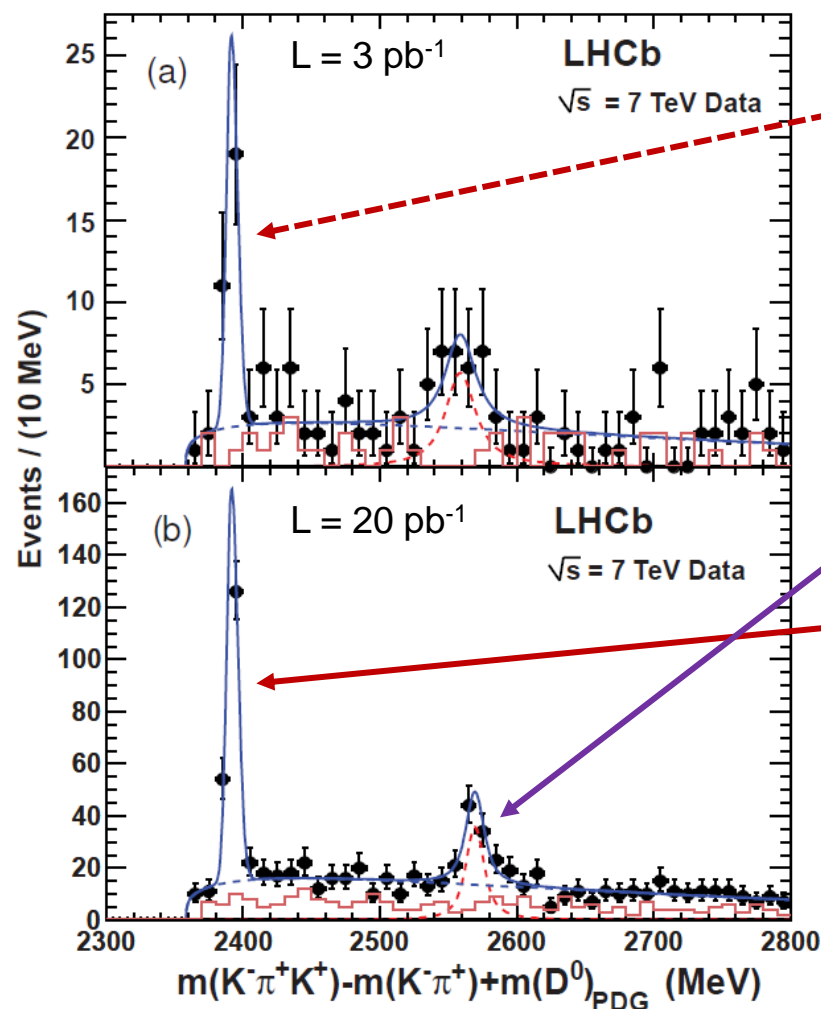
$D_{s2}^+(2573)$

>8 sigma

Events	82	\pm	17
Mass [MeVc ⁻²]	2569.4	\pm	1.6
Width [MeVc ⁻²]	12.1	\pm	4.5

(stat errors)

Phys. Lett. B 698 pp.14-20,2011
 arXiv:1102.0348



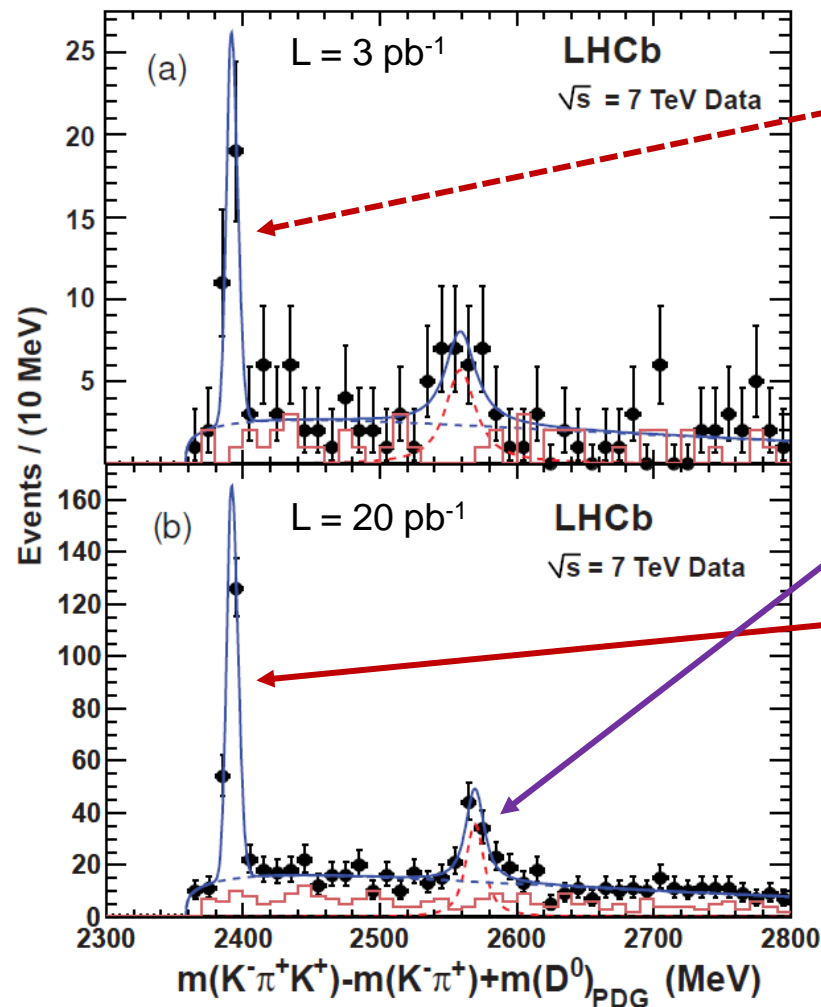
Measure BR:

$$\frac{\mathcal{B}(\overline{B}_s^0 \rightarrow D_{s1}^+ X \mu^- \overline{\nu})}{\mathcal{B}(\overline{B}_s^0 \rightarrow X \mu^- \overline{\nu})} = (5.4 \pm 1.2 \pm 0.5)\%$$

Relate both exclusive BR:

$$\frac{\mathcal{B}(\overline{B}_s^0 \rightarrow D_{s2}^{*+} X \mu^- \overline{\nu})}{\mathcal{B}(\overline{B}_s^0 \rightarrow D_{s1}^+ X \mu^- \overline{\nu})} = 0.61 \pm 0.14 \pm 0.05$$

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Measure BR:

$$\frac{\mathcal{B}(\overline{B}_s^0 \rightarrow D_{s1}^+ X \mu^- \overline{\nu})}{\mathcal{B}(\overline{B}_s^0 \rightarrow X \mu^- \overline{\nu})} = (5.4 \pm 1.2 \pm 0.5)\%$$

X

Relate both exclusive BR:

$$\frac{\mathcal{B}(\overline{B}_s^0 \rightarrow D_{s2}^{*+} X \mu^- \overline{\nu})}{\mathcal{B}(\overline{B}_s^0 \rightarrow D_{s1}^+ X \mu^- \overline{\nu})} = 0.61 \pm 0.14 \pm 0.05$$

=

Calculated exclusive BR:

$$\frac{\mathcal{B}(\overline{B}_s^0 \rightarrow D_{s2}^{*+} X \mu^- \overline{\nu})}{\mathcal{B}(\overline{B}_s^0 \rightarrow X \mu^- \overline{\nu})} = (3.3 \pm 1.0 \pm 0.4)\%$$

➤ From:

$$\frac{\mathcal{B}(\overline{B}_s^0 \rightarrow D_{s1}^+ X \mu^- \overline{\nu})}{\mathcal{B}(\overline{B}_s^0 \rightarrow X \mu^- \overline{\nu})} = (5.4 \pm 1.2 \pm 0.5)\%$$

$$\frac{\mathcal{B}(\overline{B}_s^0 \rightarrow D_{s2}^{*+} X \mu^- \overline{\nu})}{\mathcal{B}(\overline{B}_s^0 \rightarrow D_{s1}^+ X \mu^- \overline{\nu})} = 0.61 \pm 0.14 \pm 0.05$$

➤ We constrain the D** branching ratios

✓ Important for the attempt on V_{cb}

➤ We now use these measurements in our MC

✓ Important for MC-corrections of missing neutrino

∴ Important for many other SL measurements

... such as ...

