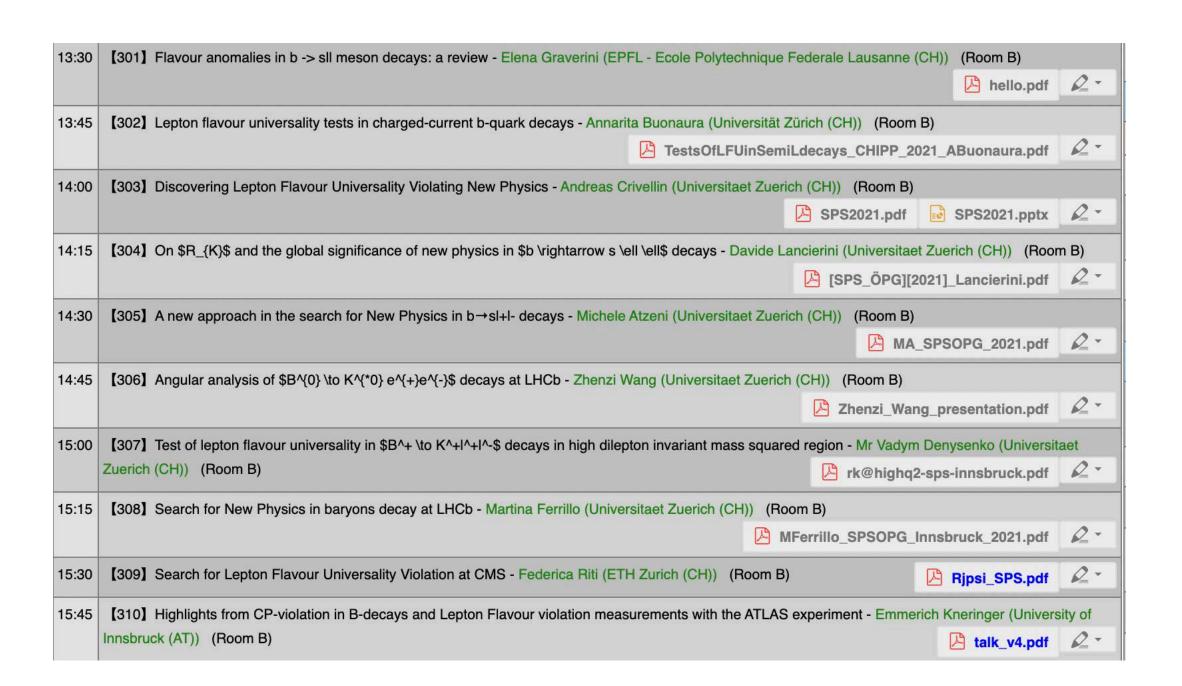
# Review of the flavour anomalies at LHCb

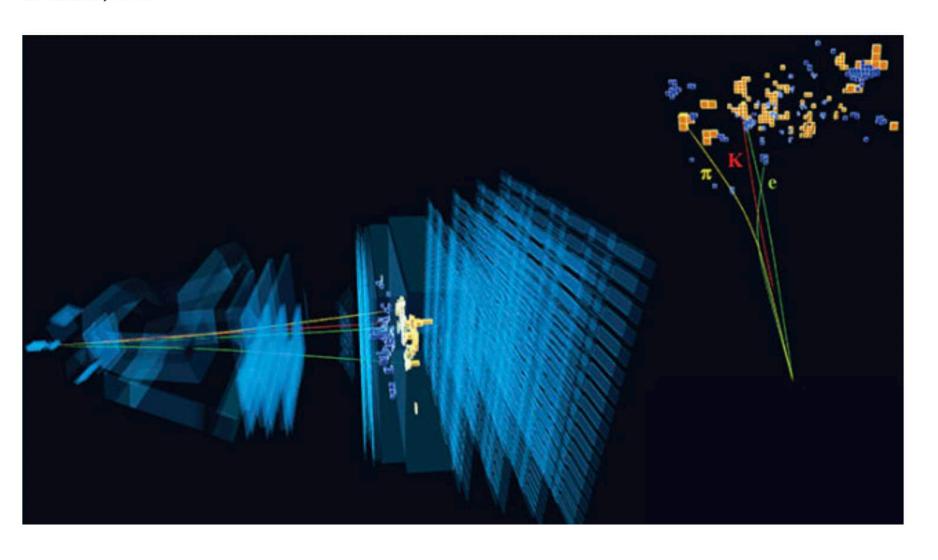
#### SPS/ÖPG joint meeting, Innsbruck



#### Intriguing new result from the LHCb experiment at CERN

The LHCb results strengthen hints of a violation of lepton flavour universality

23 MARCH, 2021



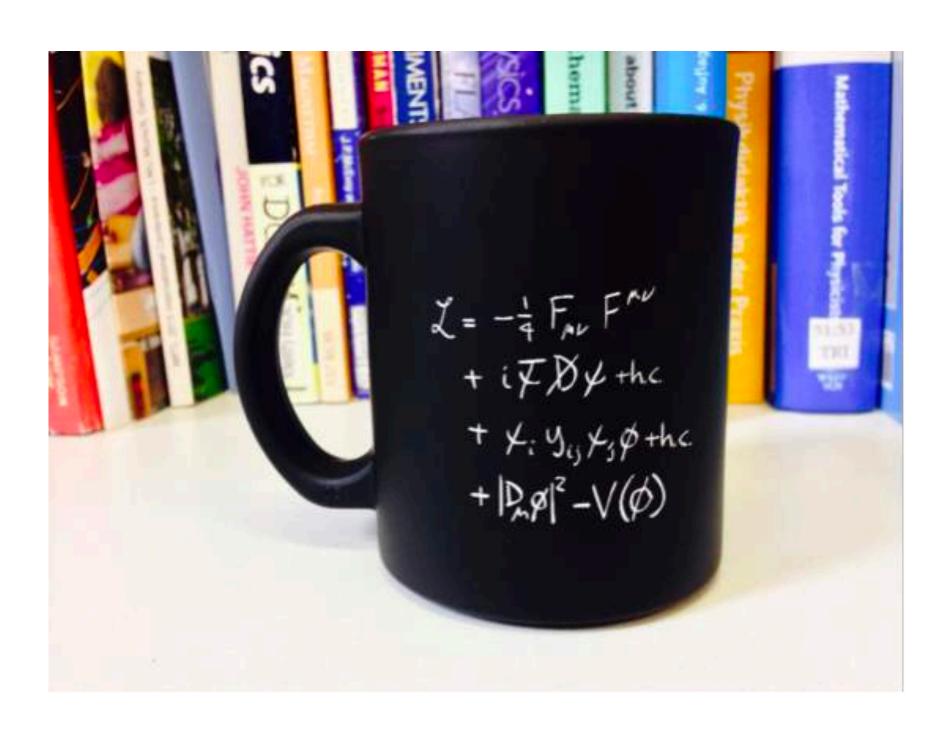
#### **Patrick Owen**



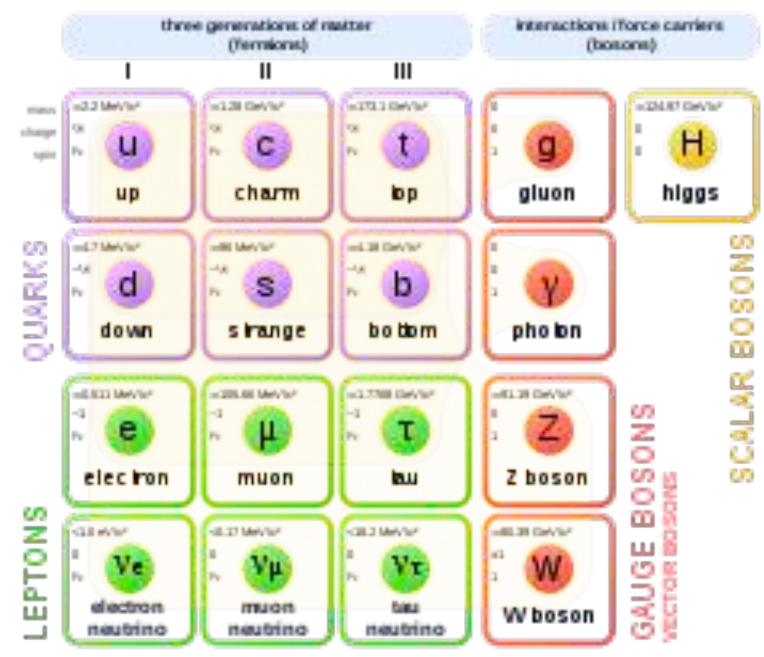


#### The Standard Model

 Vast majority of our knowledge of particle physics encoded in theory known as the Standard Model (SM).







 It describes sub-atomic interactions of fermions via the exchange of force carries (bosons).

### The dual state of the Standard Model

• The Standard Model is simultaneously a roaring success and a catastrophic failure.

