

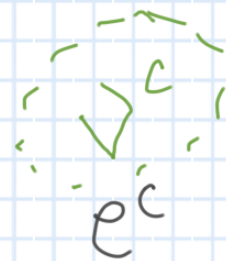
3) NEUTRINO MASSES

- NONZERO NEUTRINO MASSES IMPLY THE EXISTENCE OF NEW DEGREES OF FREEDOM BEYOND $Q u^c d^c L e^c H \oplus SU(3) \times SU(2) \times U(1)$
- WE DON'T KNOW WHAT ARE THESE NEW DEGREES OF FREEDOM. THEY CAN BE FERMIONS OR BOSONS, VERY LIGHT, VERY HEAVY, ETC.

Q) DIRAC NEUTRINOS:

$$Q = \begin{pmatrix} u \\ d \end{pmatrix} \quad \begin{matrix} u^c \\ d^c \end{matrix}$$

$$L = \begin{pmatrix} \nu \\ e \end{pmatrix}$$



$$-L \supset \lambda \bar{\nu}^c H \langle H \rangle$$

$-\frac{1}{2} \quad +\frac{1}{2} = 0$

$$m_D \nu \nu^c$$

L_D DIRAC

$\nu^c \Rightarrow [1, 1, 0]$

$SU(3) \quad SU(2)_L \quad U(1)_Y$

$$-L \supset \frac{M}{2} \nu^c \nu^c$$

ASIDE: GLOBAL SYMMETRIES OF THE SM

QUDLEH + GAUGE SYMMETRIES \Rightarrow $\begin{cases} U(1)_B & \text{ACCIDENTAL, CLASSICAL*} \\ U(1)_L & \text{GLOBAL SYMM.} \end{cases}$

e.g. $\begin{cases} Q_L(L) = +1 \times 2 & \sum Q = 0 \\ Q_L(e^c) = -1 \times 1 & \sum Q^3 = 0 \\ Q_L(\nu^c) = -1 & \sum Q = 0 \end{cases}$; $[* U(1)_{B-L} \text{ ANOMALY FREE}]$

* - WE DID NOT "ASK" FOR THEM

- IF YOU EXTEND THE SM, THERE IS NO REASON TO EXPECT THEY WILL STILL BE CONSERVED

- VIOLATED BY QUANTUM GRAVITY (?)

$$-L \ni \lambda_v L v^c H$$

$$\begin{matrix} 4 & 0 & 0 \\ \hline +1 & -1 & 0 \end{matrix}$$

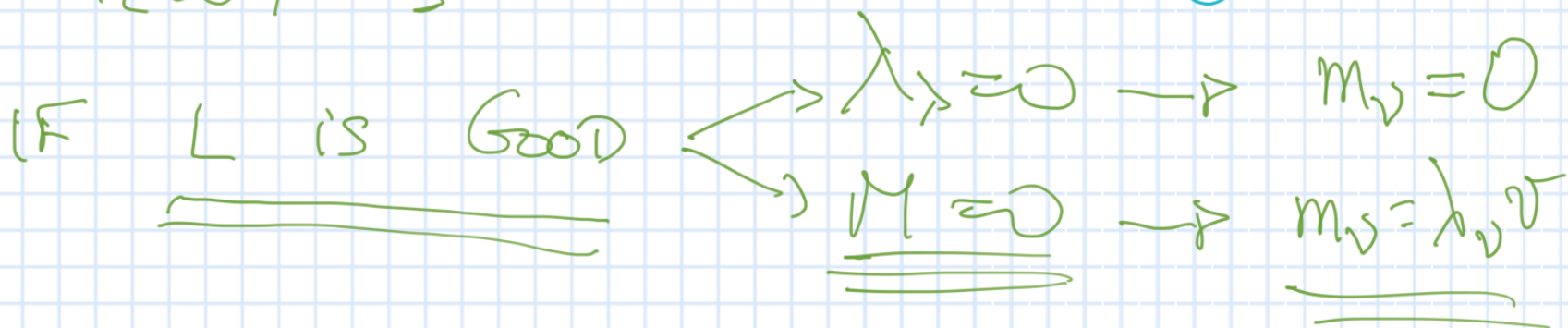
~~$$+ \begin{matrix} 0 & 0 \\ \hline M & v^e v^c \\ \hline 2 & -1 & -1 \end{matrix}$$~~

