

Dualities of two-color QCD phase diagram



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Russian
Science
Foundation



Фонд развития
теоретической физики
и математики

K.G. Klimenko, IHEP
T.G. Khunjua, University of Georgia, MSU

in the broad sense our group stems from
Department of Theoretical Physics, Moscow State University
Prof. V. Ch. Zhukovsky

details can be found in

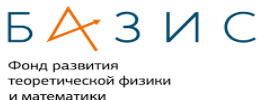
Eur.Phys.J.C 80 (2020) 10, 995 arXiv:2005.05488 [hep-ph]
JHEP 06 (2020) 148 arXiv:2003.10562 [hep-ph]
Phys.Rev. D100 (2019) no.3, 034009 arXiv: 1904.07151 [hep-ph]
JHEP 1906 (2019) 006 arXiv:1901.02855 [hep-ph]
Eur.Phys.J. C79 (2019) no.2, 151, arXiv:1812.00772 [hep-ph],
Phys.Rev. D98 (2018) no.5, 054030 arXiv:1804.01014 [hep-ph],
Phys.Rev. D97 (2018) no.5, 054036 arXiv:1710.09706 [hep-ph]
Phys.Rev. D95 (2017) no.10, 105010 arXiv:1704.01477 [hep-ph]

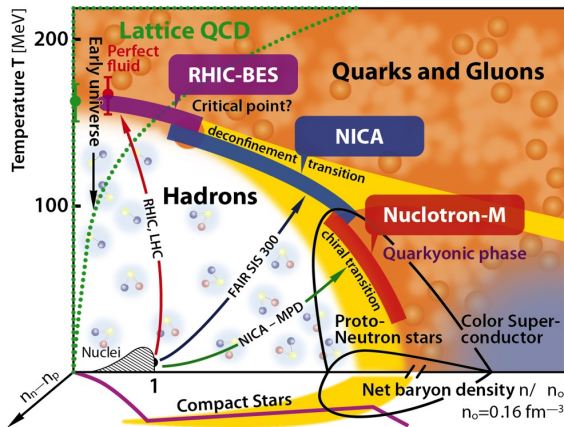
The work is supported by

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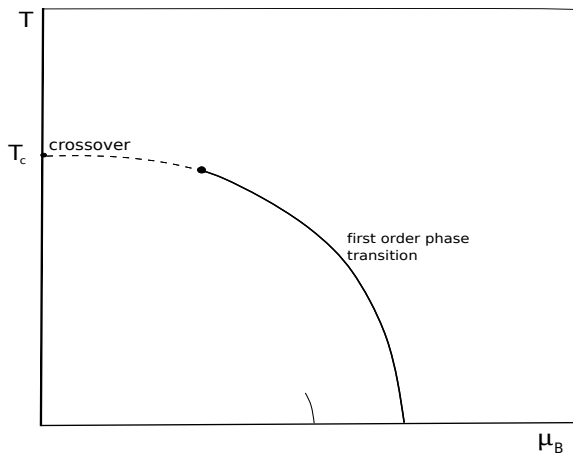
- ▶ Foundation for the Advancement of
Theoretical Physics and Mathematics





Two main phase transitions

- ▶ confinement-deconfinement
- ▶ chiral symmetry breaking phase—chiral symmetric phase

