

$$G_{\text{ms}} \equiv R_{\text{ms}} - \frac{1}{2} R_{g_{\text{ms}}} = \frac{8\pi G}{c^3} T_{\text{ms}}$$

$$S_{\text{B}} = \frac{k_{\text{B}} 4\pi G}{hc} M^2$$

$$\Psi(x) = \frac{1}{\sqrt{k_1}} (A_+ e^{ik_1 x} + A_- e^{-ik_1 x}) \quad x < 0$$

$$k_1 = \sqrt{2mE/\hbar^2}$$

$$R_{\text{ms}} - \frac{1}{2} R_{g_{\text{ms}}} + \Lambda_{g_{\text{ms}}} = \frac{8\pi G}{c^3} T_{\text{ms}}$$

Giorgio GALANTI

INAF - IASF Milano

$$H = \frac{P^2}{2m} + V(r)$$

$$P = -i\hbar\nabla$$

Re[Ψ(x)]



$$S = \frac{1}{2\pi} \int (P \dot{q} - H) dt$$

$$S = \frac{1}{4\pi G} \int R_{\mu\nu} dx^\mu dx^\nu$$

# Studying the importance of possible oscillations of photons into ALPs in pulsars

$$L = \frac{1}{4} F_{\mu\nu} F^{\mu\nu} - i\lambda F^{\mu\nu} D_\mu \lambda$$

$$H|\psi(t)\rangle = i\hbar \frac{\partial}{\partial t} |\psi(t)\rangle$$

$$\omega = \frac{\theta}{2\pi} + \frac{4\pi}{g^2}$$

$$I = \int \frac{1}{c^2} \frac{\partial^2 \psi}{\partial t^2} dt - \nabla^2 \psi + \frac{m^2 c^2}{\hbar^2} \psi = 0$$

Supervisor: Patrizia Caraveo

9th International Fermi Symposium

15th April 2021

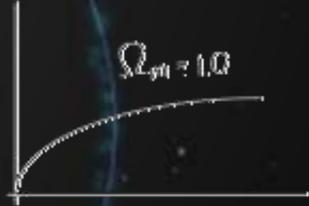
$$\frac{d}{dt} \langle A \rangle = \frac{1}{i\hbar} \langle [A, H] \rangle + \left\langle \frac{\partial A}{\partial t} \right\rangle$$

$$i\hbar \frac{\partial \psi}{\partial t} = -\frac{\hbar^2}{2m} \nabla^2 \psi + V\psi$$

$$\Delta x \Delta p \geq \frac{\hbar}{2}$$

$$A_{ij} = \frac{8\pi h \nu^3}{c^3} B_{ij}$$

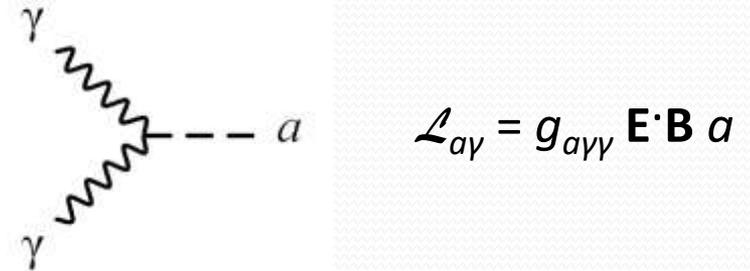
$$S = \frac{1}{2} \int d^4x \left( P^2 - \frac{P^2}{6M^2} \right)$$



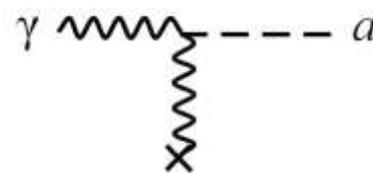
# Axion-like particles (ALPs)

- Predicted by String Theory
- Very light particles ( $m_a < 10^{-7}$  eV)
- Spin 0
- Interaction with two photons (coupling  $g_{a\gamma\gamma}$ )
- Interactions with other particles discarded
- Possible candidate for dark matter

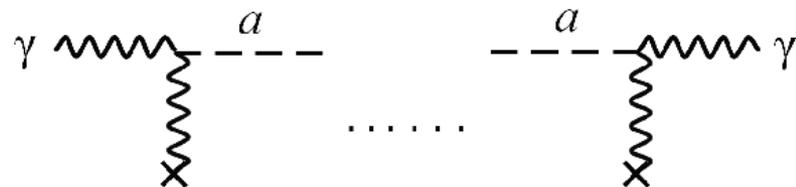
## Two photons



## In an external B field



## Photon-ALP oscillations



## QUESTION:

**Can pulsars tell us anything about axion-like particles (ALPs)?**

Possible photon-ALP conversion in the pulsar magnetic field<sup>1</sup> and/or in that of the Galaxy<sup>2</sup>?

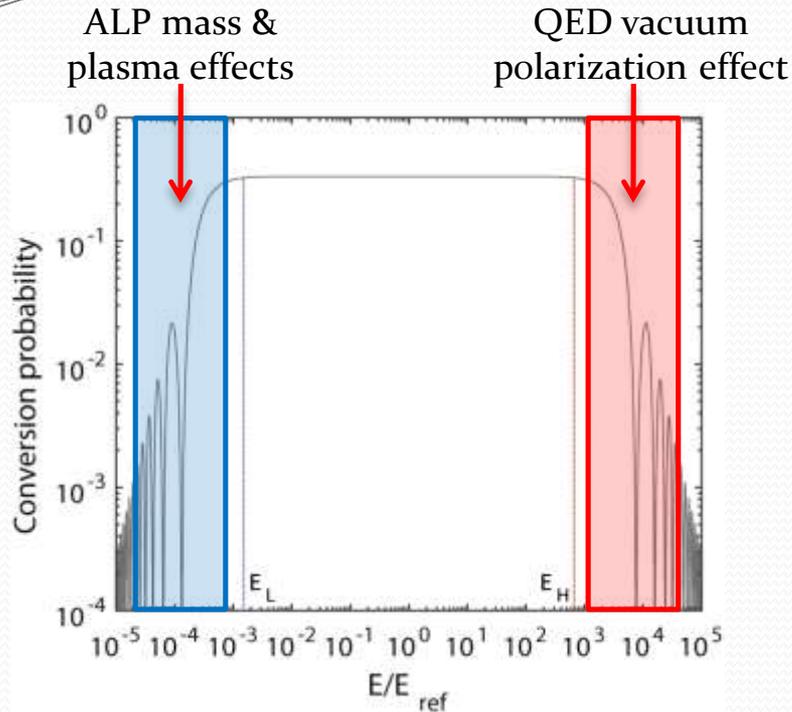
Our plan (for the future):