$$Q_L(L) = +1$$

$$Q_L(\underline{v^c}) = -1$$

$$Q_L(e^c) = -1$$

$$= 0$$



PERFECTLY GOOD MODEL

• NEW PARTICLE ν^c , VERY LIGHT FERMION

• LEPTON NUMBER IS NOT ACCIDENTAL SYM.
→ FUNDAMENTAL SYM.

$\mathcal{L} \rightarrow Q u d l e \nu^c H + \text{GAUGE SYM} + L$ IS GOOD

CONPLAINT(S)

$m_\nu = \lambda_\nu \nu$

$\hookrightarrow 100 \text{ eV}$

$\hookrightarrow 10^{-12}$

$\lambda_e \sim 10^{-5}$

$$- \mathcal{L} \supset \lambda_\nu \underline{L \nu^c H}$$

$$\lambda_\nu \sim 10^{-12} \sim 0$$

Sym.

$$Q_{\text{NEW}}(\nu^c) = +1$$

$$Q_{\text{NEW}}(Q u d L e H) = 0$$

$$Q_{\text{NEW}}(\phi) = -1$$

$$\underline{L \nu^c H \phi}$$

^

$$\langle H \rangle, \langle \phi \rangle \rightarrow$$

$$m_\nu \nu^c \nu ; m_\nu \sim \frac{v \cdot v_\phi}{\Lambda}$$

DARK SECTOR

$$\begin{aligned} \nu^c &\rightarrow L \nu^c H \\ \phi &\rightarrow \underbrace{|\phi|^2 |H|^2} \end{aligned}$$

$$v_\phi, v \ll \Lambda$$

MAJORANA MASSES

$\mathcal{L} \supset \cancel{\overbrace{m\psi\psi}^{0+0=0}}$



NO
NEW LIGHT
FIELD

violates $SU(2)$



$L = \begin{pmatrix} \nu \\ e \end{pmatrix}$

$m \cancel{\overbrace{L L}^{\text{Lorentz Singlet}}}$

$L \rightarrow -\frac{1}{2} - \frac{1}{2} = -1$

VIOLATES CHARGE INVARIANCE



$\overbrace{\begin{pmatrix} -\frac{1}{2} + \frac{1}{2} \\ (LH) (LH) \end{pmatrix}}^{\text{Lorentz Singlet}}$

GAUGE

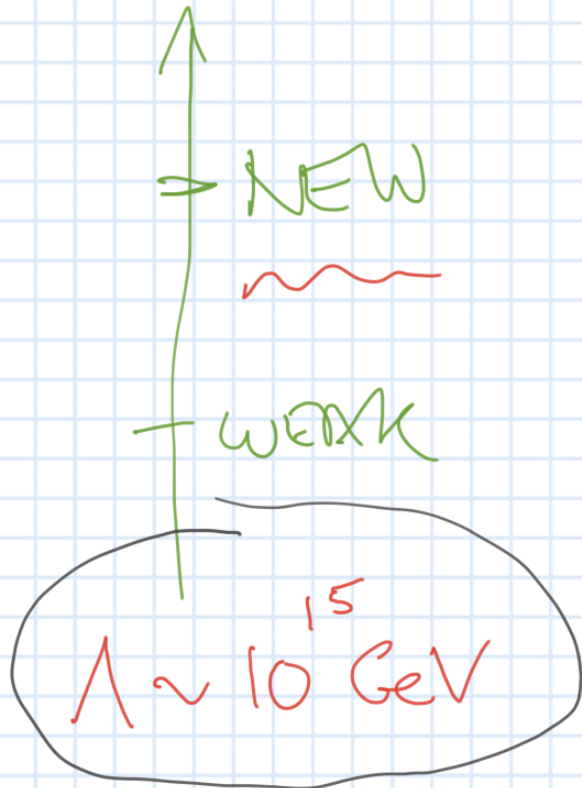
SINGLET

LORENTZ

SINGLET

$$- \mathcal{L} \supset \frac{(LH)(LH)}{\Lambda}$$

DIMENSION-5
WEINBERG (79).