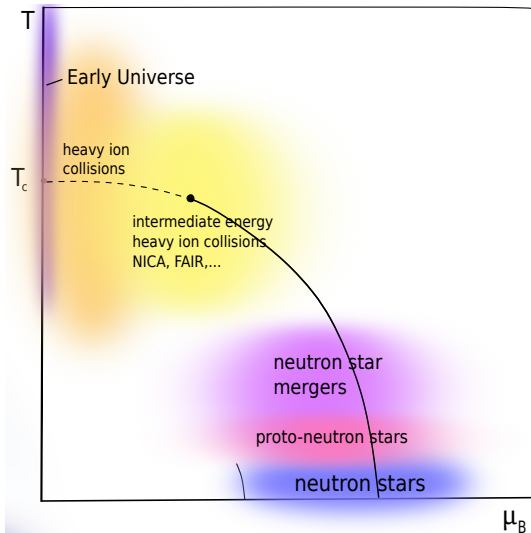


Two main phase transitions

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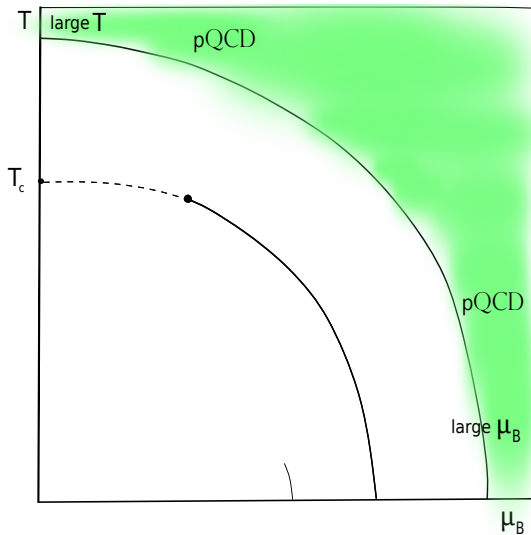
QCD at T and μ
(QCD at extreme conditions)

- ▶ Early Universe
- ▶ heavy ion collisions
- ▶ neutron stars
- ▶ proto- neutron stars
- ▶ neutron star mergers



