

# A prototype model for **quasi-Dirac** neutrinos

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[CSF, Gregoire, Tonerio, 2007.09158] [Earl, CSF, Gregoire, Tonerio, 1903.12192]



# We need New Physics

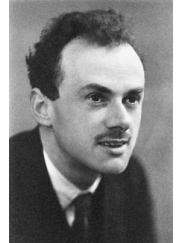
1) Neutrino mass  $m_\nu^{\text{obs}} \sim 0.05 \text{ eV}$  (**Dirac** or **Majorana**)

2) Cosmic matter-antimatter asymmetry  $Y_B^{\text{obs}} \sim 10^{-10}$

3) Dark matter  $Y_{\text{DM}}^{\text{obs}} \sim 5 \times 10^{-10} \left( \frac{1 \text{ GeV}}{m_{\text{DM}}} \right)$

$$m_M \bar{\nu}_L \nu_L^c$$

$$m_D \bar{\nu}_L \nu_R$$



Or somewhere in between

**Theoretical puzzles:** 3-generation fermions, higgs mass stability, hypercharge quantization, no strong CP violation

# Neutrino mass

At renormalized level, strictly massless

If we look further ...  $\frac{\lambda}{\Lambda} \bar{\ell}_L \tilde{H} \ell_L^c \tilde{H} \implies m_M \bar{\nu}_L \nu_L^c$

$$m_M = \frac{\lambda v^2}{\Lambda}$$



**Majorana**

- Explain lightness
- B-L violation
- L violation ( $0\nu\beta\beta$ )

# Neutrino mass

To have a **Dirac mass**, need **light degrees of freedom**

$$m_D \bar{\nu}_L \nu_R$$

- Singlet under SM
- B-L is conserved
- L is conserved (or  $\beta\beta$ )

**New symmetr(ies) needed to forbid**  $M \bar{\nu}_R^c \nu_R$

Can be realized at renormalized level  $y \bar{\ell}_L \tilde{\Phi} \nu_R$  **New Higgs doublet**

Why so **light**?  $y \langle \Phi \rangle \ll v$

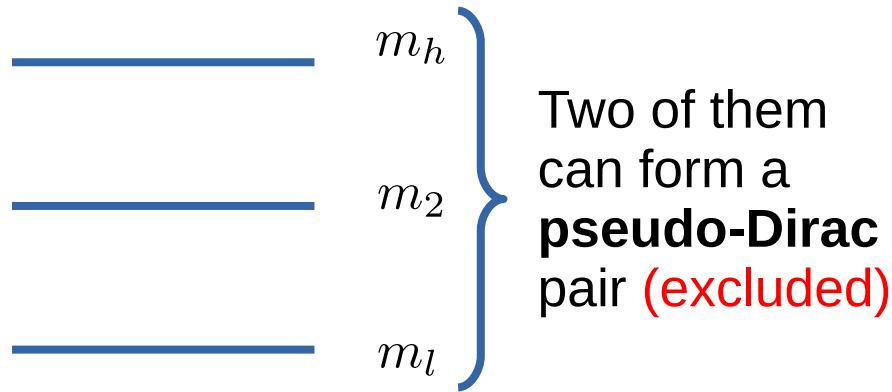
Let us look further again  $\frac{1}{\Lambda} \bar{\ell}_L \tilde{H} \nu_R \phi \implies \frac{\lambda v \langle \phi \rangle}{\Lambda} \bar{\nu}_L \nu_R$

$$\frac{1}{\Lambda} \bar{\ell}_L \tilde{H} \ell_L'^c \tilde{H}' \implies \frac{\lambda v v'}{\Lambda} \bar{\nu}_L \nu_L'^c$$

# Neutrino mass

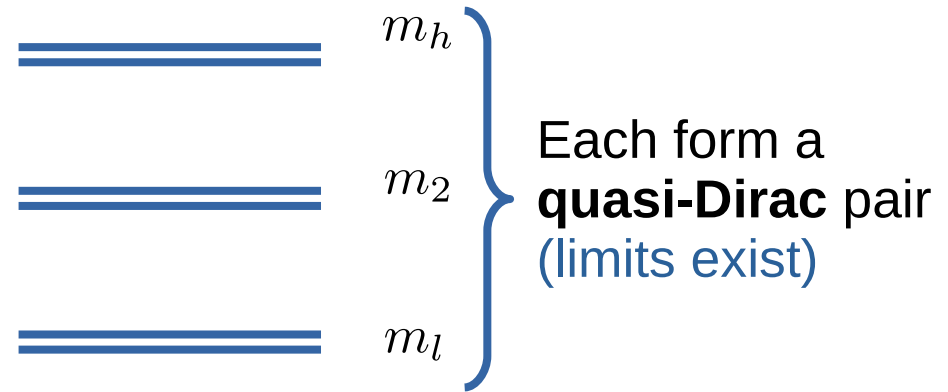
We can also have somewhere in between

Light neutrino mass spectrum (not to scale)



