

# Heavy flavour aspects in EFTs

Model-independent bounds on the  
Standard Model Effective Theory  
from Flavour Physics

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**Based on:** [Phys.Lett.B 799 \(2019\) 135062](#) (arXiv:1812.10913)

In collaboration with **Luca Silvestrini**

# UTA : A PRECISION TEST OF THE SM

From Unitarity of the CKM

$$(V^\dagger V)_{db} = 0 \Leftrightarrow \text{apex: } (\bar{\rho}, \bar{\eta})$$

Over-constrained global fit

$$\Delta m_{d,s} \Leftrightarrow B-\bar{B} \text{ mixing}$$

$$\epsilon_K \Leftrightarrow K-\bar{K} \text{ mixing}$$

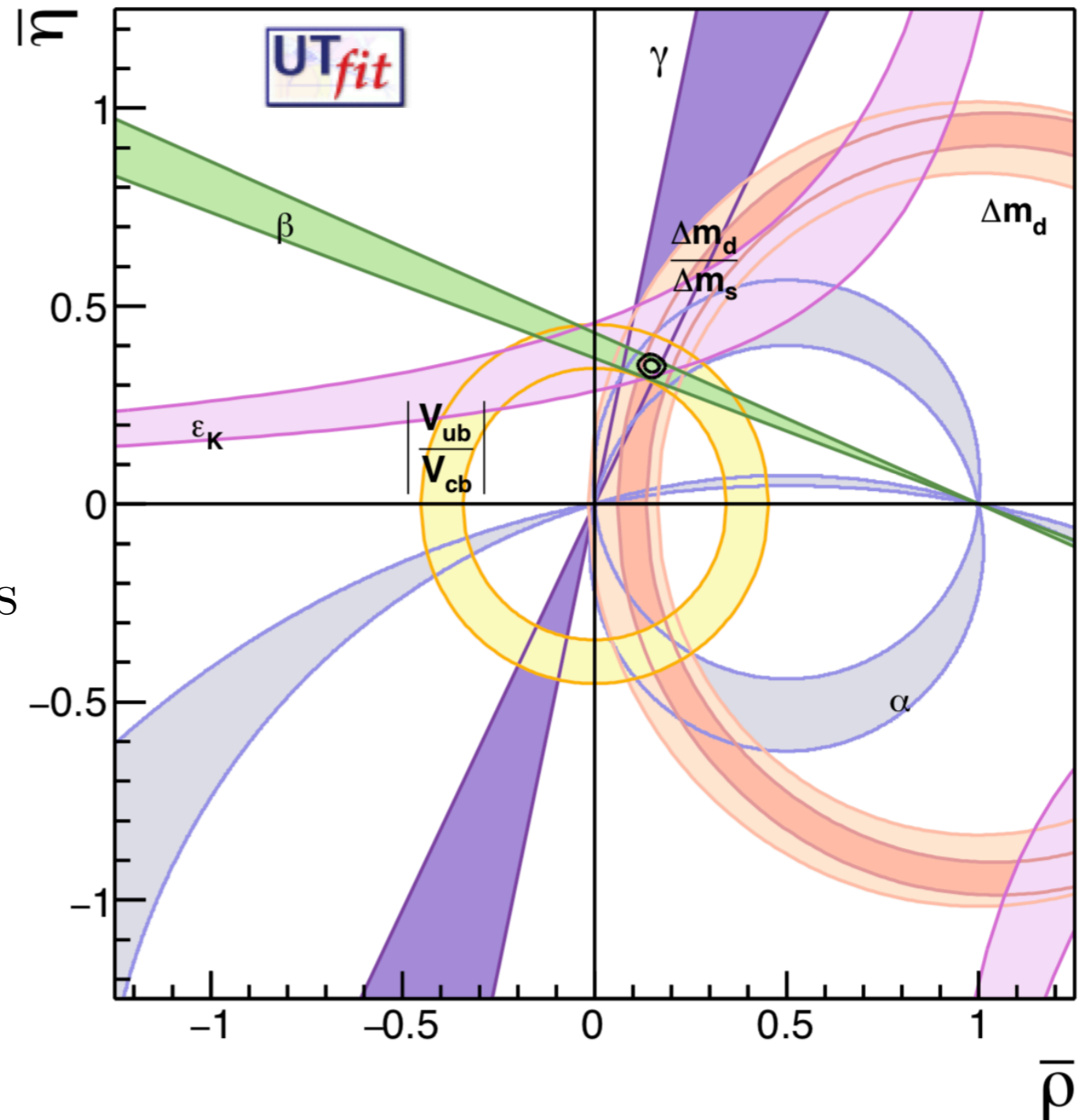
$$|V_{ub}/V_{cb}| \Leftrightarrow \text{semileptonic } B \text{ decays}$$

$$\alpha, \beta, \gamma \Leftrightarrow \text{hadronic } B \text{ decays}$$

$$\bar{\rho} = 0.148 \pm 0.013$$

$$\bar{\eta} = 0.348 \pm 0.010$$

<http://www.utfit.org/UTfit>



# WHAT ABOUT NEW PHYSICS ?

No tree-level Flavor-Changing-Neutral-Current processes (FCNCs) in the SM.

