

# Thermal shift of atomic levels in hydrogen: influence on the determination of the proton radius

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# Background

- Progress in theoretical/experimental studies of the hydrogen atom has paved the way for the determination of fundamental physical constants: Rydberg constant and proton charge radius.

The results of theoretical and experimental work are collected in CODATA:

E. Tiesinga, P.J. Mohr, D.B. Newell, B.N. Taylor, Rev. Mod. Phys. **93**, 025010 (2021)

- Most accurate measurements of transition frequencies in hydrogen:

1s – 2s - with fractional uncertainty  $4.2 \times 10^{-15}$

C.G. Parthey et al., Phys. Rev. Lett. **107**, 203001 (2011); A. Matveev et al., Phys. Rev. Lett. **110**, 230801 (2013)

1s – 3s - with fractional uncertainty  $2.5 \times 10^{-13}$

A. Grinin et al., Science **370**, 1061 (2020)

- To find the physical constants one should write a set of equations with

$$E_{nlj} = R_{\infty} \left( -\frac{1}{n^2} + f_{nlj} \left( \alpha, \frac{m_e}{m_p}, r_p \dots \right) \right) \quad (1)$$

$$E_{nlj} - E_{n'l'j'} = \Delta E_{nlj-n'l'j'}^{\text{exp}} \quad (2)$$

Since the question is about finding two variables ( $R_{\infty}$  and  $r_p$ ), it is necessary to solve two equations for independent transitions (one of them is 1s – 2s or 1s – 3s).

# Proton radius puzzle

Until recently (CODATA - 2012, 2016):

$$R_\infty = 10973731.568508(68) \text{ m}^{-1} \text{ and } r_p = 0.8759(77) \text{ fm}$$

However, the muonic hydrogen value (the Lamb shift measurement in  $\mu\text{H}$ )

$$r_p = 0.84169(66) \text{ fm}$$

R. Pohl et al., Nature **466**, 213 (2010)

$$r_p = 0.84087(39) \text{ fm}$$

A. Antognini et al., Science **339**, 417 (2013) (mixing of hyperfine sub-levels)

The problem was called "the proton radius puzzle": eH vs  $\mu\text{H}$  -  $5.6\sigma$

The problem could not be solved for about 10 years:

in  $\mu\text{H}$  very clear experiment  $\Rightarrow$  theory test

in eH clear theory  $\Rightarrow$  experiment test

# Modern status: nonresonant corrections

A. Beyer et al., Science **358**, 79 (2017):  $2s - 4p$  transition frequency measurements

$$R_\infty = 10973731.568076(96) \text{ m}^{-1} \text{ and } r_p = 0.8335(95) \text{ fm}$$

have partially solved the proton radius puzzle (coincides with  $\mu\text{H}$  experiment within the uncertainty region) by taking into account nonresonant effects (NR).

Theory: F. Low, Phys. Rev. **88**, 53 (1952)

Prof. L.N. Labzovsky for HCl: J. Phys. B **27**, L439 (1994); Physica Scripta **56**, 271 (1997)

$e\text{H}$ : Phys. Rev. Lett. **87**, 140003 (2001)

U. D. Jentschura, P. J. Mohr, Can. J. Phys. **80**, 633 (2002) (quantum interference effect - QIE)

Interfering transitions between the nearest atomic levels lead to an asymmetry of the observed line shape at a level of several tens of kilohertz.

The Fano-Voigt profile model was applied to compensate for the QIE, Science **358**, 79 (2017)

The theory is not so unambiguous and strongly depends on the experimental measurement technique

D. Solov'yev et al., J. Phys. B: **53**, 125002 (2020) and references therein.

$\mu\text{H}$  experiments are insensitive to NR corrections

The  $r_p = 0.8335(95)$  fm value is the most important step in solving the proton radius puzzle.

# Modern status: nonresonant corrections

Up to the Science **358**, 79 (2017) experiment, the QIE was considered theoretically for the one-photon scattering process in  $eH$  and  $\mu H$  atoms.

CODATA includes two-photon transitions  $2s - nS/nD$  in the  $eH$  atom.

The process-dependent NR corrections (QIE as the dominant part of them) were investigated in A. Anikin, T. Zaliialutdinov, D. Solov'yev, Phys. Rev. A **103**, 022833 (2021); JETP Letters **114**, 212 (2021), see also J. Phys. B (2021).