- Evaluate ALP hints from pulsar comparison
- Extend existing studies to the largest possible sample
- Statistical analysis to constrain ALP parameter space
- Multi-frequency analysis

[1] Perna et al. 2012

[2] Majumdar et al. 2018

# ALP-induced irregularities



- Photon-ALP conversion probability  $P_{\gamma \rightarrow a}(E, m_a, g_{a\gamma\gamma}, B)$
- Effective photon absorption
- Highlighted zones predict spectral irregularities in observational data
- Constraints on  $g_{a\gamma\gamma}$  and  $m_a$

## BLUE AREA:

- Modifications to pulsar spectra for photon-ALP conversion in the Galaxy

## RED AREA:

- Modifications to pulsar spectra for photon-ALP conversion in the source

# ALP-induced irregularities

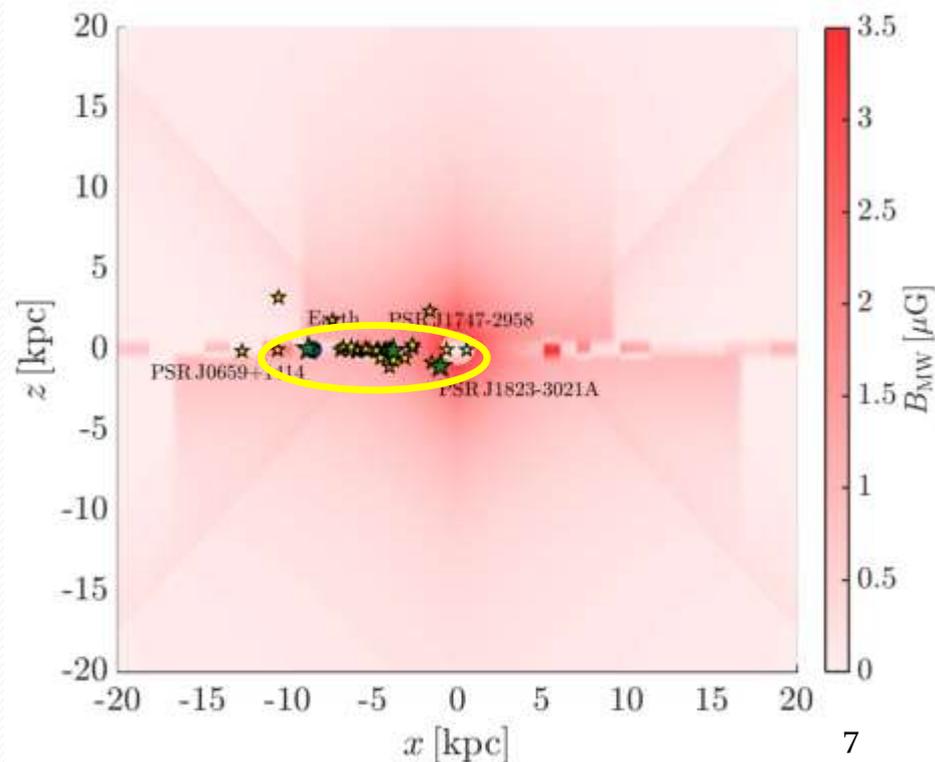
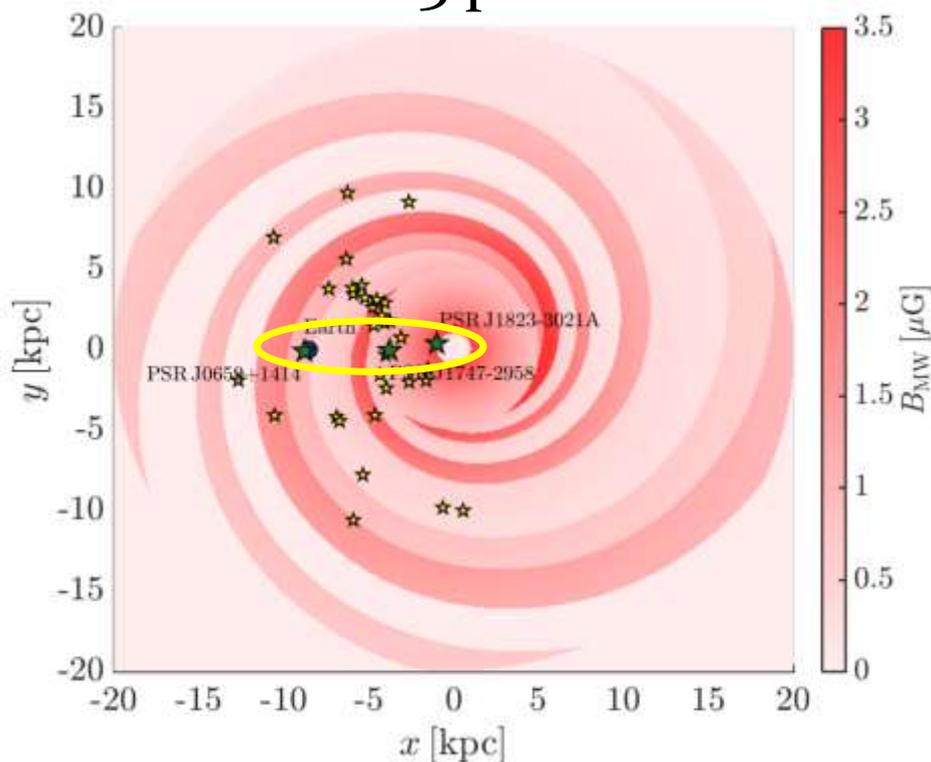
- We can calculate the photon survival probability (in the presence of photon-ALP interactions)  $P_{\gamma \rightarrow \gamma}$
- *Emitted* (intrinsic) flux  $F_{\text{em}}$  is not equal to the *observed* one  $F_{\text{obs}}$
- photon-ALP interactions  $\rightarrow F_{\text{obs}} = P_{\gamma \rightarrow \gamma} F_{\text{em}}$
- No photon-ALP interactions  $\rightarrow P_{\gamma \rightarrow \gamma} = 1, F_{\text{obs}} = F_{\text{em}}$