Furthermore, suppression of FCNC amplitudes in the SM due to GIM.

➡  $\Delta F = 2$  : excellent probe of Physics Beyond the SM!

The most general Weak Effective Hamiltonian for meson anti-meson mixing is:

[see, e.g., *M.Bona et al. JHEP 03 (2008) 049* ]

$$\mathcal{H}_{\text{eff}}^{\Delta S=2} = \sum_{i=1}^5 C_i O_i^{sd} + \sum_{i=1}^3 \tilde{C}_i \tilde{O}_i^{sd}$$

$$\mathcal{H}_{\text{eff}}^{\Delta C=2} = \sum_{i=1}^5 C_i O_i^{cu} + \sum_{i=1}^3 \tilde{C}_i \tilde{O}_i^{cu}$$

$$\mathcal{H}_{\text{eff}}^{\Delta B=2} = \sum_{i=1}^5 C_i O_i^{bq} + \sum_{i=1}^3 \tilde{C}_i \tilde{O}_i^{bq}$$

where:  $\tilde{O}_i \equiv O_i(L \leftrightarrow R)$

$$O_1^{q_i q_j} = \bar{q}_{jL}^\alpha \gamma_\mu q_{iL}^\alpha \bar{q}_{jL}^\beta \gamma^\mu q_{iL}^\beta \quad \text{SM}$$

$$O_2^{q_i q_j} = \bar{q}_{jR}^\alpha q_{iL}^\alpha \bar{q}_{jR}^\beta q_{iL}^\beta$$

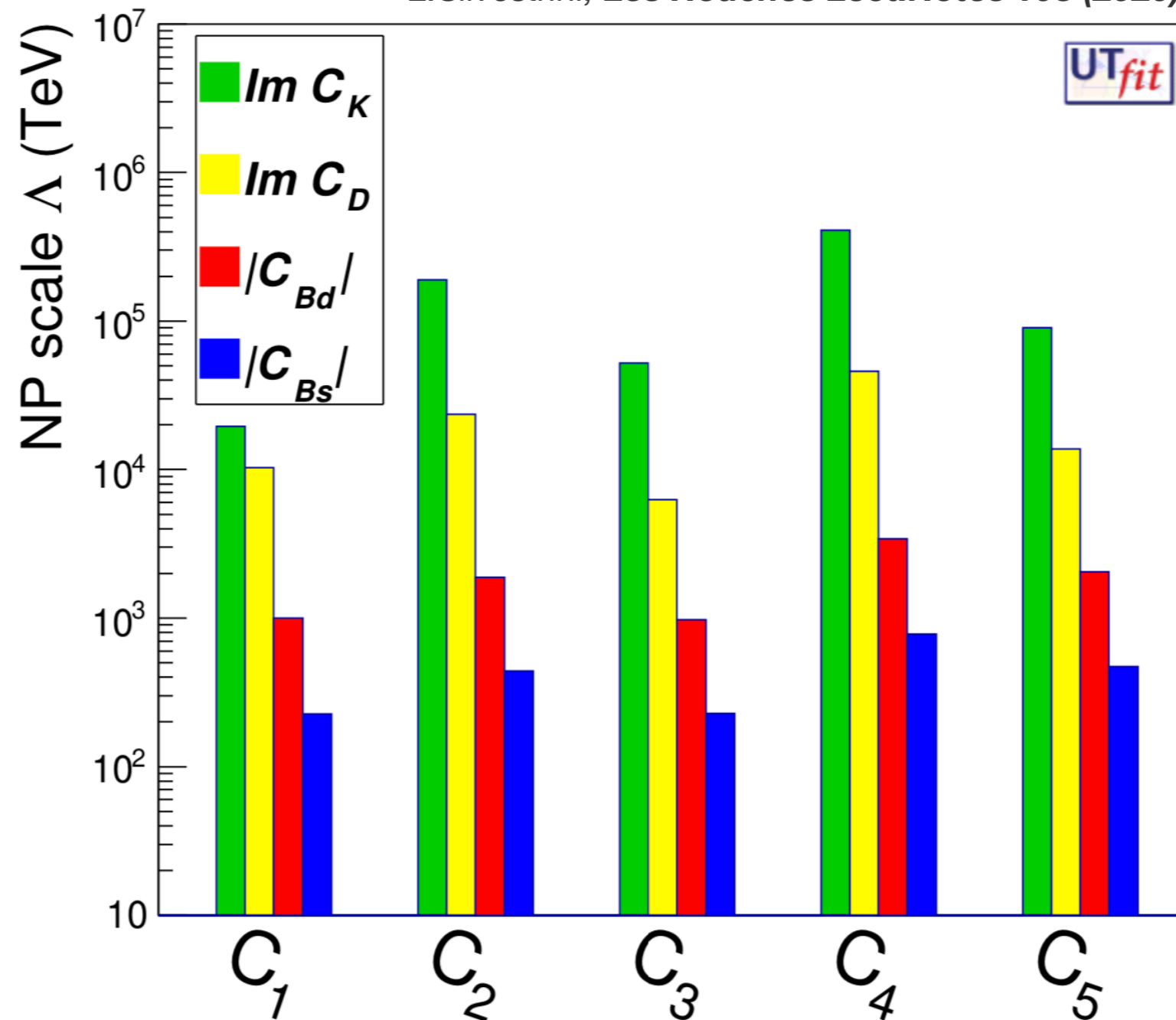
$$O_3^{q_i q_j} = \bar{q}_{jR}^\alpha q_{iL}^\beta \bar{q}_{jR}^\beta q_{iL}^\alpha$$