Flavour-specific asymmetry

Fermilab-Pub-10/114-E

Evidence for an anomalous like-sign dimuon charge asymmetry

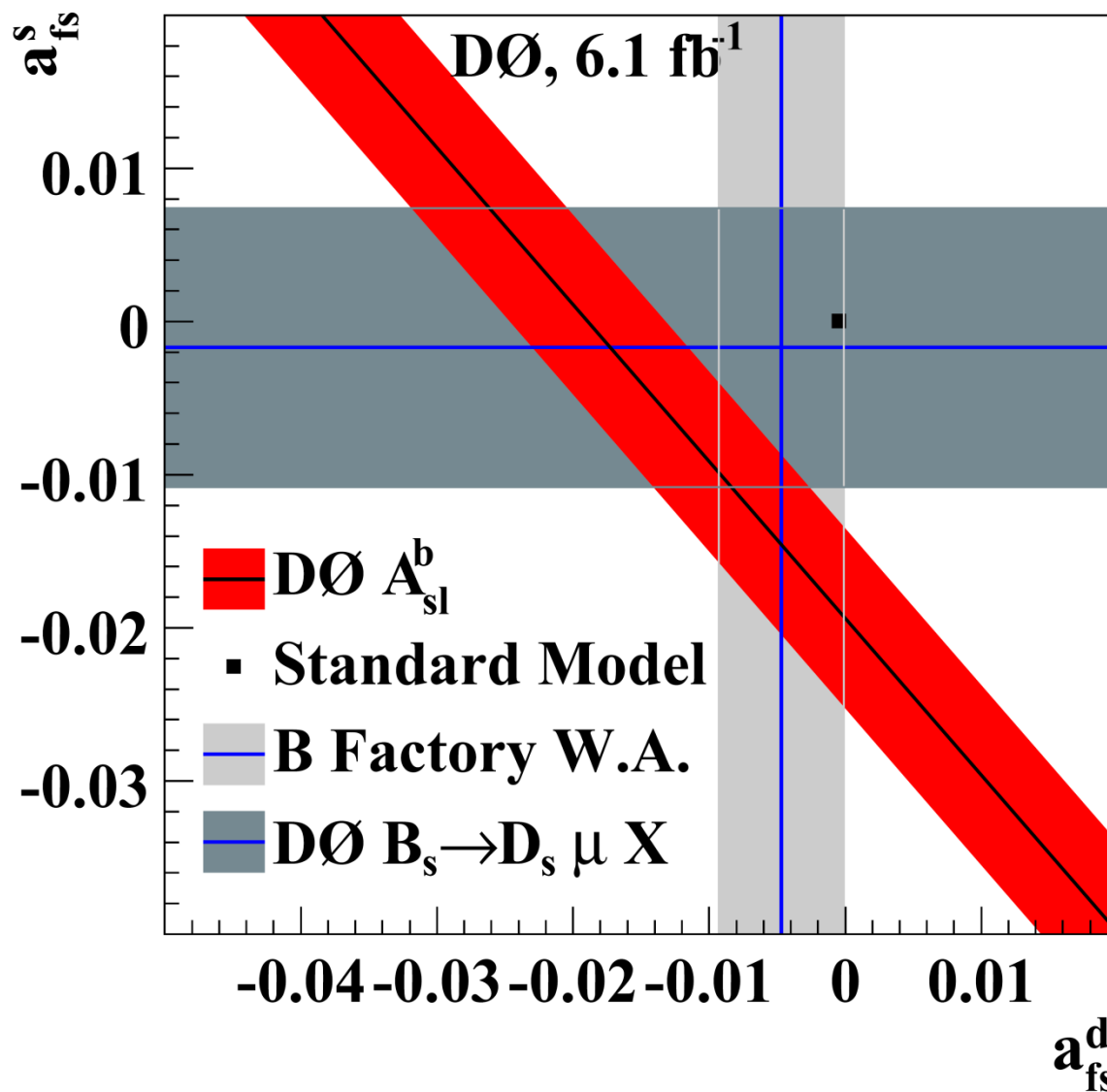
V.M. Abazov,³⁶ B. Abbott,⁷⁴ M. Abolins,⁶³ B.S. Acharya,²⁹ M. Adams,⁴⁹ T. Adams,⁴⁷ E. Aguilo,⁶ G.D. Alexeev,³⁶
 G. Alkhazov,⁴⁰ A. Alton^a,⁶² G. Alverson,⁶¹ G.A. Alves,² L.S. Ancu,³⁵ M. Aoki,⁴⁸ Y. Arnoud,¹⁴ M. Arov,⁵⁸
 A. Askew,⁴⁷ B. Åsman,⁴¹ O. Atramentov,⁶⁶ C. Avila,⁸ J. BackusMayes,⁸¹ F. Badaud,¹³ L. Bagby,⁴⁸ B. Baldin,⁴⁸
 D.V. Bandurin,⁴⁷ S. Banerjee,²⁹ E. Barberis,⁶¹ A.-F. Barfuss,¹⁵ P. Baringer,⁵⁶ J. Barreto,² J.F. Bartlett,⁴⁸
 U. Bassler,¹⁸ S. Beale,⁶ A. Bean,⁵⁶ M. Begalli,³ M. Begel,⁷² C. Belanger-Champagne,⁴¹ L. Bellantoni,⁴⁸
 J.A. Benitez,⁶³ S.B. Beri,²⁷ G. Bernardi,¹⁷ R. Bernhard,²² I. Bertram,⁴² M. Besançon,¹⁸ R. Beuselinck,⁴³

...

We measure the charge asymmetry A of like-sign dimuon events in 6.1 fb^{-1} of $p\bar{p}$ collisions recorded with the D0 detector at a center-of-mass energy $\sqrt{s} = 1.96 \text{ TeV}$ at the Fermilab Tevatron collider. From A , we extract the like-sign dimuon charge asymmetry in semileptonic b -hadron decays: $A_{\text{sl}}^b = -0.00957 \pm 0.00251 \text{ (stat)} \pm 0.00146 \text{ (syst)}$. This result differs by 3.2 standard deviations from the standard model prediction $A_{\text{sl}}^b(SM) = (-2.3_{-0.6}^{+0.5}) \times 10^{-4}$ and provides first evidence of anomalous CP-violation in the mixing of neutral B mesons.

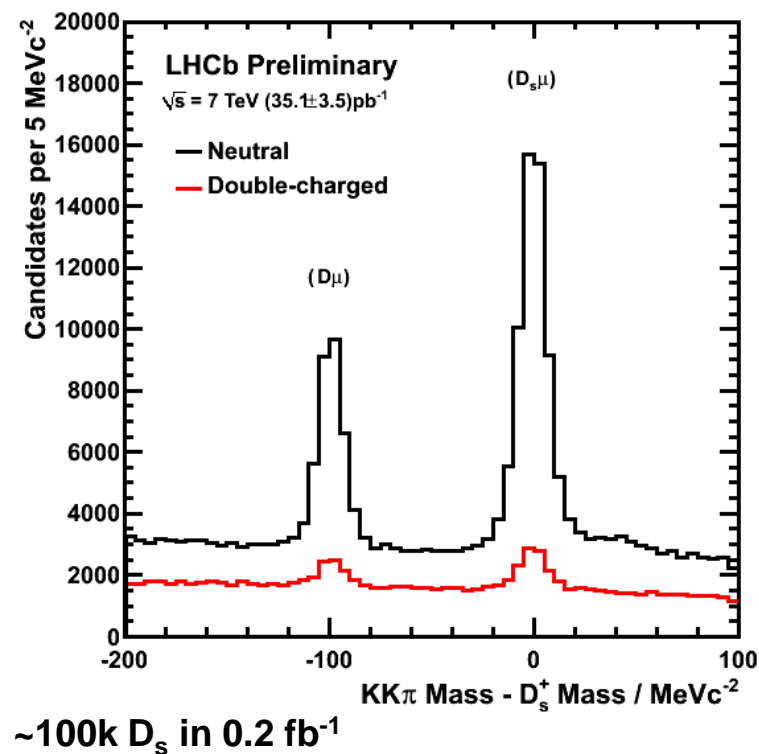
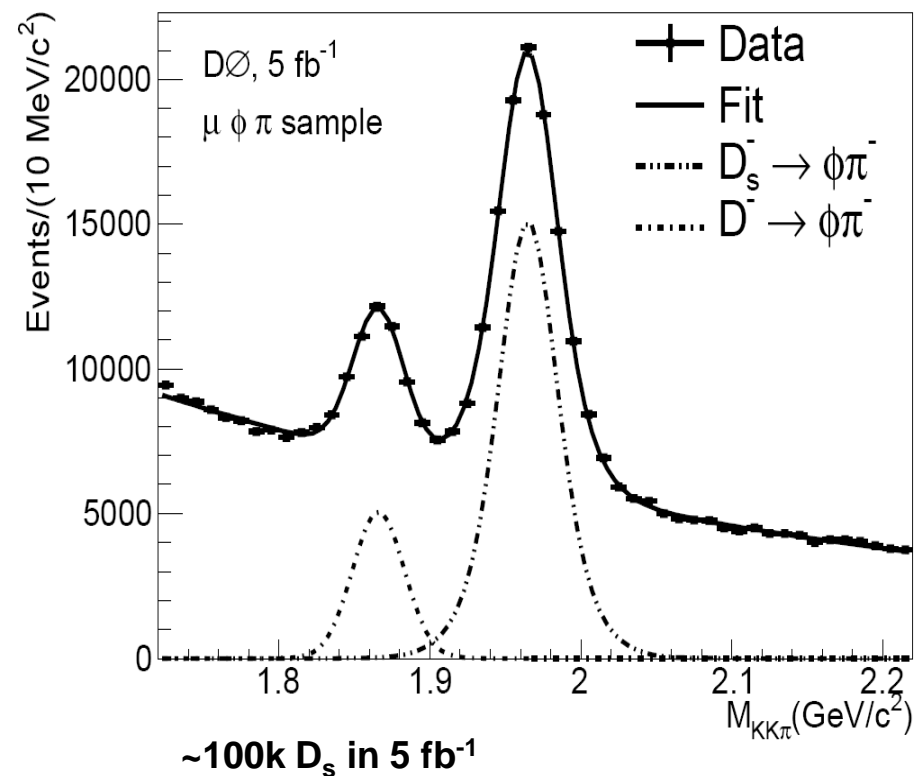
PACS numbers: 13.25.Hw; 14.40.Nd