# Photon-ALP conversion **in the source**

- Photon-ALP mixing term:  $\Delta_{a\gamma} = \frac{1}{2} g_{a\gamma\gamma} B_T$ ;  $B_T$ : **transverse**  $B$
- QED term  $\propto$  energy  $E$ :  $\Delta_{xx}^{QED} = \frac{4}{2} \left(\frac{\alpha}{45\pi}\right) \left(\frac{B_T}{B_{cr}}\right)^2 E$   $\Delta_{zz}^{QED} = \frac{7}{2} \left(\frac{\alpha}{45\pi}\right) \left(\frac{B_T}{B_{cr}}\right)^2 E$
- Dipolar magnetic field  $B$ :  $10^{11} - 10^{13}$  G in the centre
- By taking photon-ALP coupling  $g_{a\gamma\gamma} = O(10^{-11} \text{ GeV}^{-1})$
- **GeV energy band**:  $P_{\gamma \rightarrow a} < 10^{-6} \rightarrow$  **no effects**
- X-ray energy band:  $P_{\gamma \rightarrow a} = O(0.1) \rightarrow$  effects on spectra and possible correlation with pulsar emission mechanisms

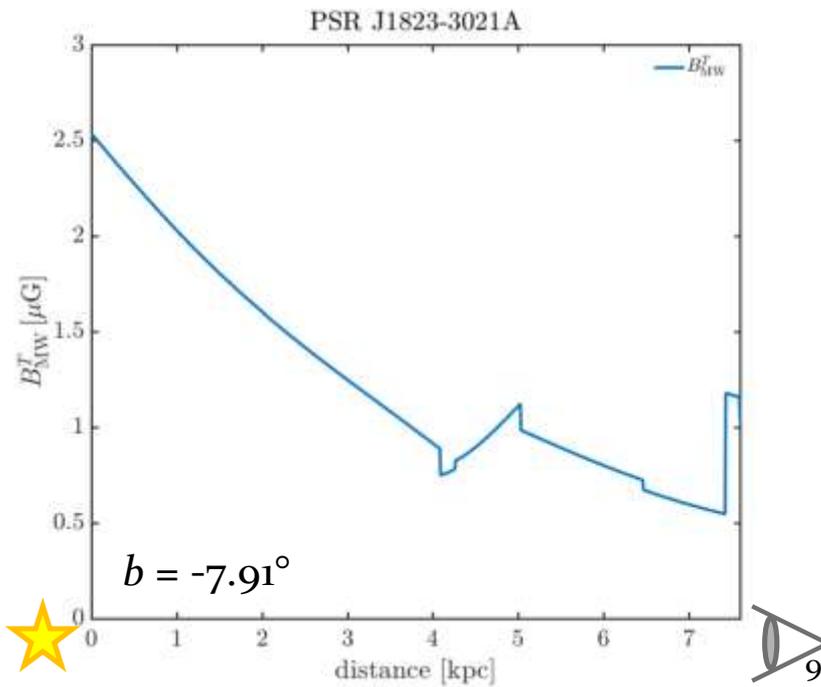
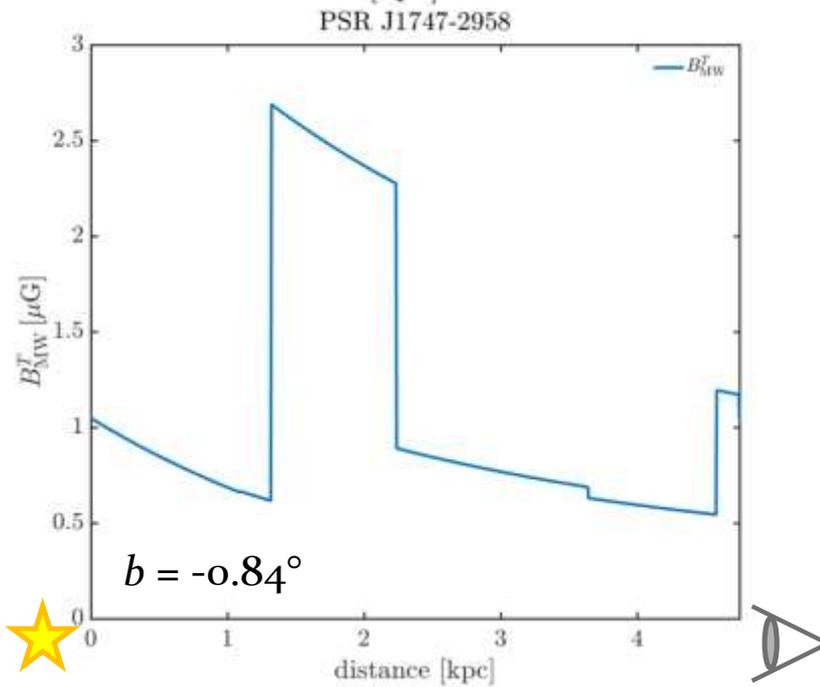
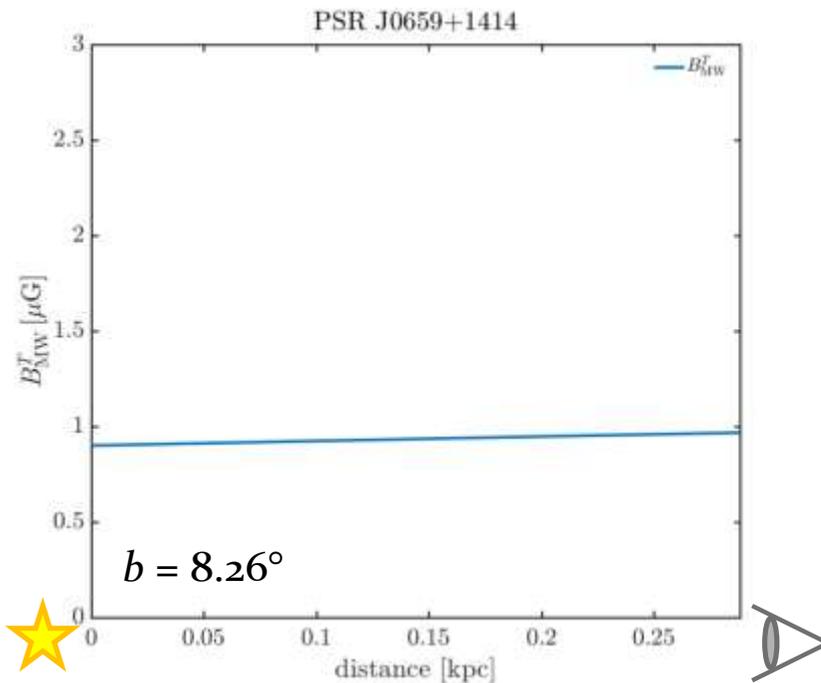
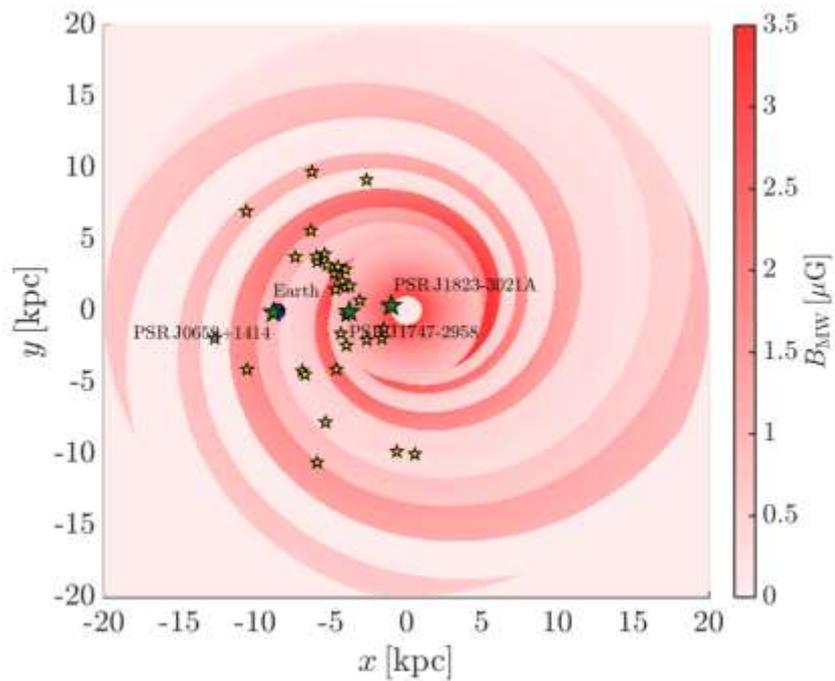
# Photon-ALP conversion **in the Galaxy**