$$O_4^{q_i q_j} = \bar{q}_{jR}^\alpha q_{iL}^\alpha \bar{q}_{jL}^\beta q_{iR}^\beta$$

$$O_5^{q_i q_j} = \bar{q}_{jR}^\alpha q_{iL}^\beta \bar{q}_{jL}^\beta q_{iR}^\alpha$$

# AMOUNT OF NP ALLOWED IN $\Delta F = 2$

L.Silvestrini, *Les Houches Lect.Notes 108 (2020)*

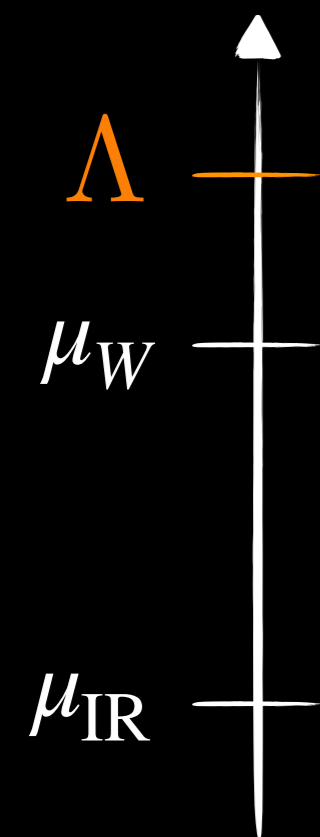


New generic source of flavor / CP violation  $\rightarrow$  high NP scale

# $\Delta F = 2$ UNDER $SU(2)_L \times U(1)_Y$

ABOVE THE EW SCALE, WE MUST EXPLOIT SM GAUGE SYMMETRY

Energy



## $\Delta F = 2$ IN THE SMEFT

$$O_{ijij}^{QQ(1)} = \bar{Q}_i \gamma_\mu Q_j \bar{Q}_i \gamma^\mu Q_j$$

$$O_{ijij}^{QQ(3)} = \bar{Q}_i \gamma_\mu \tau^A Q_j \bar{Q}_i \gamma^\mu \tau^A Q_j$$

$$O_{ijij}^{qq} = \bar{q}_i \gamma_\mu q_j \bar{q}_i \gamma^\mu q_j$$

$$O_{ijij}^{Qq(1)} = \bar{Q}_i \gamma_\mu Q_j \bar{q}_i \gamma^\mu q_j$$

$$O_{ijij}^{Qq(8)} = \bar{Q}_i \gamma_\mu T^a Q_j \bar{q}_i \gamma^\mu T^a q_j$$

CAPITAL CASE

$SU(2)_L$  Doublets

LOWER CASE

$SU(2)_L$  singlets

See, e.g.:

*J. Aebischer et al.*  
1512.02830

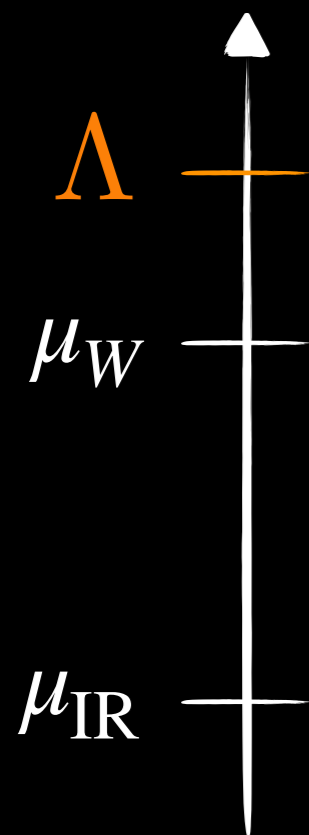
*A. Celis et al.*  
1704.04504

( $q = u, d$  stands now on for RH fields ;  $\tau^A$  Pauli matrices ;  $T^a$  color ones)

