



# NEUTRINOS: MAJORANA OR DIRAC PARTICLES?

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International Neutrino Summer School 2021



# Neutrinos: Dirac or Majorana?

## Dirac

- ❑  $\nu \neq \bar{\nu}$
- ❑ LNC

## Majorana

- ❑  $\nu = \bar{\nu}$
- ❑ LNV

## Dirac vs. Majorana: Why should we care?

- ❑ Origin of neutrino masses and mixings
- ❑ If Majorana, possible explanation for the matter-antimatter asymmetry
- ❑ Neutrinos have been full of surprises.

## Why determining this is very challenging?

- ❑ When a neutrino is ultra-relativistic, its behavior is insensitive to its nature.

## How to address Dirac vs. Majorana question?

- ❑ Find exception to the property that a Majorana neutrino behaves almost like a Dirac one.

E.g., *kaon/pion decays to heavy neutrino.*

- ❑ Find an effective way of working with non-relativistic neutrinos.

E.g., *the relic neutrinos from the Big Bang.*

- ❑ Find some process that addresses the question even though it does not involve neutrinos at all.

E.g., *Neutrinoless double beta decay.*

# Neutrinos: Dirac or Majorana?

## The most promising approach

- ❑  $0\nu\beta\beta$ : observation of  $0\nu\beta\beta$  would imply Majorana nature of neutrino.

## Other approaches

- ❑ Capture the relic neutrinos from the Big Bang on tritium

*The total capture rate is two times larger in the Majorana case compared to that in the Dirac one.*

- ❑ (Heavy) neutrino decays:  $N \rightarrow \nu + X$

Angular distribution of X in the N's rest frame: *isotropic if Majorana, almost never isotropic if Dirac.*

Physics Letters B 789 (2019) 488-495

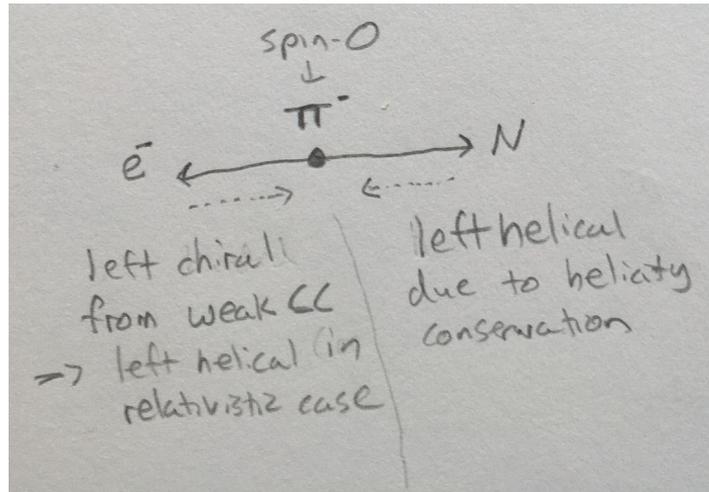
- ❑ Many more ... *E.g., LNV at LHC and ee collider.*

# Pion Decay to Heavy Neutrinos

<https://arxiv.org/abs/1805.00922>

Suppose there is a fourth neutrino mass state which has some coupling to the weak flavor states. Consider

$$\pi^- \rightarrow e^- \mathcal{N}$$
$$m_e \ll m_{\mathcal{N}} \ll m_{\pi}$$



Dirac  $\mathcal{N}$ :

- Lepton number conservation means interactions of the left-helical heavy antineutrino are suppressed