The results for the  $2s - nD$  two-photon transitions can be presented as

$$\delta_{NR} = -\frac{C_{ab}\Gamma_{nd}^2}{4C_a\Delta_{fs}} + \mathcal{O}\left(\frac{\Gamma_{nd}^4}{\Delta_{fs}^3}\right), \quad (3)$$

where  $\Gamma_{nd}$  is the level width,  $\Delta_{fs}$  is the fine splitting energy, and the coefficients  $C_{ab}$ ,  $C_a$  are calculated numerically based on the theory of angular momentum. Numerical results

	$\Delta_{fs}$ , Hz	$\Gamma_{nd}$ , Hz	$\delta_{NR}$ , Hz	Exp. unc., Hz
$4d$	$4.557026 \times 10^8$	$4.40503 \times 10^6$	$-8691.82$	$24. \times 10^3$
$6d$	$1.350231 \times 10^8$	$1.33682 \times 10^6$	$-2701.67$	$10. \times 10^3$
$8d$	$5.69628 \times 10^7$	$5.72382 \times 10^5$	$-1174.02$	$6.4 \times 10^3$
$12d$	$1.68779 \times 10^7$	$1.72261 \times 10^5$	$-358.88$	$7.0 \times 10^3$

# Thermal shift: QED theory

The influence of blackbody radiation has been well known since the early 1980s: T.F. Gallagher, W.E. Cooke, Phys. Rev. Lett. **42**, 835 (1979); J.W. Farley, W.H. Wing, Phys. Rev. A **23**, 2397 (1981) - the QM description for the Stark effect is still used.

**QED theory of BBR-induced effects:**

D. Solov'yev et al., Phys. Rev. A **92**, 022508 (2015); D. Solov'yev, Ann. Phys. **415**, 168128 (2020) The Stark shift is the same:

$$\Delta E_a^{\text{Stark}} = \frac{4e^2}{3\pi} \sum_n |\langle a | \vec{r} | n \rangle|^2 \int_0^\infty n_\beta(\omega) \frac{E_{an}\omega^3}{E_{na}^2 - \omega^2} d\omega, \quad (4)$$

$E_{na} = E_n - E_a$ ,  $n_\beta$  - the Planck distribution function.

There is an additional thermal effect:

$$\Delta E_a^\beta = -\frac{2Ze^2}{3\pi} \int_0^\infty n_\beta(\omega) \omega^2 \langle a | r^2 | a \rangle \stackrel{eH}{=} -\frac{4Ze^2\zeta(3)}{3\pi\beta^3} \frac{n_a^2}{2} \left[ 5n_a^2 + 1 - 3l_a(l_a + 1) \right] a_0^2. \quad (5)$$

Stark shift  $\sim T^4$  (deformation polarizability), thermal shift  $\sim T^3$  (orientation polarizability).

**The magnitude reaches several kHz for highly excited states at room temperature.**

# Thermal and NR corrections: proton radius

Thermal shifts combined with NR (QIE) corrections should be taken into account to determine the proton charge radius.

M. Horbatsch, E.A. Hessels, Phys. Rev. A **93**, 022513 (2016), CODATA - 2021:

$$R_\infty = 10973731.568160(21) \text{ m}^{-1} \text{ and } r_p = 0.8414(19) \text{ fm.}$$

The hyperfine structure has not been resolved for two-photon  $2s - nS/nD$  transitions.

	$r_p^{\text{HH}}$ , fm	$R_\infty^{\text{HH}}$ m <sup>-1</sup>	$r_p^{\text{HH}+\beta}$ , fm	$R_\infty^{\text{HH}+\beta}$ m <sup>-1</sup>
$1s - 2s, 2s - 8d_{3/2}$	0.8412	10973731.568152	0.8403	10973731.568142
$1s - 3s, 2s - 8d_{3/2}$	0.8407	10973731.568149	0.8396	10973731.568139
$1s - 2s, 2s - 8d_{5/2}$	0.8413	10973731.568153	0.8404	10973731.568144
$1s - 3s, 2s - 8d_{5/2}$	0.8408	10973731.568151	0.8398	10973731.568141
$1s - 2s, 2s - 12d_{3/2}$	0.8413	10973731.568152	0.8343	10973731.568082
$1s - 3s, 2s - 12d_{3/2}$	0.8407	10973731.568150	0.8328	10973731.568075
$1s - 2s, 2s - 12d_{5/2}$	0.8412	10973731.568151	0.8342	10973731.568081
$1s - 3s, 2s - 12d_{5/2}$	0.8406	10973731.568149	0.8327	10973731.568074
$1s - 2s, 2s - 4d_{5/2} - \frac{1}{4}(1s - 2s)$	0.8398	10973731.568138	0.8403	10973731.568143
$1s - 3s, 2s - 4d_{5/2} - \frac{1}{4}(1s - 2s)$	0.8398	10973731.568141	0.8403	10973731.568146
$1s - 2s, 2s - 6d_{5/2} - \frac{1}{4}(1s - 3s)$	0.8413	10973731.568153	0.8394	10973731.568134
$1s - 3s, 2s - 6d_{5/2} - \frac{1}{4}(1s - 3s)$	0.8413	10973731.568201	0.8394	10973731.568183
rms( $1s - 2s$ )	0.8410	10973731.568150	0.8383	10973731.568121
rms( $1s - 3s$ )	0.8406	10973731.568157	0.8374	10973731.568126
rms	0.8408	10973731.568153	0.8379	10973731.568123