# $\Delta F = 2$ UNDER $SU(2)_L \times U(1)_Y$

ABOVE THE EW SCALE, WE MUST EXPLOIT SM GAUGE SYMMETRY

Energy



## $\Delta F = 2$ MATCHING IN THE SMEFT

$$C_1(\mu_W) = - \left( C^{QQ^{(1)}}(\mu_W) + C^{QQ^{(3)}}(\mu_W) \right) / \Lambda^2$$

$$\tilde{C}_1(\mu_W) = -C^{qq}(\mu_W) / \Lambda^2$$

$$C_4(\mu_W) = C^{Qq^{(8)}}(\mu_W) / \Lambda^2$$

$$C_5(\mu_W) = \left( 6C^{Qq^{(1)}}(\mu_W) - C^{Qq^{(8)}}(\mu_W) \right) / (3\Lambda^2)$$

$$\Rightarrow C_2(\mu_W) = \tilde{C}_2(\mu_W) = C_3(\mu_W) = \tilde{C}_3(\mu_W) = 0$$

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1512.02830

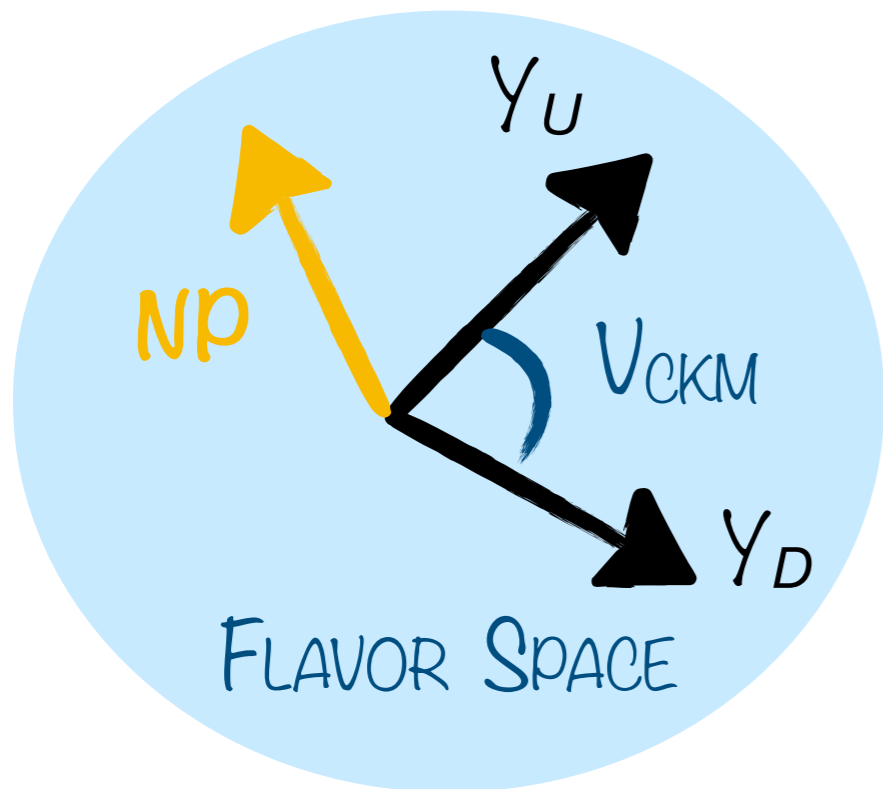
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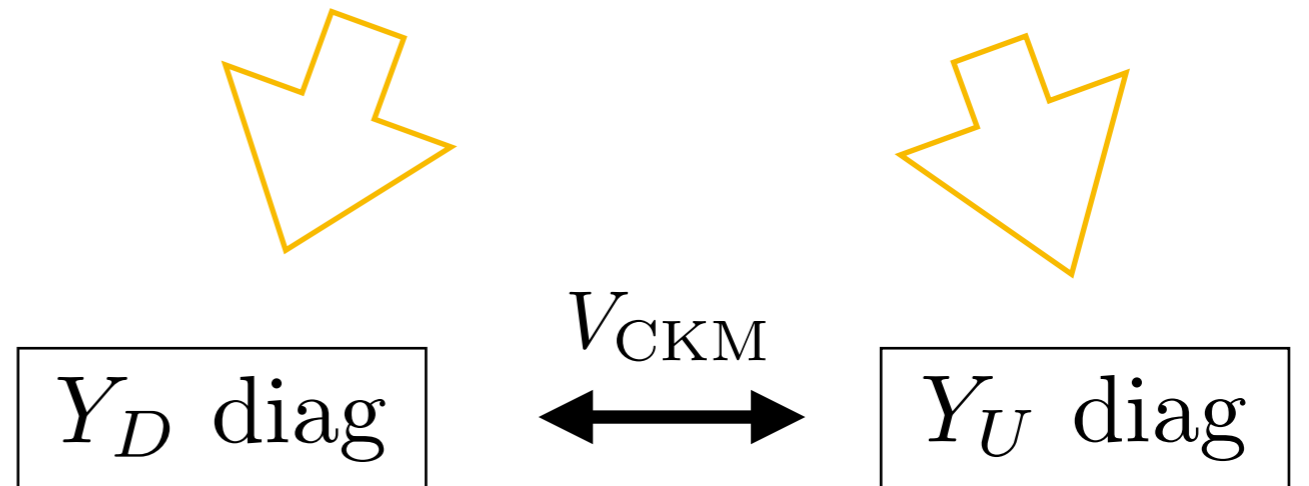
# $\Delta F = 2$ BOUNDS — A modern view