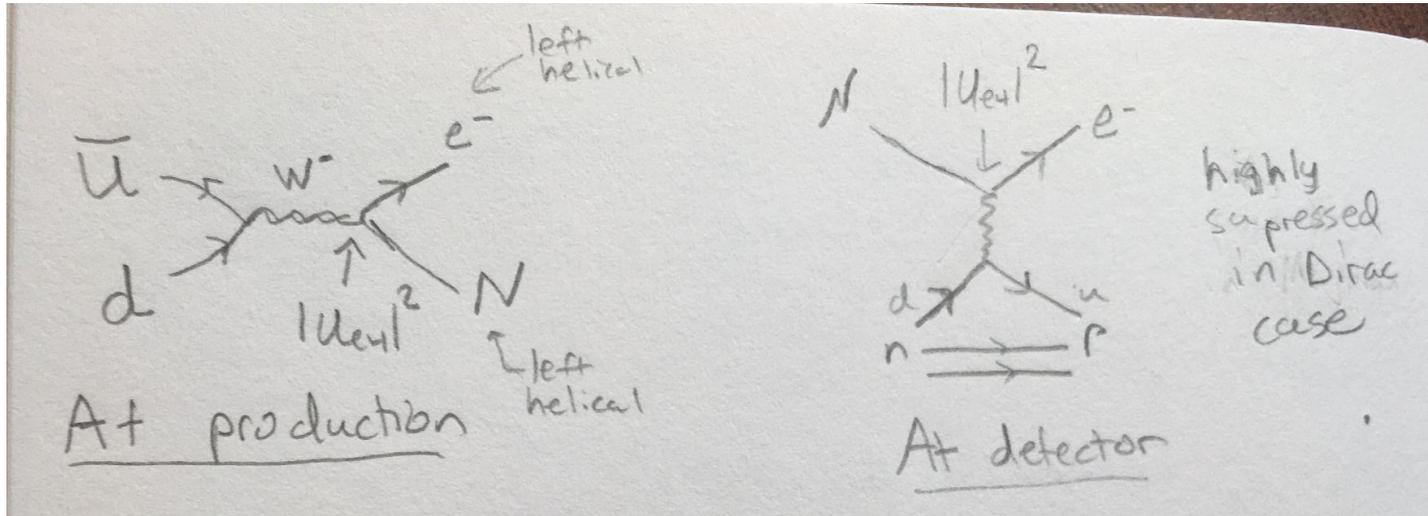
Majorana  $\mathcal{N}$ :

- Lepton number not a good symmetry
- left-helical  $\mathcal{N}$  can interact like a left-handed Dirac neutrino without helicity suppression

# Pion decay into heavy Neutrinos

If the  $N$  is Majorana in nature, we can look for an excess of charged-current quasi-elastic events from these left-handed  $N$  states in a  $\pi^-$  source

In a LArTPC these events would show up with a 1-lepton-1-proton topology, which cannot be mimicked by an antineutrino from the  $\pi^-$  source



Unfortunately suppressed by the fourth power of the new mixing

# Pion decay into heavy Neutrinos

[1] <https://arxiv.org/pdf/1905.00284v2.pdf>

[2] <https://arxiv.org/abs/1905.09694>

We can look for an excess of charged-current quasi-elastic events from these left-handed  $N$  states in an anti-neutrino beam

In a LArTPC these events would show up with a 1-lepton-1-proton topology, which cannot be mimicked by an antineutrino

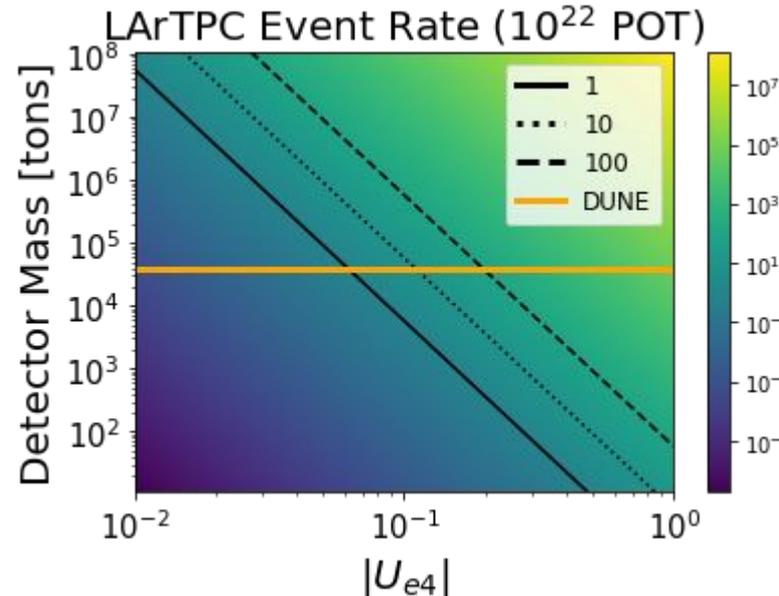
For an example flux, we consider the anti- $\bar{\nu}_e$  flux at the DUNE far detector [1]

$$\phi_{\bar{\nu}_e} \approx 10^{-12} \nu/\text{POT}/\text{cm}^2$$

$$\phi_{\mathcal{N}} = |U_{e4}|^2 \phi_{\bar{\nu}_e}$$

The CCQE cross section on LAr is [2]

$$\sigma_{\nu_e}^{\text{CCQE}} \approx 10^{-38} \text{cm}^2 \quad \sigma_{\mathcal{N}}^{\text{CCQE}} = |U_{e4}|^2 \sigma_{\bar{\nu}_e}^{\text{CCQE}}$$