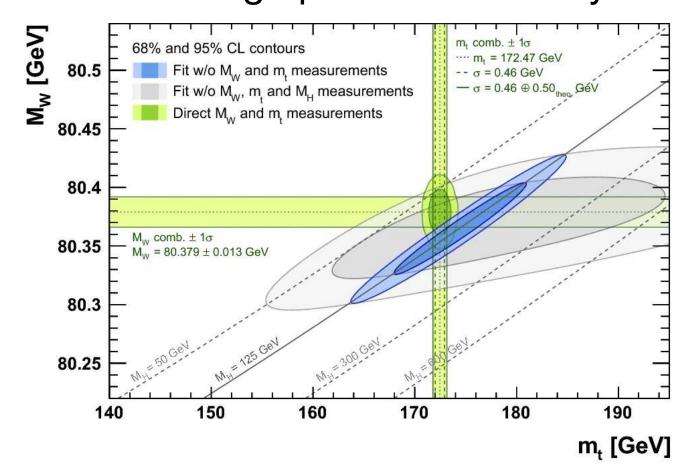
QED verified to 11 decimal places

$$a_e = 0.001\ 159\ 652\ 181\ 643(764)$$
 (th)

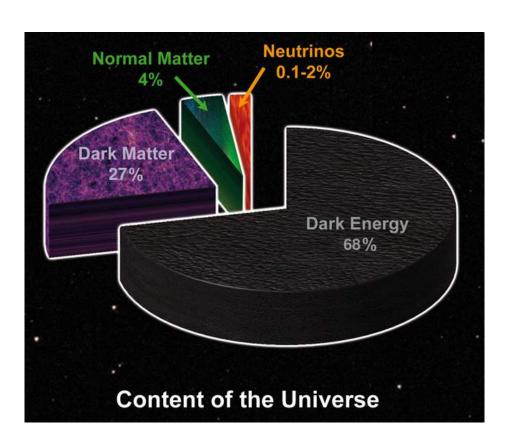
 $a_e = 0.001\ 159\ 652\ 180\ 73(28)$  (expt)

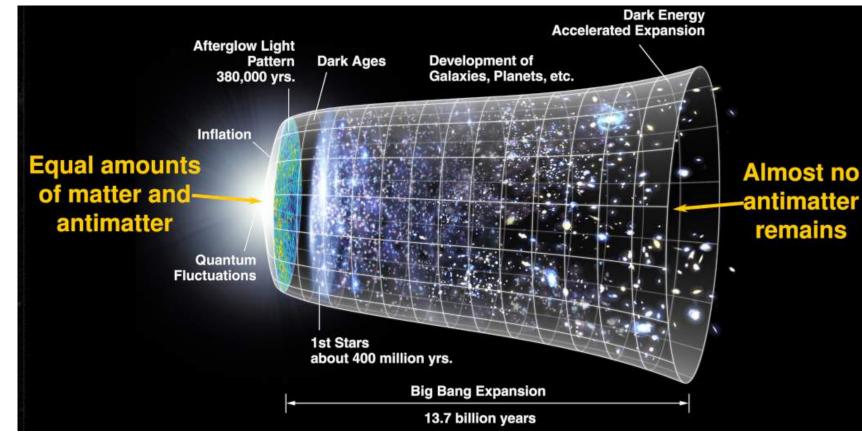
Electroweak theory still consistent under high precision scrutiny

EPJC 78, 675 (2018)



Only small fraction of universe accounted by SM particles.





The search for so-called 'new physics' is therefore highly compelling.

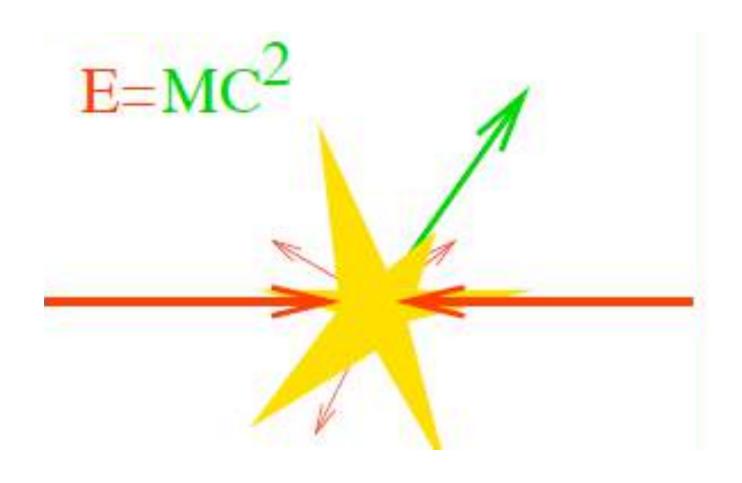
## Ways to search for new physics

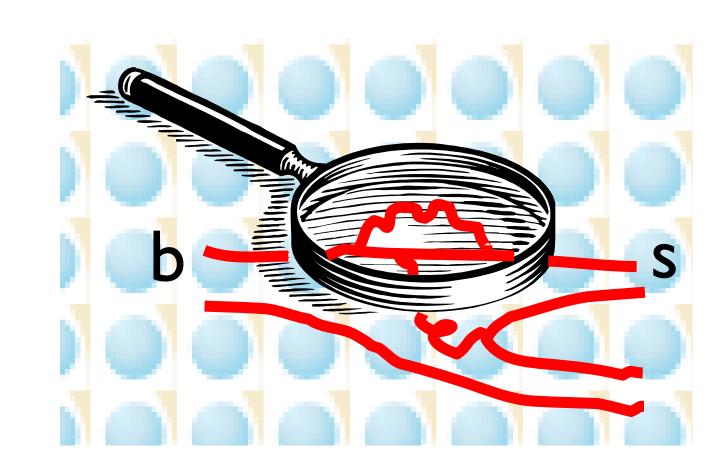
Direct detection

Indirect searches

Search for particles directly produced in high energy beams.

Measure the behaviour of SM particles and compare it to theoretical predictions.

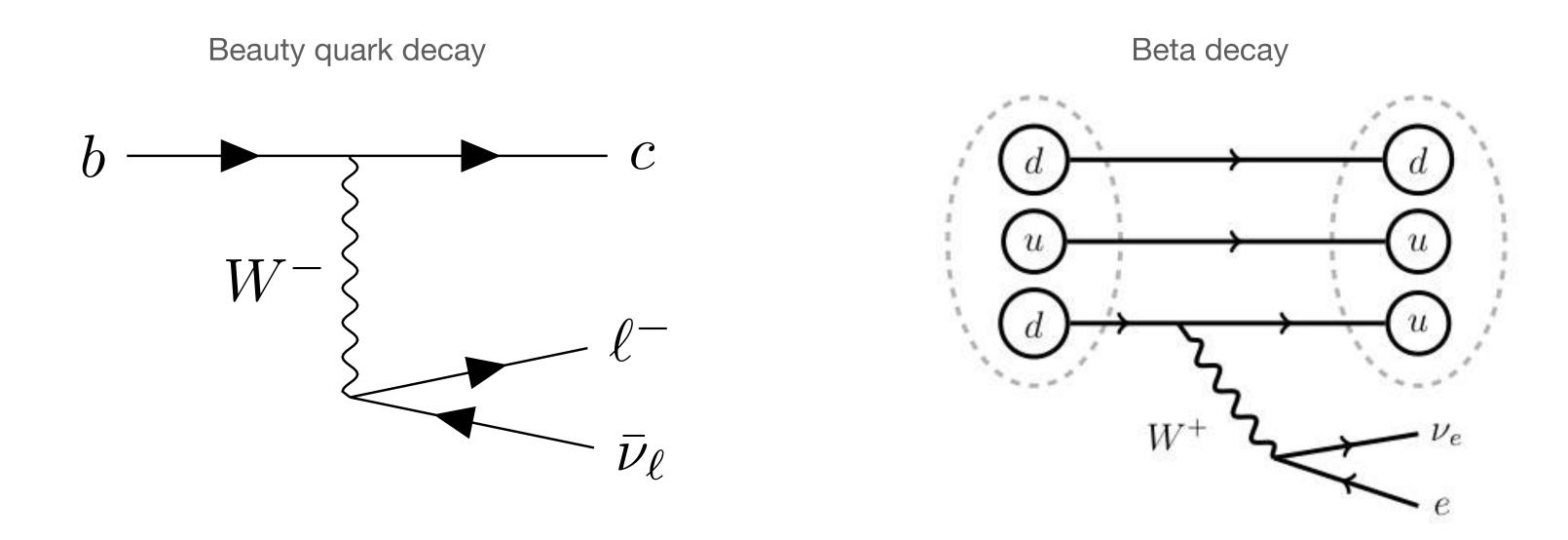




The absence for evidence of new physics in direct searches does not preclude the possibility to find it indirectly.

## New physics with beauty quarks

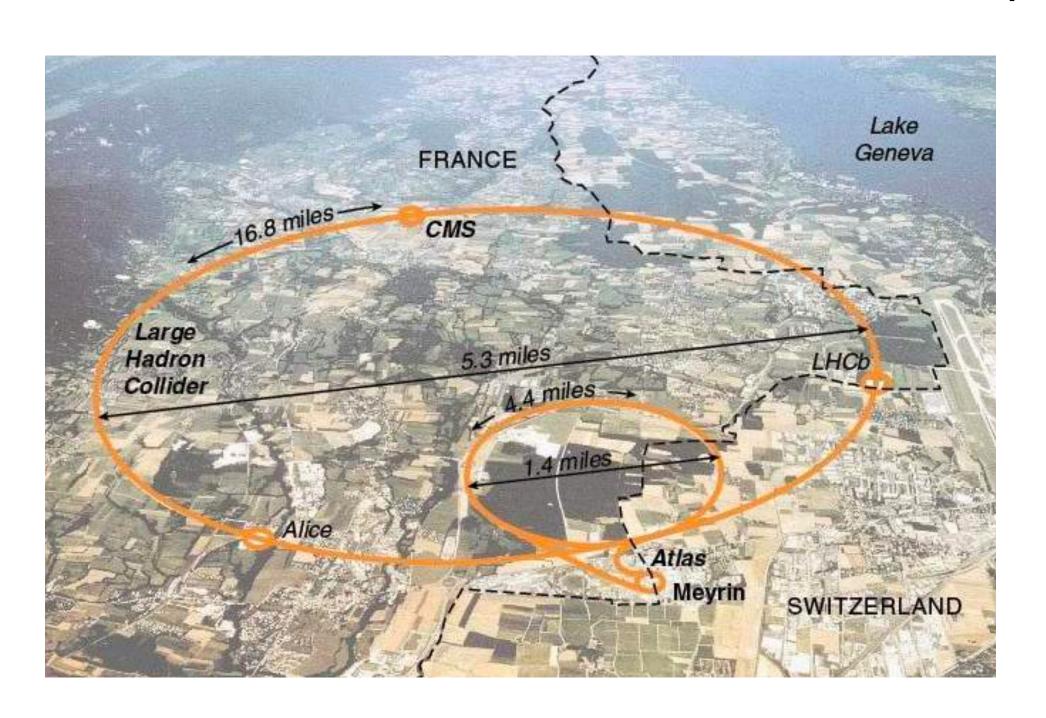
Beauty quarks decay via the weak force.