See talks from
Monday morning



Situation could
really be cleared
up by LHCb

- LHCb is reconstructing both $B_s^0 \rightarrow D_s^\mp \mu^\pm \nu_\mu$ and $B_d^0 \rightarrow D_d^\mp \mu^\pm \nu_\mu$
- LHCb is catching up with DØ very quickly



- LHC is a pp-collider, not a $p\bar{p}$ -collider
- LHCb is in the forward region
 - Can't measure the same thing as DØ
 - Need a clever new method

NB: DØ
(inclusive)

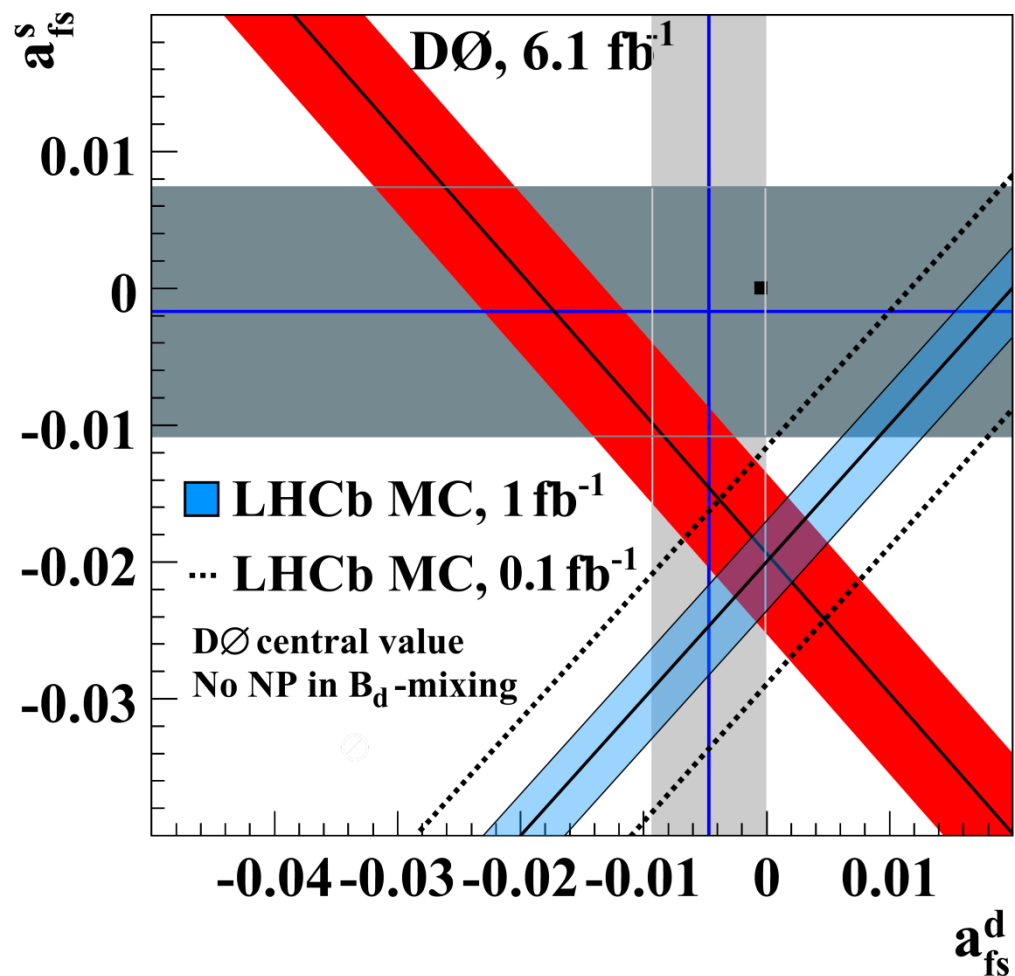
$$\left(A^b \sim \frac{a_{fs}^s + a_{fs}^d}{2} \rightarrow -(2.0 \pm 0.3) \times 10^{-4} [S.M.] \right)$$

- Subtract two asymmetries to cancel main systematics

LHCb
(subtraction)

$$\Delta A_{fs} = \frac{a_{fs}^s - a_{fs}^d}{2} \rightarrow (2.1 \pm 0.3) \times 10^{-4} [S.M.]$$

- LHCb measurement cuts at right-angles to $D\bar{D}$



NB: This is MC, scaled to real data including an estimate of systematics...

- 2010 was a fantastic year for LHCb
- Semileptonics provide an excellent probe of the SM
 - Broad physics program
 - Already competitive with only $\sim 37 \text{ pb}^{-1}$
- Already with 2010 data:
 - $b\bar{b}$ cross-section, re-discovery of B_d mixing
 - Measurement of $B_s^0 \rightarrow D_{s1,2}^{*\pm} \mu^\mp \nu X^0$
- Many other results in progress!
- In 2011 we expect 1 fb^{-1} :
 - Form factors and $V_{(u/c)b}$
 - B -mixing including ΔA_{fs}

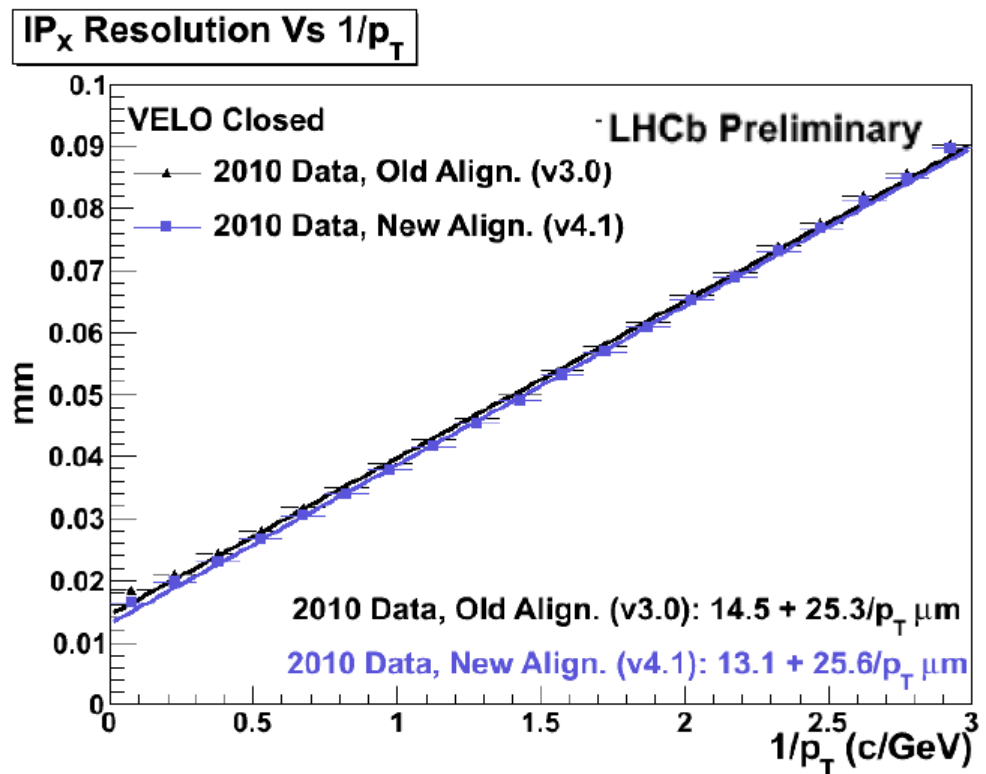
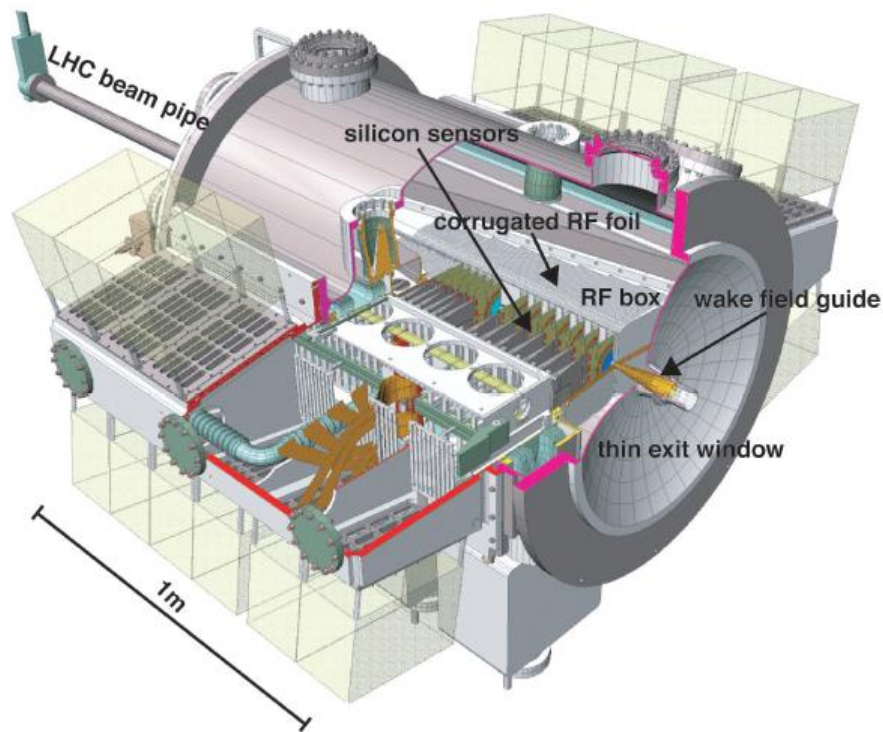
- Backups are often required