- Photon-ALP mixing term:  $\Delta_{a\gamma} = \frac{1}{2}g_{a\gamma\gamma}B_T$ ;  $B_T$ : **transverse**  $B$
- ALP mass term:  $\Delta_{aa} = -m_a^2/(2E)$ ;  $m_a$ : ALP mass
- Galactic magnetic field  $B_{MW}$  by Jansson & Farrar 2012 maps
- **Select pulsars with distance > 4 kpc** for sizable ALP effects
- Focus on 3 pulsars for now

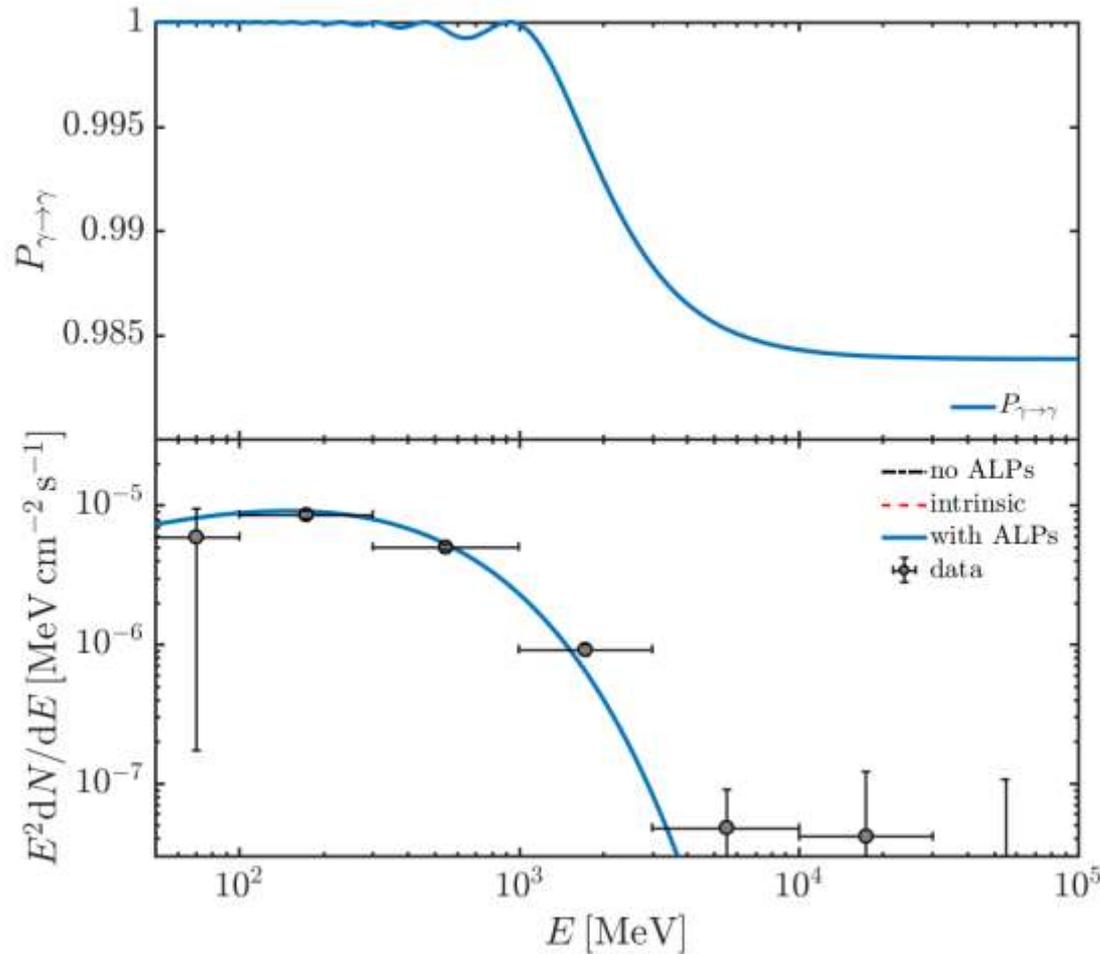


# Photon-ALP conversion in the Galaxy

- **Best pulsars for ALP studies:**
  - distance  $> 4$  kpc
  - in the Galactic plane (low Galactic latitude  $b$ )
  - located in zones where  $B_{MW}$  is high
  - better toward the centre for high **transverse**  $B_{MW}^T$  due to  $B_{MW}$  vector structure
- Spectra modified by ALPs ( $g_{a\gamma\gamma} = 0.7 \times 10^{-10} \text{ GeV}^{-1}$ ;  $m_a = 2 \times 10^{-8} \text{ eV}$ )
- Power law super exponential cutoff fit (PLECn models) in PLEC1 form  $N \exp[\Gamma \ln(E/E_o) - (E/E_{\text{cut}})^b]$  applied to *emitted* spectra (modified by  $P_{\gamma \rightarrow \gamma}$ ):
  - PSR J0659+1414 (very close, to crosscheck results)
  - PSR J1747-2958
  - PSR J1823-3021A

