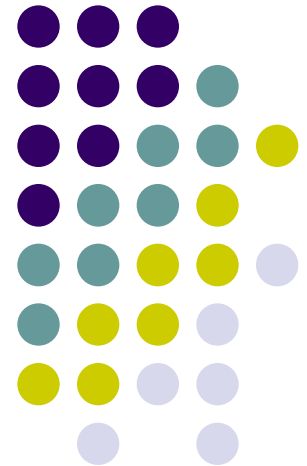




# Production cross sections of mercury and radon isotopes in complete fusion reactions with $^{36,40}\text{Ar}$ and $^{40,48}\text{Ca}$ projectiles

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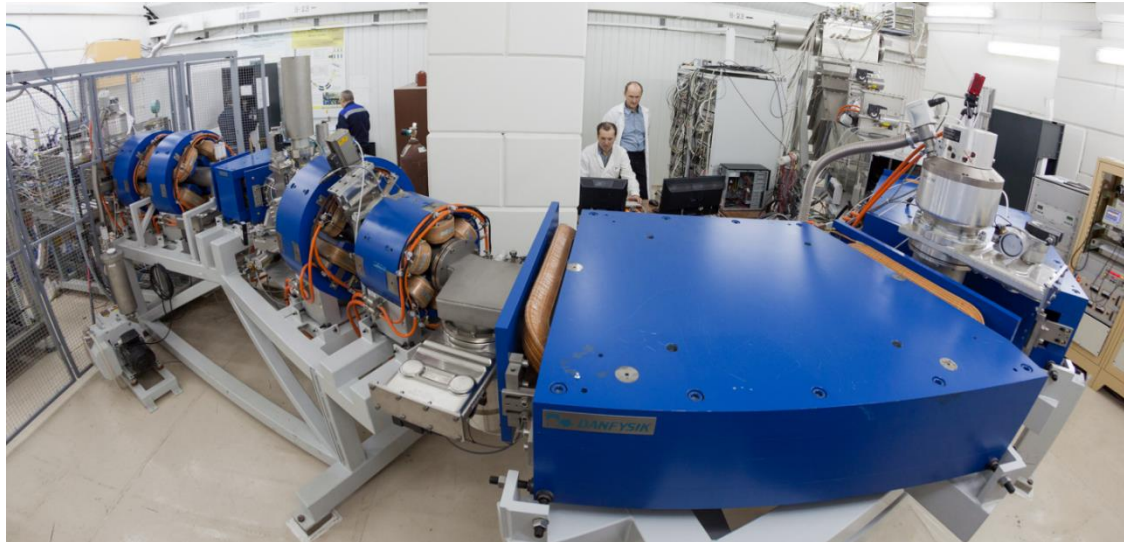
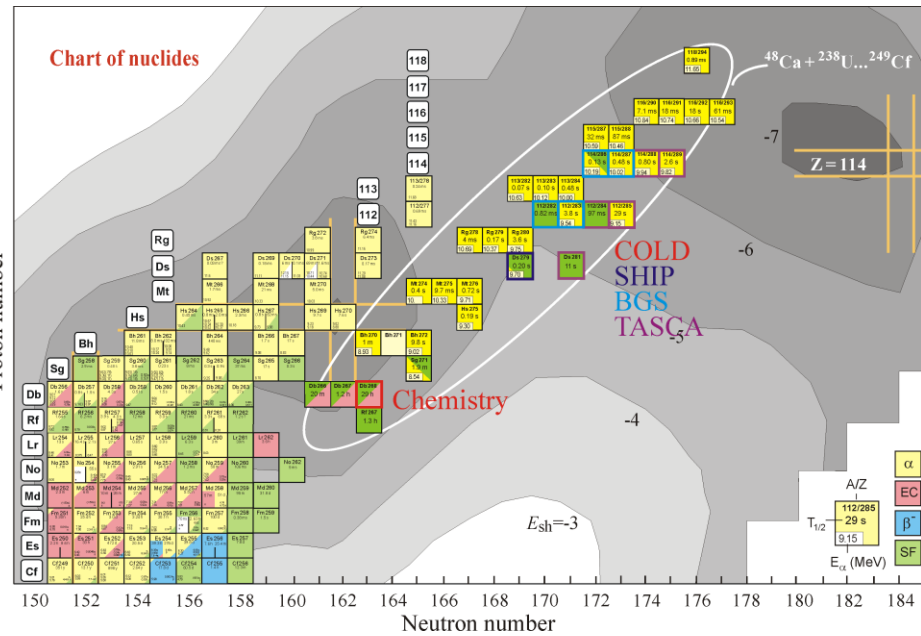
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# Synthesis of superheavy nuclei at FLNR JINR

A significant progress has been achieved in the nuclear physics experiments devoted to synthesis of new superheavy elements and study of their properties. The nuclides up to  $Z=118$  (Og) were synthesized in the reactions of doubly magic  $^{48}\text{Ca}$  ions with actinide targets.

Mass separator MASHA [1] was built in FLNR JINR, combining the **ISOL method** for the synthesis and separation of radioactive nuclei with the **classical mass spectrometry technique**.