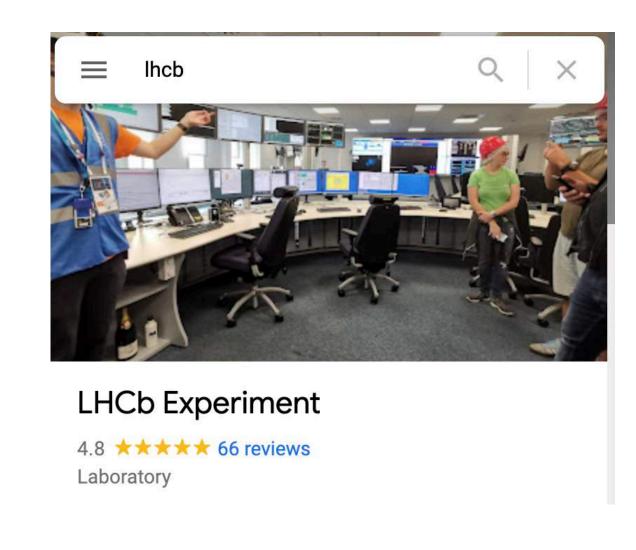


- The Wand Z bosons are over 10 times heavier than the initial decaying b-hadron, but still mediate the decay.
  - Measuring beauty quark decays can tell us about new high mass particles.
  - Such particles can change the behaviour of these decays.

## The LHCb experiment

• LHCb is the LHC's dedicated flavour physics experiment.





#### Your review

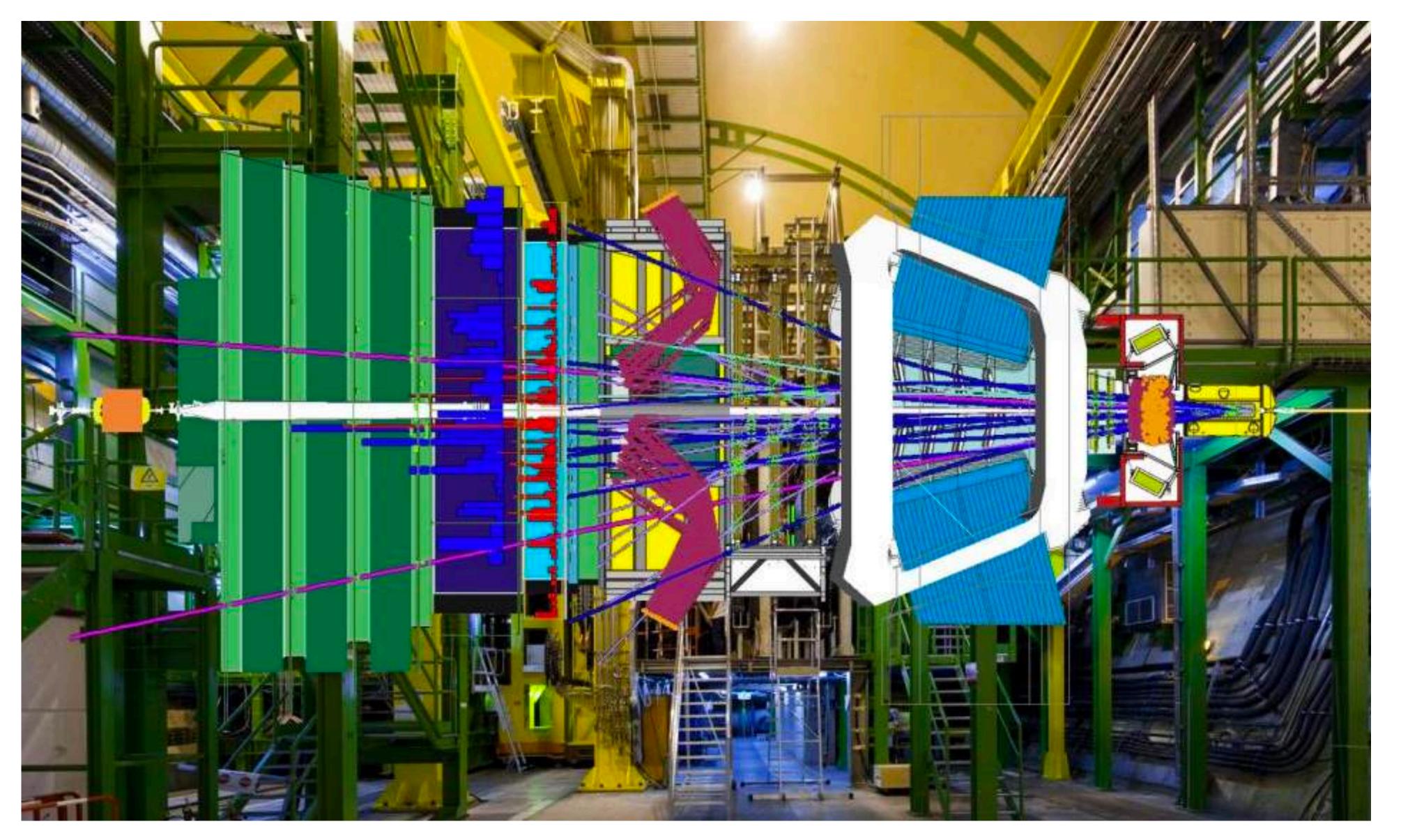


★★★★★ 2 years ago

Fantastic experiment ran by fantastic people. 10/10 would recommend to a friend if they fancied doing a bit of flavour physics

Visited in June 2019

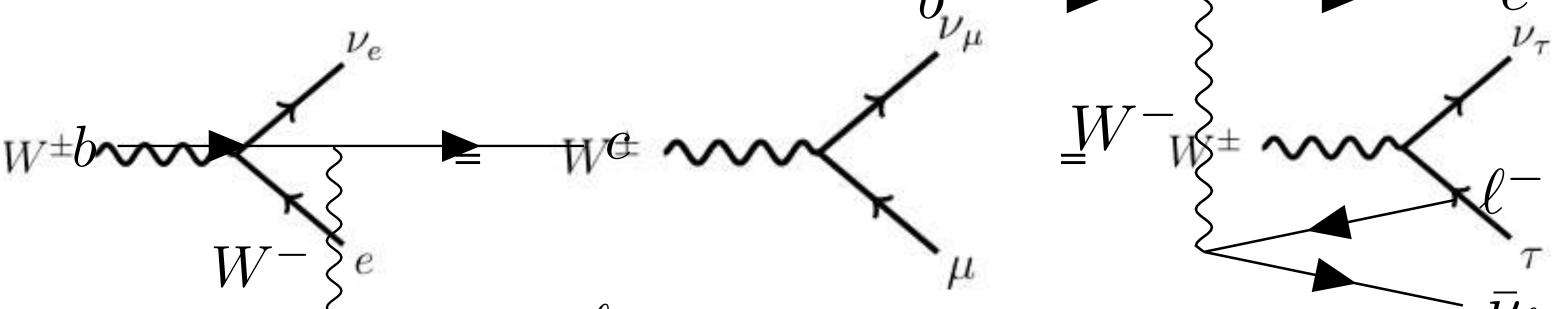
## The LHCb experiment



SPS/OPG meeting

## Lepton universality

Lepton universality is an accidental symmetry in the Standard Model,



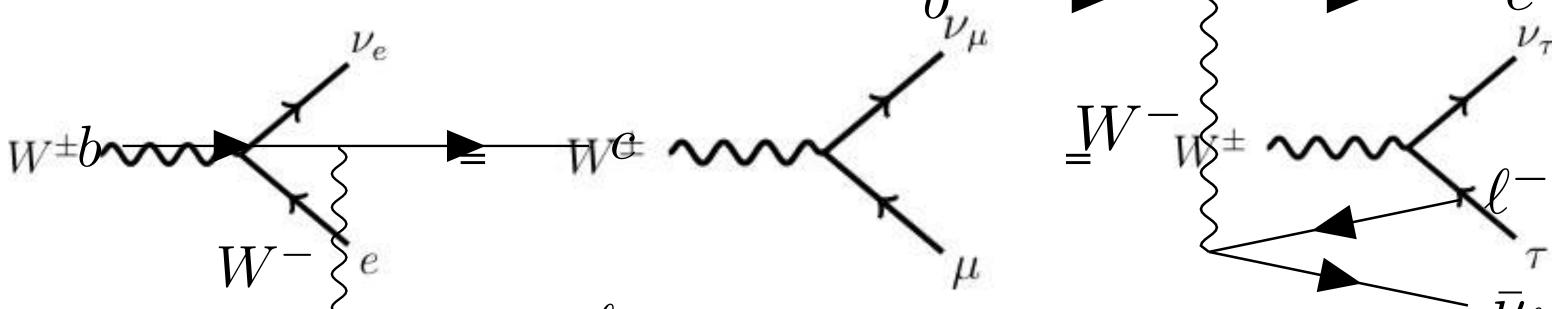
• At LHCb, we want to test it with so-called 'semileptonic decays'.