\cancel{K}

$$- \mathcal{L}(\langle H \rangle) = m_\nu \nu \nu$$

$$m_\nu = \frac{v^2}{\Lambda}$$

$$m_{\text{CHARGED}} \times E \rightarrow \text{EXPERIMENT}$$

$$\Lambda \sim 10^{15} \text{ GeV}$$



ENERGY ABOVE WHICH
WE ARE SURE THE SM
BREAKS DOWN

$$\mathcal{L} \left(\frac{1}{\sqrt{2}} G_F \right) \bar{\mu} \gamma_{\mu} P_L \nu_{\mu} \bar{\nu}_e \gamma^{\mu} P_L e$$

$\left(\frac{1}{246 \text{ GeV}} \right)^2$

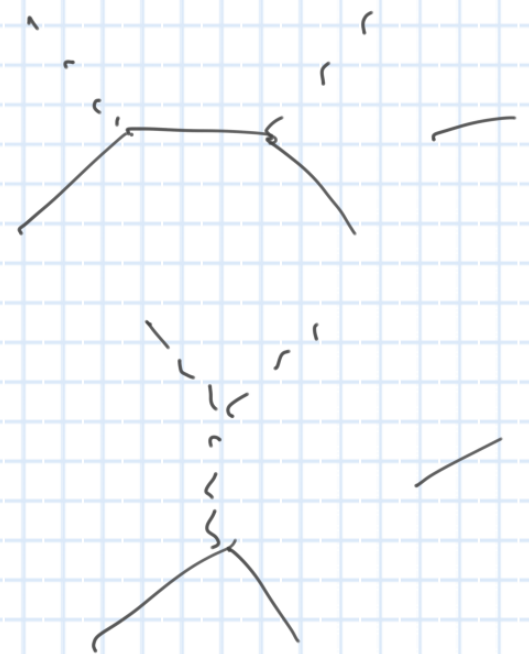
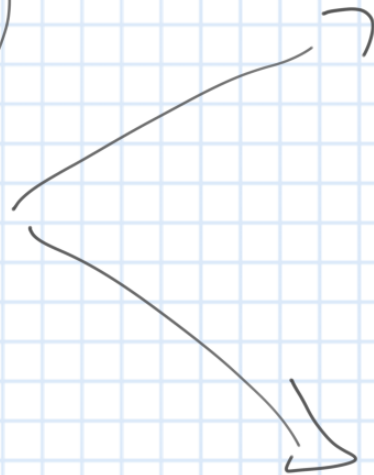
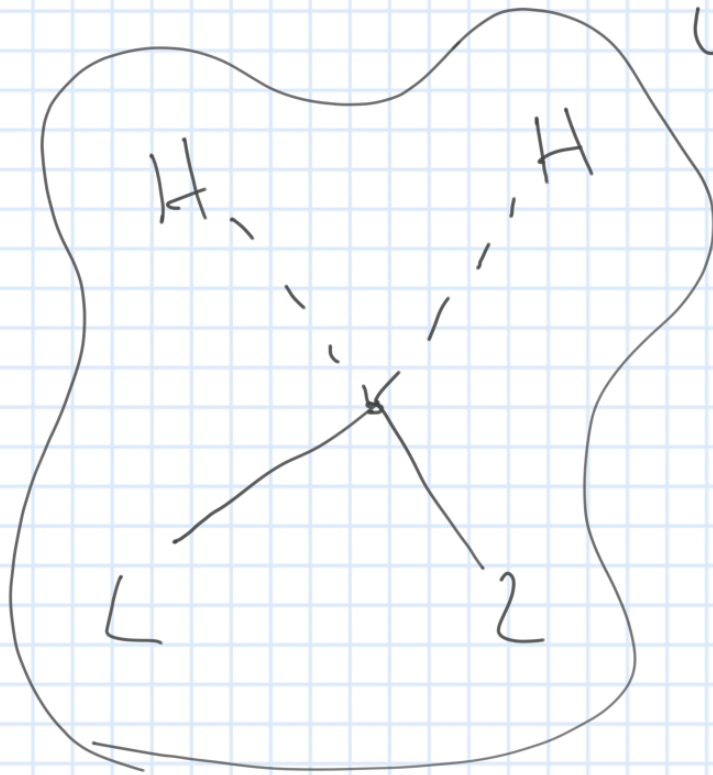
$$M_W \sim 80 \text{ GeV}$$