The life-times of some superheavy nuclei are quite long: 3.8s for  $^{283}\text{Cn}$ , 29s for  $^{285}\text{Cn}$  and 2.6s for  $^{289}\text{Fl}$ .

Cn ( $Z = 112$ ) is more volatile than its homolog Hg, while Fl ( $Z=114$ ) differs from its light homolog Pb and exhibits the properties of noble inert gases.

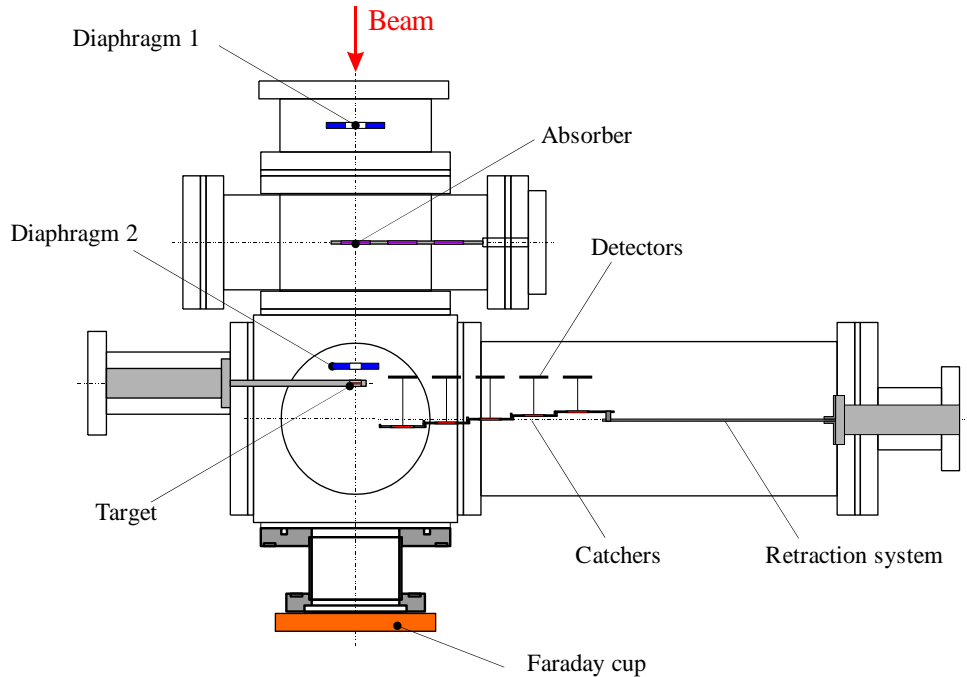
[1] A.M. Rodin *et al.*, *Inst. Exp. Techn.* Vol. 57, No. 4, pp. 386–393 (2014)

Mass Analyser of Super Heavy Atoms (MASHA)



- To verify experimentally a capability of the mass-separator MASHA to measure the masses of the volatile nuclides of Cn ( $Z=112$ ) and Fl ( $Z=114$ ), whose production cross-sections are extremely low (about several pb), *the efficiency of the mass spectrometer* is proposed to measure in the experiments with *mercury (homolog of Cn,  $Z=112$ ) and radon (noble gas as Fl,  $Z=114$ )*
- The purpose of the present work was to determine *the production absolute cross-sections of mercury and radon isotopes* produced as evaporation residues in the complete fusion reactions of  $^{36,40}\text{Ar}$  and  $^{40,48}\text{Ca}$  ions with lanthanide targets of  $^{144, 148}\text{Sm}$ ,  $^{142,144}\text{Nd}$  and  $^{166}\text{Er}$ .
- For heavy and superheavy nuclei the fusion cross-sections greatly depend on the reaction **entrance channel properties**, such as **mass asymmetry**, **charges of the interacting nuclei  $Z_1Z_2$** , **their deformations**, as well as **shell closure**, where the quasi-fission process competes with the fusion one. So, an extra motivation to carry out the experiment was to study the effect of the entrance channel on production of evaporation residues of  $^{184}\text{Hg}$  in three reactions:  $^{40}\text{Ar} + ^{144}\text{Sm}$ ,  $^{36}\text{Ar} + ^{148}\text{Sm}$ ,  $^{40}\text{Ca} + ^{144}\text{Nd}$ .
- The effect of **the neutron excess** was investigated in the reactions  $^{40}\text{Ca} + ^{144}\text{Nd} \rightarrow ^{184}\text{Hg}$  and  $^{48}\text{Ca} + ^{142}\text{Nd} \rightarrow ^{190}\text{Hg}$ .

# Experimental set-up for measurement of short-lived radon and mercury isotopes



Scheme of experiment with moving catchers.