**Majorana:** only 3 light d.o.f.\*

**Dirac:** 2 or 3 more light d.o.f.



**Majorana:** 2 or 3 more light d.o.f.

Nomenclature of [Anamiati,Fonseca, Hirsch, 1710.06249]

Sorry, **Yuber Perez** (see his talk)

# A Dirac seesaw model

Extended **gauge** symmetries  $SU(2) \times U(1) \times SU(2)' \times U(1)'$

$$-\mathcal{L} \supset M_{ab} \bar{N}_{Ra}'^c N_{Rb} + y_{\alpha a} \bar{l}_{L\alpha} \tilde{H} N_{Ra} + y'_{\alpha a} \bar{l}'_{L\alpha} \tilde{H}' N_{Ra}' + \text{H.c.}$$

$$a, b = 1, 2, \dots, N_f \quad [\text{Gu, 1209.4579}] \quad [\text{Earl, CSF, Gregoire, Toner, 1903.12192}]$$

For anomaly cancellation, quarks prime have to be introduced

Perturbative level  $L_{\text{tot}} = L - L'$

$$L_{\text{tot}}(l_{L\alpha}) = L_{\text{tot}}(N_{Ra}) = -L_{\text{tot}}(l'_{L\alpha}) = -L_{\text{tot}}(N'_{Ra})$$

Quantum level

$$\Delta_{\text{tot}} = B - B' - L_{\text{tot}}$$

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After EW & EW' breaking

$$(\bar{\nu}_L \ \bar{N}_R'^c) \begin{pmatrix} 0 & vy \\ v'y'^T & M \end{pmatrix} \begin{pmatrix} \nu_L'^c \\ N_R \end{pmatrix}$$

**Dirac seesaw formula**

$$\mathcal{M}_\nu^D \simeq -vv'yM^{-1}y'^T$$

# A Dirac seesaw model

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$$a, b = 1, 2, \dots, N_f \quad [\text{Gu, 1209.4579}] \quad [\text{Earl, CSF, Gregoire, Toner, 1903.12192}]$$

We have conserved  $\Delta_{\text{tot}} = B - B' - L_{\text{tot}}$

Can we generate a cosmic baryon asymmetry through **leptogenesis**?