$$K^+ \rightarrow \pi^0 W^+$$

$$|H|^2 (LH)^2$$

→ - Where does this come from?

- What are the particles that lead to this → UV-completion



TYPE-II SEESAW

$$(Y = -1) \quad LL$$

can be combined into a triplet

$$\begin{pmatrix} \nu \\ e \end{pmatrix} \begin{pmatrix} \nu \\ e \end{pmatrix} = \begin{pmatrix} \nu\nu \\ (\nu e + e\nu) \frac{1}{\sqrt{2}} \\ ee \end{pmatrix}$$

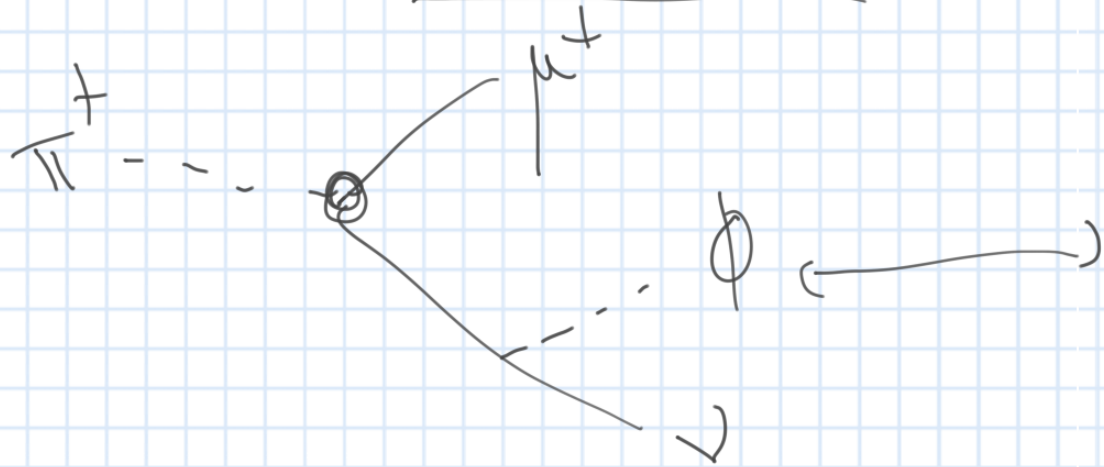
$$T = \begin{pmatrix} t^{++} \\ t^+ \\ t^0 \end{pmatrix}$$

$Y = +1$

$$\left. \begin{array}{l} \\ \\ \\ \end{array} \right\} \begin{array}{l} \boxed{Q_L(T) = -2} \\ \underline{LTL} \Rightarrow -2 \\ \langle T \rangle = v_t \ll v \\ \boxed{m_\nu = \lambda v_t} \end{array}$$