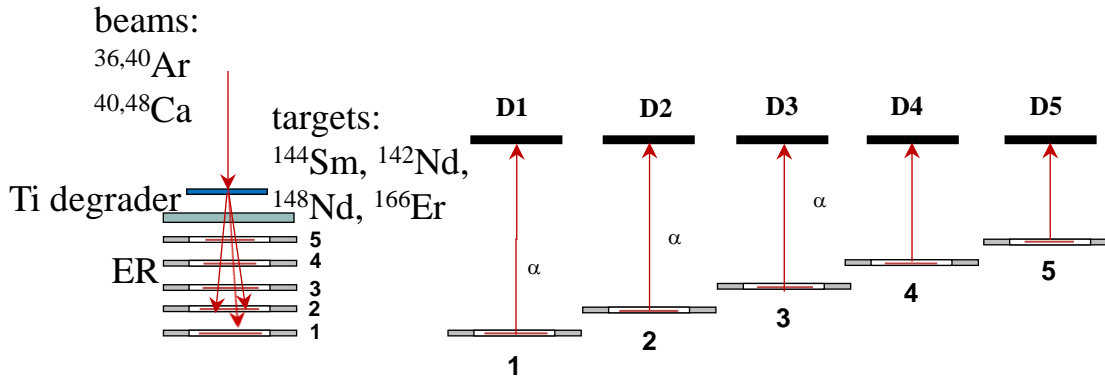
The experiments were carried out at the U400M cyclotron of the FLNR, JINR (Dubna). Beams of  $^{36}\text{Ar}$ ,  $^{40}\text{Ar}$ ,  $^{40}\text{Ca}$ , and  $^{48}\text{Ca}$  with the energies of 250-330 MeV were used. The beam intensity in experiments was  $\sim 0.2 \mu\text{A}$ . The targets were 280–330  $\mu\text{g}/\text{cm}^2$  layers of  $^{144}\text{Sm}$ ,  $^{166}\text{Er}$ ,  $^{148}\text{Sm}$ ,  $^{142}\text{Nd}$ , and  $^{144}\text{Nd}$ . Catchers: Al foils (0.8  $\mu\text{m}$ ).

**The beam interruption method** was applied for isotope identification. Geometrical efficiencies for the semiconductor detectors were 4.90; 3.95; 3.37; 2.84 and 2.43 %.



Top: pneumatic actuator with aluminum catcher foils and silicon detectors. Bottom: absorber foils in the position for registration of  $\alpha$ -particles

# Experiment with moving catcher foils



## I. accumulation mode

## II. measurement mode

In **the accumulation mode** the foils were placed one under another, so that the reaction products were stopped at different depths in dependence of their energies.

The time of moving of aluminum absorbers between their extreme positions was 0.3 s.



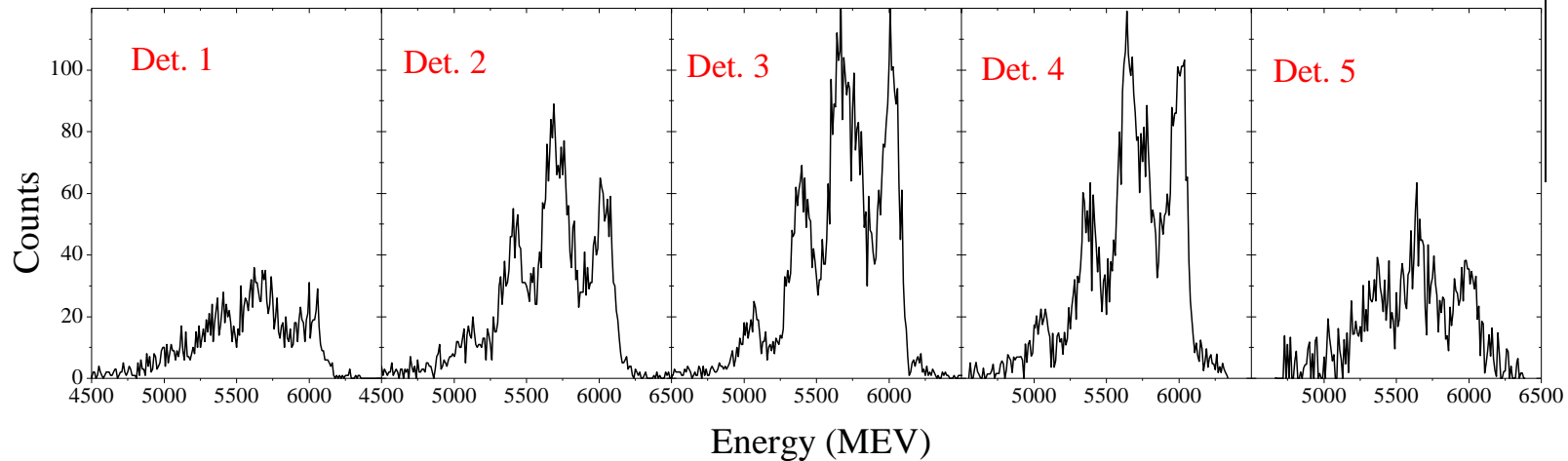
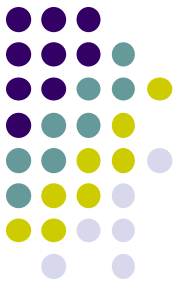
The accumulation and measurement times were chosen to be 10 s – 30 s according to the half-lives of the studied reaction products.

To register the reaction products an **advanced catcher-foil technique** was applied [D. Vermeulen, *et al.*, *Z. Phys. A* 318(1984)157].