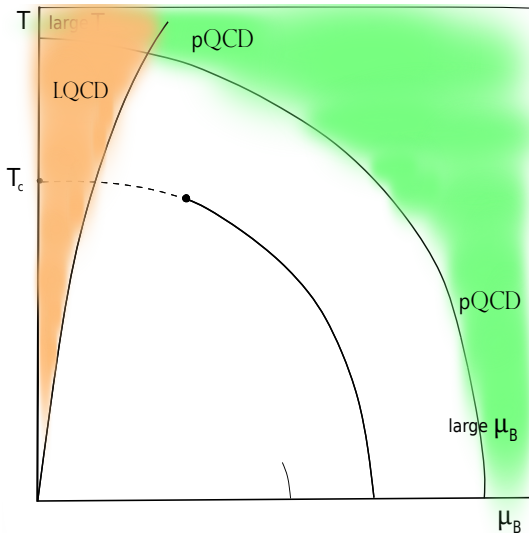
Methods of dealing with QCD

► Perturbative QCD



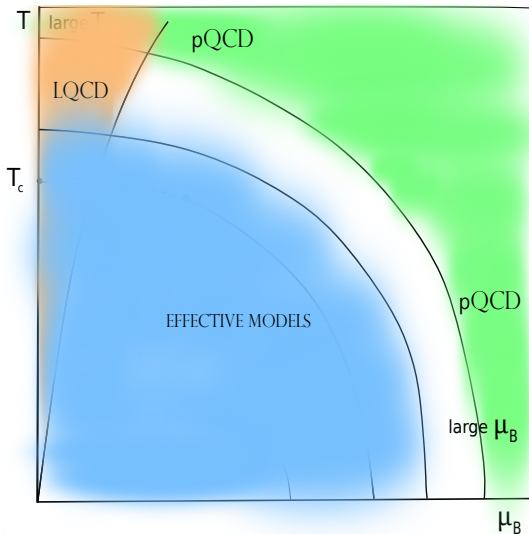
Methods of dealing with QCD

- ▶ Perturbative QCD
- ▶ First principle calculation
– lattice QCD



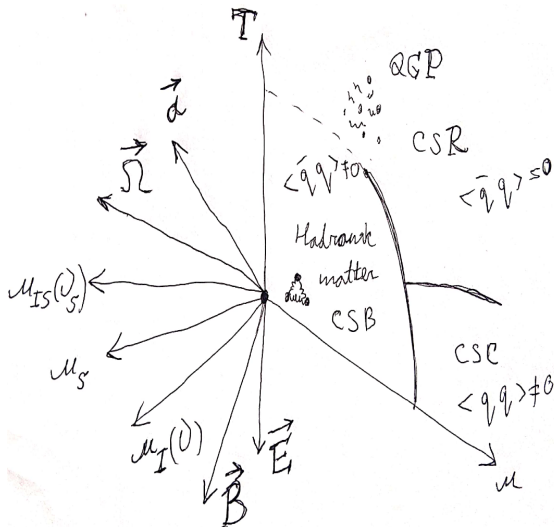
Methods of dealing with QCD

- ▶ Perturbative QCD
- ▶ First principle calculation
– lattice QCD
- ▶ Effective models
- ▶ DSE, FRG
- ▶



More than just QCD at (μ, T)

- ▶ more chemical potentials μ_i
- ▶ magnetic fields
- ▶ rotation of the system $\vec{\Omega}$
- ▶ acceleration \vec{a}
- ▶ finite size effects (finite volume and boundary conditions)



Baryon chemical potential μ_B

Allow to consider systems with non-zero baryon densities.

$$\frac{\mu_B}{3} \bar{q} \gamma^0 q = \mu \bar{q} \gamma^0 q, \quad n_B = \frac{1}{3} (n_u + n_d)$$

Baryon chemical potential μ_B

Allow to consider systems with non-zero baryon densities.

$$\frac{\mu_B}{3} \bar{q} \gamma^0 q = \mu \bar{q} \gamma^0 q, \quad n_B = \frac{1}{3} (n_u + n_d)$$

Isotopic chemical potential μ_I

Allow to consider systems with isospin imbalance ($n_n \neq n_p$).

$$\frac{\mu_I}{2} \bar{q} \gamma^0 \tau_3 q = \nu (\bar{q} \gamma^0 \tau_3 q)$$

$$n_I = n_u - n_d \quad \longleftrightarrow \quad \mu_I = \mu_u - \mu_d$$

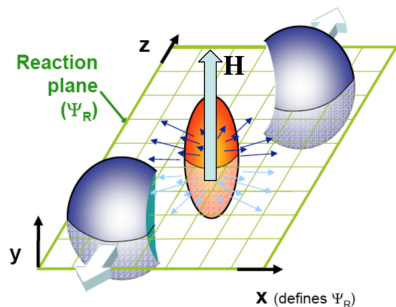
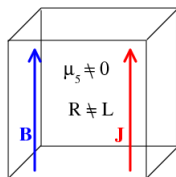
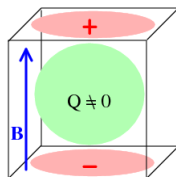
chiral (axial) chemical potential

Allow to consider systems with chiral imbalance (difference between densities of left-handed and right-handed quarks).

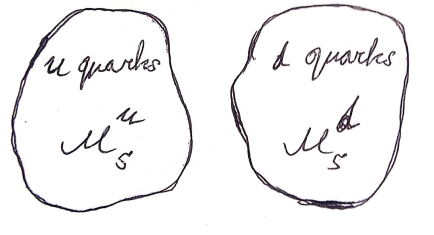
$$n_5 = n_R - n_L \quad \longleftrightarrow \quad \mu_5 = \mu_R - \mu_L$$

The corresponding term in the Lagrangian is

$$\mu_5 \bar{q} \gamma^0 \gamma^5 q$$



$$\vec{J} \sim \mu_5 \vec{B},$$



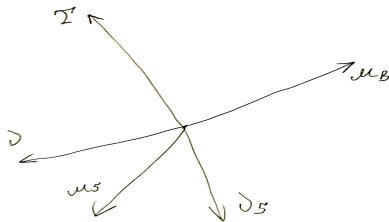
$$\mu_5^u \neq \mu_5^d \quad \text{and} \quad \mu_{I5} = \mu_5^u - \mu_5^d$$