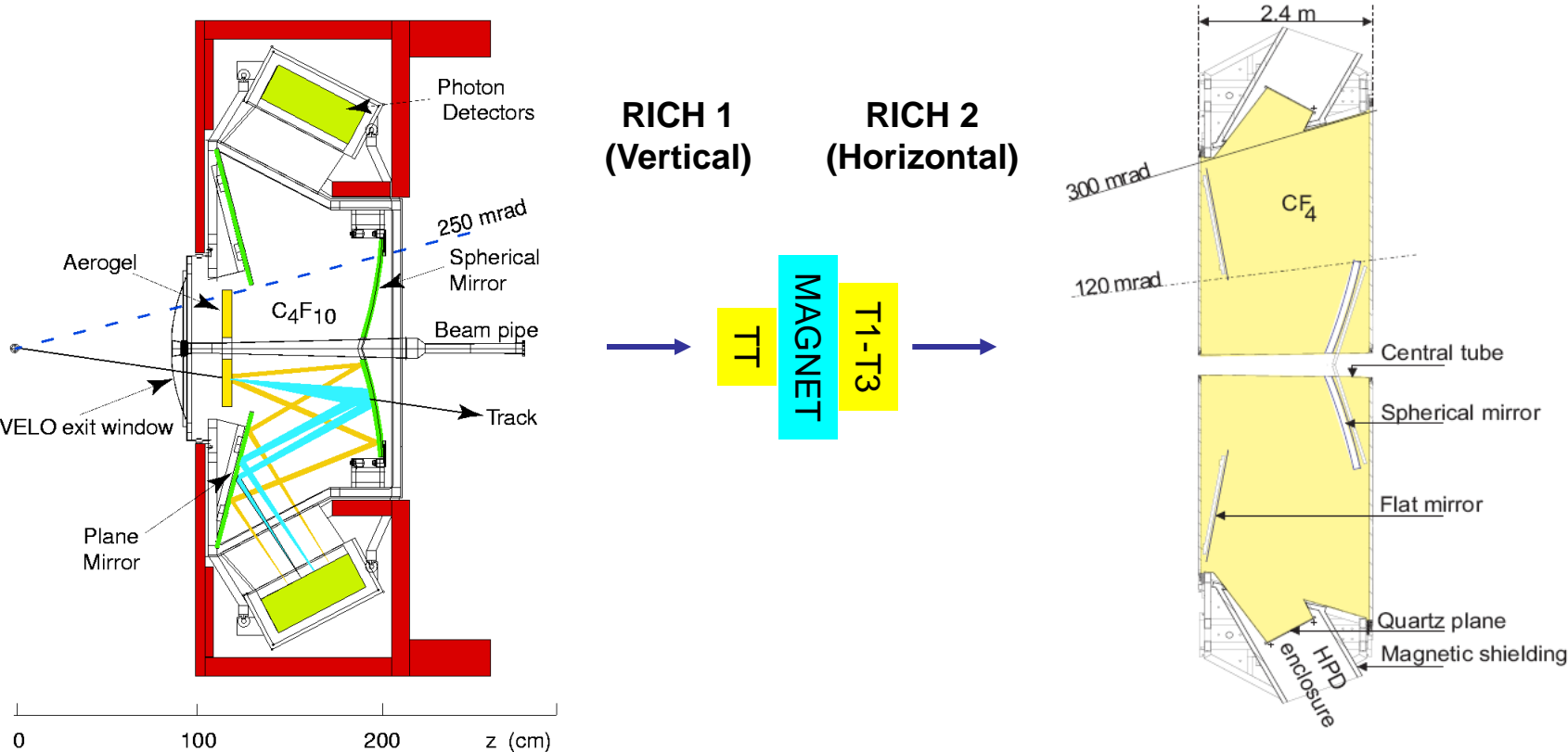
- Thanks to the LHC and the rest of the LHCb collaboration
- Specifically extra thanks to my colleagues:
 - Marina Artuso, Phillip Urquijo, Kim Vervink, Liming Zhang for direct contributions to this talk

- Theory predicts SL widths:
 - $D^{(*)*} : D \sim 3 : 1$
 - arXiv:1003.5576

- Proper Time: LHCb Velo precise down to ~50 fs



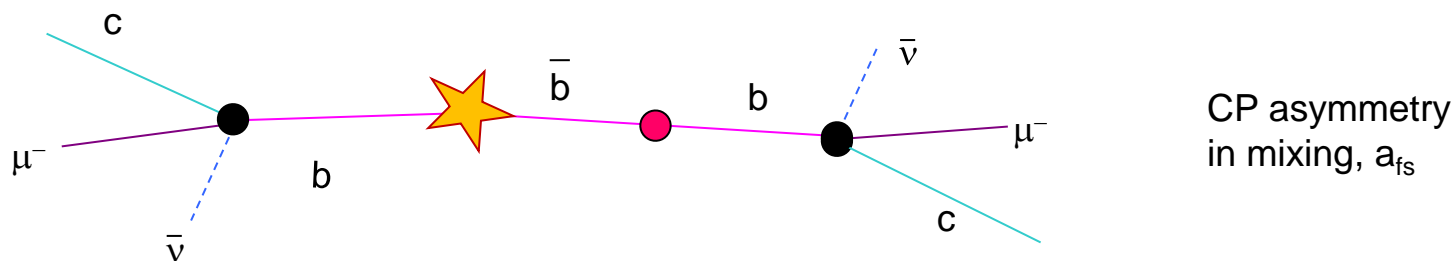
- Particle ID: separation handled by dedicated subdetectors
- Two RICHes, Calorimetry and Muon system





EVT: 49700980
RUN: 70684

- Very difficult measurement
- Observe $N(\mu^+\mu^+) \neq N(\mu^-\mu^-)$
- Flavour-specific asymmetry from B^0 -mixing in the SM:

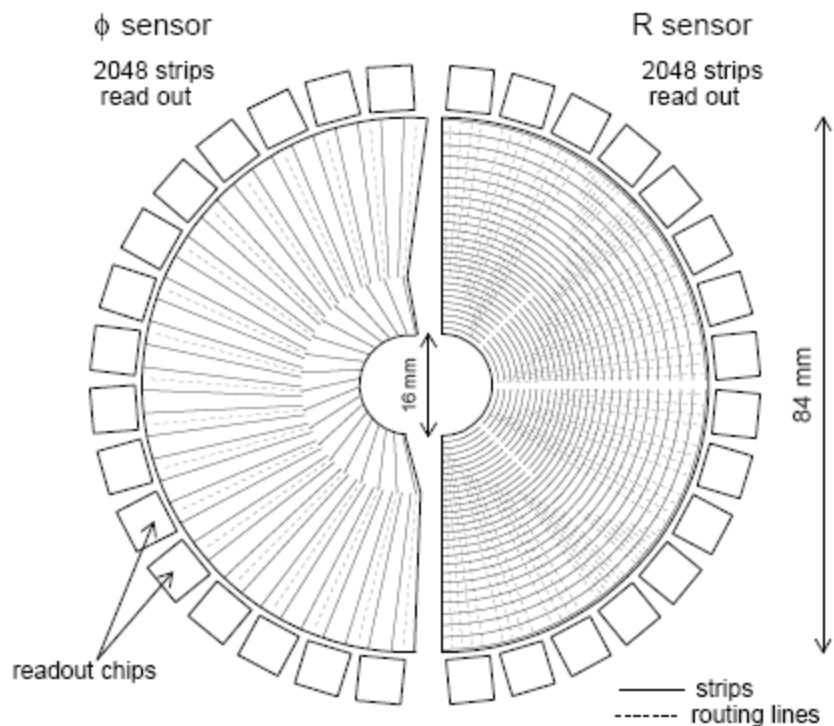
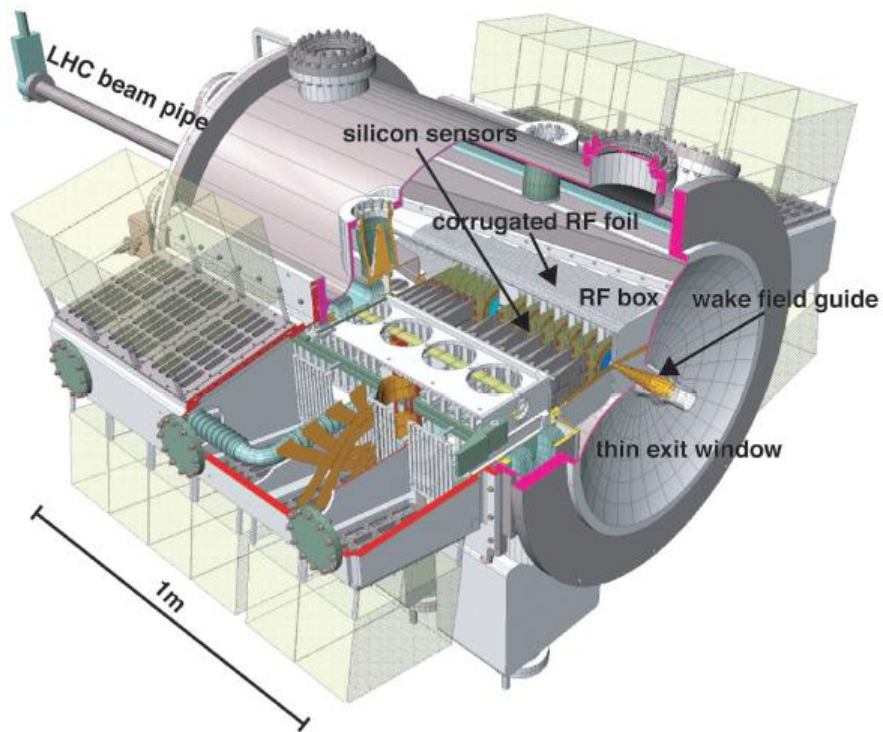