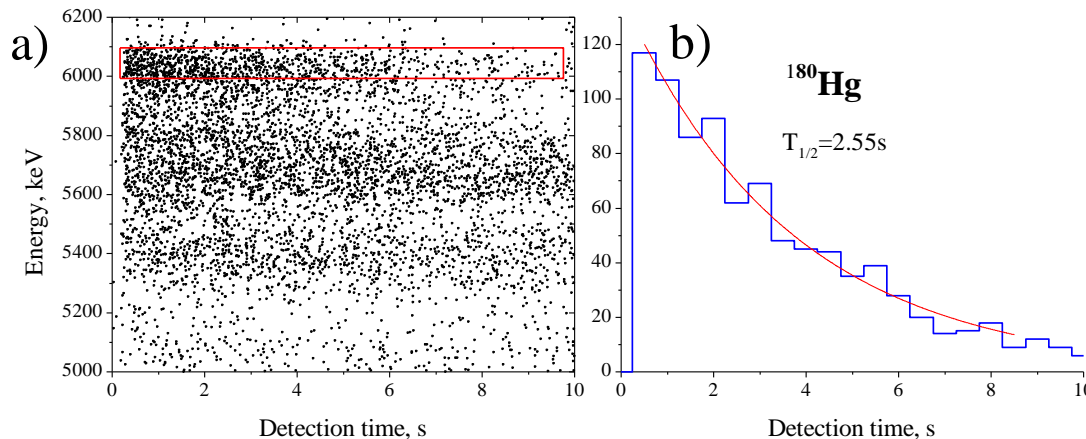
During the experiment, the catcher foils were periodically in two positions until spectra with needed statistics were accumulated.

In **the measurement mode** the foil array was extended to its full width so that each foil was placed opposite to the corresponding silicon detectors registered  $\alpha$ -decays of short-lived radon and mercury isotopes.

# Data analysis. Isotope identification



Experimental  $\alpha$ -particle spectra of the evaporation residues of the reaction  $^{40}\text{Ar}+^{144}\text{Sm}$  at the projectile energy  $E_{\text{lab}} = 180$  MeV for the foils 1-5, registered by corresponding semiconductor detectors.



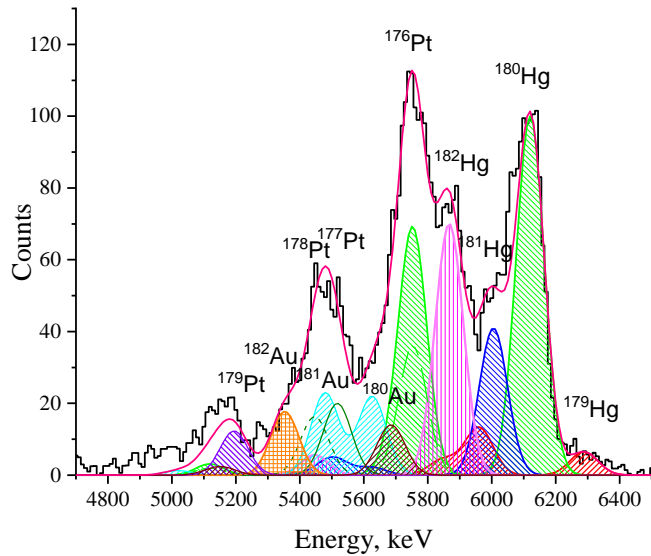
**Identification of isotopes** was performed using

1.  $\alpha$ -particle energies;
2. half-lives of isotopes;
3. analysis of excitation functions.

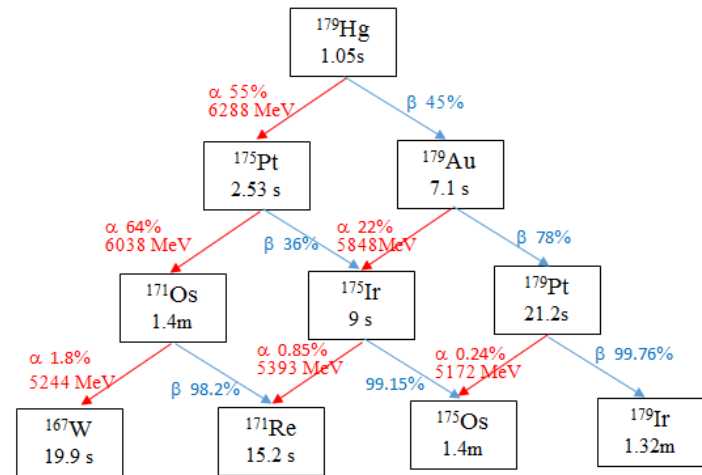
- a) two-dimensional scatter-plot of  $\alpha$ -particle energies vs. their times of registration starting from begin of measuring cycle for the reaction  $^{40}\text{Ar}+^{144}\text{Sm}$  at  $E_{\text{lab}}=187$  MeV, accumulated in third foil of catcher array;
- b) decay curve for the cut in the scatter-plot .