Term in the Lagrangian — $\frac{\mu_{I5}}{2} \bar{q} \tau_3 \gamma^0 \gamma^5 q = \nu_5 (\bar{q} \tau_3 \gamma^0 \gamma^5 q)$

$$n_{I5} = n_{u5} - n_{d5}, \quad n_{I5} \longleftrightarrow \nu_5$$

Different chemical potentials and matter content

$$\mu = \frac{\mu_B}{3}, \quad \nu = \frac{\mu_I}{2}, \quad \mu_5, \quad \nu_5 = \frac{\mu_{I5}}{2}$$



Dualities

It is not related to holography or gauge/gravity
duality

it is the dualities of the phase structures of
different systems

Dualities

Chiral symmetry breaking \iff pion condensation

Isospin imbalance \iff Chiral imbalance

The TDP

$$\Omega(T, \mu, \mu_i, \dots, \langle \bar{q}q \rangle, \dots)$$

The TDP

$$\Omega(T, \mu, \mu_i, \dots, \langle \bar{q}q \rangle, \dots)$$

$$\Omega(T, \mu, \nu, \nu_5, \dots, M, \pi, \dots)$$

The TDP

$$\Omega(T, \mu, \mu_i, \dots, \langle \bar{q}q \rangle, \dots) \qquad \Omega(T, \mu, \nu, \nu_5, \dots, M, \pi, \dots)$$

The TDP (phase daigram) is invariant under
Interchange of - condensates - matter content

$$\Omega(M, \pi, \nu, \nu_5)$$

$$M \longleftrightarrow \pi, \qquad \nu \longleftrightarrow \nu_5$$

$$\Omega(M, \pi, \nu, \nu_5) = \Omega(\pi, M, \nu_5, \nu)$$

$\mu_B \neq 0$ impossible on lattice but if $\mu_B = 0$

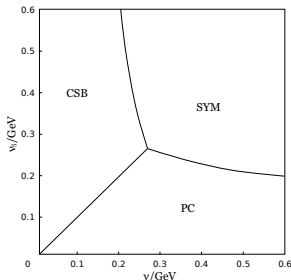
Duality is shown to take place in particular cases in lattice QCD

► **QCD at μ_5** — (μ_5, T)

V. Braguta, A. Kotov et al, ITEP lattice group

► **QCD at μ_I** — (μ_I, T)

G. Endrodi, B. Brandt et al, Emmy Noether junior research group, Goethe-University Frankfurt, Institute for Theoretical Physics ()



Two colour QCD case

QC_2D

There are a lot similarities:

- ▶ similar phase transitions:

*confinement/deconfinement, chiral symmetry
breaking/restoration at large T and μ*

- ▶ A lot of physical quantities coincide up to few dozens percent

*Critical temperature $T_c/\sqrt{\sigma}$, topological susceptibility
 $\chi^{\frac{1}{4}}/\sqrt{\sigma}$ shear viscosity η/s*

There are **no sign problem** in SU(2) case

$$(Det(D(\mu)))^\dagger = Det(D(\mu))$$

and lattice simulations at non-zero baryon density are possible

It is a great playground for studying dense matter

Phase diagram of QC_2D

Condensates and phases

$$M = \langle \sigma(x) \rangle \sim \langle \bar{q}q \rangle,$$

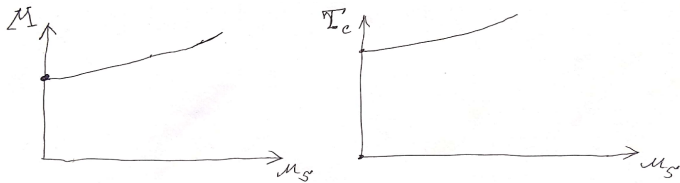
CSB phase: $M \neq 0$,

$$\pi_1 = \langle \pi_1(x) \rangle = \langle \bar{q}\gamma^5\tau_1q \rangle,$$

PC phase: $\pi_1 \neq 0$,

$$\Delta = \langle \Delta(x) \rangle = \langle qq \rangle = \langle q^T C \gamma^5 \sigma_2 \tau_2 q \rangle,$$