- In the standard model a_{fs} is almost negligible

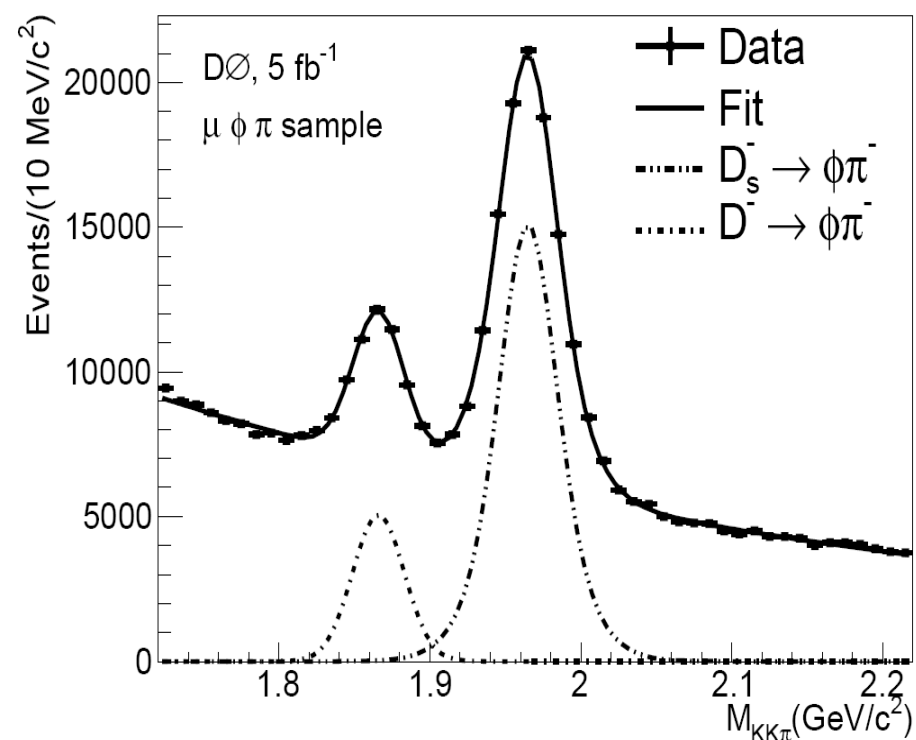
$$A^b \approx \frac{a_{fs}^s + a_{fs}^d}{2} \qquad SM = (-2.0 \pm 0.3) \times 10^{-4} \qquad D\emptyset \approx (-1 \pm 0.3)\%$$



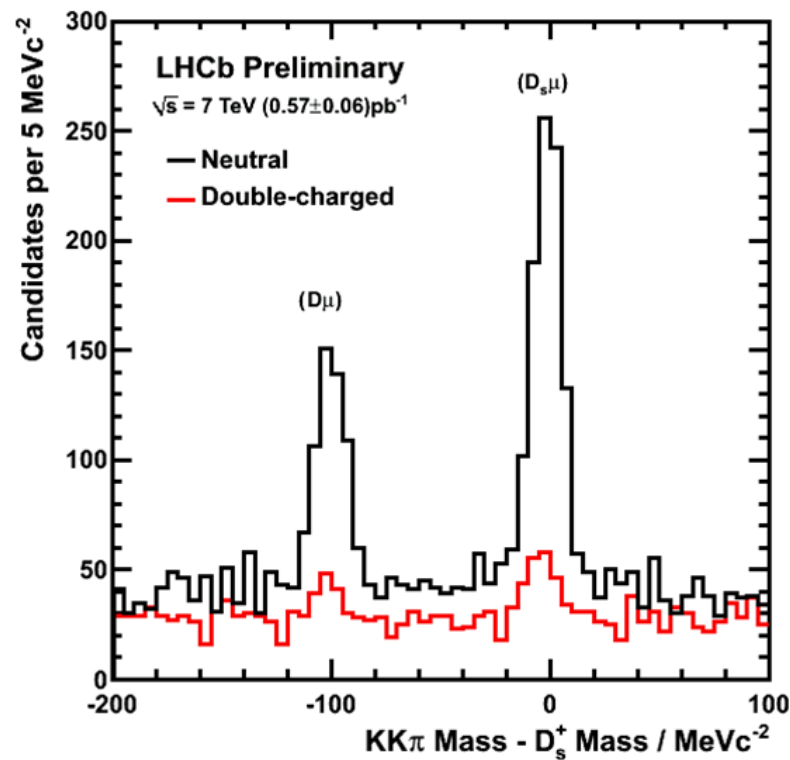
- Proper Time: LHCb Velo precise down to ~50 fs



- Our forte: exclusive, reconstructed, b -decays
- In particular, time-dependent measurements




~100k D_s in 5 fb^{-1}



Estimate 100k D_s in 0.1 fb^{-1}

$$A_{fs}^q(t) = \frac{\Gamma(f) - \Gamma(\bar{f})}{\Gamma(f) + \Gamma(\bar{f})}$$

$$A_{fs}^q(t) = \frac{a_{fs}^q}{2} \left[- \left(\frac{a_{fs}^q}{2} \right) \frac{\cos(\Delta m_q t)}{\cosh(\Delta \Gamma_q t / 2)} \right]$$



 $10^{-3} \rightarrow 10^{-5}$

$$A_{fs}^q(t) = \frac{\Gamma(f) - \Gamma(\bar{f})}{\Gamma(f) + \Gamma(\bar{f})}$$

$$A_{fs}^q(t) = \frac{a_{fs}^q}{2} \left[-\frac{\delta_c^q}{2} - \left(\frac{a_{fs}^q}{2} + \frac{\delta_p^q}{2} \right) \frac{\cos(\Delta m_q t)}{\cosh(\Delta \Gamma_q t / 2)} + \frac{\delta_b^q}{2} \left(\frac{B}{S} \right)^q \right]$$