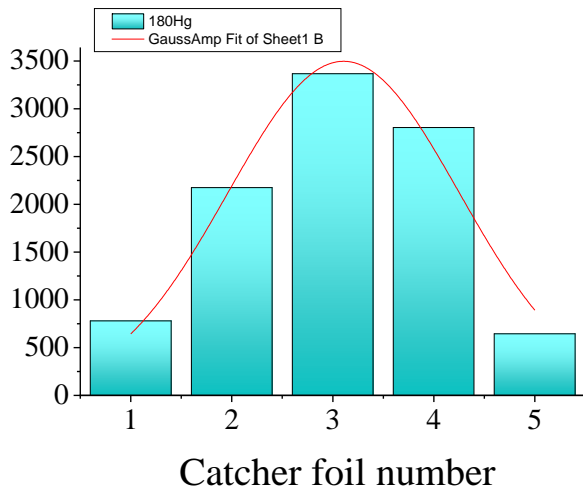
# CROSS SECTION CALCULATIONS



Experimental spectrum of  $\alpha$ -particle energy from the fourth catcher foil for the reaction  $^{40}\text{Ar} + ^{144}\text{Sm}$  at  $E_{\text{lab}} = 187$  MeV. The yields of different isotopes with their daughter and grand-daughter nuclei are shown by colour lines.



Decay chain of  $^{179}\text{Hg}$  with  $\alpha$  and  $\beta$  branchings for each isotope in per cent.

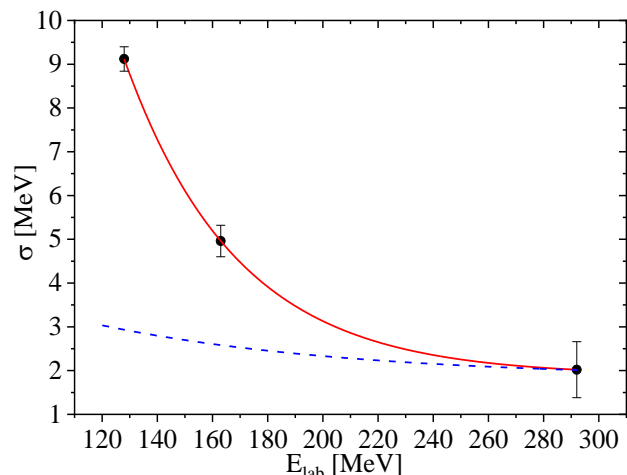
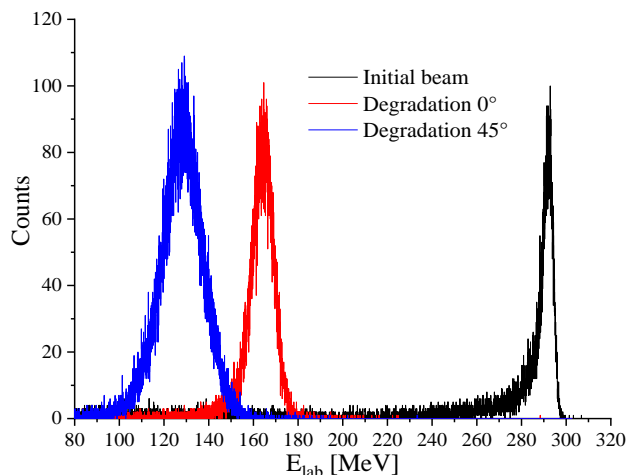


Yields of isotopes  $^{180}\text{Hg}$  accumulated by catcher foils 1-5.

$$\sigma = \frac{N_{det}^{\alpha} \lambda (1 - \exp(-\lambda T))}{g \beta \epsilon_{\alpha} I (1 - \exp(-\lambda t_B)) \exp(-\lambda t_m) (1 - \exp(-\lambda t_D)) t_D}$$

$N_{det}^{\alpha}$  - number of  $\alpha$ -particles registered during one run  
 T - the duration of one measurement cycle;  $T = t_B + 2t_m + t_D$ ,  
 $t_B$  - accumulation time,  
 $t_D$  - detection time,  
 $t_m$  - the move time of the retraction system,  
 $g$  - geometrical efficiency of the detector,  
 $\beta = dSr$  ( $d$  - target thickness,  $S$  - target area,  $r$  - target nucleus density),  
 $\lambda$  - decay constant of studied isotope  
 I - integral

# Corrections of excitation functions for beam energy spread



Left: measured energy spectrum of the  $^{36}\text{Ar}$  beam without target and Ni foil (black), with Ni foil (10,3  $\mu\text{m}$ ,  $0^\circ$  and target, and after rotating the Ni foil by  $45^\circ$  (blue). Right: the dependence of the  $^{40}\text{Ar}$  beam energy spread in the middle of the target (black squares) and its exponential fit (red line). A blue dashed line – SRIM[2] calculations.