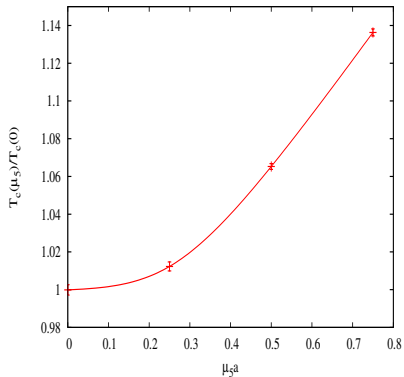
BSF phase: $\Delta \neq 0$.

QCD at non-zero μ_5 

catalysis of CSB by chiral imbalance:

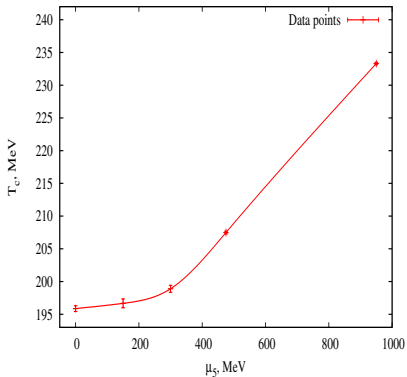
- ▶ increase of $\langle \bar{q}q \rangle$ as μ_5 increases
- ▶ increase of critical temperature T_c of chiral phase transition (crossover) as μ_5 increases

SU(2)



V. Braguta, A. Kotov et al, *JHEP* 1506, 094
(2015), PoS LATTICE 2014, 235 (2015)

SU(3)



V. Braguta, A. Kotov et al, *Phys. Rev. D* 93,
034509 (2016), arXiv:1512.05873 [hep-lat]

Instead of chiral symmetry

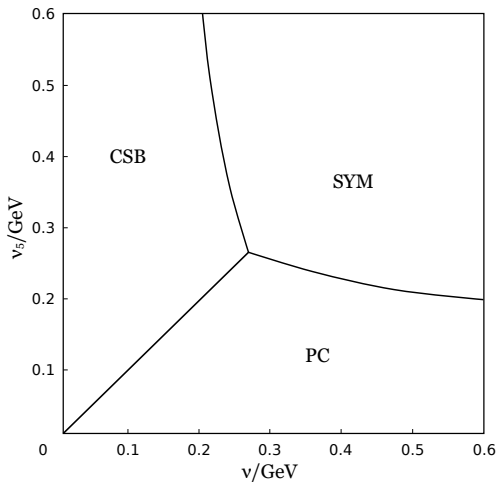
$$SU_L(2) \times SU_R(2)$$

one has Pauli-Gursey flavor symmetry

$$SU(4)$$

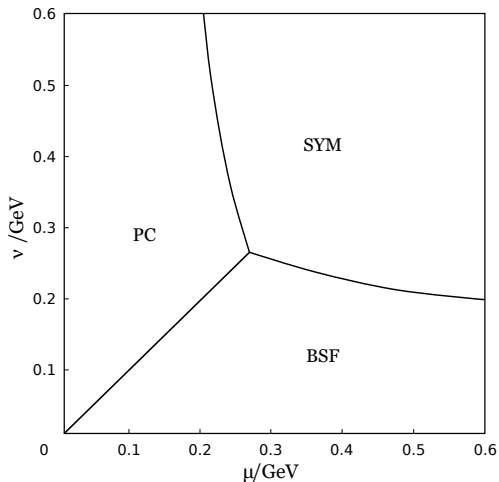
Two colour NJL model

$$L = \bar{q} \left[i\hat{\partial} - m_0 \right] q + H \left[(\bar{q}q)^2 + (\bar{q}i\gamma^5 \vec{\tau}q)^2 + (\bar{q}i\gamma^5 \sigma_2 \tau_2 q^c) (\bar{q}^c i\gamma^5 \sigma_2 \tau_2 q) \right]$$