$$U(3)_Q \times U(3)_u \times U(3)_d$$



**Q:** In which **basis** we are defining NP?

$$\mathcal{L}_{SMEFT} = \mathcal{L}_{SM} + \sum_{i,d>4} \frac{C_i \mathcal{O}_i^{(d)}}{\Lambda_{NP}^{d-4}}$$



*Basis where down-quark Yukawa matrix is diagonal*

*Basis where up-quark Yukawa matrix is diagonal*

Orientation in Flavor space imprints NP phenomenology: *2 extremes at hand.*

**Important point, since in the SMEFT up and down sectors are correlated!**



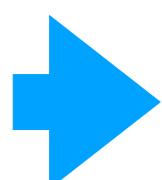
# $\Delta F = 2$ BOUNDS — A modern view

Let's take for instance:  $O_{1111}^{QQ} = (\bar{Q}_1 \gamma_\mu Q_1)^2$

→  $\Delta F = 0$ , still subject to Flavor constraints!

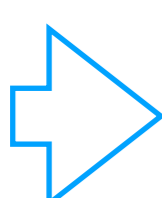
- 1) If aligned with  $Y_U$ :  $d_{L,i} \rightarrow V_{1j} d_{L,j} \rightarrow K - \bar{K}$
- 2) If aligned with  $Y_D$ :  $u_{1,L} \rightarrow V_{1j}^\dagger u_{L,j} \rightarrow D - \bar{D}$

*going to mass basis*

  $\Lambda_{\text{NP}}^{QQ1111} \gtrsim$

- 1) 415 TeV
- 2) 267 TeV

Note that misalignment of NP in Flavor space is not relevant for fully right-handed quark operators.

  $O_{ijkl}^{uu} = \bar{u}_i \gamma_\mu u_j \bar{u}_k \gamma^\mu u_l$   
 well constrained only in 1212  
*(independent of flavor alignment)*