BREAK \hookrightarrow SPONTANEOUS \rightarrow GOLDSTONE BOSON

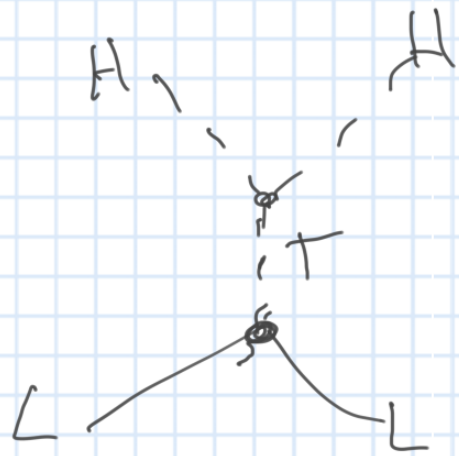
(MASSLESS
SCALAR)



[MAJONON]

$$m_T \gg m_H$$

$$U \supset \frac{\lambda}{4} \left(H^\dagger H \right)^2$$



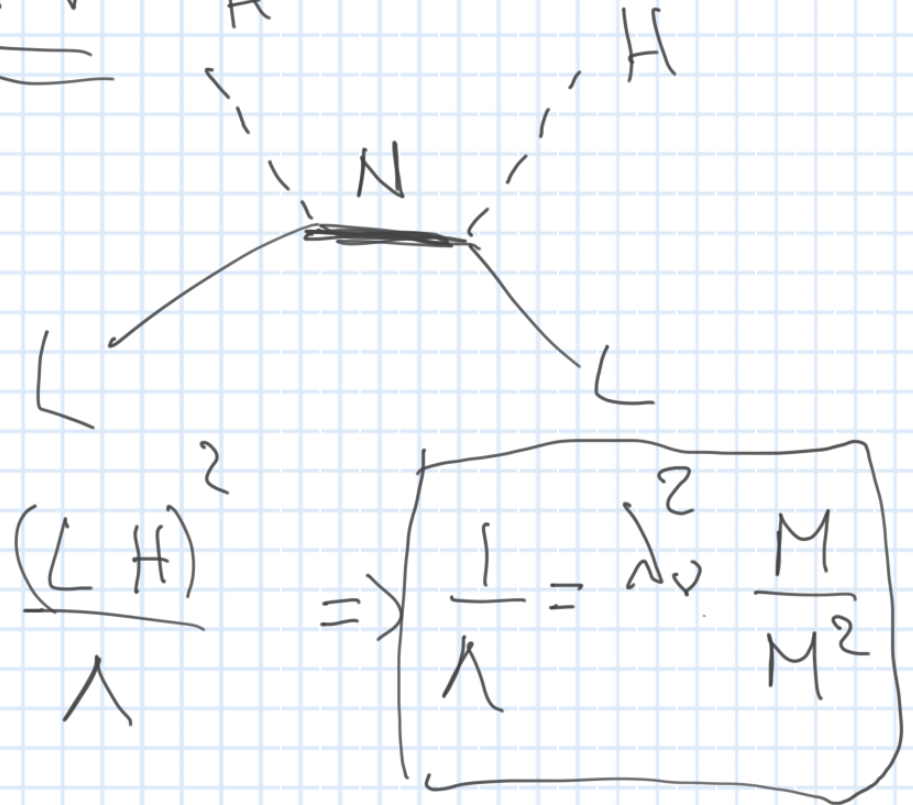
$$(LH)(LH) \frac{\lambda \kappa}{m_T^2}$$

$$\frac{1}{\Lambda} \sim \frac{\lambda \kappa}{m_T^2}$$

TYPE-I SEESAW

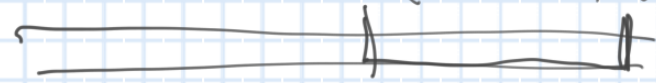
$$- \mathcal{L} \supset \lambda_\nu L \nu^c H + \frac{M}{2} \nu^e \nu^e H$$