# Thermal and NR corrections: proton radius

Proton charge radius,  $r_p$ , and rydberg constant,  $R_\infty$ , using data of M. Horbatsch, E.A. Hessels, Phys. Rev. A **93**, 022513 (2016) and A. Grinin et al., Science **370**, 1061 (2020), thermal effects and NR corrections.

	$r_p^{\text{HH}+\beta}$ , fm	$R_\infty^{\text{HH}+\beta}$ m <sup>-1</sup>	$r_p^{\text{HH}+\beta+\text{NR}}$ , fm	$R_\infty^{\text{HH}+\beta+\text{NR}}$ m <sup>-1</sup>
1s - 2s, 2s - 8d <sub>3/2</sub>	0.8403	10973731.568142	0.8410	10973731.568150
1s - 3s, 2s - 8d <sub>3/2</sub>	0.8396	10973731.568139	0.8405	10973731.568147
1s - 2s, 2s - 8d <sub>5/2</sub>	0.8404	10973731.568144	0.8409	10973731.568149
1s - 3s, 2s - 8d <sub>5/2</sub>	0.8398	10973731.568141	0.8403	10973731.568146
1s - 2s, 2s - 12d <sub>3/2</sub>	0.8343	10973731.568082	0.8345	10973731.568084
1s - 3s, 2s - 12d <sub>3/2</sub>	0.8328	10973731.568075	0.8330	10973731.568077
1s - 2s, 2s - 12d <sub>5/2</sub>	0.8342	10973731.568081	0.8344	10973731.568082
1s - 3s, 2s - 12d <sub>5/2</sub>	0.8327	10973731.568074	0.8328	10973731.568075
1s - 2s, 2s - 4d <sub>5/2</sub> - $\frac{1}{4}(1s - 2s)$	0.8403	10973731.568143	0.8481	10973731.568223
1s - 3s, 2s - 4d <sub>5/2</sub> - $\frac{1}{4}(1s - 2s)$ ,	0.8403	10973731.568146	0.8481	10973731.568220
1s - 2s, 2s - 6d <sub>5/2</sub> - $\frac{1}{4}(1s - 3s)$	0.8394	10973731.568134	0.8414	10973731.568154
1s - 3s, 2s - 6d <sub>5/2</sub> - $\frac{1}{4}(1s - 3s)$	0.8394	10973731.568183	0.8414	10973731.568202
rms(1s - 2s)	0.8383	10973731.568121	0.8401	10973731.568140
rms(1s - 3s)	0.8374	10973731.568126	0.8394	10973731.568145
rms	0.8379	10973731.568123	0.8397	10973731.568142

# Conclusions

Proton radius puzzle solved according to CODATA-2021

$$r_p = 0.8414(19) \text{ fm}$$

- $\mu\text{H}$  experiment leads to **stable** value of  $r_p$
- Most significant contribution arises from the hyperfine splitting in the eH atom
  - $r_p = 0.8759(77) \text{ fm} \rightarrow r_p = 0.8414(19) \text{ fm}$
  - The latter is the same as  $\mu\text{H}$  results:  $r_p = 0.84169(66), 0.84087(39) \text{ fm}$
- Nonresonant corrections in the eH atom: QIE as the most significant part
- Further experimental improvements intersect with the need for an accurate analysis of NR corrections for two-photon transitions
- Thermal environment:
  - The thermal Stark shift is well known and included into the analysis
  - There are other thermal-induced energy shifts as well
- Further experimental improvements intersect with the need for an accurate analysis of the thermal shifts affecting the last digits of the proton charge radius

# Thank you for the attention

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