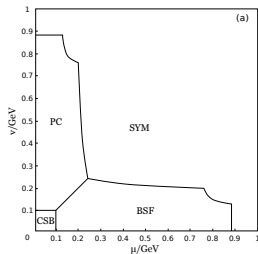


*Isospin imbalance ν
generates pion condensation*

*Chiral imbalance ν_5
generates chiral symmetry
breaking*



*Baryon density μ generates
diquark condensation*

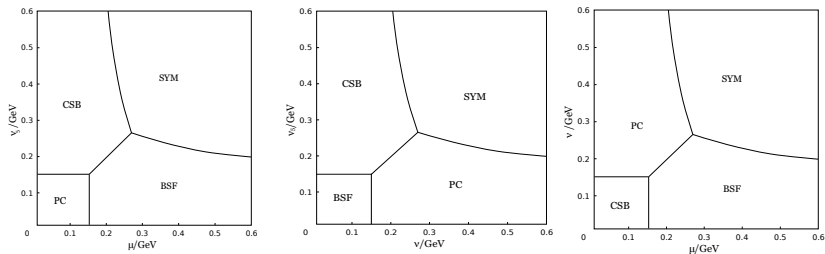


$$(a) \quad \mathcal{D}_1 : \quad \mu \longleftrightarrow \nu, \quad \pi_1 \longleftrightarrow |\Delta|, \quad PC \longleftrightarrow BSF$$

J. Andersen, T. Brauner, D. T. Son, M. Stephanov, J. Kogut, ...

$$(b) \quad \mathcal{D}_3 : \quad \nu \longleftrightarrow \nu_5, \quad M \longleftrightarrow \pi_1, \quad PC \longleftrightarrow CSB$$

$$(c) \quad \mathcal{D}_2 : \quad \mu \longleftrightarrow \nu_5, \quad M \longleftrightarrow |\Delta|, \quad CSB \longleftrightarrow BSF$$



- Phase diagram is **highly symmetric** due to **dualities**

The **whole phase diagram**, including diquark condensation, **in two color case** can be obtained from the results of behaviour of CSB and PC phases .

(μ, ν, ν_5) phase diagram is highly symmetric due to dualities

and intermingled by dualities

$$(\mu, \nu, \nu_5 \mid \mu_5)$$

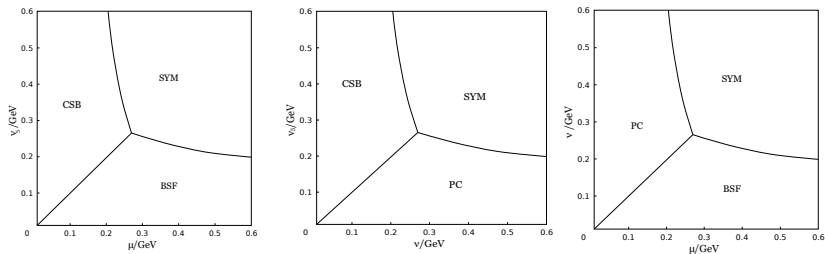
μ_5 stands alone from other μ, ν, ν_5 in QC₂D

μ_5 respects the dualities but it does not transform under them

μ_5 deforms the (μ, ν, ν_5) phase diagram

The influence of μ_5 is constrained by dual properties

- ▶ Chameleon nature of chiral imbalance μ_5 is also a consequence of dualities
 - ▶ The feature of μ_5 of being universal catalizer is a consequence of dualities as well
-

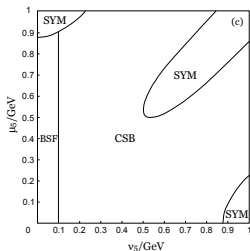
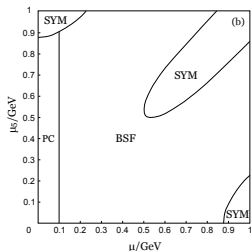
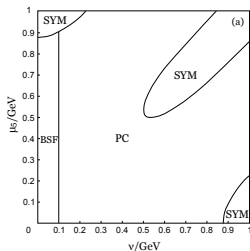


