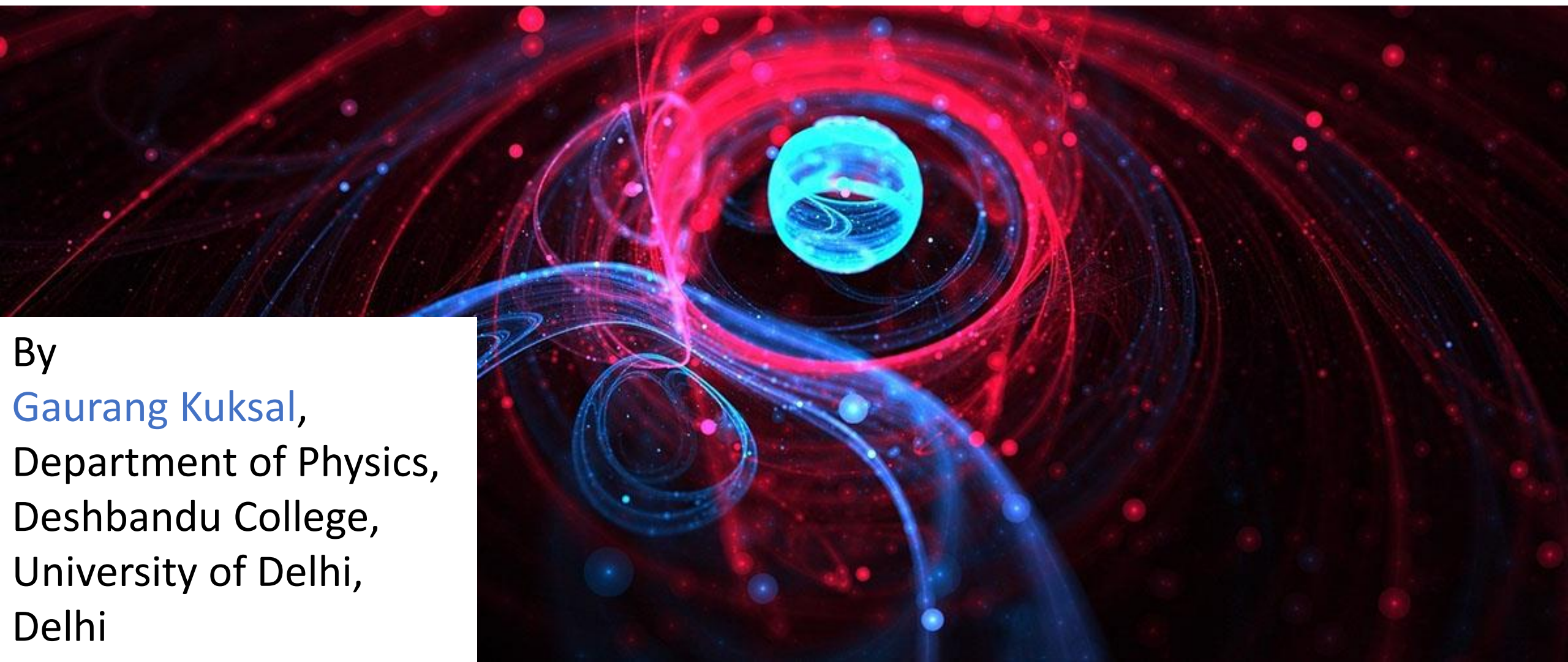


EMISSION OF TWO GAMMA RAY PHOTON FROM QUARK-GLUON PLASMA WITH CHEMICAL POTENTIAL

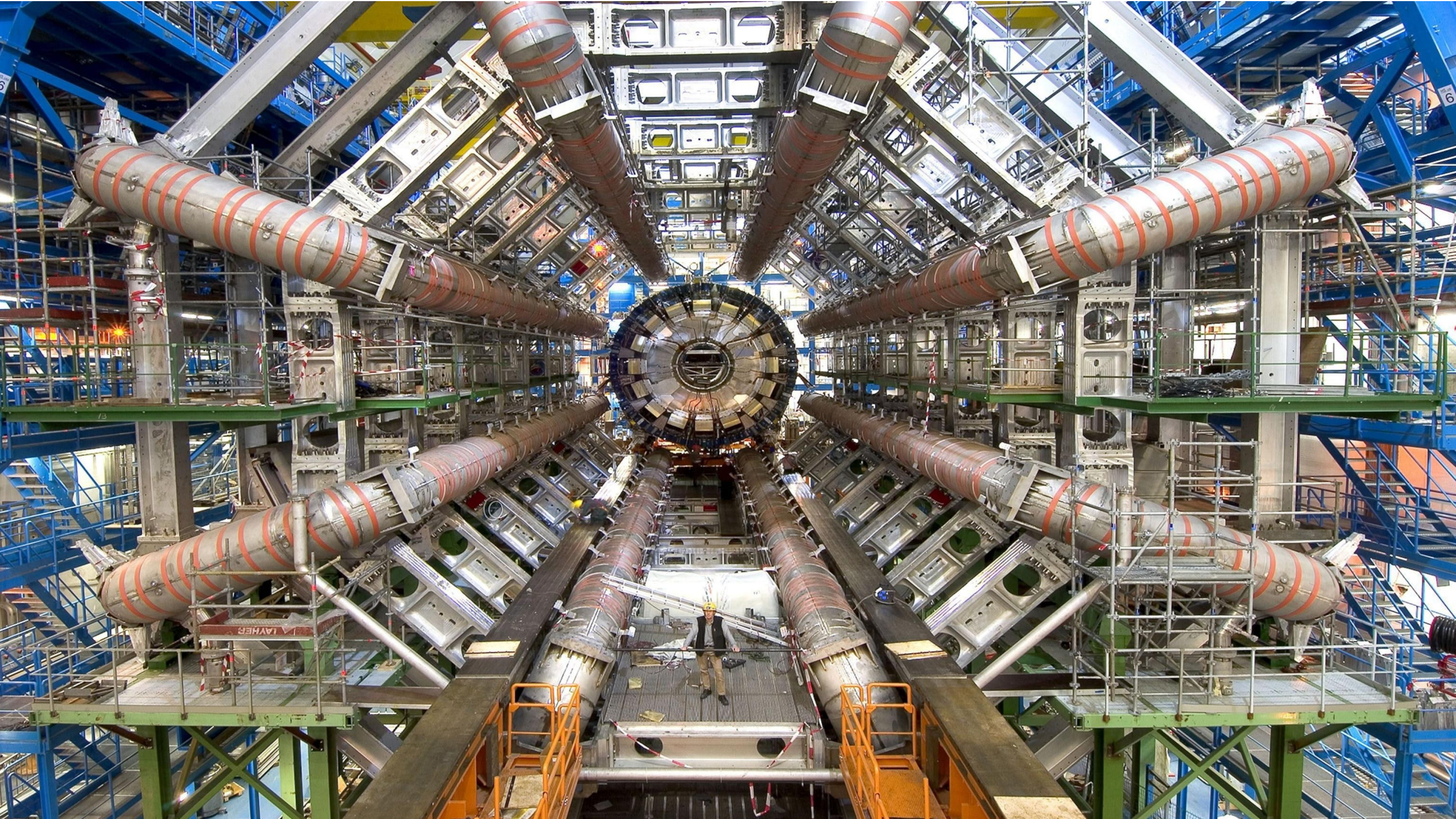


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Delhi

INTRODUCTION



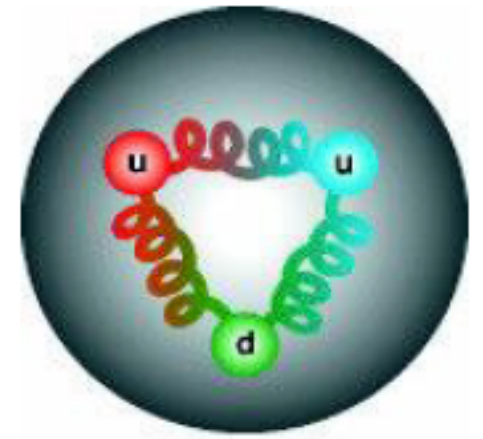
<div> <div>QUARKS</div> <div>LEPTONS</div> </div>	<div> <div>mass → $\approx 2.3 \text{ MeV}/c^2$</div> <div>charge → $2/3$</div> <div>spin → $1/2$</div> </div> <div> <div>u</div> <div>up</div> </div>	<div> <div>mass → $\approx 1.275 \text{ GeV}/c^2$</div> <div>charge → $2/3$</div> <div>spin → $1/2$</div> </div> <div> <div>c</div> <div>charm</div> </div>	<div> <div>mass → $\approx 173.07 \text{ GeV}/c^2$</div> <div>charge → $2/3$</div> <div>spin → $1/2$</div> </div> <div> <div>t</div> <div>top</div> </div>	<div> <div>0</div> <div>0</div> <div>1</div> </div> <div> <div>g</div> <div>gluon</div> </div>	<div> <div>mass → $\approx 126 \text{ GeV}/c^2$</div> <div>charge → 0</div> <div>spin → 0</div> </div> <div> <div>H</div> <div>Higgs boson</div> </div>
	<div> <div>mass → $\approx 4.8 \text{ MeV}/c^2$</div> <div>charge → $-1/3$</div> <div>spin → $1/2$</div> </div> <div> <div>d</div> <div>down</div> </div>	<div> <div>mass → $\approx 95 \text{ MeV}/c^2$</div> <div>charge → $-1/3$</div> <div>spin → $1/2$</div> </div> <div> <div>s</div> <div>strange</div> </div>	<div> <div>mass → $\approx 4.18 \text{ GeV}/c^2$</div> <div>charge → $-1/3$</div> <div>spin → $1/2$</div> </div> <div> <div>b</div> <div>bottom</div> </div>	<div> <div>0</div> <div>0</div> <div>1</div> </div> <div> <div>γ</div> <div>photon</div> </div>	
	<div> <div>mass → $0.511 \text{ MeV}/c^2$</div> <div>charge → -1</div> <div>spin → $1/2$</div> </div> <div> <div>e</div> <div>electron</div> </div>	<div> <div>mass → $105.7 \text{ MeV}/c^2$</div> <div>charge → -1</div> <div>spin → $1/2$</div> </div> <div> <div>μ</div> <div>muon</div> </div>	<div> <div>mass → $1.777 \text{ GeV}/c^2$</div> <div>charge → -1</div> <div>spin → $1/2$</div> </div> <div> <div>τ</div> <div>tau</div> </div>	<div> <div>mass → $91.2 \text{ GeV}/c^2$</div> <div>charge → 0</div> <div>spin → 1</div> </div> <div> <div>Z</div> <div>Z boson</div> </div>	
	<div> <div>mass → $< 2.2 \text{ eV}/c^2$</div> <div>charge → 0</div> <div>spin → $1/2$</div> </div> <div> <div>ν_e</div> <div>electron neutrino</div> </div>	<div> <div>mass → $< 0.17 \text{ MeV}/c^2$</div> <div>charge → 0</div> <div>spin → $1/2$</div> </div> <div> <div>ν_μ</div> <div>muon neutrino</div> </div>	<div> <div>mass → $< 15.5 \text{ MeV}/c^2$</div> <div>charge → 0</div> <div>spin → $1/2$</div> </div> <div> <div>ν_τ</div> <div>tau neutrino</div> </div>	<div> <div>mass → $80.4 \text{ GeV}/c^2$</div> <div>charge → ± 1</div> <div>spin → 1</div> </div> <div> <div>W</div> <div>W boson</div> </div>	<div>GAUGE BOSONS</div>



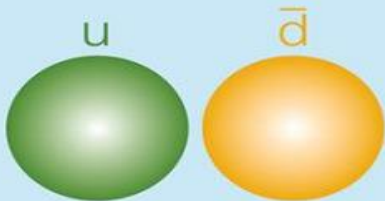
What is Quark-Gluon Plasma?

At room temperature, quarks and gluons are always confined inside **colorless** objects (**hadrons**):