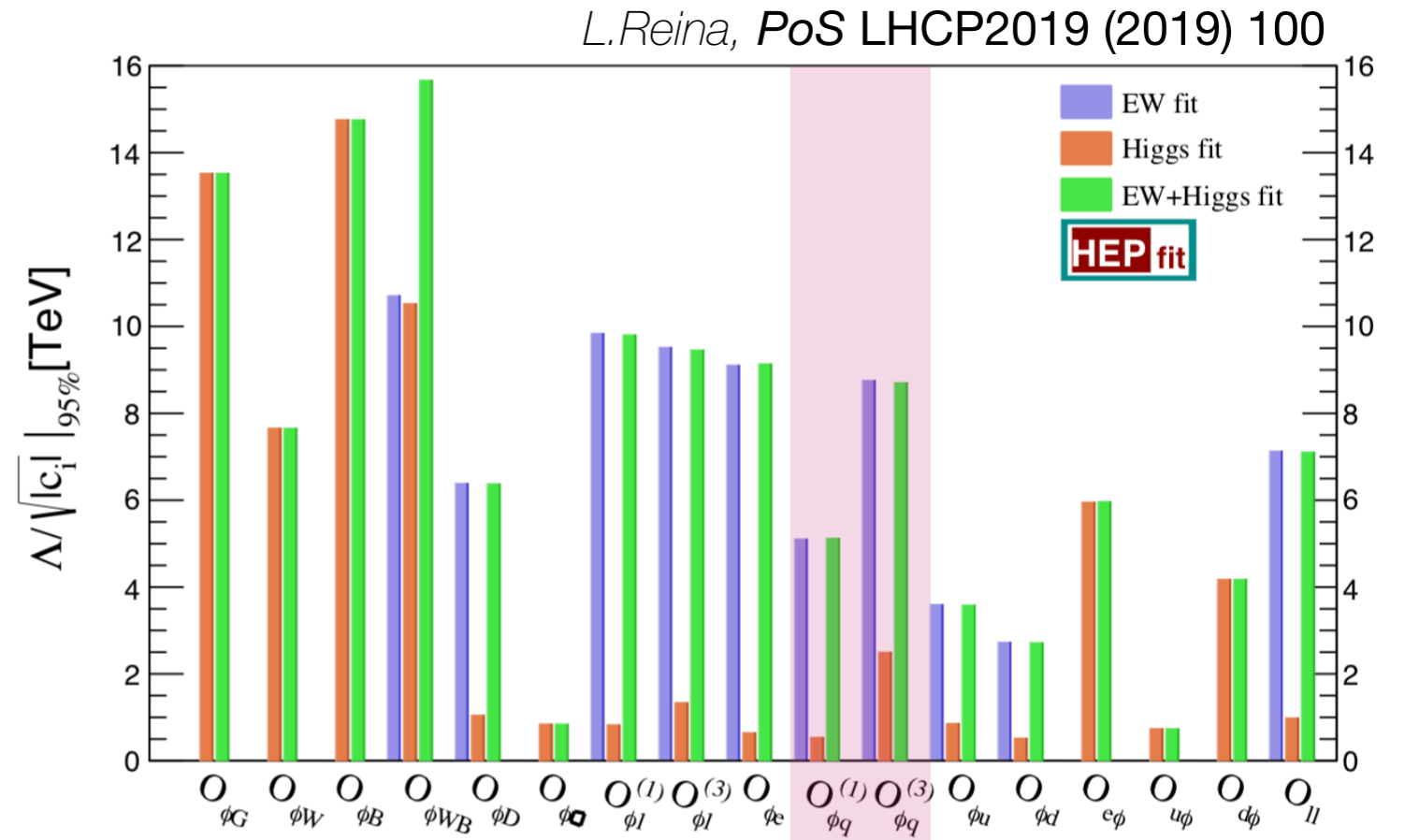
$ijkl$	$C_{ijkl}^{QQ(1,3)}$ [TeV <sup>-2</sup> ]		$C_{ijkl}^{dd}$ [TeV <sup>-2</sup> ]
	$Y_D$ diag	$Y_U$ diag	$Y_{D,U}$ diag
1111	$5.8^\diamond 10^{-6}$	$1.4^\square 10^{-5}$	$\emptyset$
1112	$(7.0^\diamond, 0.19^\diamond) 10^{-7}$	$(17^\square, 0.051^\square) 10^{-7}$	$(3.2^\square, 1.3^\square) 10^{-3}$
1113	$(15^\diamond, 0.44^\diamond) 10^{-6}$	$(39^\Delta, 0.12^\square) 10^{-6}$	$(1.4^\Delta, 1.2^\Delta)$
1122	$2.9^\diamond 10^{-6}$	$6.8^\square 10^{-6}$	$\emptyset$
1123	$(5.6^\diamond, 2.3^\diamond) 10^{-6}$	$(1.5^\square, 3.4^\square) 10^{-6}$	$(\emptyset, \emptyset)$
1133	$1.3^\diamond 10^{-4}$	$3.6^\square 10^{-5}$	$\emptyset$
1212	$(31^\diamond, 0.22^\square) 10^{-8}$	$(28^\diamond, 0.25^\square) 10^{-8}$	$(65^\square, 0.22^\square) 10^{-8}$
1213	$(3.5^\diamond, 0.1^\diamond) 10^{-6}$	$(15^\square, 0.3^\square) 10^{-7}$	$(16^\square, 0.3^\square) 10^{-4}$
1221	$2.9^\diamond 10^{-6}$	$6.8^\square 10^{-6}$	$\emptyset$
1222	$(7.0^\diamond, 0.19^\diamond) 10^{-7}$	$(17^\square, 0.05^\square) 10^{-7}$	$(3.2^\square, 1.3^\square) 10^{-3}$
1223	$(15^\diamond, 0.44^\diamond) 10^{-6}$	$(39^\diamond, 0.12^\diamond) 10^{-6}$	$(\emptyset, \emptyset)$
1231	$(5.6^\diamond, 2.3^\diamond) 10^{-6}$	$(1.5^\square, 3.4^\square) 10^{-6}$	$(\emptyset, \emptyset)$
1232	$(1.3^\diamond, 2.2^\diamond) 10^{-6}$	$(3.7^\square, 1.4^\square) 10^{-7}$	$(7.2^\square, 2.9^\square) 10^{-3}$
1233	$(3.1^\diamond, 6.2^\diamond) 10^{-5}$	$(8.9^\square, 3.4^\square) 10^{-6}$	$(\emptyset, \emptyset)$
1313	$(1.1^\Delta, 0.90^\Delta) 10^{-6}$	$(1.1^\Delta, 0.95^\Delta) 10^{-6}$	$(1.0^\Delta, 0.88^\Delta) 10^{-6}$
1322	$(15^\diamond, 0.44^\diamond) 10^{-6}$	$(39^\diamond, 0.12^\diamond) 10^{-6}$	$(\emptyset, \emptyset)$
1323	$(3.3^\diamond, 0.1^\diamond) 10^{-4}$	$(2.4^\Delta, 2.1^\Delta) 10^{-6}$	$(6.3^\Delta, 5.3^\Delta) 10^{-1}$
1331	$1.3^\diamond 10^{-4}$	$3.6^\square 10^{-5}$	$\emptyset$
1332	$(3.1^\diamond, 6.2^\diamond) 10^{-5}$	$(8.9^\square, 3.4^\square) 10^{-6}$	$(\emptyset, \emptyset)$
1333	$(6.7^\diamond, 15^\diamond) 10^{-4}$	$(6.9^\Delta, 5.8^\Delta) 10^{-5}$	$(1.4^\Delta, 1.2^\Delta)$
2222	$5.8^\diamond 10^{-6}$	$1.4^\square 10^{-5}$	$\emptyset$
2223	$(5.7^\diamond, 2.3^\diamond) 10^{-6}$	$(1.5^\square, 3.4^\square) 10^{-6}$	$(3.0^\nabla, 0.98^\nabla) 10^{-1}$
2233	$1.3^\diamond 10^{-4}$	$3.6^\square 10^{-5}$	$\emptyset$
2323	$(2.2^\nabla, 0.75^\nabla) 10^{-5}$	$(2.1^\Delta, 0.80^\Delta) 10^{-5}$	$(2.2^\nabla, 0.74^\nabla) 10^{-5}$
2332	$1.3^\diamond 10^{-4}$	$3.6^\square 10^{-5}$	$\emptyset$
2333	$(2.2^\diamond, 2.8^\diamond) 10^{-3}$	$(2.8^\nabla, 0.94^\nabla) 10^{-4}$	$(3.0^\nabla, 0.98^\nabla) 10^{-1}$
3333	$4.2^\diamond 10^{-1}$	$1.3^\nabla 10^{-2}$	$\emptyset$