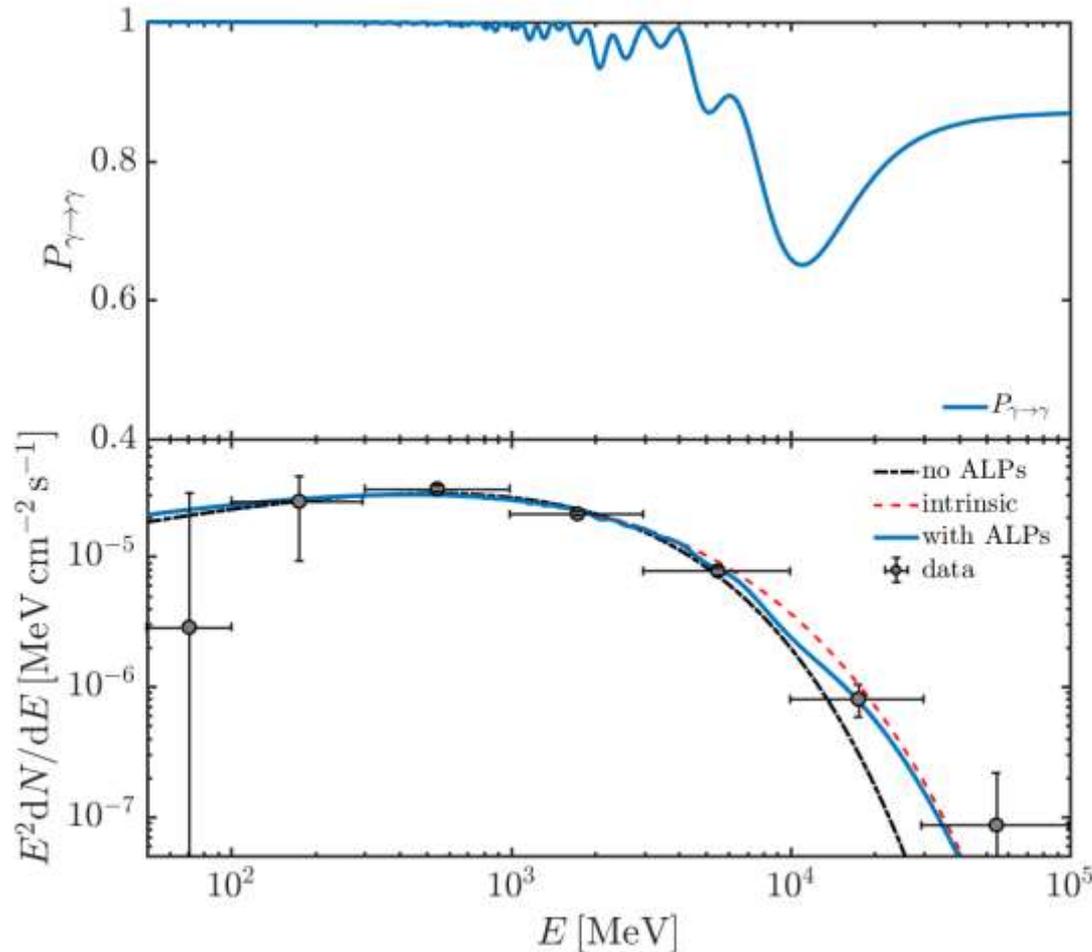


# PSR J0659+1414



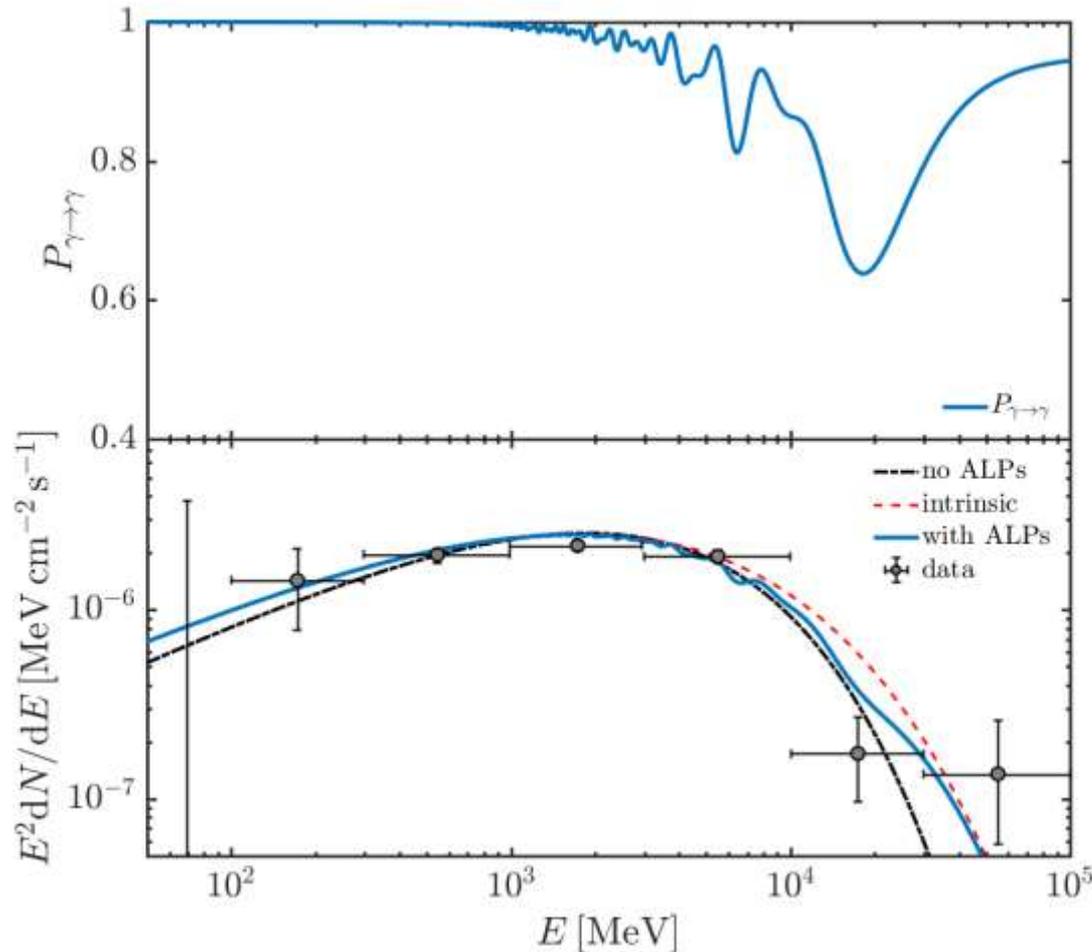
- No sizeable ALP effects (as expected)
- Fit unchanged
- Data and fits without ALPs  $\rightarrow$  from *Fermi*-LAT repository