protons, neutrons, pions,

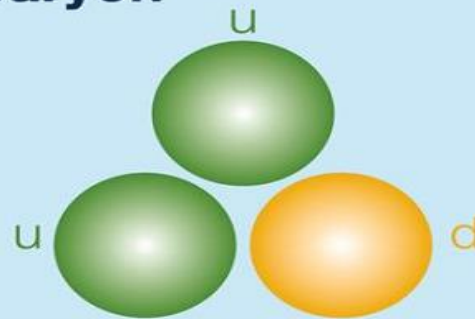


meson



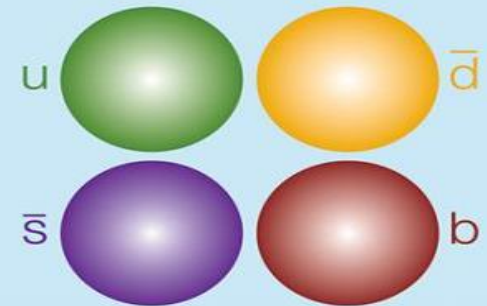
Mesons are made
of two quarks

baryon



Baryons are made
of three quarks

tetraquark



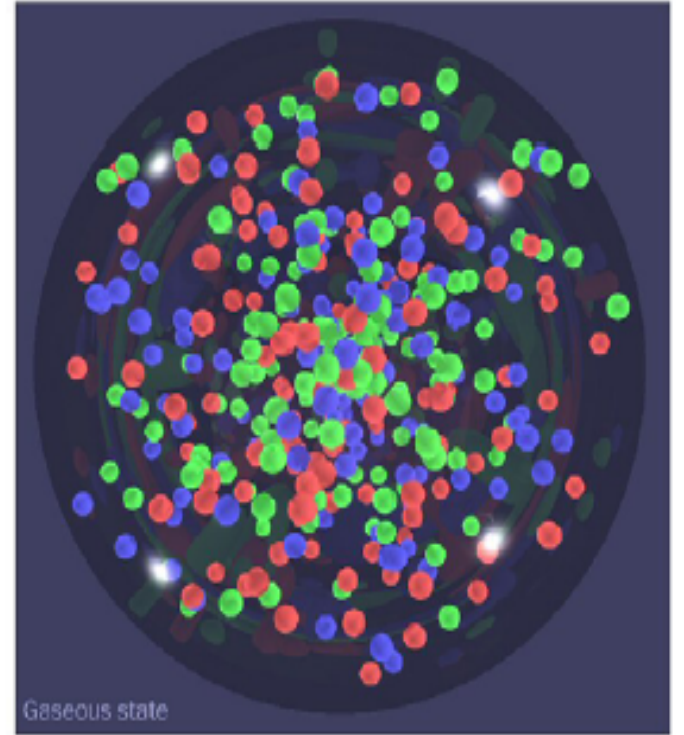
Tetraquarks are
made of four quarks

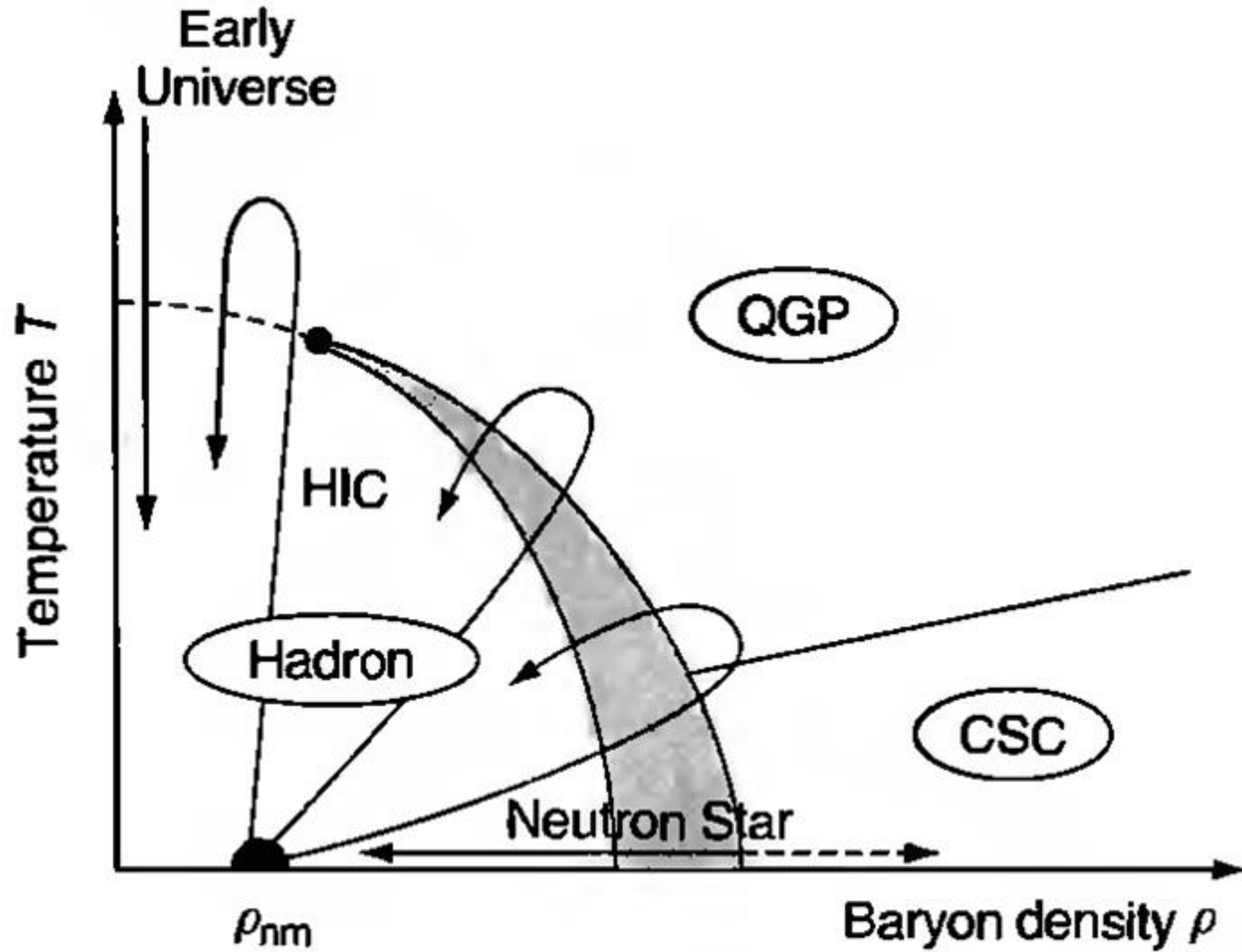
Very high temperature (asymptotic freedom):

- Interactions become weak
- quarks and gluons **deconfined**
- **Quark-gluon plasma (QGP)**

Infinitely high temperature:

QGP may behave like an **ideal gas**.