 $b \longrightarrow \text{Charged current: BF} \sim 10^{-2}$   $t \longrightarrow c$   $\gamma/Z \longrightarrow \ell^+$   $\ell^ R_{K^{(*)}} = \frac{\mathcal{B}(\text{Egure 1:}K^{(*)}\mu^+\mu^-)}{\mathcal{B}(B \to K^{(*)}e^+e^-)} \bar{\nu}_{\ell}$ 

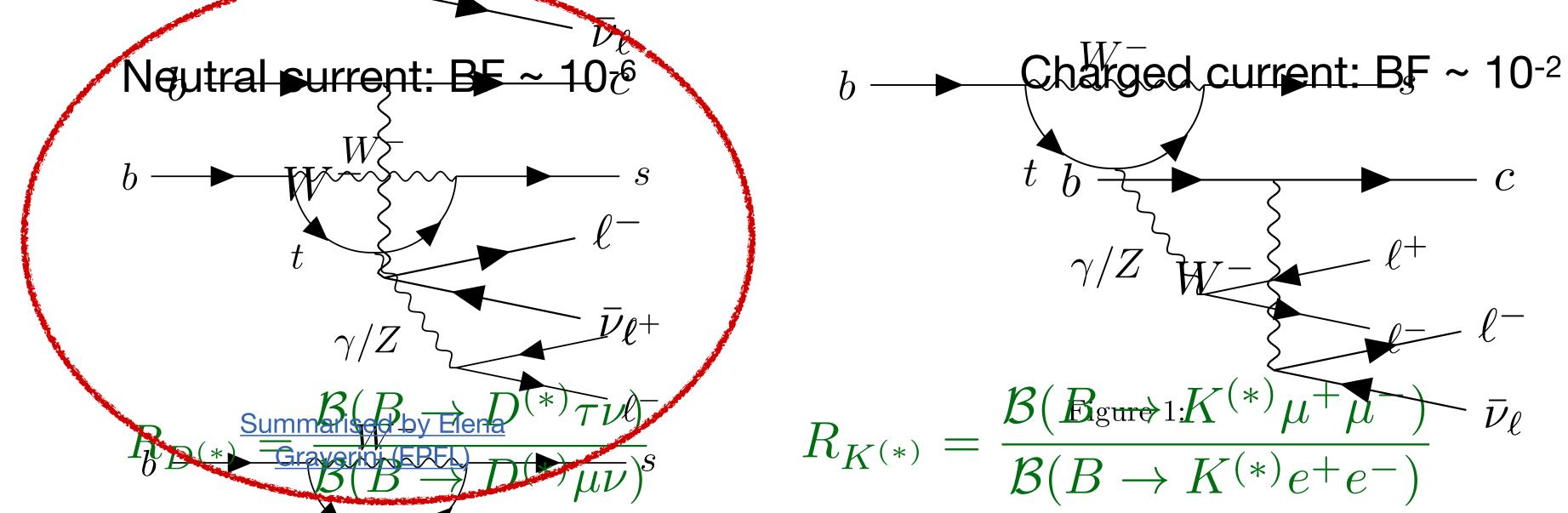
• Compare the decay probabilities (BF) involving different charged Jepton types  $\ell$ .

## Lepton universality

Lepton universality is an accidental symmetry in the Standard Model,



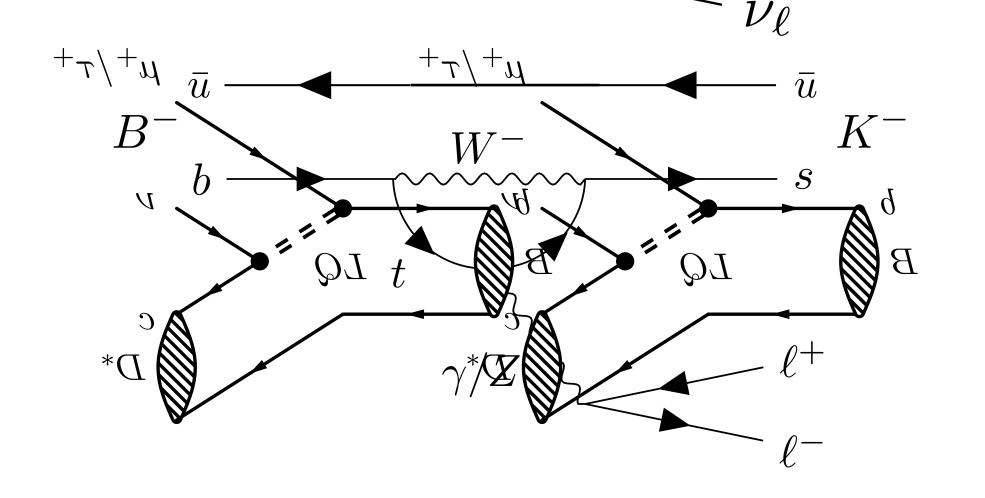
• At LHCb, we want to test it with so-called 'semileptonic decays'.



• Compare the decay probabilities (BF) involving different charge  $\sqrt[4]{lepton types} \ell$ .

<u>Discussed by Davide</u> <u>Lancierini (UZH)</u>

Quarks never appear by themselves, detect bound states as mesons (qq) or baryons (qqq).

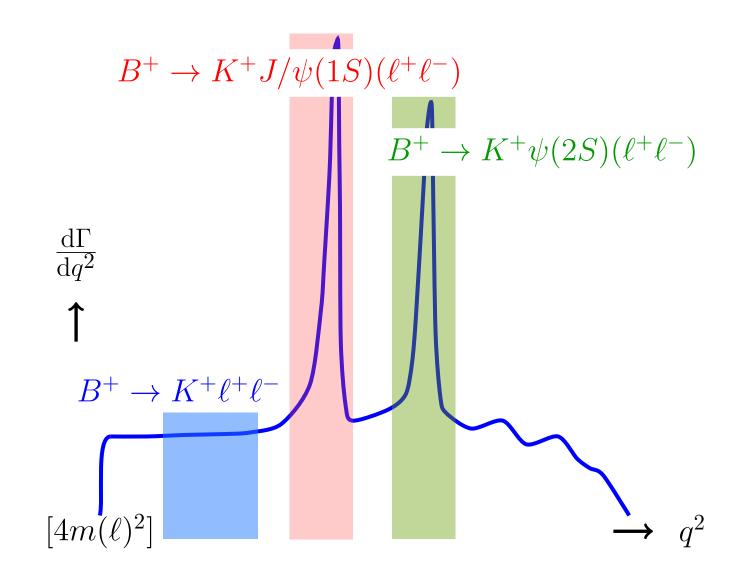


 Take ratio of decay probabilities involving muons and electrons to form R<sub>K</sub>.

$$R_K = \frac{\mathcal{B}(B^- \to K^- \mu^+ \mu^-)}{\mathcal{B}(B^- \to K^- e^+ e^-)}$$

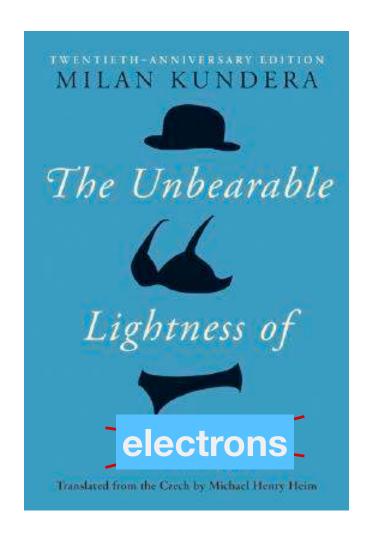
- Important kinematic quantity: q<sup>2</sup>=m<sub>μμ</sub><sup>2</sup>.
- Select signal in the low q<sup>2</sup> region.
- Resonant enhancement from  $c\bar{c}$  mesons.

$oldsymbol{J}/\psi(1S)$ Decay Modes		Screenshot from PDG	Expand all decays		
	Mode	Fraction $(\Gamma_i  /  \Gamma)$	Scale Factor/ Conf. Level	P(MeV/c)	)
$\Gamma_5$	$e^+e^-$	$(5.971 \pm 0.032)\%$		1548	~
$\Gamma_7$	$\mu^+\mu^-$	$(5.961 \pm 0.033)\%$		1545	~

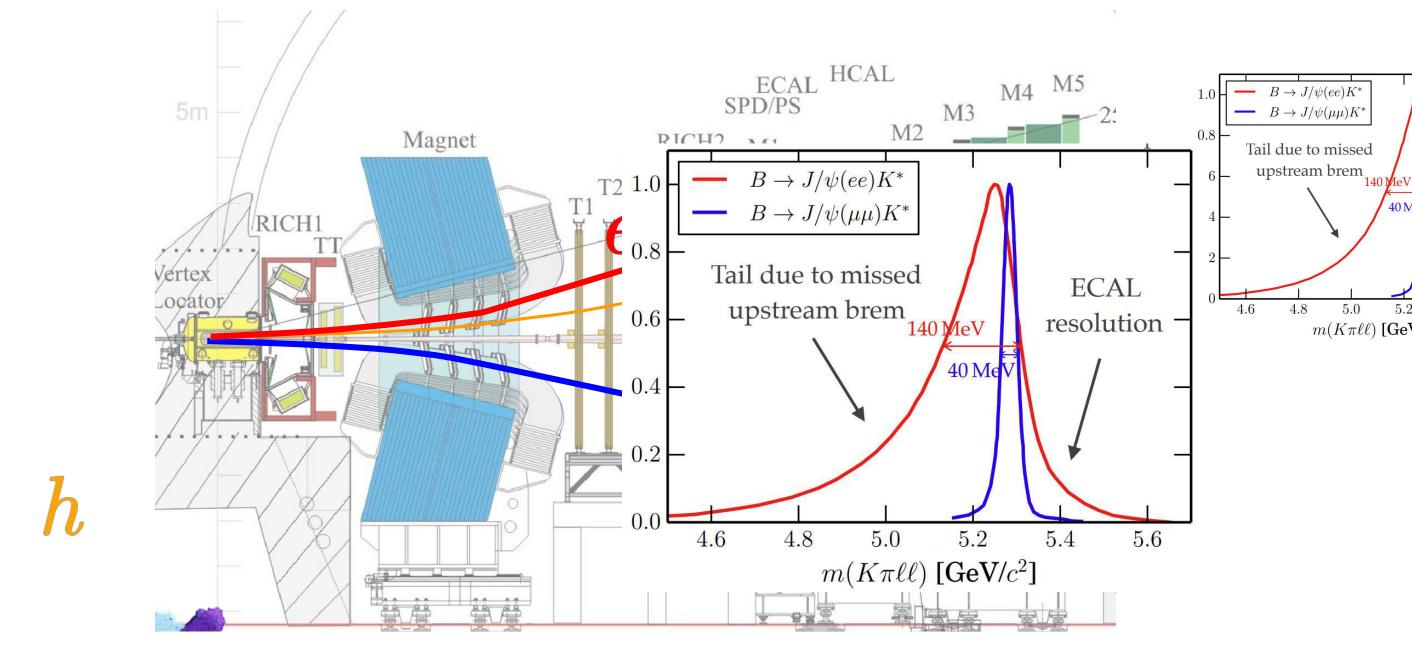


## The unbearable lightness of electrons

Electrons are 200 times lighter than muons
 —> undergo bremsstrahlung more often.