# PSR J1747-2958



- ALP effects for  $E > 1 \text{ GeV}$
- Tentative new heuristic fit “with ALPs” resulting in
  - $\Gamma = 1.61$
  - $E_{\text{cut}} = 811.7 \text{ MeV}$
  - $b = 0.56$
- Small increase of  $\Gamma$ , small decrease of  $E_{\text{cut}}$  and  $b$

# PSR J1823-3021A

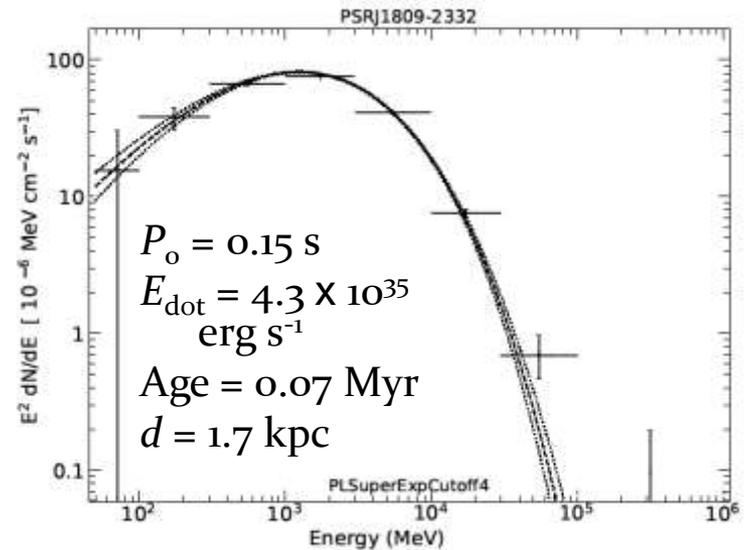
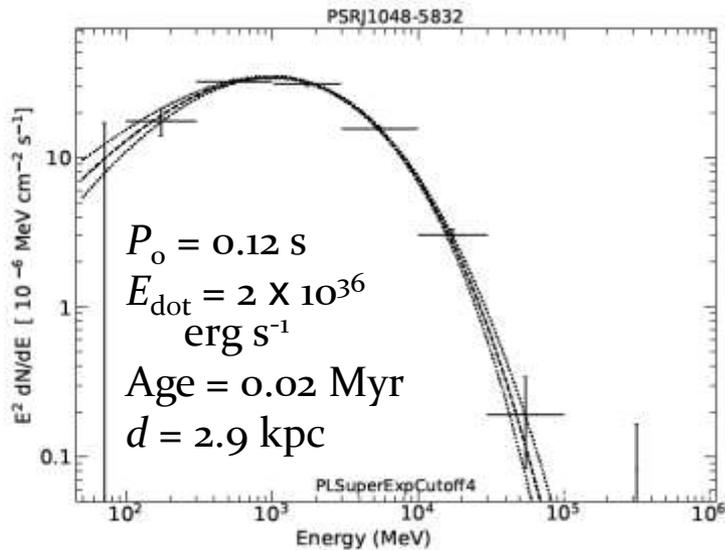
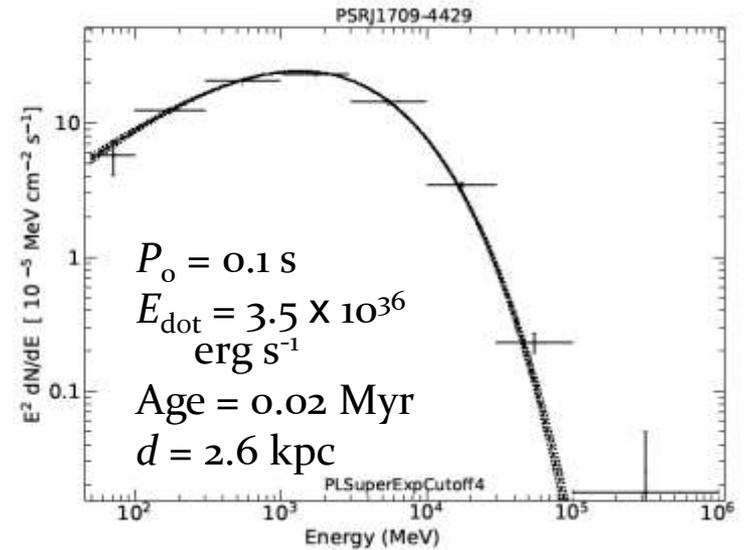
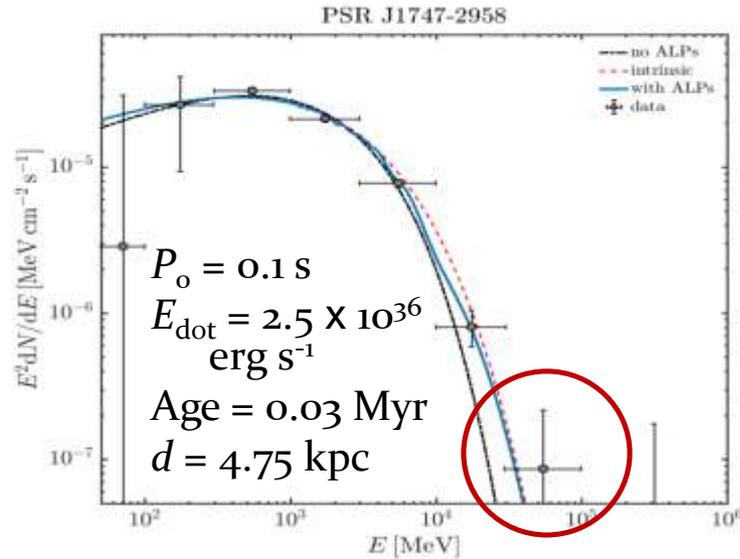


- ALP effects for  $E > 1 \text{ GeV}$
- Tentative new heuristic fit “with ALPs” resulting in
  - $\Gamma = 1.33$
  - $E_{\text{cut}} = 1100.8 \text{ MeV}$
  - $b = 0.53$
- Small increase of  $\Gamma$ , small decrease of  $E_{\text{cut}}$  and  $b$