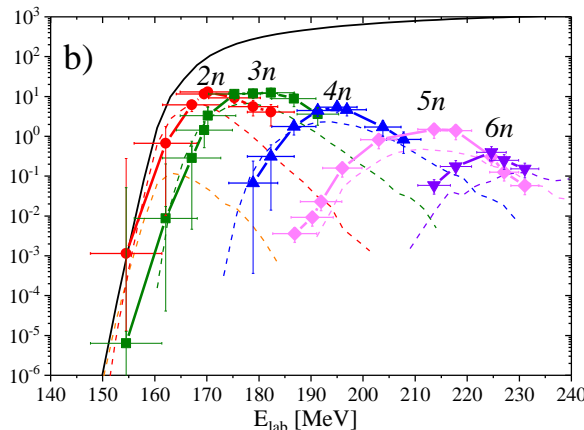
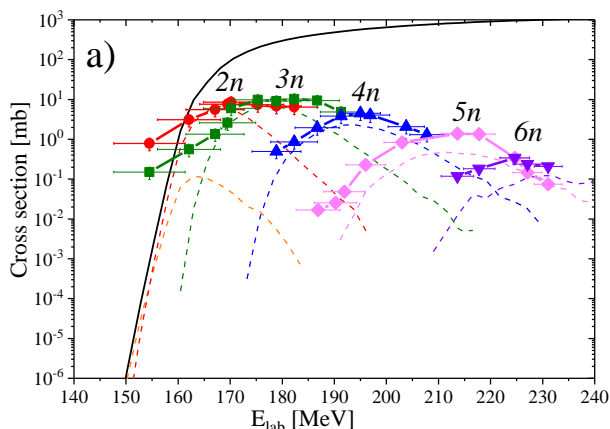
Gold deconvolution method [1]

$$D = \frac{1}{\sigma\sqrt{2\pi}} \cdot e^{-\frac{1}{2}\left(\frac{E-E_c}{\sigma}\right)^2}$$

$$\sigma = B \cdot e^{-\frac{E_c}{t}} + \sigma_0,$$

$$h^{ij} = C^j \cdot e^{-2\left(\frac{E^i - E_c^j}{\sigma}\right)^2}$$

$$C^j = \frac{1}{\sum_i e^{-2\left(\frac{E^i - E_c^j}{\sigma}\right)^2}}$$



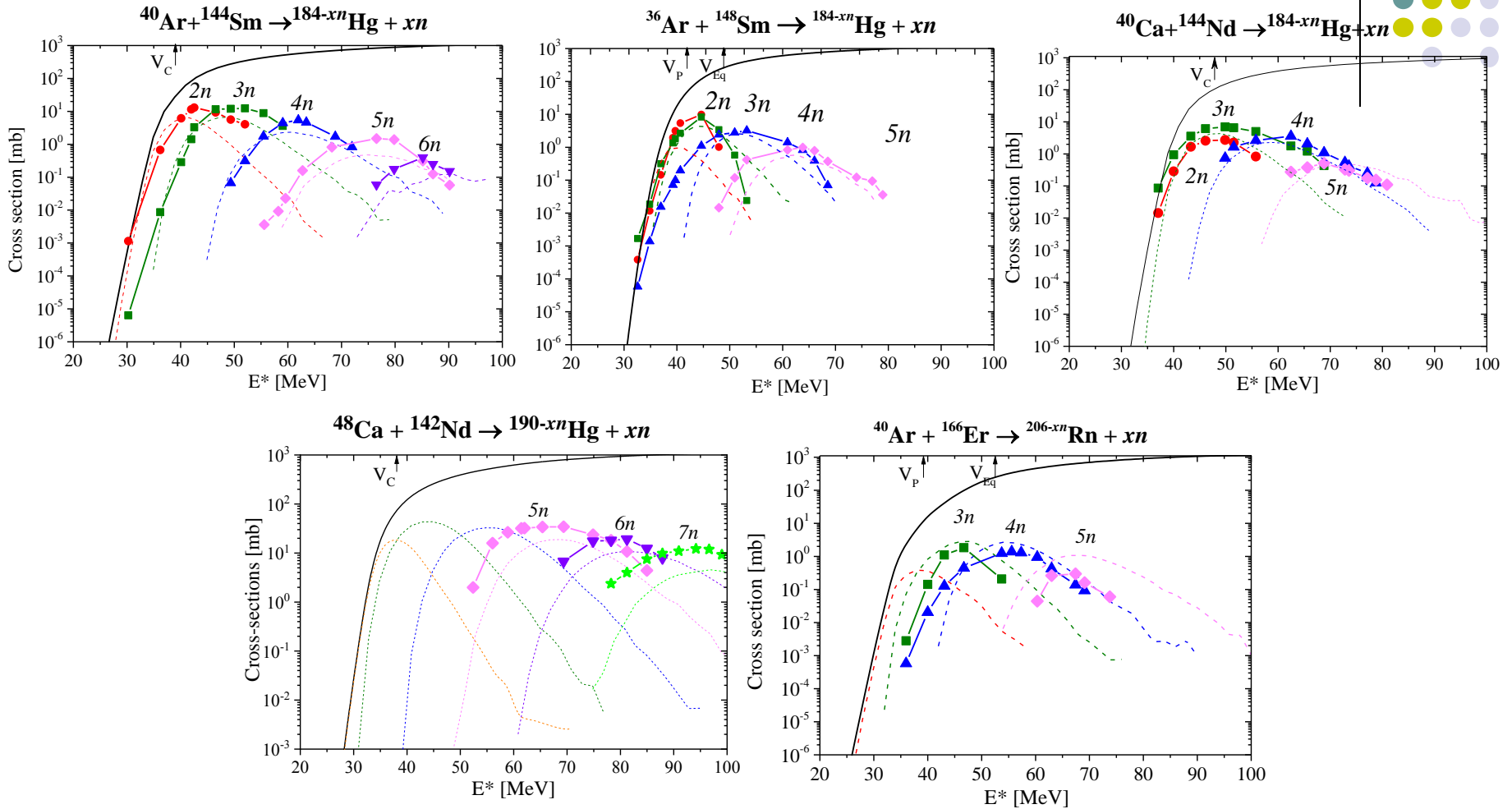
Cross sections of xn-evaporation channels of the reaction  $^{40}\text{Ar} + ^{144}\text{Sm}$  measured by using the catcher foil method: a) before and b) after deconvolution. Dashed curves are the cross sections calculated by the coupled-channel model nrv.