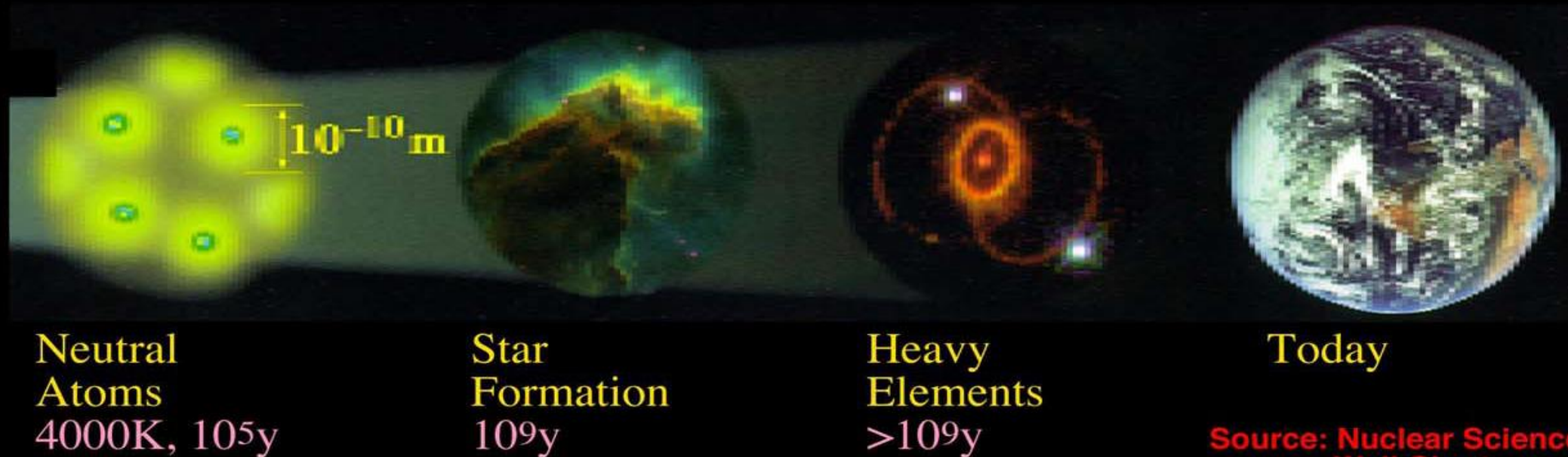
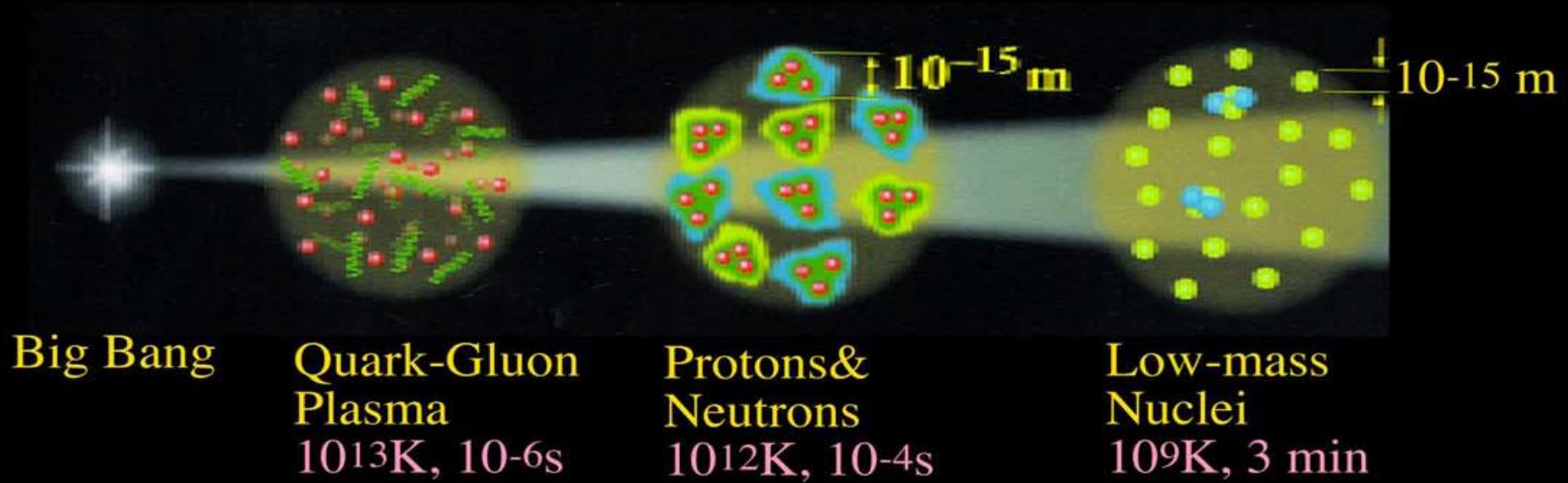