$$M \gg m_H$$



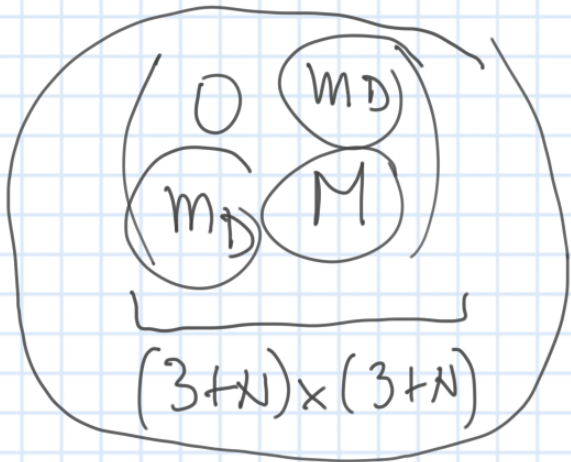
$$-L \supset \lambda_\nu \overline{LH\nu^c} + \frac{M}{2} \nu^c \nu^c$$

$$-L \supset m_D \nu \nu^c + \frac{M}{2} \nu^c \nu^c = \frac{1}{2} (\nu \nu^c) \begin{pmatrix} 0 & m_D \\ m_D & M \end{pmatrix} \begin{pmatrix} \nu \\ \nu^c \end{pmatrix}$$

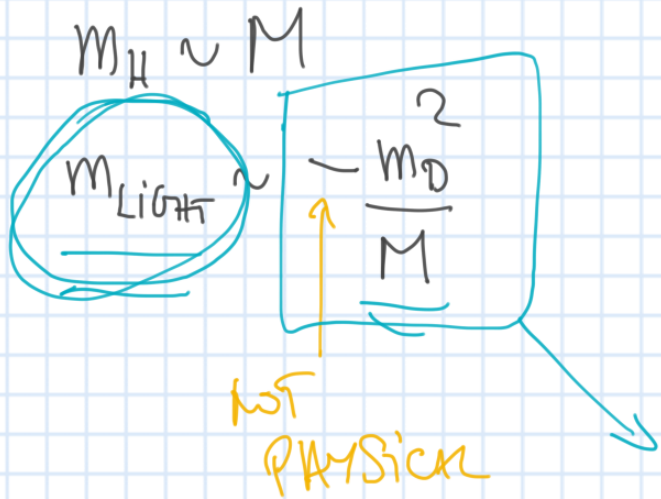


MASS
MATRIX

EXPLAINS NEUTRINO MASSES
MAJORANA, NEW FERMIONS!



$$M \gg m_D$$



$N = \# \text{ OF } \nu^c \text{-FIELDS } (N \geq 2)$

$$\theta_{\text{QS}}^2 \sim \frac{m_{\text{LIGHT}}}{m_H}$$

$$\begin{aligned} \cancel{\nu_H} &\sim \cancel{\nu^c} + \cancel{\epsilon \nu} \\ \cancel{\nu_{\text{LIGHT}}} &\sim \cancel{\nu} - \epsilon \cancel{\nu^c} \end{aligned}$$

$$\epsilon \sim \frac{m_D}{M}$$

$$\sim \sqrt{\frac{m_{\text{LIGHT}}}{m_H}}$$

SEESAW

$$\rightarrow \underline{\underline{M \ll m_D}}$$

$$m_{\pm} = \underline{\underline{m_D \pm M}}$$

PSEUDO-DIRAC REGIME

$$\Delta \underline{\underline{\Delta m^2 \sim 4Mm_D}}$$

$$\nu_+ = \frac{1}{\sqrt{2}} (\nu + \nu^c)$$

$$\nu_- = \frac{1}{\sqrt{2}} (\nu - \nu^c)$$

(SOLAR)

$$\lesssim 10^{-12} \text{ eV}^2$$

$$- \mathcal{L} \supset \underbrace{\lambda_\nu L \nu^c H}_m + \frac{M}{2} \nu^c \nu^c$$

λ_ν, M
DATA (NEUTRINO DATA)