**Yes!** 
$$\sum_{\alpha} Y_{B/3-L_{\alpha}} - \sum_{\alpha} Y_{B'/3-L'_{\alpha}} - \sum_a Y_{\Delta N_a} = 0$$

$$\sum_{\alpha} Y_{B/3-L_{\alpha}} = \sum_{\alpha} Y_{B'/3-L'_{\alpha}} \neq 0$$

}

$$Y_B = \kappa Y_{B-L}$$

$$Y'_B = \kappa' Y_{B-L}$$

**Asymmetric  
Dark Matter**

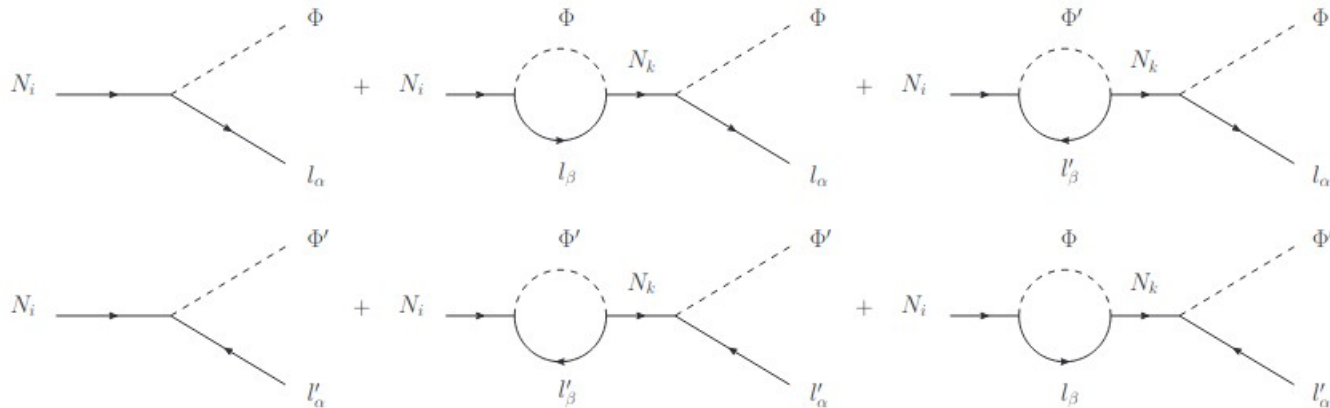


# A Dirac seesaw model: leptogenesis

From heavy N decays     [Gu, 1209.4579] [Earl, CSF, Gregoire, Toner, 1903.12192]

**CPT:** 
$$\sum_{\alpha} [\Gamma(N_a \rightarrow \ell_{\alpha} H) - \Gamma(\bar{N}_a \rightarrow \bar{\ell}_{\alpha} H^*)] = \sum_{\alpha} [\Gamma(\bar{N}_a \rightarrow \ell'_{\alpha} H') - \Gamma(N_a \rightarrow \bar{\ell}'_{\alpha} H'^*)]$$

$$\sum_{\alpha} \epsilon_{a\alpha} = \sum_{\alpha} \epsilon'_{a\alpha}$$



$$\Phi = H, \Phi' = H'$$

$$i, k, a = 1, 2, \dots, N_f$$

Require  $N_f \geq 2$

# A Dirac seesaw model: leptogenesis

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$$\sum_{\alpha} \epsilon_{a\alpha} = \sum_{\alpha} \epsilon'_{a\alpha}$$

We obtain analogous of **Davidson-Ibarra bound**

$$\left| \sum_{\alpha} \epsilon_{a\alpha} \right| \leq \frac{M_1 |m_3^2 - m_1^2|}{16\pi(m_3 + m_1)vv'}$$

**High scale leptogenesis**      $M_1 \gtrsim 10^6 \text{ GeV}$

# A quasi-Dirac seesaw model

Introducing small violation of  $\Delta_{\text{tot}} = B - B' - L_{\text{tot}}$  [CSF, Gregoire, Toner, 2007.09158]

$$-\mathcal{L} \supset \frac{1}{2} m_{ab} \bar{N}_{Ra}^c N_{Rb} + \frac{1}{2} m'_{ab} \bar{N}'_{Ra}{}^c N'_{Rb} + \tilde{y}_{\alpha a} \bar{l}_{L\alpha} \tilde{H} N'_{Ra} + \tilde{y}'_{\alpha a} \bar{l}'_{L\alpha} \tilde{H}' N_{Ra} + \text{H.c.}$$

$$m, m' \ll M, \quad \tilde{y}, \tilde{y}' \ll y, y'$$

Can arise from spontaneously broken

$$U(1)_{\Delta_{\text{tot}}}$$

$$\frac{\phi^2}{\Lambda} \bar{N}_R^c N_R$$

$$\frac{\phi^{*2}}{\Lambda} \bar{N}'_R{}^c N'_R$$

$$\frac{\phi^{*2}}{\Lambda^2} \bar{l}_L \tilde{H} N'_R$$

$$\frac{\phi^2}{\Lambda^2} \bar{l}'_L \tilde{H}' N_R$$

Even if no “soft” breaking  $m, m' = 0$

$$\delta m \sim \frac{(y \tilde{y}^\dagger + \tilde{y}' y'^T) M}{8\pi^2}$$

$$\delta m' \sim \frac{M(\tilde{y}'^\dagger y' + y^T \tilde{y})}{8\pi^2}$$

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Many new parameters!

We can consider some symmetries [Lee & Yang, Phys.Rev. 104 (1956) 254-258]

[Berezinsky, Narayan, Vissani, hep-ph/0210204] [Berezhiani, hep-ph/0312335]

$$Z_{2M} : \Psi_{L,R}(t, \vec{x}) \leftrightarrow \Psi'_{R,L}(t, -\vec{x})$$

$$M = M^\dagger \quad m^* = m' \quad y^* = y' \quad \tilde{y}^* = \tilde{y}'$$

$$Z_{2D} : \Psi_{L,R} \leftrightarrow \Psi'_{L,R}$$

$$M = M^T \quad m = m' \quad y = y' \quad \tilde{y} = \tilde{y}'$$

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Heavy quasi-Dirac N

$$\mathcal{M} = \begin{pmatrix} m & M^T \\ M & m' \end{pmatrix}$$

$$M_a^{\mp 2} = M_a^2 \mp \delta M_a^2 \quad \delta M_a^2 \sim M_a^2 \times \frac{|m' + m^*|}{M_a}$$