$D$  – response function of a beam energy spreading;  
 $E_c$  – beam energy in the middle of the target;  
 $\sigma$  – energy variance;  
 $h^{ij}$  – matrix elements of response function  
 $C^j$  – normalization constant

1. M. Morhac *et al.*, Nucl. Instr. Meth. A 401, 385-408 (1997)  
 2. J.F. Ziegler *et al.*, Nucl. Instr. Meth. B, 268, p. 1818-1823 (2010)



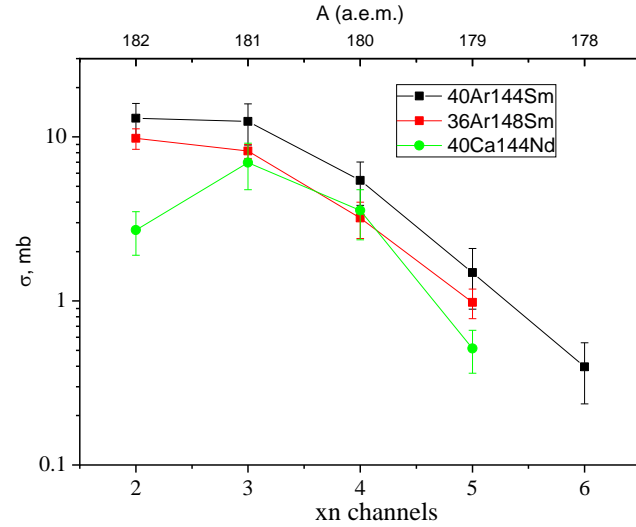
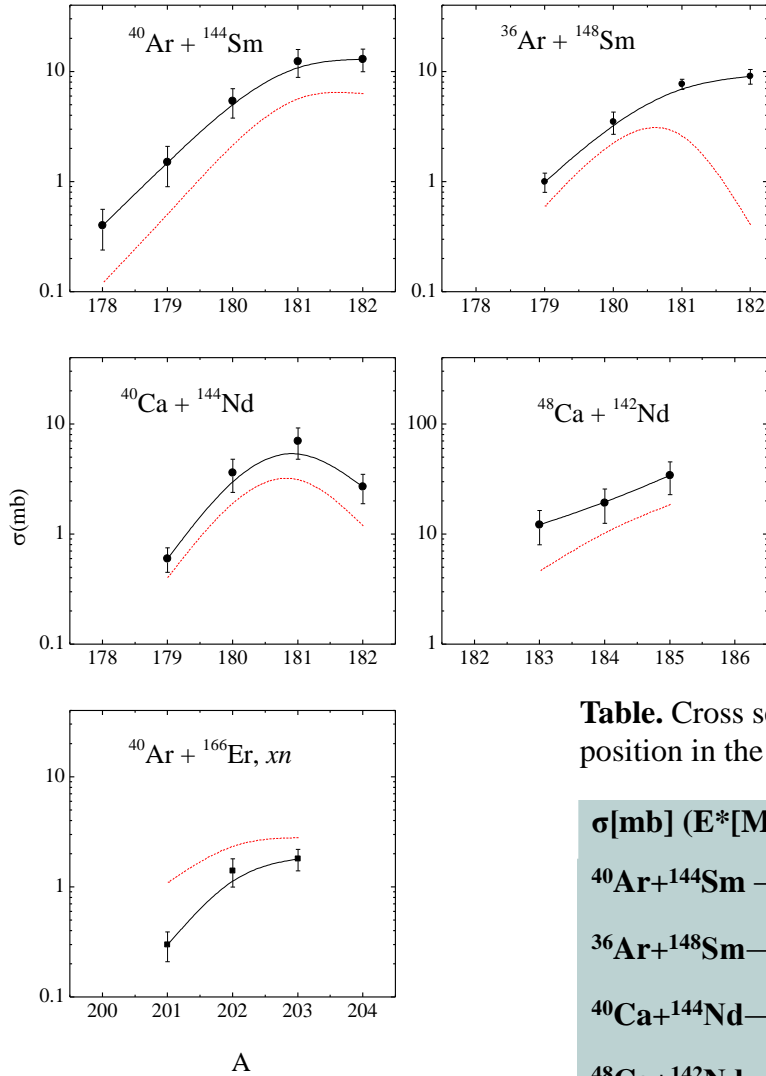
# Excitation functions of the reactions $^{40}\text{Ar}+^{144}\text{Sm}$ , $^{36}\text{Ar}+^{148}\text{Sm}$ , $^{40}\text{Ca}+^{144}\text{Nd}$ , $^{48}\text{Ca}+^{142}\text{Nd}$ and $^{40}\text{Ar}+^{166}\text{Er}$