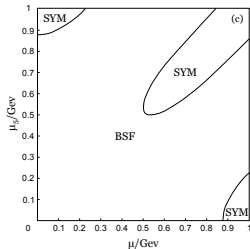
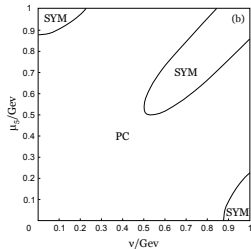
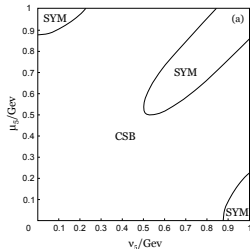
For the moderate values the phase diagram is quite concise

- ▶ Baryon density $\mu \iff$ diquark condensation
- ▶ Isospin imbalance $\nu \iff$ pion condensation
- ▶ Chiral imbalance $\nu_5 \iff$ chiral symmetry breaking

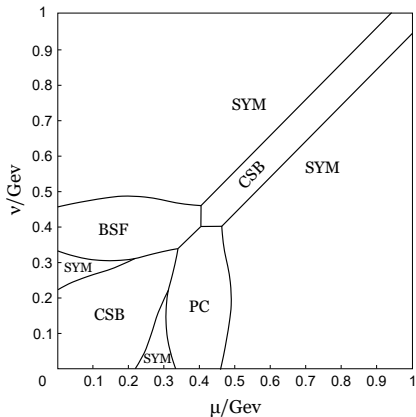
Chiral imbalance μ_5 catalyzes all the phenomena

Catalysis of CSB phenomenon was known as in $SU(3)$ as well as in $SU(2)$ cases

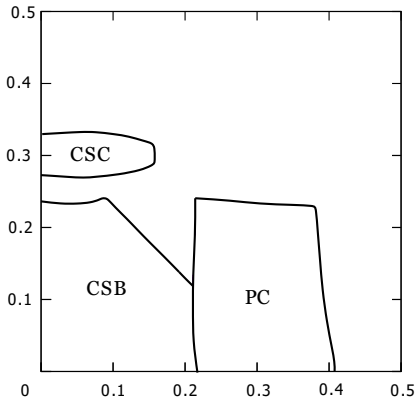


Chameleon nature of chiral imbalance μ_5 μ_5 mimics other chemical potentials μ , ν , ν_5 

Chiral imbalance μ_5 leads to several rather peculiar phases in the system, e. g. the **diquark condensation** in the region of the phase diagram at $\mu = 0$



Chiral imbalance μ_5 leads to the **diquark condensation** in the region of the phase diagram at $\mu = 0$ in three color case



Dualities \mathcal{D}_1 , \mathcal{D}_2 and \mathcal{D}_3 were found in

- In the framework of NJL model
 - In the mean field approximation
-

Dualities are connected with Pauli-Gursey group

Dualities were found in

- In the framework of NJL model
beyond mean field
- In QC_2D non-perturbatively (at the level of
Lagrangian)

Duality \mathcal{D} is a remnant of chiral symmetry

Duality was found in

- ▶ In the framework of NJL model beyond mean field or at all orders of N_c approximation
 - ▶ In QCD non-perturbatively (at the level of Lagrangian)
-

- ▶ $(\mu_B, \mu_I, \nu_5, \mu_5)$ phase diagram was studied in two color color case
- ▶ It was shown that there exist dualities in QCD and QC_2D
Richer structure of Dualities in the two colour case
- ▶ There have been shown ideas how dualities can be used
Duality is not just entertaining mathematical property but an instrument with very high predictivity power
- ▶ Dualities have been shown non-perturbatively in the two colour case
- ▶ Duality has been shown non-perturbatively in QCD