Resonant leptogenesis from each quasi-Dirac pair See **Juan Racker's** talk

CP violation for leptogenesis is bounded

$$|\epsilon_a| \lesssim \frac{\max(|\tilde{y}|, |\tilde{y}'|)}{\max(|y|, |y'|)}$$

Nonzero for  
each  
generation

# A quasi-Dirac seesaw model

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Light quasi-Dirac  $\nu$

$$\mathcal{M}_\nu \simeq -\mathcal{M}_D \mathcal{M}^{-1} \mathcal{M}_D^T \quad \mathcal{M}_D = \begin{pmatrix} \nu y & \nu \tilde{y} \\ f \tilde{y}' & f y' \end{pmatrix}$$

$$m_a^\mp = m_a \mp \delta m_a \quad m_a \sim \frac{|yy'|vv'}{M_a} \quad \delta m_a \sim m_a \times \frac{|y^* \tilde{y}^* v^2 + y' \tilde{y}' v'^2|}{|yy'|vv'}$$

Quasi-Dirac seesaw mechanism

# A quasi-Dirac seesaw model

Introducing small violation of  $\Delta_{\text{tot}} = B - B' - L_{\text{tot}}$  [CSF, Gregoire, Toner, 2007.09158]

$$-\mathcal{L} \supset \frac{1}{2} m_{ab} \bar{N}_{Ra}^c N_{Rb} + \frac{1}{2} m'_{ab} \bar{N}'_{Ra}{}^c N'_{Rb} + \tilde{y}_{\alpha a} \bar{l}_{L\alpha} \tilde{H} N'_{Ra} + \tilde{y}'_{\alpha a} \bar{l}'_{L\alpha} \tilde{H}' N_{Ra} + \text{H.c.}$$

We can connect “high” and low scale observables

1) Resonant leptogenesis ends by  $T > 132 \text{ GeV}$   $M_a, |\epsilon_a| \lesssim \frac{\delta m_a}{2m_a}$

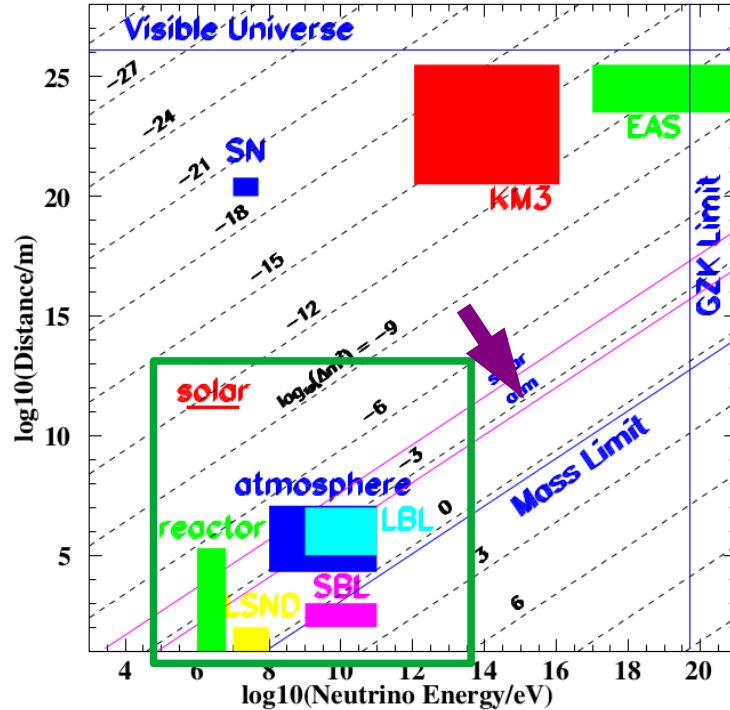
$$Y_B^{\text{obs}} \sim 10^{-10} \quad |\epsilon_a| \gtrsim 10^{-6} \implies \delta m_a \gtrsim 10^{-7} \text{ eV} \left( \frac{m_a}{0.05 \text{ eV}} \right)$$