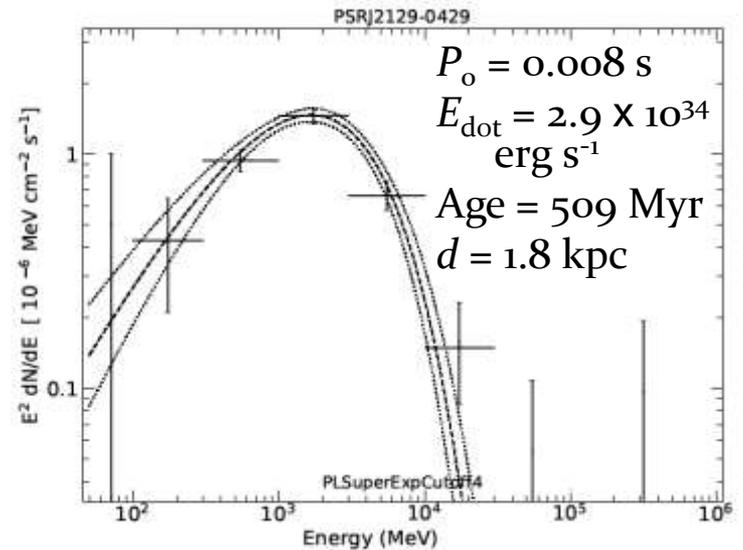
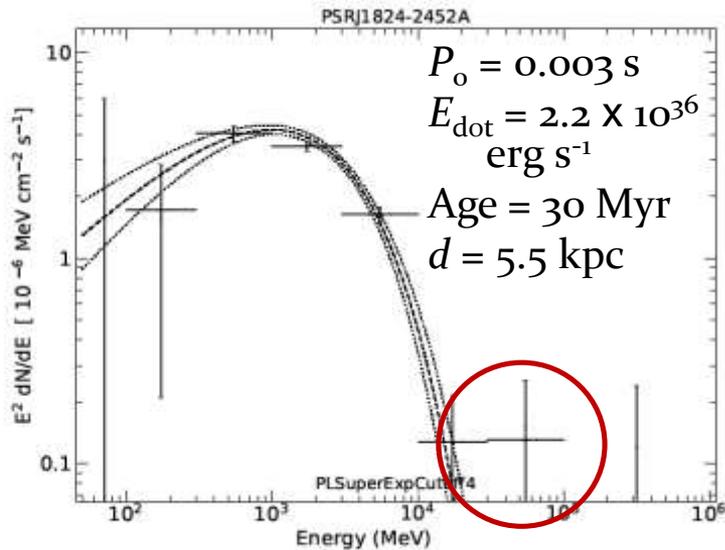
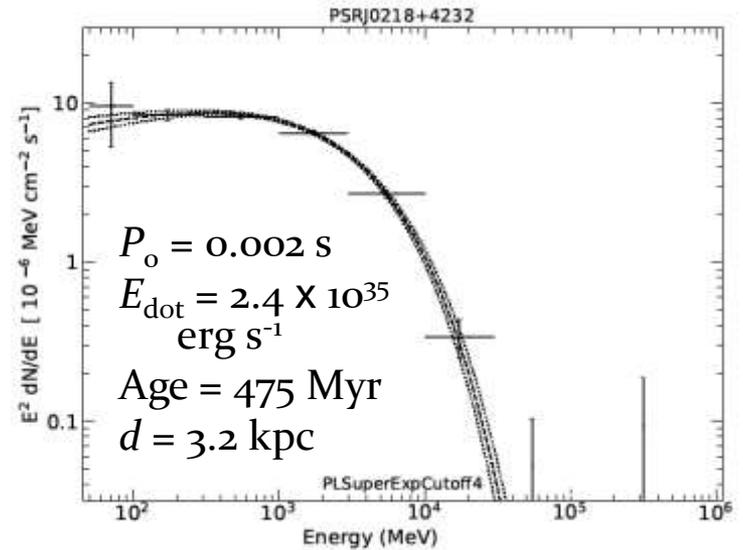
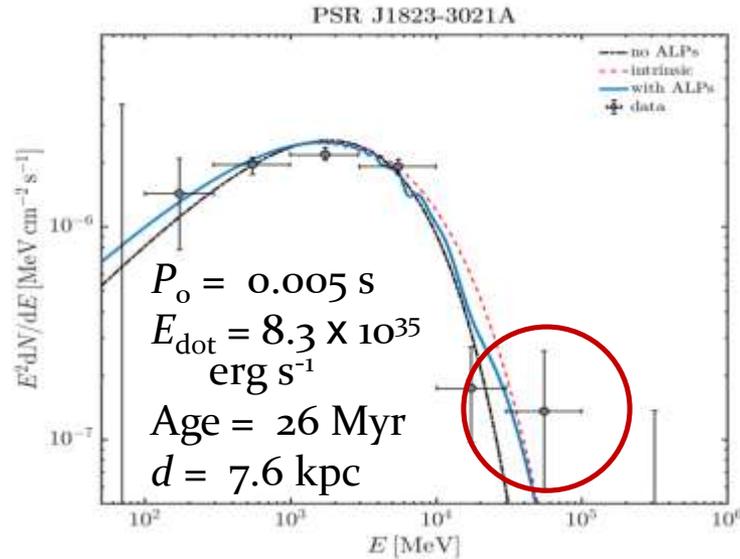
# Photon-ALP conversion in the Galaxy

- **Distance is the key parameter**
- Comparison with pulsars similar to those considered but at *different distances* (similar *emitted* spectra are expected)
- **Hint on the presence of ALPs** if:
  - Comparing similar pulsars at different distances should yield different spectral shapes as a function of distance
  - **farther pulsars should have spectra more modified than closer ones**
- Without such correlation → hint against ALPs

# PSR J1747-2958



# PSR J1823-3021A



# Conclusions (hints so far...)

- In the GeV energy range what matters for photon-ALP conversion is the *transverse* component of  $B_{\text{MW}}$  only
- **Best pulsars for ALP studies: distant, in the Galactic plane, towards the centre**
- Possible improvement of the spectral fits (to be confirmed with a more ROBUST analysis)
- Possible existence of ALP hint from the behavior of pulsar with similar properties but at different distances (to be investigated further)
- $g_{a\gamma\gamma} = O(10^{-10} \text{ GeV}^{-1})$  close to the upper limit of CAST,  $m_a = O(10^{-8} \text{ eV})$

