

Non-statistical nature of fragments' spin distributions in binary nuclear fission

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Since for spontaneous and low-energy induced fission, compound fissile nuclei and primary fission fragments in the vicinity of the scission point are in cold nonequilibrium states [1], when constructing the spin distributions of these fragments, it is necessary to take into account [2,3] only zero transverse bending- and wriggling-vibrations of the indicated fissile nuclei. Expressing the normalized distribution function of $W(\mathbf{J}_1, \mathbf{J}_2)$ fission fragments over spins \mathbf{J}_1 and \mathbf{J}_2 in terms of the product of the squared moduli of the wave functions of zero bending- and wriggling-vibrations, one can obtain [4]:

$$W(\mathbf{J}_1, \mathbf{J}_2) = \frac{4J_1 J_2}{\pi C_b C_w} \exp \left[-\frac{1}{2} \left(\frac{1}{C_b} + \frac{1}{C_w} \right) (J_1^2 + J_2^2) + \left(\frac{1}{C_b} - \frac{1}{C_w} \right) J_1 J_2 \cos \phi \right], \quad (1)$$

where ϕ ($0 \leq \phi \leq 2\pi$) is the angle between the two-dimensional spin vectors of fragments \mathbf{J}_1 and \mathbf{J}_2 lying in plane XY . By integrating in (1) over variables J_2 and ϕ , one can obtain [4] the normalized distribution of spin J_1 of the first fission fragment and estimate the average value \bar{J}_1 of spin J_1 :

$$W(J_1) = \frac{4J_1}{C_b + C_w} \exp \left[-\frac{2J_1^2}{C_b + C_w} \right], \quad \bar{J}_1 = \int_0^\infty J_1 W(J_1) dJ_1 = \frac{1}{2} \sqrt{\frac{\pi}{2}} (C_b + C_w)^{1/2}. \quad (2)$$

For a fissile nucleus ^{236}U at values [4] of parameters $M_w = 1.6 \cdot 10^6 \text{ MeV} \cdot \text{Fm}^2 \cdot \text{s}^2$; $M_b = 2.0 \cdot 10^6 \text{ MeV} \cdot \text{Fm}^2 \cdot \text{s}^2$; $K_w = 295 \text{ MeV} \cdot \text{rad}^{-2}$; $K_b = 52 \text{ MeV} \cdot \text{rad}^{-2}$; $\hbar\omega_w = 2.3 \text{ MeV}$; $\hbar\omega_b = 0.9 \text{ MeV}$; $C_w = 132\hbar^2$ and $C_b = 57\hbar^2$, it follows that the energies of vibrational quanta $\hbar\omega_w$ and coefficients C_w for wriggling-vibrations turn out to be noticeably larger than those for bending-vibrations. This means that the main contribution to \bar{J}_1 (2) comes from wriggling vibrations. Then the calculated value $\bar{J}_1 = 8.6$ correlates well with the experimental [5] average values of the spins of fission fragments $\bar{J}_1 = 7-9$.

This means that the spin distribution of fission fragments is determined with a good degree of accuracy by taking into account zero wriggling and bending vibrations of a composite fissile system. This confirms the assumption [6] about the inequality of the statistical Gibbs distribution with temperature T for the spin distribution of fragments, which is used in [1].

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Primary authors: LYUBASHEVSKY, Dmitrii (Voronezh State University); KADMENSKY, Stanislav (Voronezh State University)

Presenter: LYUBASHEVSKY, Dmitrii (Voronezh State University)

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