$10^{-3} \rightarrow 10^{-5}$ 10^{-2} 10^{-2} 10^{-3}

➤ Polluting asymmetries are much larger than a_{fs}

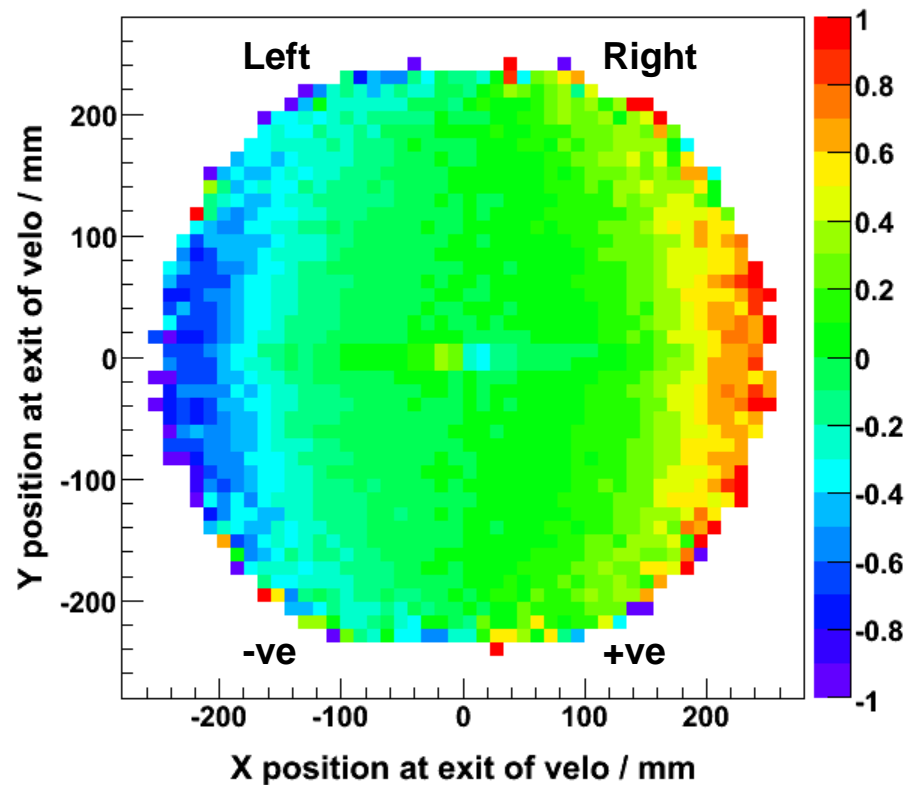
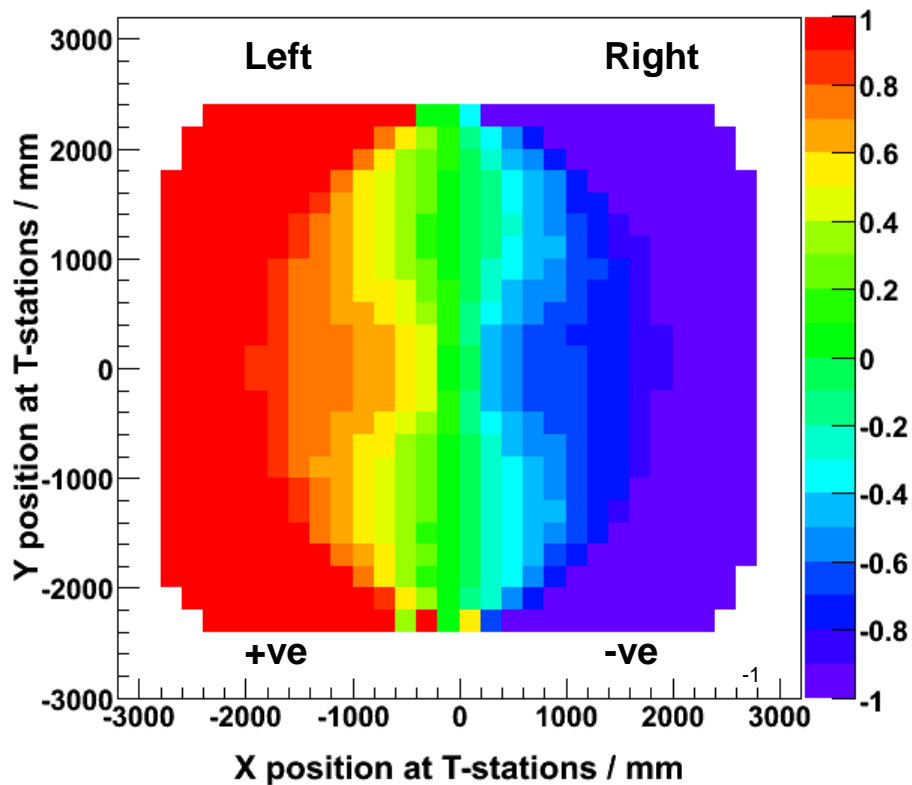
- Detector asymmetry δ_c $\sim(10^{-2})$
- Production asymmetry δ_p $\sim(10^{-2})$
- Background asymmetry δ_b $\sim(10^{-3})$

$$\delta_c = \frac{\varepsilon(\bar{f}_i)}{\varepsilon(f_i)} - 1$$

$$\delta_p = \frac{N(\bar{I}_0)}{N(I_0)} - 1$$

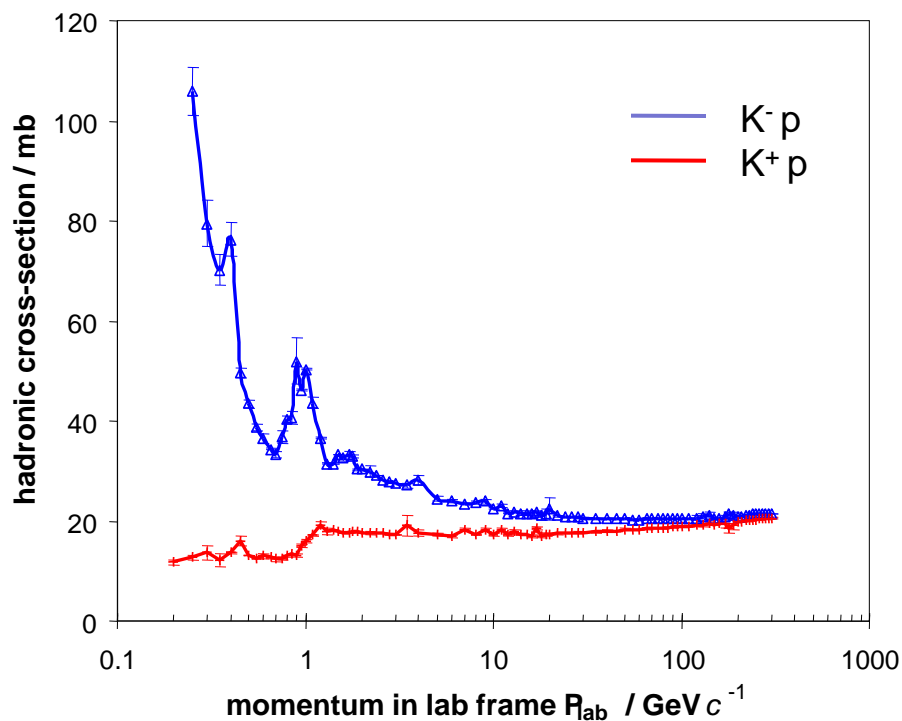
$$\delta_b = \frac{\bar{B}/\bar{S}}{B/S} - 1$$

Asymmetry from Long Muon Tracks Reconstructed in MC

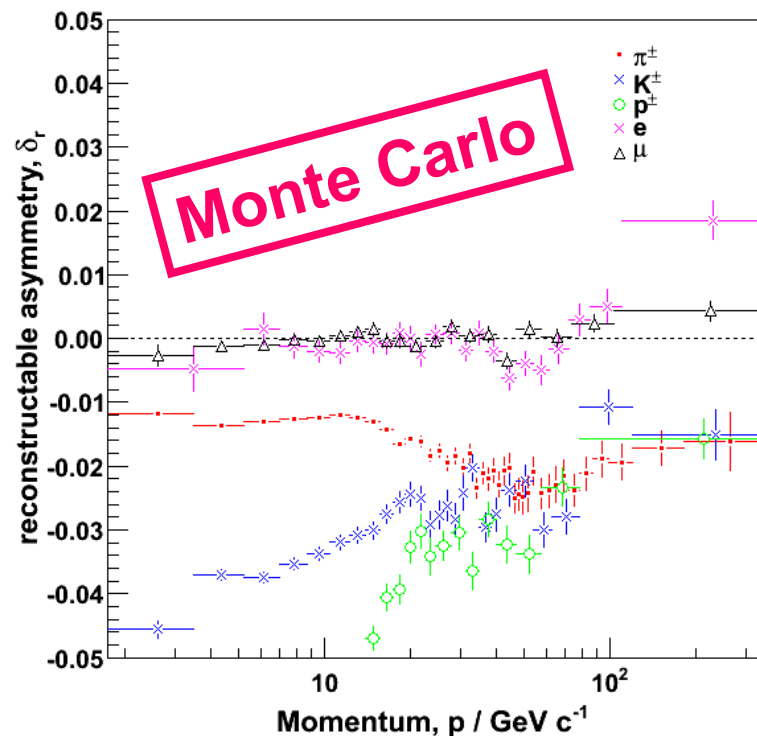


- Magnet divides +/- charge, allowing +/- asymmetry
- by reversing magnet in D0: δ_c reduced from 3% \rightarrow $\sim 0.1\%$

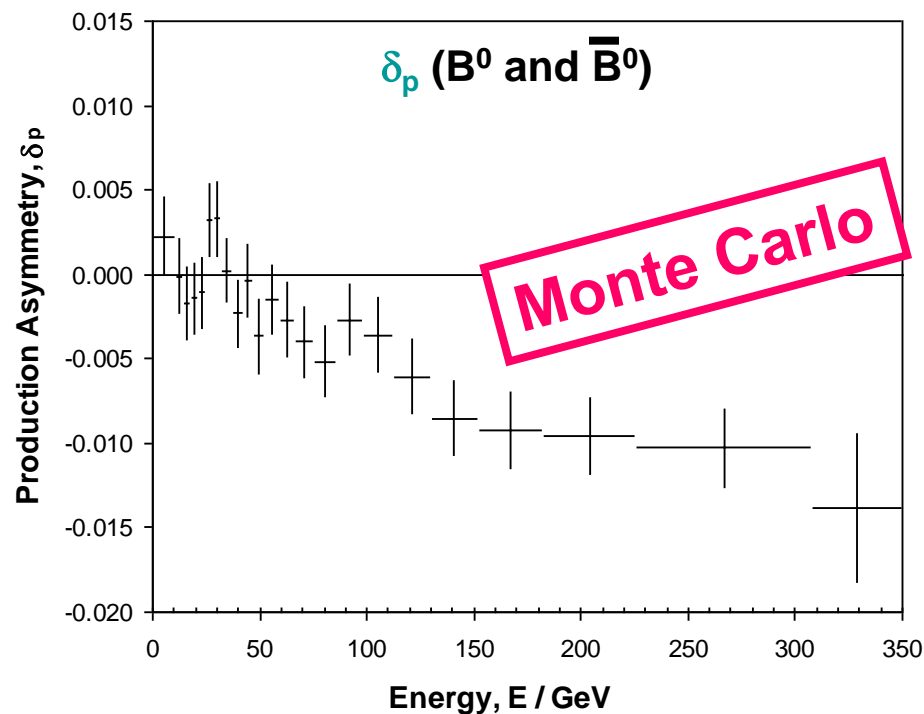
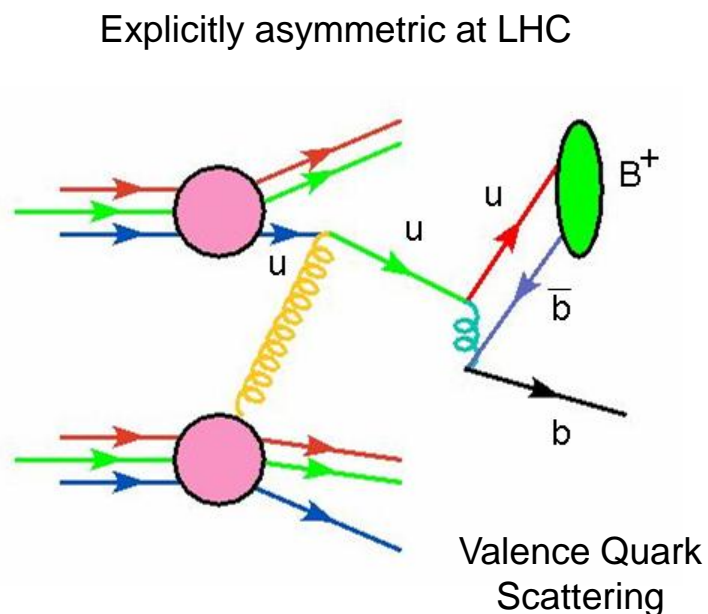
Kaon interaction cross-section