# Neutrino Magnetic Moment

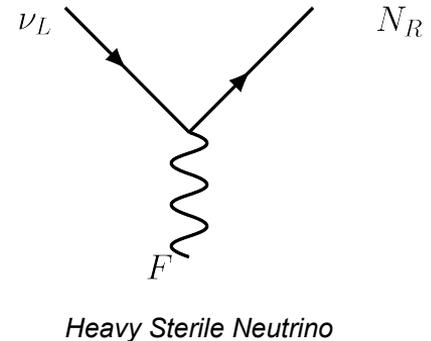
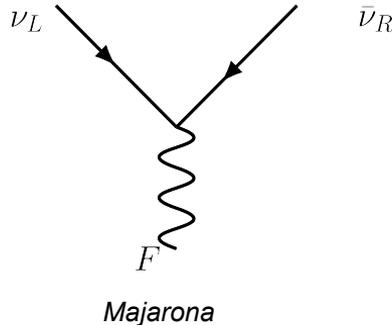
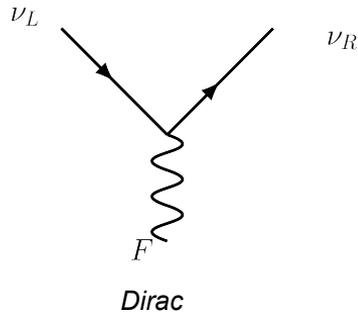
- The Lagrangian describing neutrino magnetic moments is:

$$\mathcal{L} = \mu_{\nu}^{ij} (\nu_i \sigma_{\mu\nu} \nu_j F^{\mu\nu}) + h.c. \quad (\text{Majorana}), \quad \text{or} \quad \mathcal{L} = \mu_{\nu}^{ij} (\bar{\nu}_i \sigma_{\mu\nu} \nu_j F^{\mu\nu}) + h.c. \quad (\text{Dirac}),$$

In the case of Majorana neutrinos,  $\mu_{\nu}^{ij}$  is constrained to be antisymmetric under CPT invariance. Whereas Dirac neutrino magnetic moments can have diagonal or transition moments.

Standard model predicts  $O[10^{-19}]$  Bohr Magnetron.

The final state neutrino is sterile in the Dirac case. It is possible to write a lepton number violating transition magnetic moment in the Majorana case. We can in principle detect these transition moments.



# Detecting large neutrino magnetic moments.

We can look at regions of high magnetic fields to .

1. Solar neutrinos could undergo spin flavour precession. This don't see the appropriate antineutrinos.
2. Supernova neutrino oscillations could be affected

Interestingly, mass of neutrino forces stronger constraints on Dirac moments.

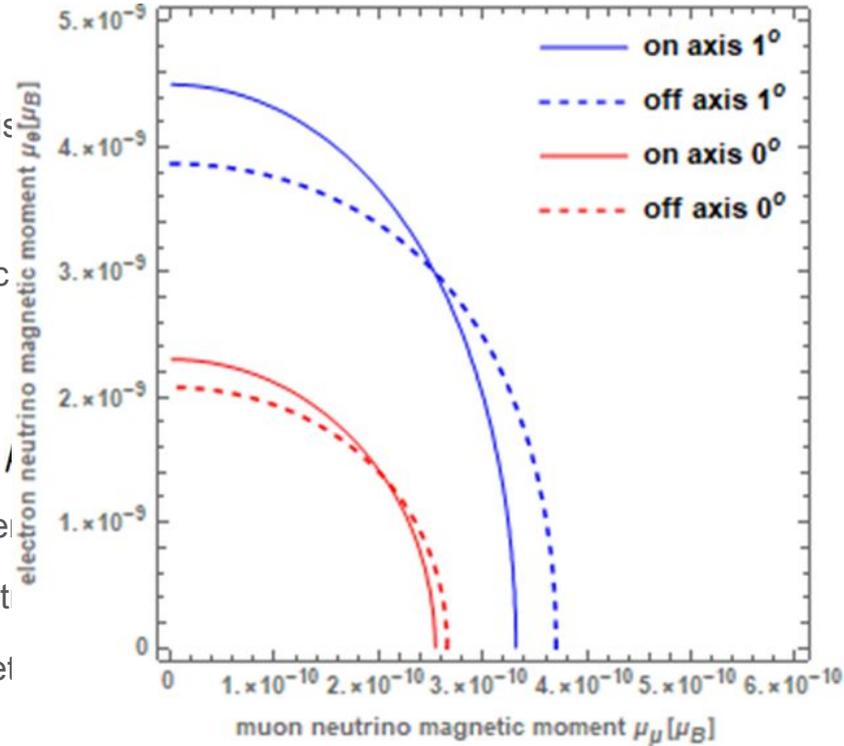
Neutrino electron scattering:

$$\frac{d\sigma}{dT}(\nu_j e \rightarrow \nu_i e)_{em} = \mu^2 \frac{\pi\alpha^2}{E_\nu m_e^2} \left( \frac{E_\nu}{T} - 1 \right).$$

We can look for low energy recoils at dark matter detectors like XENON.