QGP and the Universe



History of the Universe

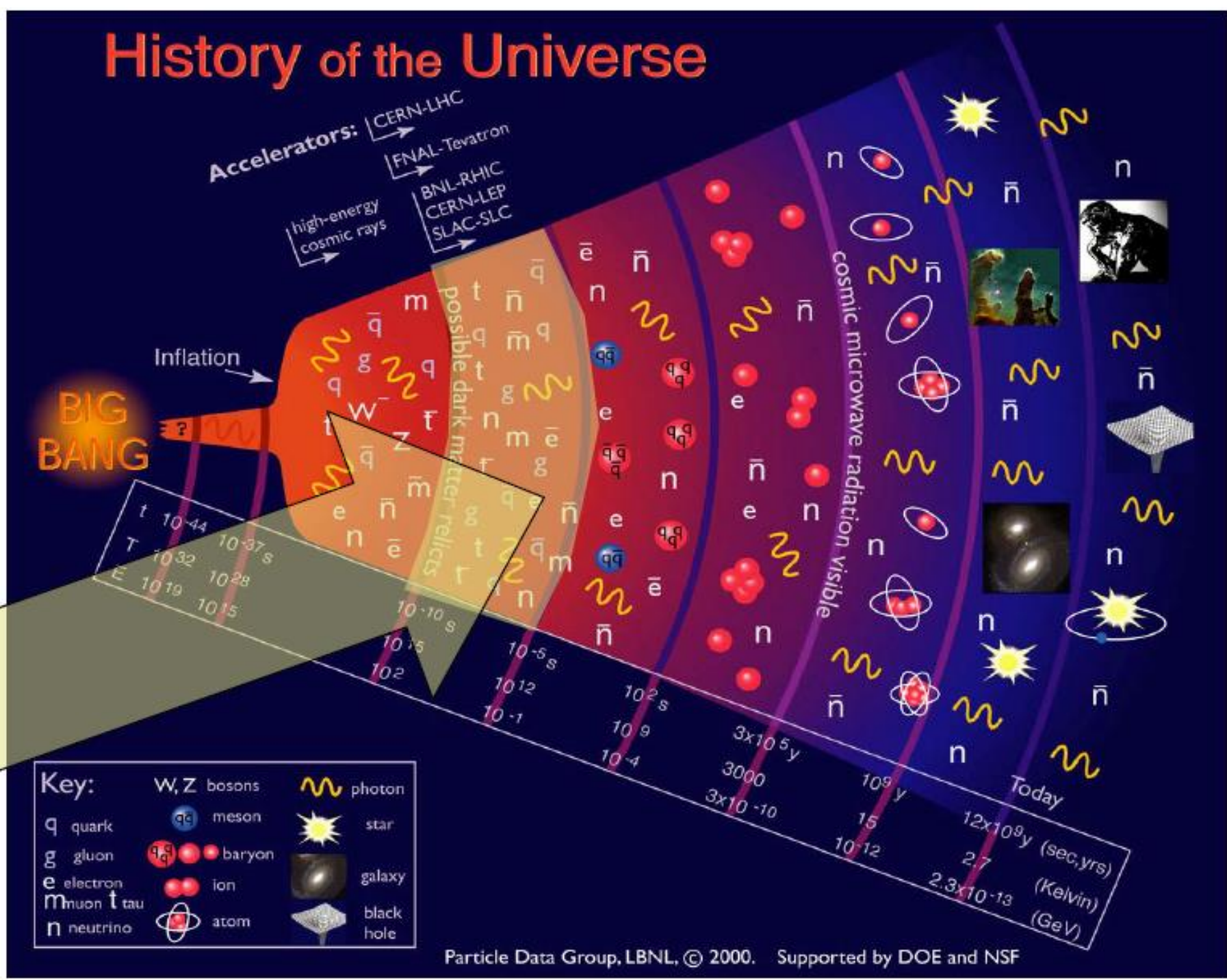


Source: Nuclear Science
Wall Chart

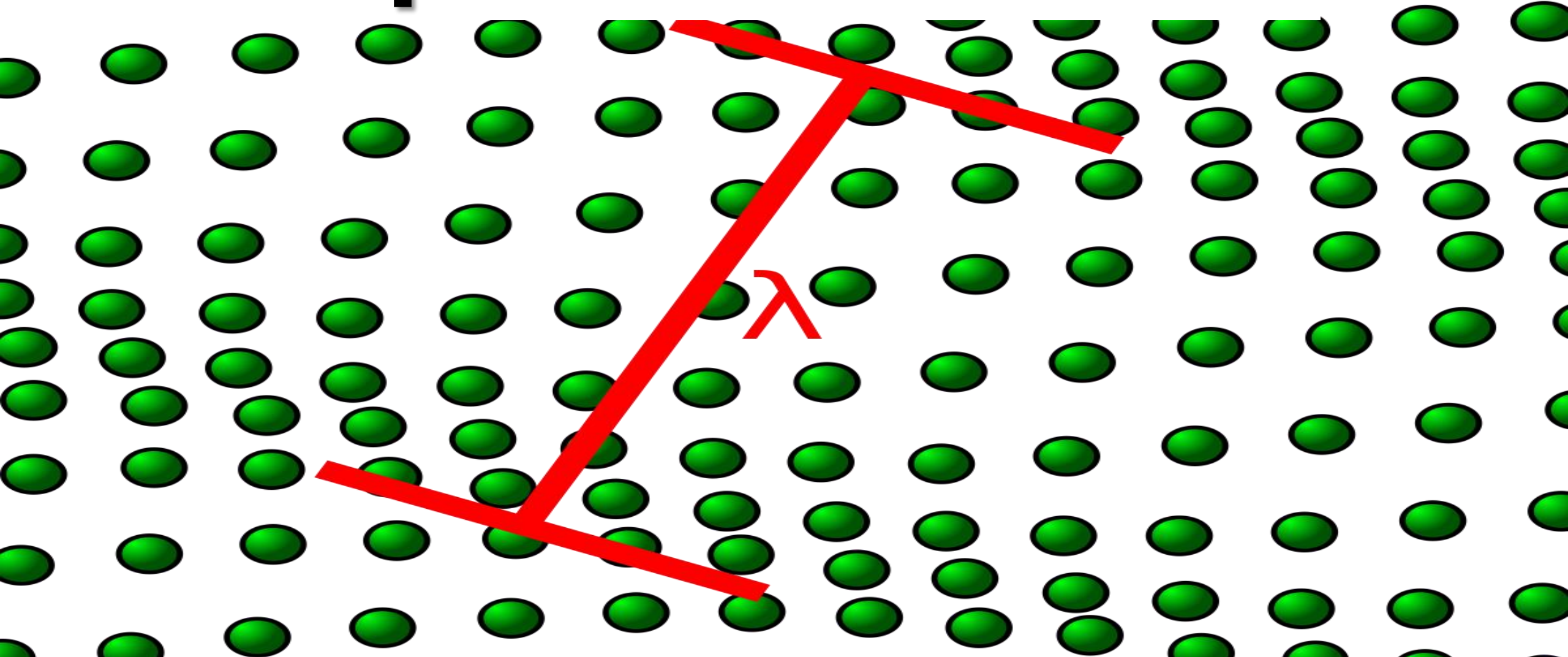
History of the Universe

- **Energy density**
 - $1 \text{ GeV/fm}^3 \rightarrow$
 - $1.8 \times 10^{15} \text{ g/cm}^3$
- **Temperature**
 - $170 \text{ MeV} \rightarrow$
 - $2.0 \times 10^{12} \text{ K}$

Conditions that prevailed $\approx 10 \mu\text{s}$ after the Big Bang



Quasi-particle model



Proposed by Golviznin and Satz and then by Peshier et. al. ,