# $\Delta F = 2$ BOUNDS — RGE driven

$ij$	$C_{ij}^{HQ^{(1,3)}} [\text{TeV}^{-2}]$	
	$Y_D$ diag	$Y_U$ diag
11	$\emptyset$	$4.1^\square 10^{-3}$
12	$(8.9^\square, 3.8^\square) 10^{-4}$	$(9.9^\square, 3.8^\square) 10^{-4}$
13	$(7.4^\Delta, 6.3^\Delta) 10^{-3}$	$(7.6^\Delta, 6.4^\Delta) 10^{-3}$
22	$\emptyset$	$4.1^\square 10^{-3}$
23	$(3.0^\nabla, 1.0^\nabla) 10^{-2}$	$(3.1^\nabla, 1.0^\nabla) 10^{-2}$
33	$\emptyset$	$7.3^\Delta 10^{-1}$

$\Lambda_{\text{NP}} \gtrsim 15 \text{ TeV}$



$ijkl$	$C_{ijkl}^{LeQu} [\text{TeV}^{-2}]$	$C_{ijkl}^{LedQ} [\text{TeV}^{-2}]$
	$Y_D$ diag	$Y_U$ diag
2221	$(5.1^\diamond, 1.6^\diamond) 10^{-1}$	$(4.2^\square, 0.13^\square) 10^{-1}$
2222	$(22^\diamond, 6.8^\diamond) 10^{-1}$	$(18^\square, 0.58^\square) 10^{-1}$
2223	$(\emptyset, \emptyset)$	$(4.3^\square, 1.6^\square)$
3321	$(3.0^\diamond, 0.93^\diamond) 10^{-2}$	$(24^\square, 0.8^\square) 10^{-3}$
3322	$(1.3^\diamond, 0.4^\diamond) 10^{-1}$	$(10^\square, 0.34^\square) 10^{-2}$
3323	$(3.1^\diamond, 3.6^\diamond)$	$(2.5^\square, 0.9^\square) 10^{-1}$
3331	$(\emptyset, 9.5^\diamond)$	$(8.5^\Delta, 11^\Delta)$
3332	$(\emptyset, \emptyset)$	$(\emptyset, 8.9^\nabla)$

$$O_{ijkl}^{LeQu} \equiv \bar{L}_i e_j \epsilon \bar{Q}_k u_l$$

mixes into  $\Delta F = 2$  operators in the up sector via  $Y_U Y_L \rightarrow$  constraints from D-D $\bar{}$  mixing only

$$O_{ijkl}^{LedQ} \equiv \bar{L}_i e_j \bar{d}_k Q_l$$

mixes into  $\Delta F = 2$  operators in the down sector via  $Y_D Y_L \rightarrow$  constraints from K-K $\bar{}$  & B-B $\bar{}$  mixing only