Cross sections (symbols) of  $xn$ -evaporation channels of the reactions measured by using the catcher foil method. Dashed curves are the cross sections calculated by the coupled-channel model nrv.  $V_c$  are the Coulomb barriers for spherical interacting nuclei,  $V_p$  and  $V_{Eq}$  are the Coulomb barriers for polar and equatorial mutual orientation of deformed ones, respectively.

<http://nrv.jinr.ru/nrv> A.V. Karpov *et al.*, *Physics of Atomic Nuclei* 2016, V. 79, p. 749–761.

# Results



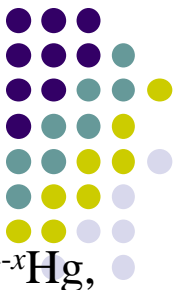
The maximum values of the xn-channels of the reactions leading to the  $^{184}\text{Hg}$  formation.

**Table.** Cross section maximum values in separate channels of complete fusion reactions and their position in the excitation energy scale.

$\sigma[\text{mb}] (E^*[\text{MeV}])$	2n	3n	4n	5n	6n	7n
$^{40}\text{Ar}+^{144}\text{Sm} \rightarrow ^{184}\text{Hg}$	13±3 (42.5±3.3)	12.4±3.5 (52±2.6)	5.42±1.6 (62±2.3)	1.49±0.6 (76.5±2.2)	0.39±0.16 (85±2.1)	
$^{36}\text{Ar}+^{148}\text{Sm} \rightarrow ^{184}\text{Hg}$	9.8±1.4 (48.9±4.9)	8.3±0.8 (48.9±4.8)	3.2±0.8 (60.7±3.2)	0.98±0.2 (75.2±2.3)		
$^{40}\text{Ca}+^{144}\text{Nd} \rightarrow ^{184}\text{Hg}$	2.7±0.8 (49.9±5.7)	7.1±2.2 (46.1±5.7)	3.6±1.2 (62.5±4.1)	0.5±0.15 (68.8±3.8)		
$^{48}\text{Ca}+^{142}\text{Nd} \rightarrow ^{190}\text{Hg}$				34.4±11.3 (69.4±2.6)	19.2±6.6 (81.2±2.1)	12.2±4.2 (94.1±2)
$^{40}\text{Ar}+^{166}\text{Er} \rightarrow ^{206}\text{Rn}$		1.8±0.38 (46.7±3)	1.4±0.41 (55.6±2.4)	0.3±0.09 (67.4±2)		

The comparison of experimental (solid lines) and theoretical [nrv] (dash lines) maxima of excitation functions for xn-channels of the studied reactions.

# Conclusion



Evaporation residue cross sections in the reactions  $^{144}\text{Sm}(^{40}\text{Ar},xn)^{184-x}\text{Hg}$ ,  $^{148}\text{Sm}(^{36}\text{Ar},xn)^{184-x}\text{Hg}$ ,  $^{144}\text{Nd}(^{40}\text{Ca},xn)^{184-x}\text{Hg}$ ,  $^{166}\text{Er}(^{40}\text{Ar},xn)^{206-x}\text{Rn}$ , and  $^{142}\text{Nd}(^{48}\text{Ca},xn)^{190-x}\text{Hg}$  have been measured at FLNR JINR (Dubna). The improved catcher foil method was applied for the isotope identification.

The results for the reactions  $^{144}\text{Sm}(^{40}\text{Ar},xn)^{184-x}\text{Hg}$ ,  $^{144}\text{Nd}(^{40}\text{Ca},xn)^{184-x}\text{Hg}$ ,  $^{166}\text{Er}(^{40}\text{Ar},xn)^{206-x}\text{Rn}$ , and  $^{142}\text{Nd}(^{48}\text{Ca},xn)^{190-x}\text{Hg}$  were obtained for the first time. By comparison of the  $^{40}\text{Ar} + ^{144}\text{Sm}$ ,  $^{36}\text{Ar} + ^{148}\text{Sm}$  and  $^{40}\text{Ca} + ^{144}\text{Nd}$  reactions leading to the same compound nucleus, the effect of the Coulomb barrier value on the position of the  $2n$  channel maximum cross section was investigated.

Slightly higher values of experimentally determined cross sections in comparison with coupled-channel calculated ones were shown for all reactions leading to mercury isotopes, as opposed to the reaction  $^{40}\text{Ar} + ^{166}\text{Er}$  leading to  $^{206}\text{Rn}$ . The latter could be caused by a rather high values of  $1pxn$ -evaporation channel cross sections in this reaction.



The largest Mendeleev Table in Eurasia on the wall of the Archimedes swimming pool was inaugurated in July 2021 to the JINR 65-year anniversary.

**Thank you  
for your attention!**

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