V.M.Bannur, Agotiya, Chandra to explain the EOS of QGP
(Quasiparticles are ubiquitous in finite-temperature calculations H
A Weldon).

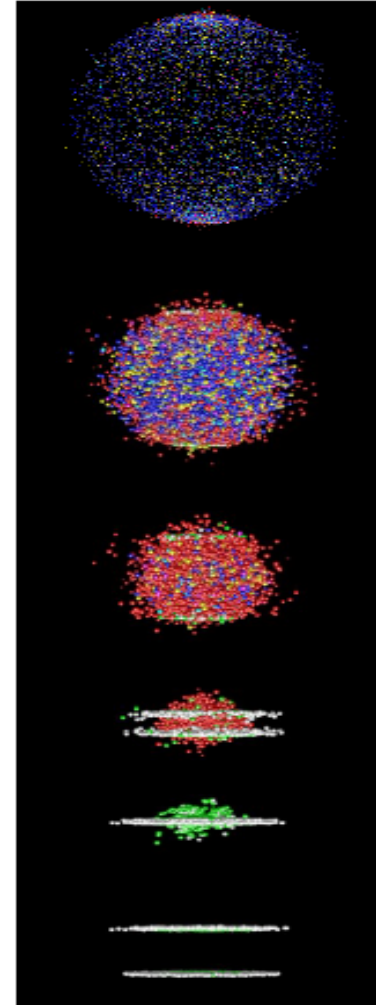
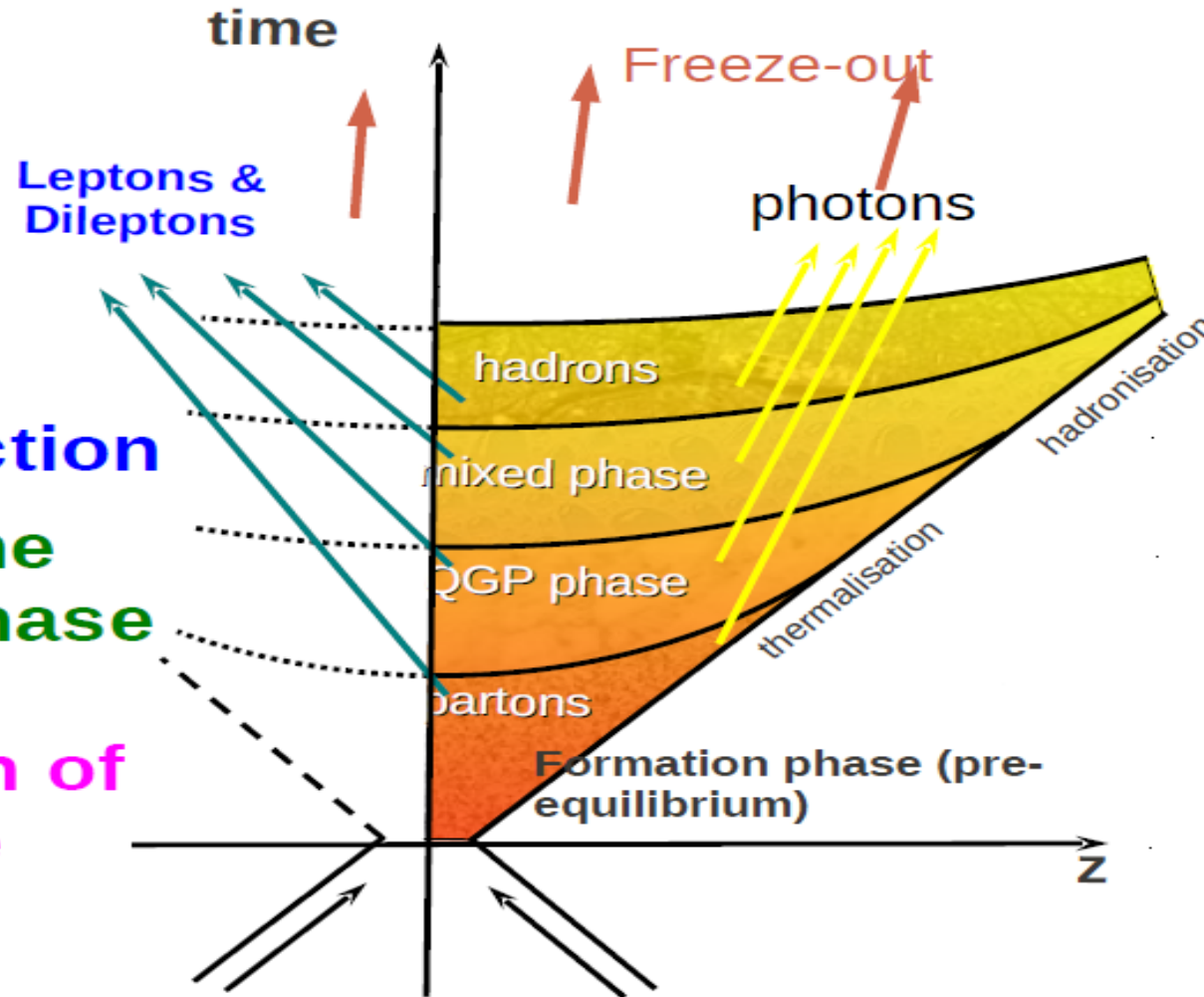
- Quarks and gluons in QGP are not bare quarks and gluons, but they are quasi-particles.
- Quasiparticles are the thermal excitations of the interacting quarks and gluons.
- Due to collective behaviour in QGP, massless partons acquire masses equal to their respective plasma frequencies and become quasi-partons.
- QGP is made up of non-interacting quasi-partons.