# SMEFT: A FLAVORFUL SUMMARY

*Phys.Lett.B 799 (2019) 135062 (arXiv:1812.10913)*

## SMEFT RGE

$O_{jk}^{HQ(1[3])}$ $(H^\dagger i \overleftrightarrow{D}_\mu [A] H) (\bar{Q}_j \gamma^\mu [\tau^A] Q_k)$	$O_{jjkl}^{LedQ}$ $(\bar{L}_j e_j) (\bar{d}_k Q_l)$	$O_{jjkl}^{LeQu}$ $(\bar{L}_j e_j) i\tau^2 (\bar{Q}_k u_l)$	$O_{jklm}^{ud(1[8])}$ $(\bar{u}_j \gamma_\mu [T^a] u_k) (\bar{d}_l \gamma^\mu [T^a] d_m)$	$O_{jklm}^{QuQd(1[8])}$ $(\bar{Q}_j \gamma_\mu [T^a] u_k) i\tau^2 (\bar{Q}_l \gamma^\mu [T^a] d_m)$
$O_{jklm}^{QQ(1[3])}$ $(\bar{Q}_j \gamma_\mu [\tau^A] Q_k) (\bar{Q}_l \gamma^\mu [\tau^A] Q_m)$	$O_{jklm}^{uu}$ $(\bar{u}_j \gamma_\mu u_k) (\bar{u}_l \gamma^\mu u_m)$	$O_{jklm}^{dd}$ $(\bar{d}_j \gamma_\mu d_k) (\bar{d}_l \gamma^\mu d_m)$	$O_{jklm}^{Qd(1[8])}$ $(\bar{Q}_j \gamma_\mu [T^a] Q_k) (\bar{d}_l \gamma^\mu [T^a] d_m)$	$O_{jklm}^{Qu(1[8])}$ $(\bar{Q}_j \gamma_\mu [T^a] Q_k) (\bar{u}_l \gamma^\mu [T^a] u_m)$

*poorly constrained*

FLAVOR MISALIGNMENT

Interesting future directions ahead:

—> **IMPACT OF OTHER FLAVOR MEASUREMENTS** (e.g.,  $\Delta F = 1$  FCNCs)

*See recent work in: 2003.05432 (R.Aoude et al.) , 2005.05366 (D.A.Faroughy et al.) ,  
2009.07276 (J.Aebischer et al.) , 2101.07273 (S.Bruggisser et al.)*

—> **INTERPLAY WITH DI-BOSON SEARCHES / EW PRECISION / HIGGS / TOP PHYSICS**

***See previous talks by David Marzocca and by Gudrun Hiller!***

**BACK-UP**

# $\Delta F = 2$ BOUNDS — A modern view

Leading-log RGE effects may be also relevant ...

Let's compare bounds in the SMEFT on  $\Delta F = 1$  operators that run into  $\Delta F = 2$  via RGE:

$$\mathcal{A}_{\text{NP}}^{\Delta F=2} \lesssim \mathcal{A}_{\text{SM}}^{\Delta F=2} \varepsilon_{\Delta F=2} \quad \Bigg| \quad \mathcal{A}_{\text{NP}}^{\Delta F=1} \lesssim \mathcal{A}_{\text{SM}}^{\Delta F=1} \varepsilon_{\Delta F=1}$$

$\mathcal{A}$  being short-distance amplitude,  $\varepsilon_{\Delta F=i}$  amount of NP allowed in  $\Delta F = i$

$$\mathcal{A}_{\text{SM}}^{\Delta F=2} \sim \frac{(Y_{\text{SM}} Y_{\text{SM}}^\dagger)^2}{\Lambda_{\text{EW}}^2}, \quad \mathcal{A}_{\text{NP}}^{\Delta F=2} \sim \frac{C_{\text{NP}}^{\Delta F=2}}{\Lambda_{\text{NP}}^2} \quad \Bigg| \quad \mathcal{A}_{\text{SM}}^{\Delta F=1} \sim \frac{Y_{\text{SM}} Y_{\text{SM}}^\dagger}{\Lambda_{\text{EW}}^2}, \quad \mathcal{A}_{\text{NP}}^{\Delta F=1} \sim \frac{C_{\text{NP}}^{\Delta F=1}}{\Lambda_{\text{NP}}^2}$$

$\Delta F = 2$  BOUND ON NEW PHYSICS SCALE IS STRONGER IFF:

$$\frac{C_{\text{NP}}^{\Delta F=2} \Lambda_{\text{EW}}^2}{(Y_{\text{SM}} Y_{\text{SM}}^\dagger)^2 \varepsilon_{\Delta F=2}} \gtrsim \frac{C_{\text{NP}}^{\Delta F=1} \Lambda_{\text{EW}}^2}{Y_{\text{SM}} Y_{\text{SM}}^\dagger \varepsilon_{\Delta F=1}}$$

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RGE in the SMEFT implies:

$$C_{\text{NP}}^{\Delta F=2} = Y_{\text{SM}} Y_{\text{SM}}^\dagger C_{\text{NP}}^{\Delta F=1} \mathcal{R} \equiv \frac{1}{16\pi^2} \log \left( \frac{\Lambda_{\text{NP}}}{\Lambda_{\text{EW}}} \right) \sim \mathcal{O}(\%)$$

THEN,  $\Delta F = 2$  BOUND RELEVANT IFF

$$\mathcal{R} \gtrsim \varepsilon_{\Delta F=2} / \varepsilon_{\Delta F=1}$$