

## A fresh field-theoretic calculation of the deuteron magnetic and quadrupole moments

We continue our explorations [1] of the electromagnetic properties of the deuteron with help of the method of unitary clothing transformations (UCT) [2,3]. It is the case, where one has to deal with the matrix elements  $\langle \mathbf{P}', M' | J^\mu(0) | \mathbf{P}, M \rangle$  (to be definite in the lab. frame). Here the operator  $J^\mu(0)$  is the Nöther current density  $J^\mu(x)$  at the point  $x = (t, \mathbf{x}) = 0$ , sandwiched between the eigenstates of a "strong" field Hamiltonian  $H$ , viz., the deuteron states  $|\mathbf{P}, M\rangle$ . Latter belong to the two-clothed-nucleon subspace with the Hamiltonian  $H = P^0 = K_F + K_I$  and the boost operator  $\mathbf{B} = \mathbf{B}_F + \mathbf{B}_I$ , where free parts  $K_F$  and  $\mathbf{B}_F$  are  $\sim b_c^\dagger b_c$  and interactions  $K_I$  and  $\mathbf{B}_I$  are  $\sim b_c^\dagger b_c^\dagger b_c b_c$ . Further, we use the expansion in the  $R$ -commutators

$$J^\mu(0) = W J_c^\mu(0) W^\dagger = J_c^\mu(0) + [R, J_c^\mu(0)] + \frac{1}{2} [R, [R, J_c^\mu(0)]] + \dots,$$

where  $J_c^\mu(0)$  is the initial current in which the bare operators  $\{\alpha\}$  are replaced by the clothed ones  $\{\alpha_c\}$  and  $W = \exp R$  the corresponding UCT. In its turn, the operator being sandwiched between the two-clothed-nucleon states contributes as  $J^\mu(0) = J_{one-body}^\mu + J_{two-body}^\mu$ . By keeping only the one-body contribution we arrive to certain off-energy-shell extrapolation of the so-called relativistic impulse approximation (RIA) in the theory of e.m. interactions with nuclei (bound systems). Of course, the RIA results [1] should be corrected including more complex mechanisms of e-d scattering (see other our contribution). Starting from the operator of the magnetic dipole moment  $\mu = \frac{1}{2} \int d\mathbf{x} \mathbf{x} \times \mathbf{J}(x)$  (reminiscent of the Biot-Savart formula from magnetostatics) one can show that the magnetic dipole moment of the deuteron being defined after Sachs as  $z$ -component of the matrix elements between narrow wave packets of the vector  $\mu$  for the stretched configuration, looks as

$$\mu_d = \lim_{q \rightarrow 0} \left[ -\frac{i}{2} \text{curl}_q \langle \frac{q}{2}; 1 | \mathbf{J}(0) | -\frac{q}{2}; 1 \rangle^z \right],$$

where the matrix elements  $\langle \frac{q}{2}; M_J | \mathbf{J}(0) | -\frac{q}{2}; M_J \rangle$  ( $M_J = (\pm 1, 0)$  projection of the total angular momentum) determine the corresponding current in the Breit frame. In this way the deuteron magnetic dipole moment can be expressed as

$$\mu_d = \frac{1}{m_d} \langle \mathbf{0}; 1 | \frac{1}{2} [\mathbf{B} \times \mathbf{J}(0)]^z | \mathbf{0}; 1 \rangle$$

with the deuteron state  $|\mathbf{0}; 1\rangle$  in the rest frame. In parallel, considering interaction energy of the system with the charge density  $\rho(x) = J^0(x)$  in static external electric field and expanding it in the Cartesian electric moments one encounters the quadrupole moment tensor  $Q_{ij} = \int d\mathbf{x} [3x_i x_j - \delta_{ij} \mathbf{x}^2] \rho(\mathbf{x})$  ( $i, j = 1(x), 2(y), 3(z)$ ). Then repeating the same trick with wave packets one gets the matrix elements

$$\langle JM_J' | Q_{ij} | JM_J \rangle = - \lim_{q \rightarrow 0} \left[ \left\{ 3 \frac{\partial^2}{\partial q_i \partial q_j} - \delta_{ij} \frac{\partial^2}{\partial q_i^2} \right\} \langle \frac{q}{2} | \rho(0) | -\frac{q}{2} \rangle \right]$$

to introduce electric quadrupole moment  $Q = \langle JJ | Q_{33} | JJ \rangle$ . These formulae have been a departure point for our preceding RIA calculations [1]. Here we will show our results with the meson exchange currents and boost ( $\mathbf{B}_I$  part) contributions included.

### References

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