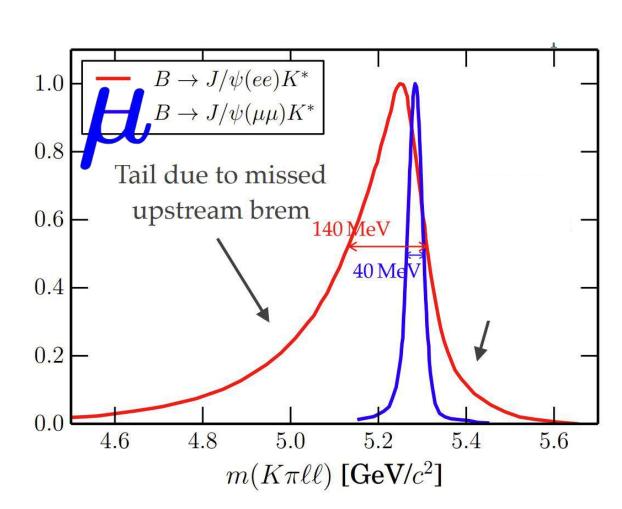


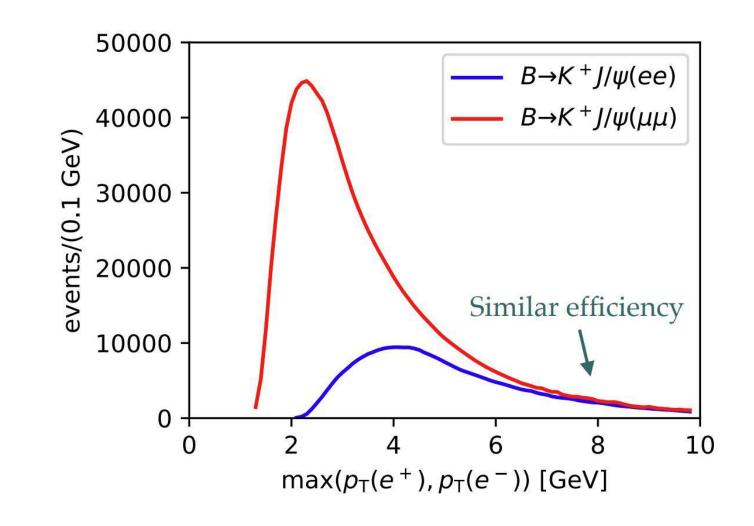
e



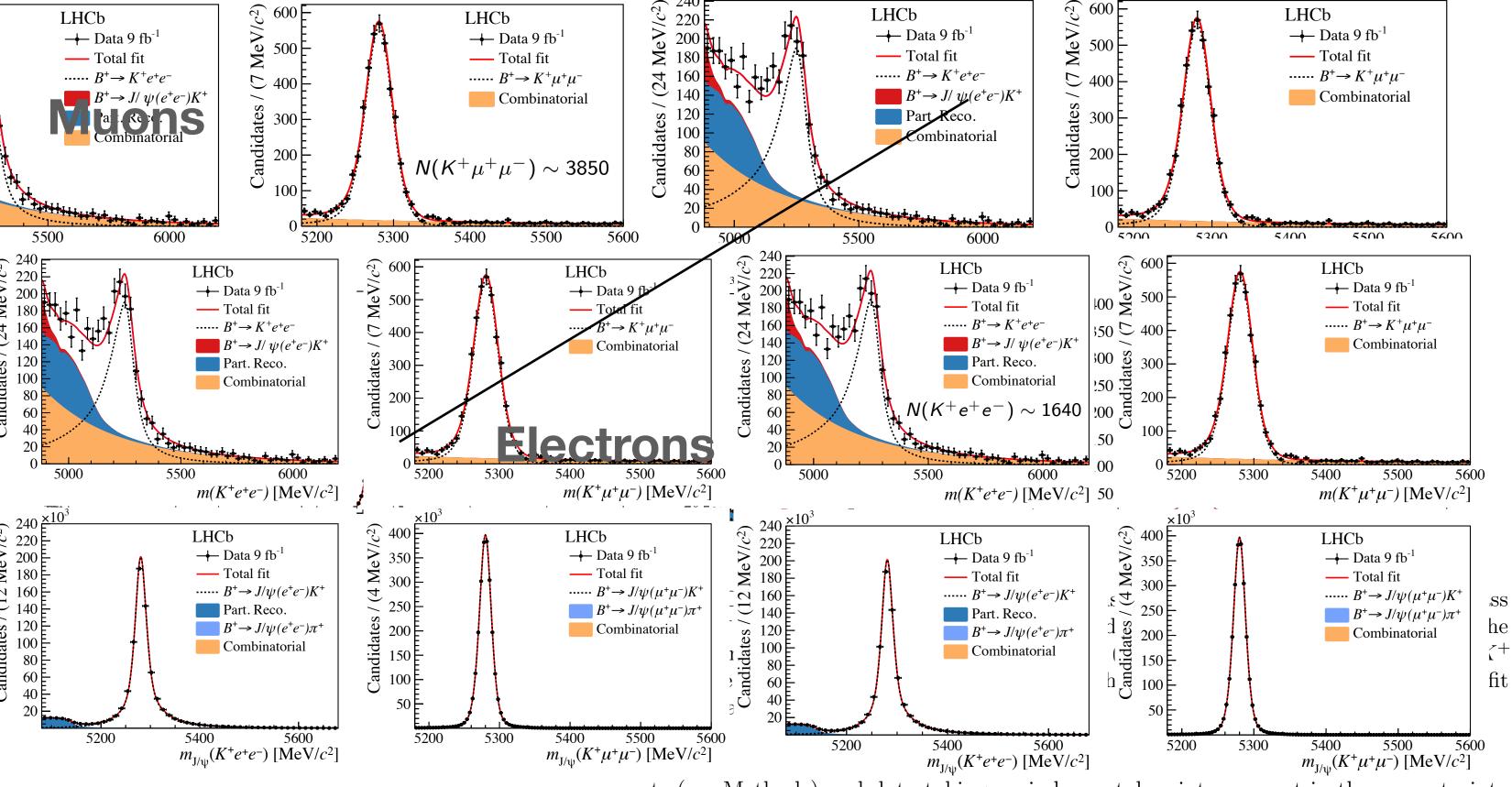
Credit: M. Atzeni

- Two effects from this:
  - Worse mass resolution for electrons.
- Worse efficiency for electrons.





$$R_K = \frac{\mathcal{B}(B^- \to K^- \mu^+ \mu^-)}{\mathcal{B}(B^- \to K^- e^+ e^-)} = \frac{N(B^- \to K^- \mu^+ \mu^-)}{N(B^- \to K^- e^+ e^-)} \frac{\epsilon_{e^+ e^-}}{\epsilon_{\mu^+ \mu^-}}$$



Candidate invariant mass distributions. Distribution of the entrant (see Methods) and data-taking periods are taken into account in these constraints.  $(K^+\ell^+\ell^-)$  for candidates with (left) electron and (right) muon pairs in the mass distribution to the uncertainty is then determined by scanning the nonresonant  $B^+ \to K^+\ell^+\ell^-$  signal channels and (bottom) resonant  $B^+ \to K^+\ell^-\ell^-$  fit projection is superimposed. In the resonant-mode distributions with the efficiencies fixed to their fitted values.

The determination of the  $r_{J/\psi}$  ratio requires control of the relative selection efficiencies

Patrick Owen for the resonant electron and muon modes, and does not therefore benefit from the sole (see Methods) and data-taking periods are taken into account in the scale of the relative selection efficiencies.

Patrick Owen for the resonant electron and muon modes, and does not therefore benefit from the scale of the relative selection efficiencies.

$$R_K = \frac{\mathcal{B}(B^- \to K^- \mu^+ \mu^-)}{\mathcal{B}(B^- \to K^- e^+ e^-)} = \frac{N(B^- \to K^- \mu^+ \mu^-)}{N(B^- \to K^- e^+ e^-)} = \frac{N(B^- \to K^- \mu^+ \mu^-)}{N(B^- \to K^- e^+ e^-)} = \frac{N(B^- \to K^- \mu^+ \mu^-)}{N(B^- \to K^- e^+ e^-)} = \frac{N(B^- \to K^- \mu^+ \mu^-)}{N(B^- \to K^- e^+ e^-)} = \frac{N(B^- \to K^- \mu^+ \mu^-)}{N(B^- \to K^- e^+ e^-)} = \frac{N(B^- \to K^- \mu^+ \mu^-)}{N(B^- \to K^- e^+ e^-)} = \frac{N(B^- \to K^- \mu^+ \mu^-)}{N(B^- \to K^- e^+ e^-)} = \frac{N(B^- \to K^- \mu^+ \mu^-)}{N(B^- \to K^- e^+ e^-)} = \frac{N(B^- \to K^- \mu^+ \mu^-)}{N(B^- \to K^- e^+ e^-)} = \frac{N(B^- \to K^- \mu^+ \mu^-)}{N(B^- \to K^- e^+ e^-)} = \frac{N(B^- \to K^- \mu^+ \mu^-)}{N(B^- \to K^- e^+ e^-)} = \frac{N(B^- \to K^- \mu^+ \mu^-)}{N(B^- \to K^- e^+ e^-)} = \frac{N(B^- \to K^- \mu^+ \mu^-)}{N(B^- \to K^- e^+ e^-)} = \frac{N(B^- \to K^- \mu^+ \mu^-)}{N(B^- \to K^- e^+ e^-)} = \frac{N(B^- \to K^- \mu^+ \mu^-)}{N(B^- \to K^- e^+ e^-)} = \frac{N(B^- \to K^- \mu^+ \mu^-)}{N(B^- \to K^- e^+ e^-)} = \frac{N(B^- \to K^- \mu^+ \mu^-)}{N(B^- \to K^- e^+ e^-)} = \frac{N(B^- \to K^- \mu^+ \mu^-)}{N(B^- \to K^- e^+ e^-)} = \frac{N(B^- \to K^- \mu^+ \mu^-)}{N(B^- \to K^- e^+ e^-)} = \frac{N(B^- \to K^- \mu^+ \mu^-)}{N(B^- \to K^- e^+ e^-)} = \frac{N(B^- \to K^- \mu^+ \mu^-)}{N(B^- \to K^- e^+ e^-)} = \frac{N(B^- \to K^- \mu^+ \mu^-)}{N(B^- \to K^- e^+ e^-)} = \frac{N(B^- \to K^- \mu^+ \mu^-)}{N(B^- \to K^- e^+ e^-)} = \frac{N(B^- \to K^- \mu^+ \mu^-)}{N(B^- \to K^- e^+ e^-)} = \frac{N(B^- \to K^- \mu^+ \mu^-)}{N(B^- \to K^- e^+ e^-)} = \frac{N(B^- \to K^- \mu^+ \mu^-)}{N(B^- \to K^- e^+ e^-)} = \frac{N(B^- \to K^- \mu^+ \mu^-)}{N(B^- \to K^- e^+ e^-)} = \frac{N(B^- \to K^- \mu^+ \mu^-)}{N(B^- \to K^- e^+ e^-)} = \frac{N(B^- \to K^- \mu^+ \mu^-)}{N(B^- \to K^- e^+ e^-)} = \frac{N(B^- \to K^- \mu^+ \mu^-)}{N(B^- \to K^- e^+ e^-)} = \frac{N(B^- \to K^- \mu^+ \mu^-)}{N(B^- \to K^- e^+ e^-)} = \frac{N(B^- \to K^- \mu^+ \mu^-)}{N(B^- \to K^- e^- \mu^-)} = \frac{N(B^- \to K^- \mu^+ \mu^-)}{N(B^- \to K^- e^- \mu^-)} = \frac{N(B^- \to K^- \mu^+ \mu^-)}{N(B^- \to K^- e^- \mu^-)} = \frac{N(B^- \to K^- \mu^-)}{N(B^- \to K^- e^- \mu^-)} = \frac{N(B^- \to K^- \mu^-)}{N(B^- \to K^- e^- \mu^-)} = \frac{N(B^- \to K^- \mu^-)}{N(B^- \to K^- e^- \mu^-)} = \frac{N(B^- \to K^- \mu^-)}{N(B^- \to K^- \mu^-)} = \frac{N(B^- \to K^- \mu^-)}{$$