Resultant charge asymmetry (MC)



- Matter detector → hadronic interactions are asymmetric
- Dominant systematic at order 1%



- LHC is a proton-proton collider: not CP-symmetric
- LHCb is at high rapidity where production asymmm. are largest
- There is never a simple control channel to measure δ_p

- Take B_s/B_d with the same final states ($f=KK\pi\mu$)

$$\Gamma(f) = N e^{-\Gamma t} \left[(1 + x_1) \cosh\left(\frac{\Delta\Gamma t}{2}\right) + (x_2 + x_3) \cos(\Delta m t) \right]$$

$$\Gamma(\bar{f}) = N e^{-\Gamma t} \left[(1 - x_1) \cosh\left(\frac{\Delta\Gamma t}{2}\right) + (x_2 - x_3) \cos(\Delta m t) \right]$$

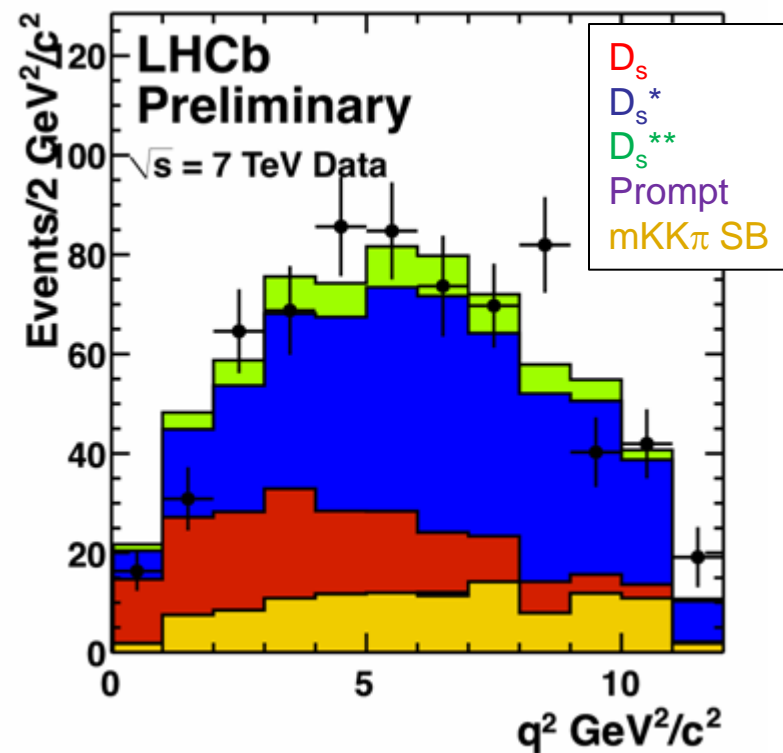
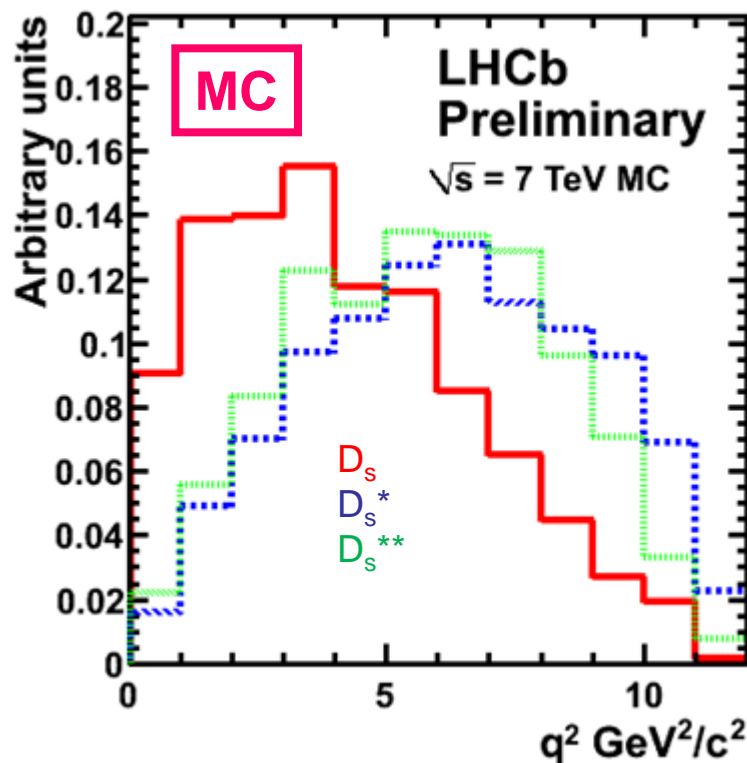
where: $x_1 = A_c + a_{fs}$ $x_2 = 2A_c A_p$ $x_3 = 2A_p - a_{fs}$

- All **production asymmetry** is in x_2/x_3 , just throw it away
- Measure the **difference** between B_s and B_d

$$\Delta A_{fs}^{s,d} = \frac{x_1^s - x_1^d}{2} = \frac{a_{fs}^s - a_{fs}^d}{2}$$

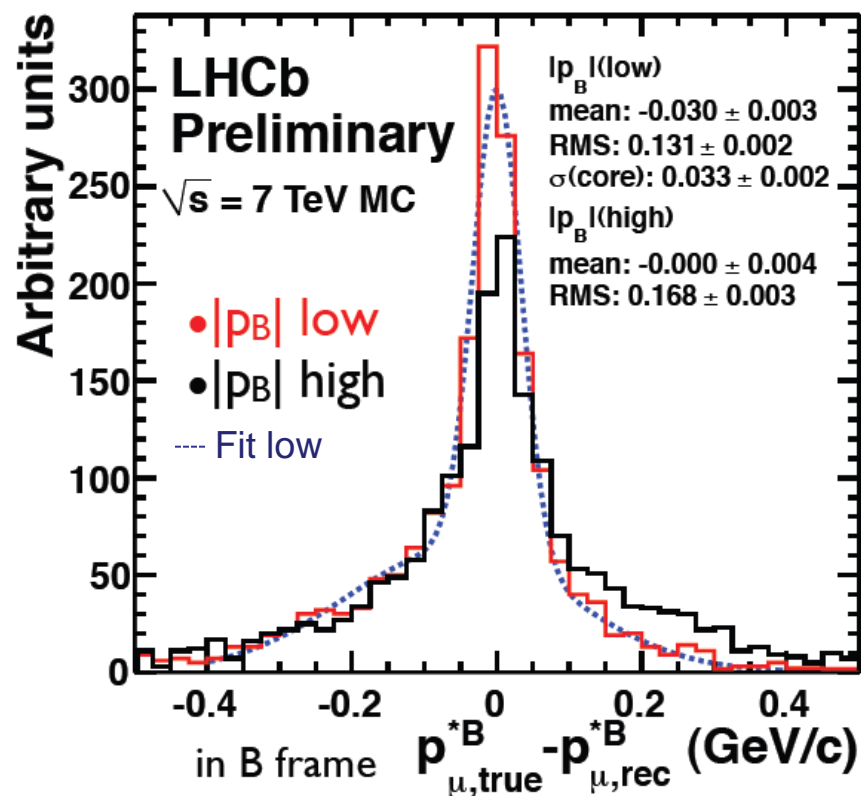
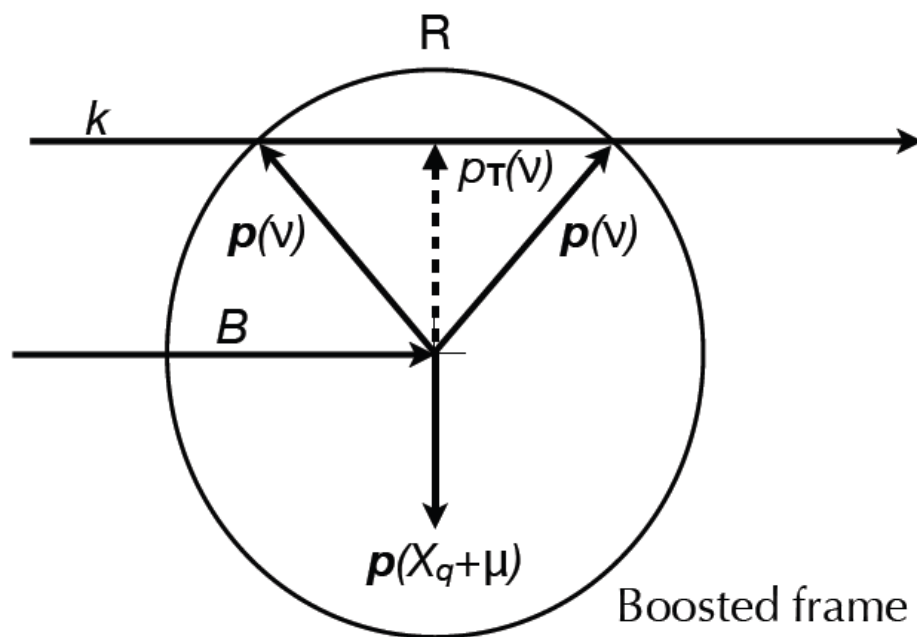
$$SM = \left(+2.5^{+0.5}_{-0.6} \right) \times 10^{-4}$$