# Future perspectives

- Possible ON-peak analysis to consider pulsar phase (thanks David)
- Perform a likelihood analysis to test the best ALP parameter space ( $g_{a\gamma\gamma}$ ,  $m_a$ ) and significance of the model
- Find hints for ALPs or constrain ALP parameter space by considering all known pulsars
- Continue analysis of the “ALP behavior” of similar pulsars
- Study the ALP effects in the X-ray band:
  - Modeling photon-ALP conversion in the source
  - Different pulsar  $B$  fields (dipolar, multipolar, ...)
  - Different emission models
- Study of possible links between GeV and X-ray band

$$G_{rms} \equiv R_{rms} - \frac{1}{2} R_{g_{rms}} = \frac{8\pi G}{c^3} T_{rms}$$

$$S_{\text{eff}} = \frac{k_B 4\pi G}{hc} M^2$$

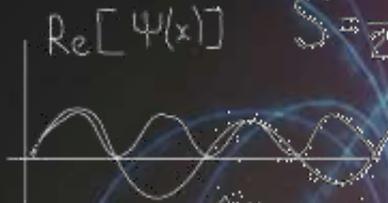
$$\psi(x) = \frac{1}{\sqrt{K}} (A_+ e^{ikx} + A_- e^{-ikx}) \quad x < 0$$

$$k_1 = \sqrt{2mE/\hbar^2}$$

$$R_{rms} - \frac{1}{2} R_{g_{rms}} + \Lambda_{g_{rms}} = \frac{8\pi G}{c^3} T_{rms}$$

$$H = \frac{P \cdot P}{2m} + V(r)$$

$$P = -i\hbar \nabla$$



$$\sigma = \frac{2\pi \hbar^2 \mu^2}{T^2 (1 - \mu^2)}$$

$$S = \frac{1}{2\pi} \int P \sqrt{1 - g_{\mu\nu}^2} dx$$



$$S = \frac{e^2 k A}{4\hbar c}$$

$$L = \text{tr} \left[ \frac{1}{g} F_{\mu\nu} F^{\mu\nu} - i\lambda \Gamma^{\mu} D_{\mu} \lambda \right]$$

$$H|\psi(t)\rangle = i\hbar \frac{\partial}{\partial t} |\psi(t)\rangle$$

Thank you

$$r = \frac{\theta}{2\pi} + \frac{4\pi}{g^2}$$

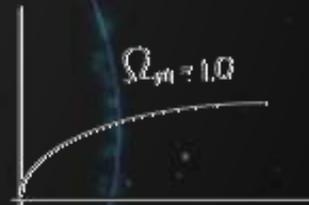
$$I = \int e^{-ax^2/2} dx = \sqrt{\frac{2\pi}{a}}$$

$$E^2 = p^2 c^2 + m^2 c^4$$

$$\frac{1}{c^2} \frac{\partial^2}{\partial t^2} \psi - \nabla^2 \psi + \frac{m^2 c^2}{\hbar^2} \psi = 0$$

$$p = \hbar k = \frac{\hbar m v}{c} = \frac{\hbar}{\lambda}$$

$$S = \frac{1}{2} \int d^4x \left( P^2 - \frac{P^2}{6M^2} \right)$$



$$A_{ij} = \frac{8\pi \hbar v^2}{c^3} B_{ij}$$

$$S_{gr} = \langle \dot{P} | \dot{S} | \dot{P} \rangle$$

$$dY = e^{-\int_0^s V(X(\tau)) d\tau} (X, s) \frac{\partial u}{\partial X} dW$$

$$\frac{d}{dt} \langle A \rangle = \frac{1}{i\hbar} \langle [\hat{A}, \hat{H}] \rangle + \left\langle \frac{\partial \hat{A}}{\partial t} \right\rangle$$

$$i\hbar \frac{\partial}{\partial t} \psi = -\frac{\hbar^2}{2} \sum_{n=0}^{\infty} \frac{1}{m_n} \nabla_n^2 \psi + V\psi$$

$$\Delta x \Delta p \geq \frac{\hbar}{2}$$

$$G_{rms} \equiv R_{rms} - \frac{1}{2} R_{g_{rms}} = \frac{8\pi G}{c^3} T_{rms}$$

$$S_{\text{class}} = \frac{k_B 4\pi G}{hc} M^2$$

$$\psi(x) = \frac{1}{\sqrt{K}} (A_- e^{ikx} + A_+ e^{-ikx}) \quad x < 0$$

$$k_1 = \sqrt{2mE/\hbar^2}$$

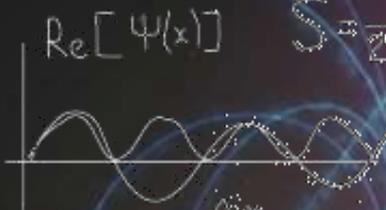
$$R_{rms} - \frac{1}{2} R_{g_{rms}} + \Lambda_{g_{rms}} = \frac{8\pi G}{c^3} T_{rms}$$

$$\sigma = \frac{2\pi \hbar^2 \mu^2}{T^2 (2\mu - \mu)}$$



$$H = \frac{P^2}{2m} + V(r)$$

$$P = -i\hbar \nabla$$



$$S = \frac{1}{2\pi} \int P dx - g_{\text{class}} x$$

$$S = \frac{e^2 k A}{4\hbar c}$$

$$L = \text{tr} \left[ \frac{1}{g} F_{\mu\nu} F^{\mu\nu} - i\lambda \Gamma^{\mu} D_{\mu} \lambda \right]$$

$$H|\psi(t)\rangle = i\hbar \frac{\partial}{\partial t} |\psi(t)\rangle$$

$e^-$

$$\frac{\delta(k_+ k_-)}{k^2}$$

$$E = mc^2$$

$$E^2 = (pc)^2 + (mc^2)^2$$

$$r = \frac{\theta}{2\pi} + \frac{4\pi}{g^2}$$

$$I = \int e^{-ax^2/2} dx = \sqrt{\frac{2\pi}{a}}$$

$$E^2 = p^2 c^2 + m^2 c^4$$

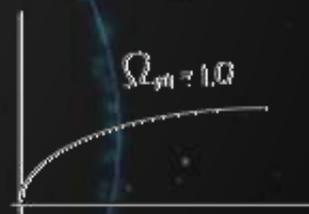
$$\frac{1}{c^2} \frac{\partial^2}{\partial t^2} \psi - \nabla^2 \psi + \frac{m^2 c^2}{\hbar^2} \psi = 0$$

$$p = \hbar k = \frac{\hbar m v}{c} = \frac{\hbar}{\lambda}$$

$$A_{ij} = \frac{8\pi \hbar v^2}{c^3} B_{ij}$$

$$S_{\text{gr}} = \langle \mathcal{R} | \mathcal{S} | \mathcal{I} \rangle$$

$$S = \frac{1}{2} \int d^4x \left( P^2 + \frac{P^2}{6M^2} \right)$$



$$dY = e^{-\int_0^s V(X(\tau)) d\tau} (X, s) \frac{\partial u}{\partial X} dW$$

$$\frac{d}{dt} \langle A \rangle = \frac{1}{i\hbar} \langle [\hat{A}, \hat{H}] \rangle + \left\langle \frac{\partial \hat{A}}{\partial t} \right\rangle$$

$$i\hbar \frac{\partial}{\partial t} \psi = -\frac{\hbar^2}{2} \sum_{n=0}^{\infty} \frac{1}{m_n} \nabla_n^2 \psi + V\psi$$

$$\Delta x \Delta p \geq \frac{\hbar}{2}$$