 $m_{//\psi}(K^+\mu^+\mu^-) [\text{MeV}/c^2]$ 

e 2: Candidate invariant mass distributions. Distribution of the invariant mass distributions. Distribution of the invariant mass distributions. Distribution of the invariant mass distributions.  $K^{+}\ell^{-}$ ) is the determined by scanning the scanning rs. The fit projection is superimposed. In the resonant-mode distributions switch fithe efficiencies fixed to their fitted values. onents are too small to be visible. The determination of the  $r_{J/\psi}$  ratio requires control of the relative selection efficiencies

120

for the resonant electron and muon modes, and does not therefore benefit from the Patrick Owen 169 s (see Methods) and data-taking periods are taken into account on the constraints systematic effects in the double ratio used to measure  $R_K$ . Given the scale

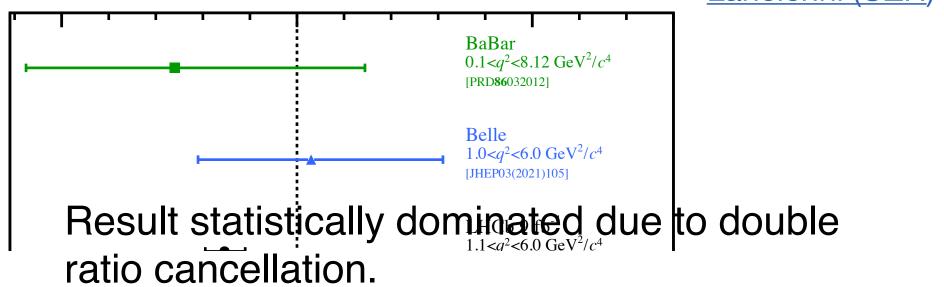
 $m_{\rm J/\psi}(K^+e^+e^-) \, [{\rm MeV}/c^2]$ 

#### R<sub>K</sub> with full Run 1 and Run 2 LHCb data

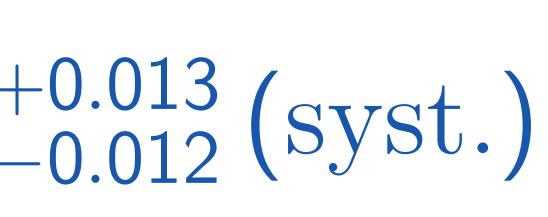
Discussed by Davide Lancierini (UZH)

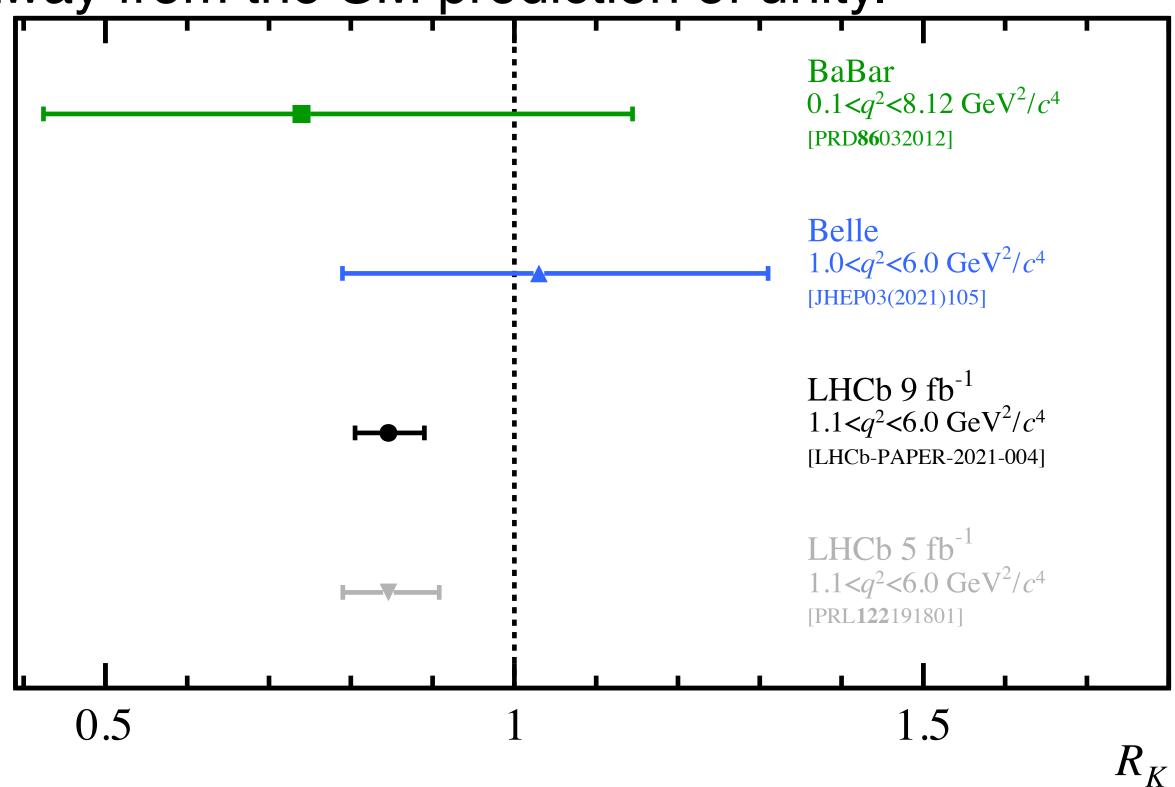
The measured value of  $R_K$  is:

$$P_{11} = 0.046 + 0.042 (stat) + 0.013 (stat)$$



2 LHCb data
3.1 standard deviations away from the SM prediction of unity.



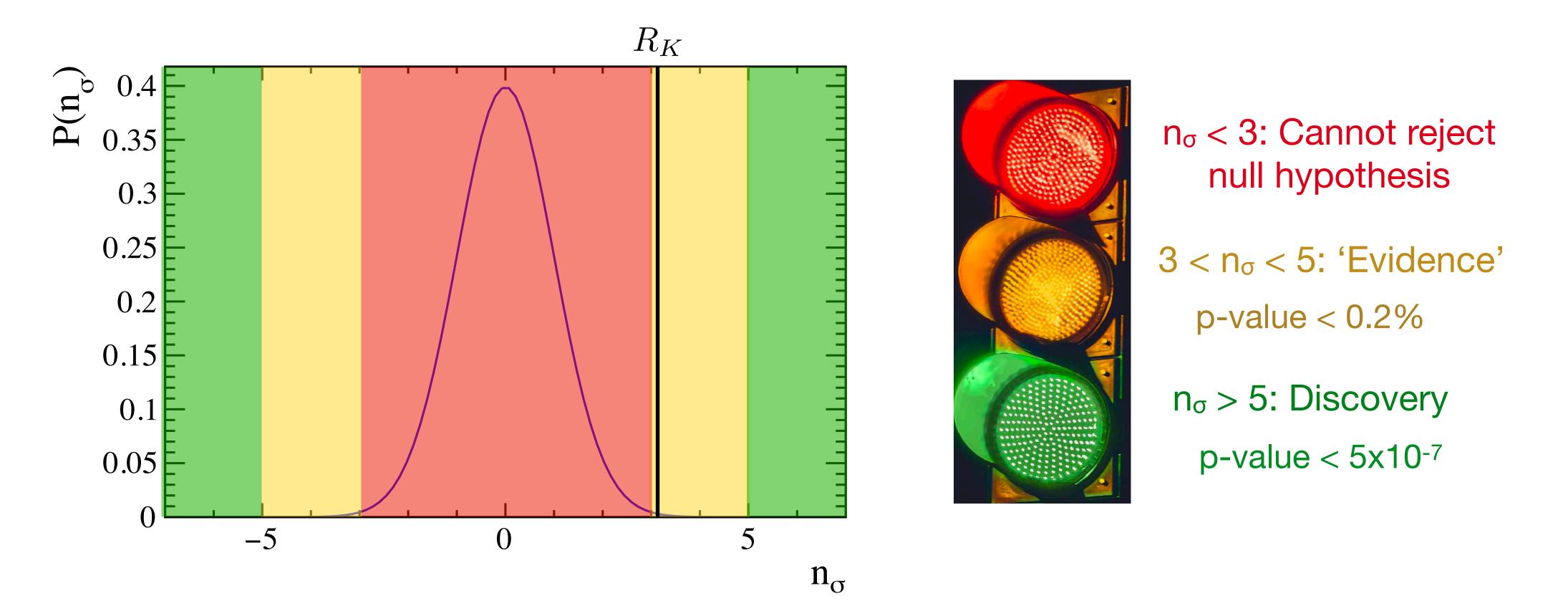


nodel

rigger & kinematics

## Discovery convention in particle physics

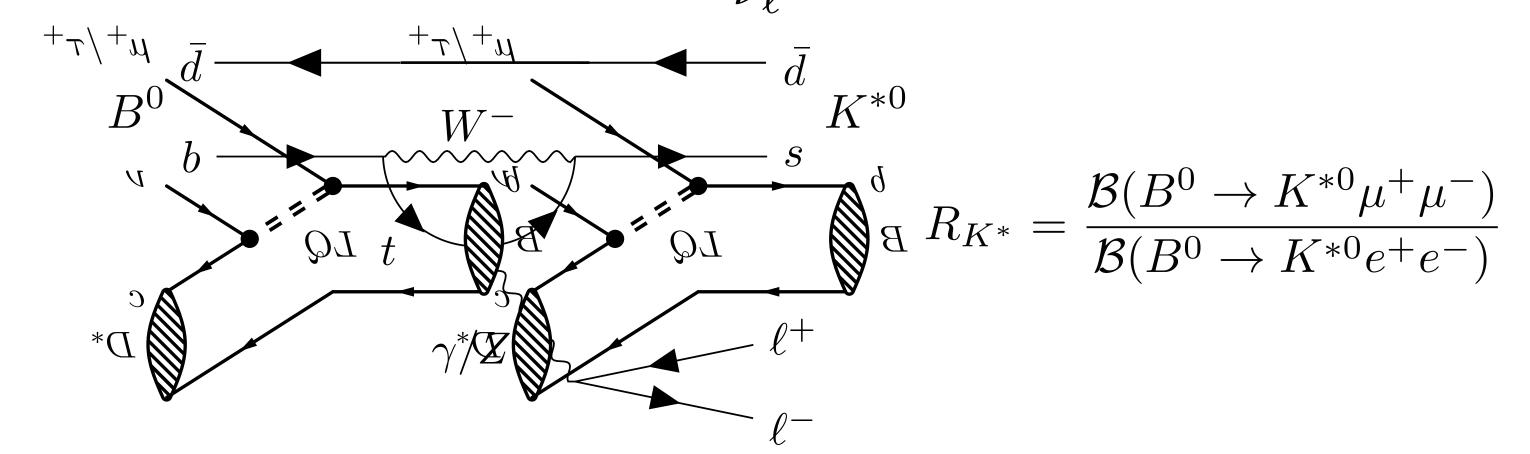
Comparing the likelihood of the null hypothesis to an alternative one.

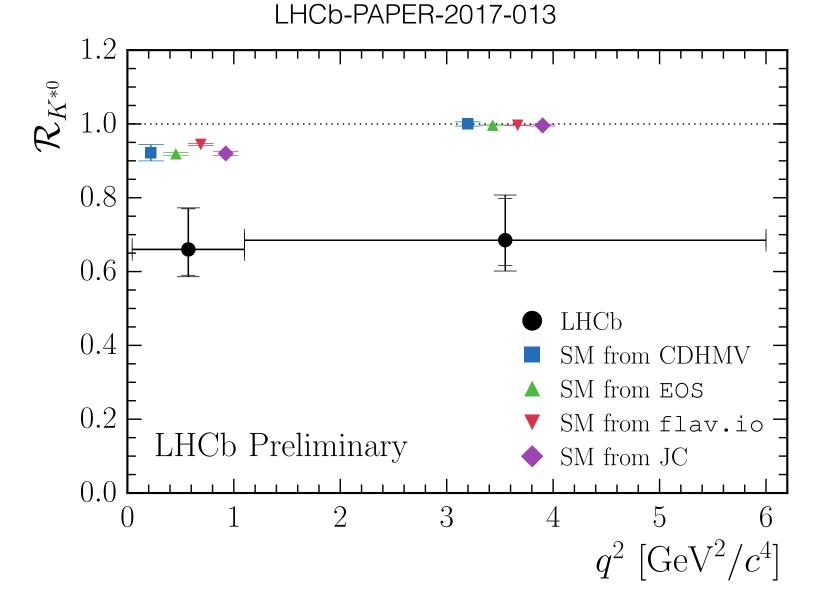


• A significance of 3.1σ is therefore conventionally 'evidence' for lepton universality violation.

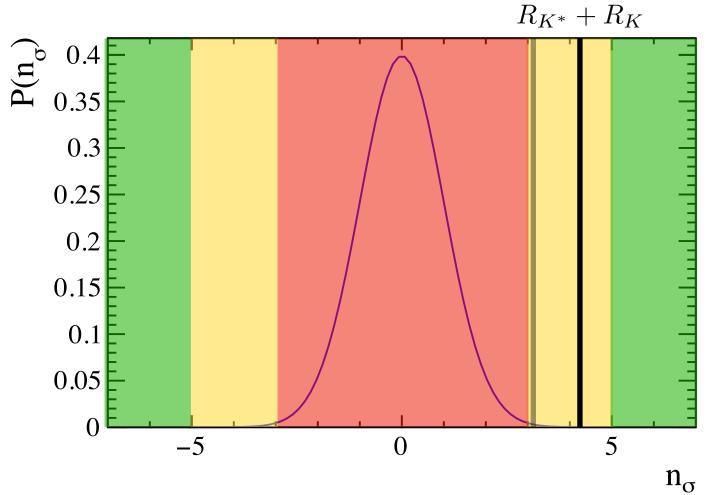
## The lepton universality ratio Rk\*

• We have also tested lepton universality in a closely related mode,  $B^0 \to K^{*0} \ell^+ \ell^-$ 





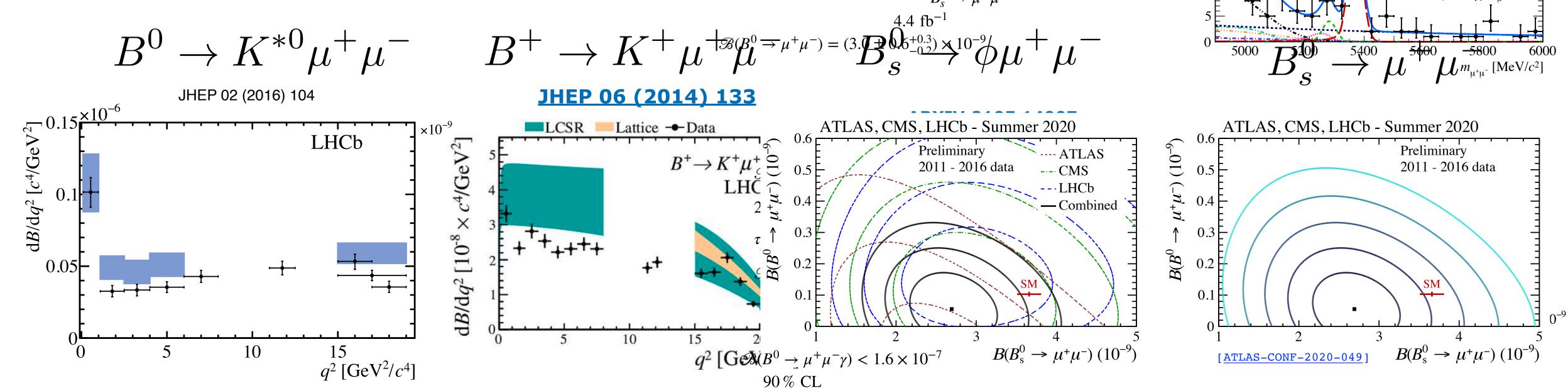
- Also see deficit of muons compared to electrons in two q² regions. Significances 2.2-2.4σ.
- Combined significance ~4.2σ.
- Is there anything else we can measure?



 $\mathcal{R}_{K^{*0}}$ 

## The branching fraction

Measure the individual branching fractions and compare with SM pred



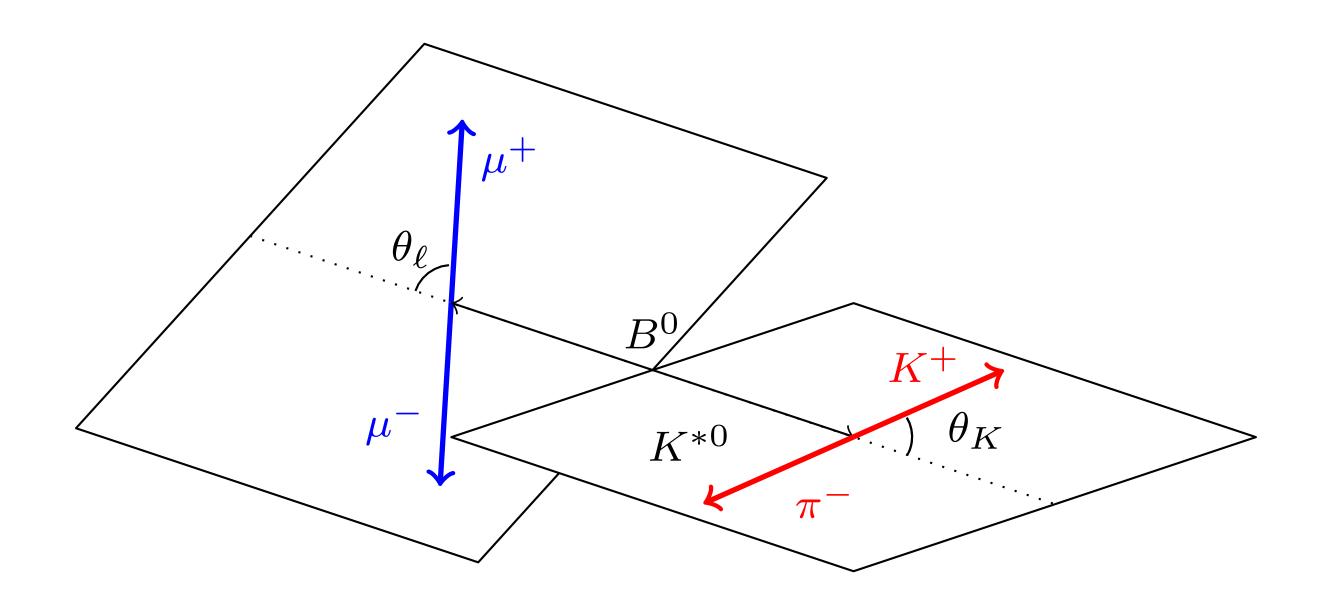
- Theoretical uncertainty now large but a coherent picture emerges: Deficit of muons.
- Could new physics also change the angular distribution as well as the decay rate?