Best terrestrial constraints from solar neutrinos, but they can't constrain large magnetic moments.

We can also place constraints on muon and electron flavor magnetic moments.



# Neutrino Magnetic Moment

Neutrino masses imply that neutrinos have non-zero electromagnetic dipole moments. The nature of  $\mu_{\nu}$  will depend on whether the neutrinos are Majorana or Dirac fermions.

$$\mathcal{L} = \mu_{\nu}^{ij} (\nu_i \sigma_{\mu\nu} \nu_j F^{\mu\nu}) + h.c. \quad (\text{Majorana})$$

*anti-symmetric by CPT invariance  $\square \nu$  and  $\bar{\nu}$   
must have magnetic moments of equal magnitude*

$$\mathcal{L} = \mu_{\nu}^{ij} (\bar{\nu}_i \sigma_{\mu\nu} \nu_j F^{\mu\nu}) + h.c. \quad (\text{Dirac})$$

Majorana have only transition moments,  
Dirac have diagonal and transition moments.

Effective neutrino magnetic moment measured by an experiment:

$$\mu_{\nu}^2(\nu_{\alpha}, L, E) = \sum_j \sum_{kk'} U_{\alpha k}^* U_{\alpha k'} e^{-i\Delta m_{kk'}^2 L/2E} \mu_{jk} \mu_{jk'}$$

	Dirac	Majorana
$\mu_{ii}$	$\frac{3eG_F m_i}{8\sqrt{2}\pi^2} \approx 3.2 \times 10^{-19} \left(\frac{m_i}{1 \text{ eV}}\right) \mu_B$	0
$\mu_{ij}, i \neq j$	$\sim 4 \times 10^{-23} \mu_B$	$\sim 2 \times \mu_{ij}^D$

*Measuring a magnetic moment different than what predicted for Dirac neutrinos would indicate BSM physics!*

*But it would be tricky to distinguish between Majorana and Dirac, since it is difficult to distinguish between diagonal and transition moments.*

# Neutrino Magnetic Moment

## Triangle Inequalities for Majorana-Neutrino Magnetic Moments

Example: muon neutrino beam interacting with a target.

$$\nu_\mu e^- \rightarrow \nu_X e^-$$

*Final flavour not known*

Results in an “effective” magnetic moment that considers the superposition of flavour states for the outgoing neutrino:

$$|\mu_{\nu_\mu}| \equiv \sqrt{|\mu_{e\mu}|^2 + |\mu_{\tau\mu}|^2}$$

$$\sqrt{|\mu_{e\mu}|^2 + |\mu_{\tau\mu}|^2} (\bar{\nu}_X^c \sigma_{\alpha\beta} \nu_\mu F^{\alpha\beta}),$$

where

$$\bar{\nu}_X^c \equiv \frac{(\mu_{e\mu} \bar{\nu}_e^c + \mu_{\tau\mu} \bar{\nu}_\tau^c)}{\sqrt{|\mu_{e\mu}|^2 + |\mu_{\tau\mu}|^2}}.$$

*Which comes from the vertex*

If  $|\mu_{\nu_\tau}|^2 \leq |\mu_{\nu_e}|^2 + |\mu_{\nu_\mu}|^2$ , neutrinos are Majorana.

*Due to the antisymmetry of  $\mu_{ij}$  in the Majorana case.*

Practically, the most promising approach to checking these inequalities is to improve the limits on  $\mu_{\nu_\tau}$ , which is possible at the proposed SHiP experiment at CERN.

Upper limits from accelerator experiments:

$$\mu_{\nu_e}^{\text{eff}} < 1.1 \cdot 10^{-9} \mu_B \quad \text{LAMPF}$$

$$\mu_{\nu_\mu}^{\text{eff}} < 6.8 \cdot 10^{-10} \mu_B \quad \text{LSND}$$

$$\mu_{\nu_\tau}^{\text{eff}} < 3.9 \cdot 10^{-7} \mu_B \quad \text{DONUT}$$

# Summary and Outlook

Neutrinoless double beta decay is, currently, the most promising way to probe the nature of neutrinos - Dirac or Majorana.

There are many other ways to probe Majorana vs. Dirac nature such as look for interactions/decays with heavy neutrinos, measure the cosmic neutrino background, measure the (tau-)neutrino magnetic moment, etc.

- Need experiments with impressive energy resolution, better sensitivities, large statistics, etc...