- Choose a set of form-factors for the LHCb MC (HQET2/ISGW2)
- Set the relative branching ratios from $D/D^*/D^{**}$, $\sim 2.5:\sim 6:\sim 1$



- 0.8 pb^{-1} examined so far, should become powerful in 50 pb^{-1}

- Use kinematics to calculate the neutrino momentum
 - Two ambiguous solutions, the lower momentum has lower error



Flavour-specific asymmetry

... a smoking gun for new physics??

1. $p\bar{p}$ -interactions within a symmetric experiment
2. Correct all experimental biases (magnets, mis-id ...)
3. Observe $N(\mu^+\mu^+) \neq N(\mu^-\mu^-)$
4. In the SM, the favoured way to make charge asymmetry is if:

$$b\bar{b} \longrightarrow \mu^+\mu^+ \neq b\bar{b} \longrightarrow \mu^-\mu^-$$

5. Which comes from B^0 -mixing:

$$b\bar{b} \Rightarrow \bar{B}^0 B^0 \sim \bar{B}^0 \bar{B}^0 \rightarrow \mu^+\mu^+ X \neq b\bar{b} \Rightarrow \bar{B}^0 B^0 \sim B^0 B^0 \rightarrow \mu^-\mu^- X$$

➤ In the standard model it is almost negligible

$$A^b \approx \frac{a_{fs}^s + a_{fs}^d}{2} \quad SM = (-2.0 \pm 0.3) \times 10^{-4} \quad D\emptyset \approx (-1 \pm 0.3)\%$$

- a_{fs} is very sensitive to new physics (NP) even if:
 - ❖ Tree-level processes are SM-dominated
 - ❖ SM flavour structure
 - ❖ Unitary CKM
- With very weird scenarios (like *leptoquarks*)
 - Probe NP mixing, interference and/or decays
- Usual formula is modified:

$$a^{SM} \approx \text{Im} \left\{ \frac{\Gamma_{12}^{SM}}{M_{12}^{SM}} \right\}$$

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 - Probe NP mixing, interference and/or decays
- If we allow a single NP phase in the mixing Θ

$$a^{NP} \approx \text{Im} \left\{ \frac{\Gamma_{12}^{SM}}{M_{12}^{SM}} \right\} \cos \Theta - \text{Re} \left\{ \frac{\Gamma_{12}^{SM}}{M_{12}^{SM}} \right\} \sin \Theta$$

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 - (first part is just the SM value)

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- With very weird scenarios (like *leptoquarks*)
 - Probe NP mixing, interference and/or decays
- If we allow a single NP phase in the mixing Θ
 - (first part is just the SM value)

$$a^{NP} \approx 2.1 \times 10^{-5} \cos \Theta + 4.0 \times 10^{-3} \sin \Theta$$

- Up to **200-times** the SM!!! [[[... still... < DØ measurement]]]

Flavour-specific asymmetry

At LHCb

- At the LHC we have extra complications in the measurement
- Polluting asymmetries, which are all much larger than a_{fs}
 - Production asymmetry $\delta_p \sim (10^{-2})$
 - Detector asymmetry $\delta_c \sim (10^{-2})$
 - Background asymmetry $\delta_b \sim (10^{-3})$
- Use a, time-dependent, untagged, simultaneous fit to $B_s + B_d$
- Subtract two asymmetries to eliminate detector component

$$\Delta A_{fs} = \frac{a_{fs}^s - a_{fs}^d}{2} = (2.1 \pm 0.3) \times 10^{-4}$$

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NB: DØ

$$\left(A^b \sim \frac{a_{fs}^s + a_{fs}^d}{2} \sim -(2.0 \pm 0.3) \times 10^{-4} \right)$$

$$A_{fs}^q(t) = \frac{\Gamma(f) - \Gamma(\bar{f})}{\Gamma(f) + \Gamma(\bar{f})}$$

$$A_{fs}^q(t) = \frac{a_{fs}^q}{2}$$



$10^{-3} \rightarrow 10^{-5}$

$$-\left(\frac{a_{fs}^q}{2}\right) \frac{\cos(\Delta m_q t)}{\cosh(\Delta \Gamma_q t / 2)}$$

$$A_{fs}^q(t) = \frac{\Gamma(f) - \Gamma(\bar{f})}{\Gamma(f) + \Gamma(\bar{f})}$$

$$A_{fs}^q(t) = \frac{a_{fs}^q}{2} \left[-\frac{\delta_c^q}{2} - \left(\frac{a_{fs}^q}{2} + \frac{\delta_p^q}{2} \right) \frac{\cos(\Delta m_q t)}{\cosh(\Delta \Gamma_q t / 2)} + \frac{\delta_b^q}{2} \left(\frac{B}{S} \right)^q \right]$$

$10^{-3} \rightarrow 10^{-5}$ 10^{-2} 10^{-2} 10^{-3}

➤ Polluting asymmetries are much larger than a_{fs}

- Detector asymmetry δ_c $\sim(10^{-2})$
- Production asymmetry δ_p $\sim(10^{-2})$
- Background asymmetry δ_b $\sim(10^{-3})$

$$\delta_c = \frac{\varepsilon(\bar{f}_i)}{\varepsilon(f_i)} - 1$$

$$\delta_p = \frac{N(\bar{I}_0)}{N(I_0)} - 1$$

$$\delta_b = \frac{\bar{B}/\bar{S}}{B/S} - 1$$

- We measure time-dependent decay rates:

$$\begin{aligned}\Gamma(f) &= N e^{-\Gamma t} (1 + A_c) \left[(1 + A_{fs}) \cosh\left(\frac{\Delta\Gamma t}{2}\right) + (2A_p - A_{fs}) \cos(\Delta m t) \right] \\ &\rightarrow N e^{-\Gamma t} \left[(1 + A_c + A_{fs}) \cosh\left(\frac{\Delta\Gamma t}{2}\right) + (2A_p - A_{fs} + 2A_p A_c) \cos(\Delta m t) \right] \\ \Gamma(\bar{f}) &\rightarrow N e^{-\Gamma t} \left[(1 - A_c - A_{fs}) \cosh\left(\frac{\Delta\Gamma t}{2}\right) + (A_{fs} - 2A_p + 2A_p A_c) \cos(\Delta m t) \right]\end{aligned}$$

- A_c , A_p and A_{fs} are correlated and cannot be separately fitted
- First, reparameterise

- Just to make it easier to see what we're doing...

$$\Gamma(f) = N e^{-\Gamma t} \left[(1 + x_1) \cosh\left(\frac{\Delta\Gamma t}{2}\right) + (x_2 + x_3) \cos(\Delta m t) \right]$$

$$\Gamma(\bar{f}) = N e^{-\Gamma t} \left[(1 - x_1) \cosh\left(\frac{\Delta\Gamma t}{2}\right) + (x_2 - x_3) \cos(\Delta m t) \right]$$

where: $x_1 = A_c + a_{fs}$ $x_2 = 2A_c A_p$ $x_3 = 2A_p - a_{fs}$

- **production asymmetry** is an *initial state asymmetry*
- Changes the mixing amplitude, does not change the physics
- Fit for x_1 independently, which now only has **detector asym**

- Take B_s/B_d with the same final states ($f=KK\pi\mu$)

$$\Gamma(f) = N e^{-\Gamma t} \left[(1 + x_1) \cosh\left(\frac{\Delta\Gamma t}{2}\right) + (x_2 + x_3) \cos(\Delta m t) \right]$$

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$$SM = \left(+2.5^{+0.5}_{-0.6} \right) \times 10^{-4}$$

- MC sensitivities, Real data yields and systematics
 - 0.1 fb^{-1} $\sigma \sim 5 \times 10^{-3}$... First result (2011)
 - 1.0 fb^{-1} $\sigma \sim 2 \times 10^{-3}$... 5σ observation? (2012/2013)