OUR WORK



Electromagnetic radiations

- ➔ Why studying electromagnetic radiations ?
- ★ Electromagnetic probes do not interact strongly
- ★ Suffer little or no final state interaction
- ★ Give access to the hot and dense phase of the reaction
- ★ Carry information of the system at the time of their production



Direct photons

Photons in the QGP-phase are mainly produced via the following processes:

$$q + \bar{q} \rightarrow \gamma + g$$

Annihilation + Compton effect

$$g + q \rightarrow \gamma + q$$

$$q + \bar{q} \rightarrow \gamma + \gamma$$

Mainly produced in QGP state

We have to deal with background problems because of hadronic decays into photons, especially by

$$\pi^0 \rightarrow \gamma + \gamma \quad ; \quad \pi^+ + \pi^- \rightarrow \gamma + \gamma \quad \text{and} \quad \eta \rightarrow \gamma + \gamma (\text{or } 3\pi^0)$$

In principle, these photons can be distinguished by invariant mass analysis. But this is very difficult, if 100 neutral pions or more are produced in the reaction.

The diphoton production rate is usually comes from the processes:



The number of diphotons emitted per space-time volume per invariant mass in the quark gluon plasma phase is expressed as:

$$\frac{dN_{\gamma\gamma}}{d^4x dM} = \frac{1}{(2\pi)^4} (M^2 - 4m_q^2)^{1/2} M^2 \lambda_q^2 T^2 \sigma_{q\bar{q} \rightarrow \gamma\gamma} (M^2) \left(\frac{M}{T}\right) K_1\left(\frac{M}{T}\right)$$

$$\text{where, } \lambda_q = e^{\frac{\mu}{T}}$$

Diphoton mass spectrum can be written as:

$$\frac{dN_{\gamma\gamma}}{dM^2 dy} = \frac{1}{2(2\pi)^4} \times \pi R^2 M^3 \int_{\tau_i}^{\tau_f} \sigma_{q\bar{q} \rightarrow \gamma\gamma} (M^2) \left(1 - 4m_q^2 / M^2\right)^{1/2} T(\tau) K_1\left(\frac{M}{T}\right) \tau d\tau$$

Electromagnetic radiations

I) Photons:

Testing pQCD

Thermal photons

II) Low Mass Region:

Vector mesons in medium

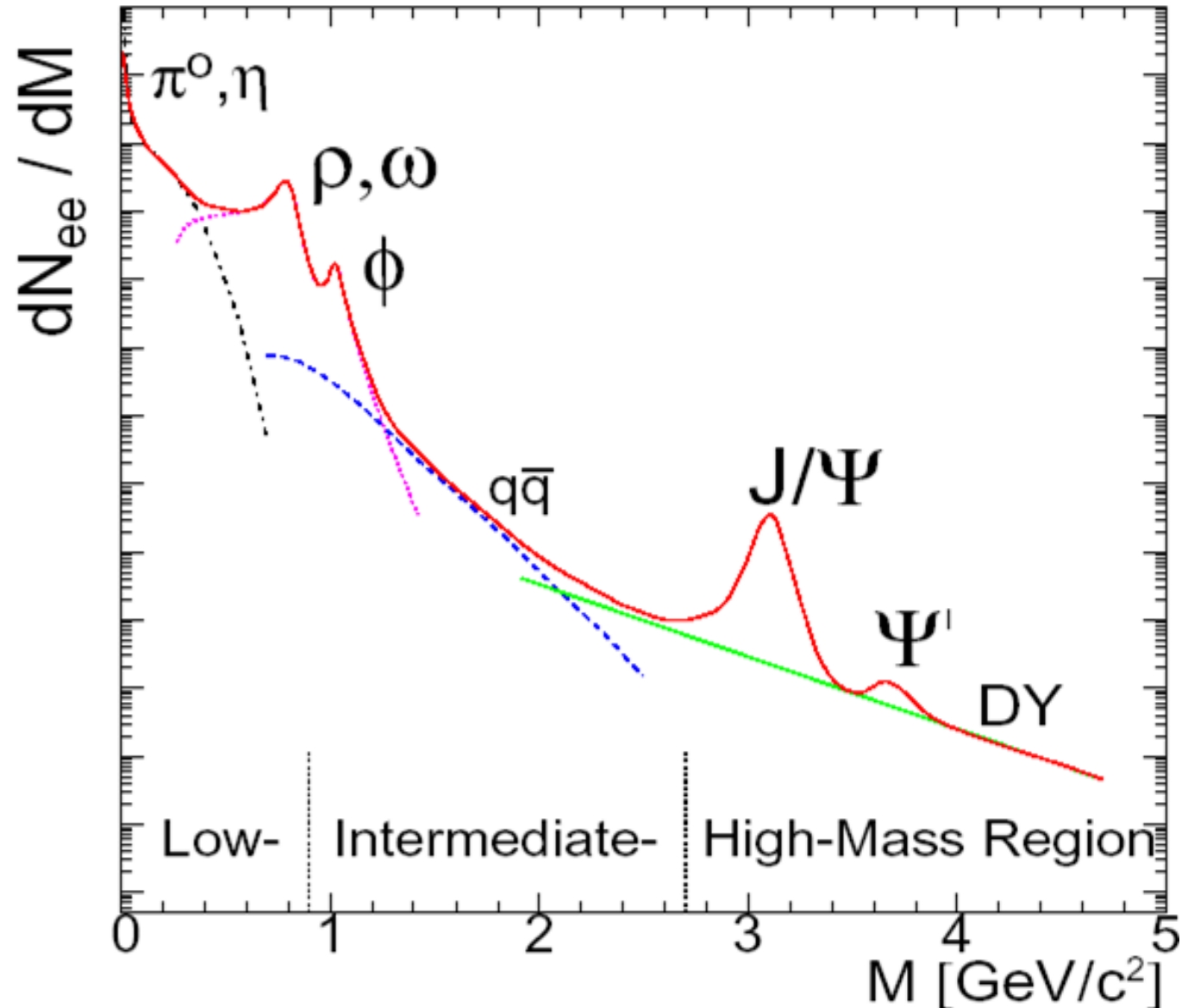
III) Intermediate Mass Region:

Thermal QGP

Open charm

IV) High Mass Region:

Heavy quarks resonances

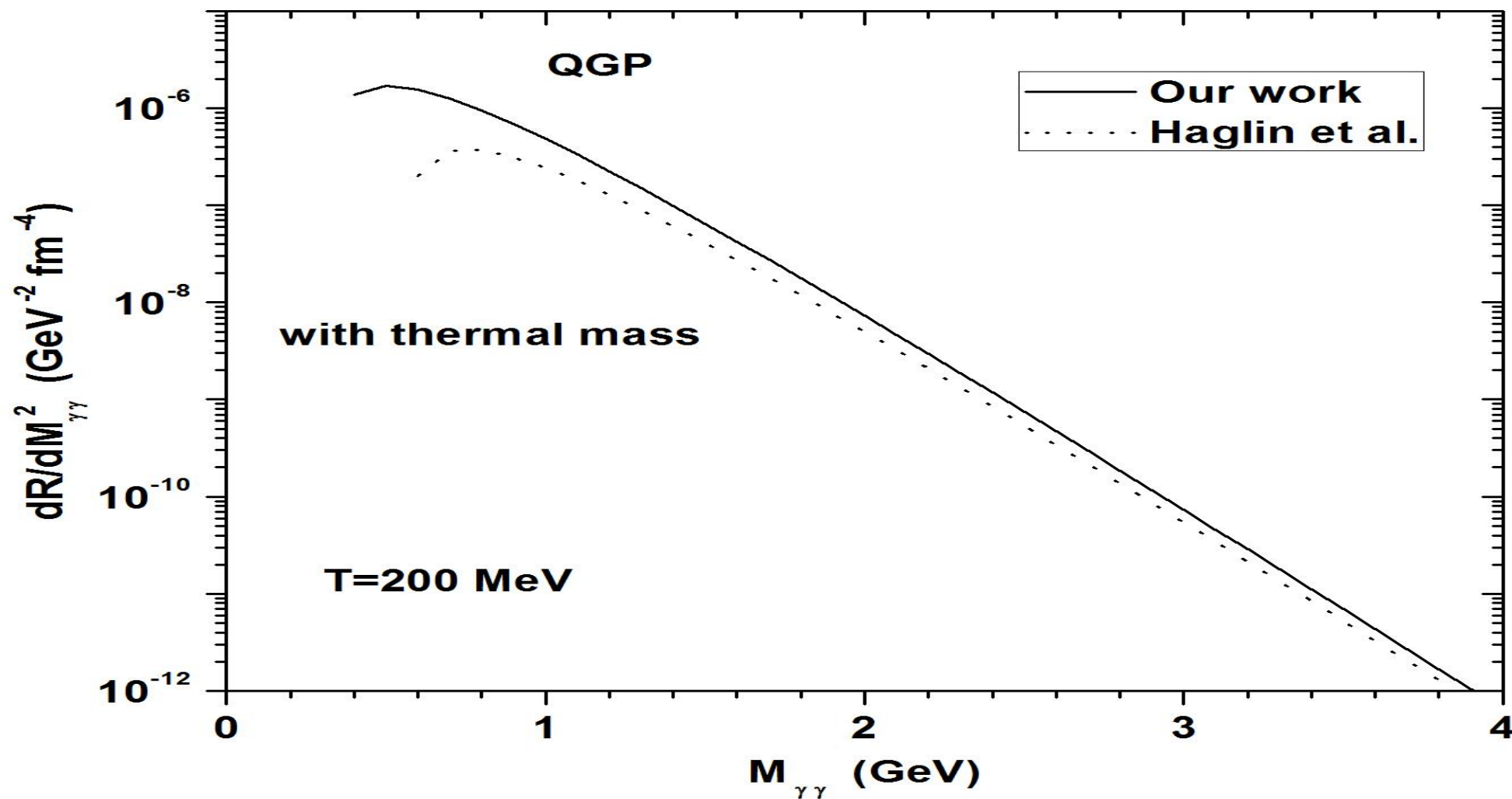


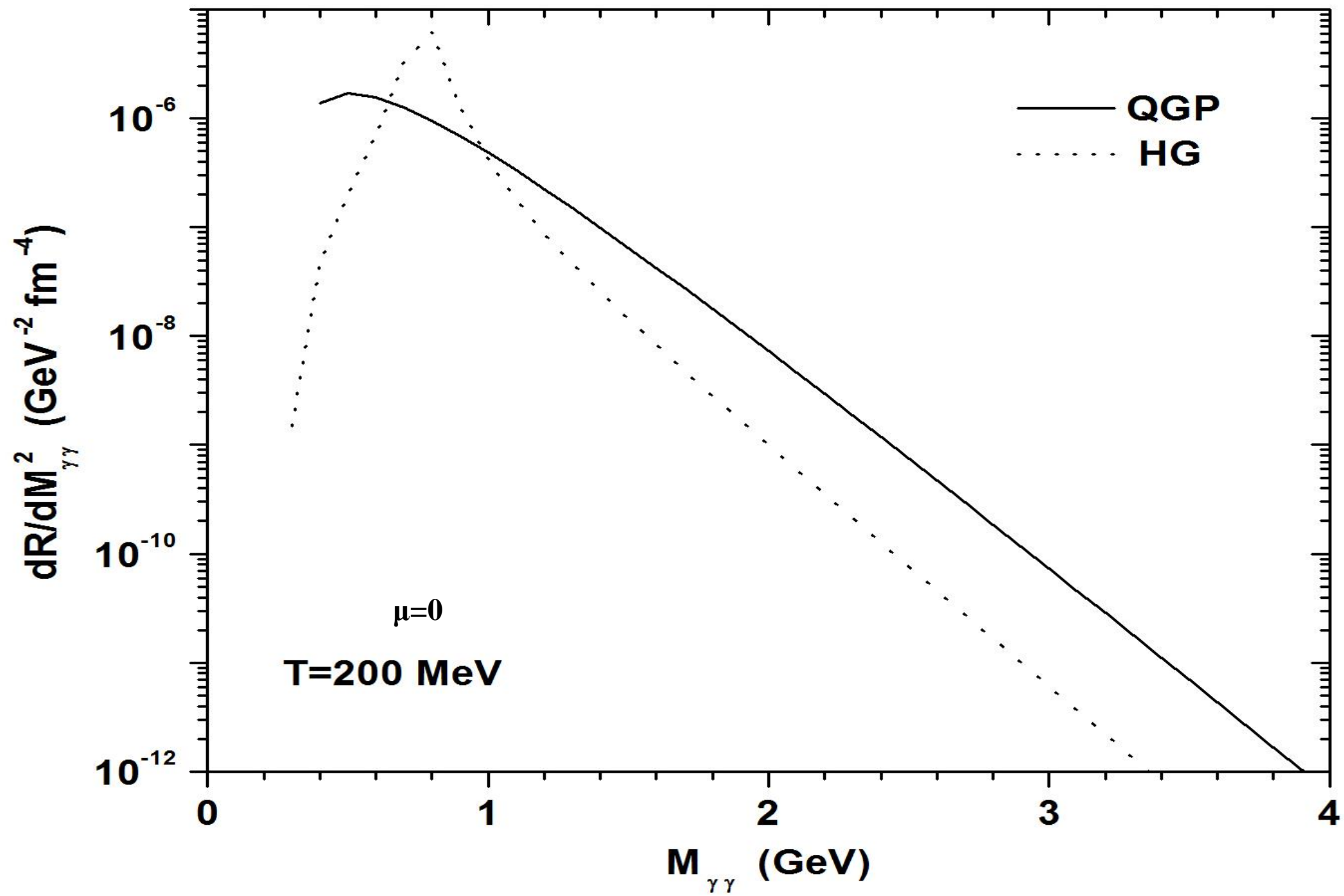
RESULTS



Diphoton mass spectrum :

$\mu=0$





Effect of chemical potential

All our previous calculations were based on the assumption of vanishing chemical potential for the quarks.

The changes in result by introducing finite chemical potentials, by making following change

from $-\frac{E}{T}$ to $\frac{\mu - E}{T}$

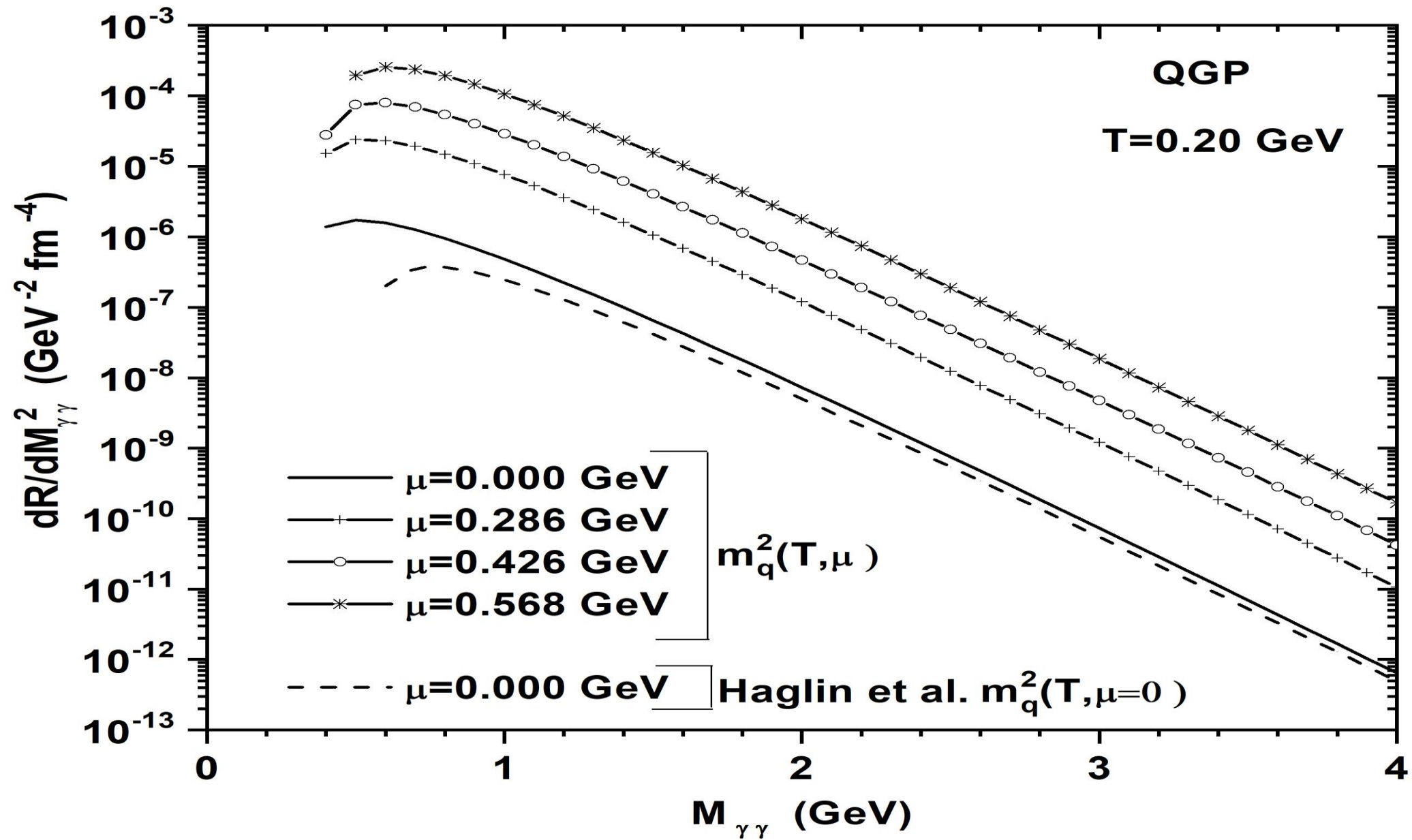
where 'μ' is the potential.

Finite value of quark mass depend on temperature and chemical potential (Quasi-Particle Model):

$$m_q^2(T, \mu) = \gamma_q g^2 T^2 \left[1 + \frac{\mu^2}{\pi^2 T^2} \right]$$

where, μ is the chemical potential.

*Use: Peshier et al., Bannur et al., Kumar et al.,
Srivastava et al. and many more.*



Summary

- . Two photon decays using a phenomenological model emitted from a quark-gluon plasma at high temperature and chemical potential is studied.**
- . We used a similar idea of our earlier calculations of photon production rate in QGP and hadronic state at RHIC and LHC. For the phenomenology of heavy-ion collisions, we present emission spectra by taking into account the parametrization factors in thermal dependent quark mass.**
- . The quasi-model of quark mass is used in such a way that the diphoton emission is reviewed and significant with quark mass dependent on temperature and chemical potential. This indicates that quark mass plays a pivoted role in the diphoton production rate.**

THANK YOU