2) Mass splitting that can be observed at low scale

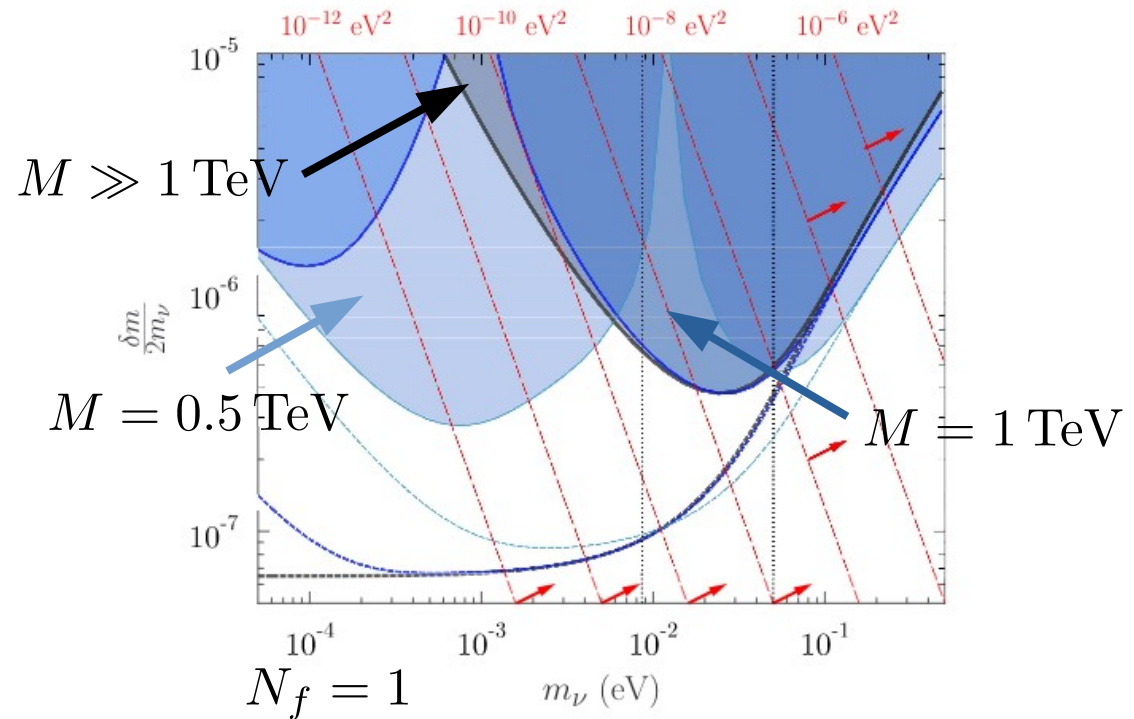
$$\varepsilon_a^2 \equiv m_a^{+2} - m_a^{-2} = 4m_a \delta m_a \gtrsim 2 \times 10^{-8} \text{ eV}^2$$

# A quasi-Dirac seesaw model

Overlapping parameter space with **neutrino oscillation experiments**



[Beacom et al., hep-ph/0307151]



[CSF, Gregoire, Toner, 2007.09158]



# A quasi-Dirac seesaw model

On-going effort to identify all the viable parameter space for experimental searches

## 1) Start with Dirac seesaw parametrization

$$y = \frac{i}{\sqrt{vv'}} U_\nu^* D_{\sqrt{m}} X D_{\sqrt{M}} \quad y' = \frac{i}{\sqrt{vv'}} V_\nu^* D_{\sqrt{m}} (X^{-1})^T D_{\sqrt{M}}$$

$$Z_{2M} : y = y'^* \implies V_\nu = -U_\nu^*, \quad X^{-1,\dagger} = X \quad \text{Complex orthogonal}$$

$$Z_{2D} : y = y' \implies V_\nu = U_\nu, \quad X^{-1,T} = X \quad \text{Hermitian}$$

$$2) \text{ Generate } \tilde{y}, \tilde{y}' \ll y, y', \quad \boxed{\delta m, \delta m'} \lesssim m, m' \ll M \quad \text{Radiative contributions}$$

$$3) \text{ Verify } Y_B^{\text{obs}} \sim 10^{-10}$$

$$4) \text{ Verify if consistent with experimental constraints } \delta m_a, \theta_k, \delta_k \quad \begin{matrix} (a = 1, 2, 3) \\ (k = 1, \dots, 9) \end{matrix} \quad 17$$

\*Previous results apply for each generation if  $m, m', M$  are all diagonal

# Outlook

- Light neutrinos can be **quasi-Dirac** with a large parameter space to be explored
- We present a “prototype” model with **quasi-Dirac seesaw mechanism**
- “**Baryogenesis miracle**”: parameter space consistent with leptogenesis can be probed in neutrino oscillation experiments
- Inputs are welcome: how to constraint the parameter space with neutrino oscillation experiments or other observables