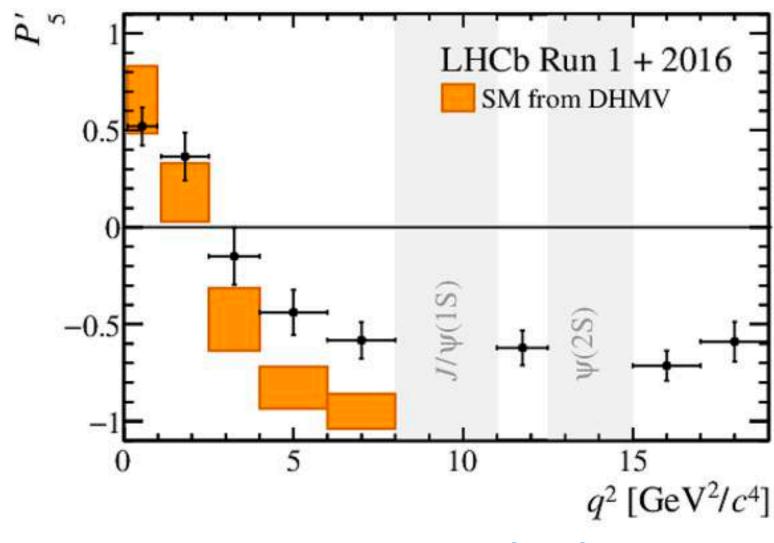
LHCb

BDT > 0.5

## The $B \to K^* \mu^+ \mu^-$ angular analysis

• In addition to decay rates, compare angular distribution and compare with SM.





PHYS. REV. LETT. 125 (2020) 011802

- Discrepancy also seen in the angular distribution.
- Are these anomalies consistent with each other and what is the combined significance?

## Effective field theory

We

ries (EFTs).

• Similarly to the  $\beta$ -decay we can integrate out the heavy The idea is a generalisation of Fermi's theory of weak decays: 1:  $\beta$ -de

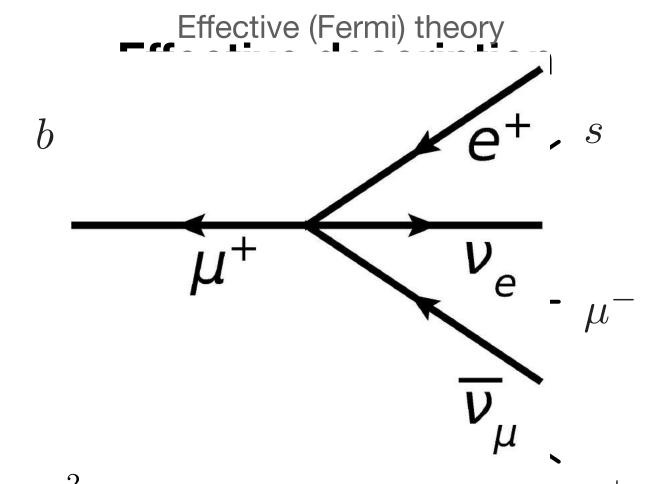
 $Marsh \epsilon$ ark lang been in tion of t

d electro

at the fo

e inclusic

Full theory
Full theory



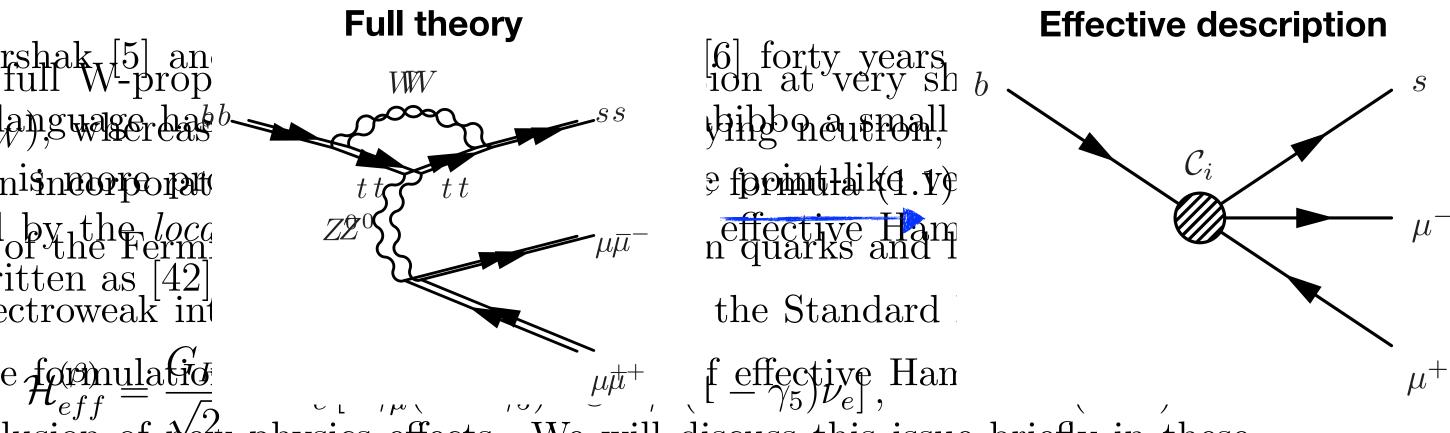
 $\mathcal{M} = \frac{g_L^2}{2} \bar{x}(k_3) \bar{\sigma}_\rho x(p) \frac{1}{g^2 - m^2} \bar{x}(k_1) \bar{\sigma}_\rho y(k_2) \qquad \mathcal{M} \approx -\frac{g_L^2}{2m^2} [\bar{x}(k_3) \bar{\sigma}_\rho x(p)] [\bar{x}(k_1) \bar{\sigma}_\rho y(k_2)] / [1 + \mathcal{O}(q^2/m_W^2)]$ • Model independent description in effective field theory

• This results in an effective hamiltonian, written as a combination of the short encoding info of the short distance physics

O<sub>i</sub> four-fermion operators

## EFTs in Reavy flavour physics

utron The part at the quark level in the full (a) and effective (b) theory.



lusion of New physics effects. We will discuss this issue briefly in these

miliar Familiar operators:
Fermi theor

 $\mathcal{H}_{ef} \; \mathcal{O}_7^{(\prime)} = \qquad \qquad \gamma$  he local oper  $S_{L(R)}$ 

(a)

 $b \to s\ell^+\ell^-$ 

 $\mathcal{O}_{9,10}^{(\prime)} = \begin{cases} b_{L,R} \\ \mathcal{O}_{9,10}^{(\prime)} = \\ s_{L(R)} \end{cases}$ 

 $\mathcal{D}_{S,P}^{(\prime)} = S_{L(R)}$ 

O<sub>7</sub> gives a long distance contribution to  $b^{R,L}$ >sll via the photon.

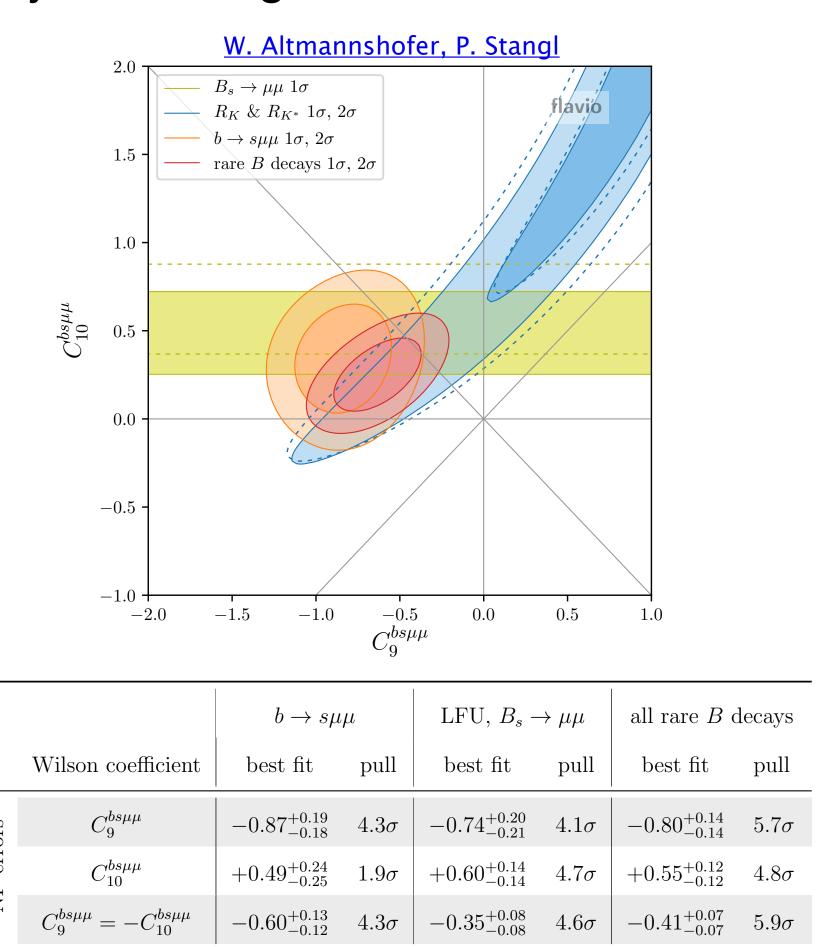
O<sub>9,10</sub> can be different for different lepton flavours.

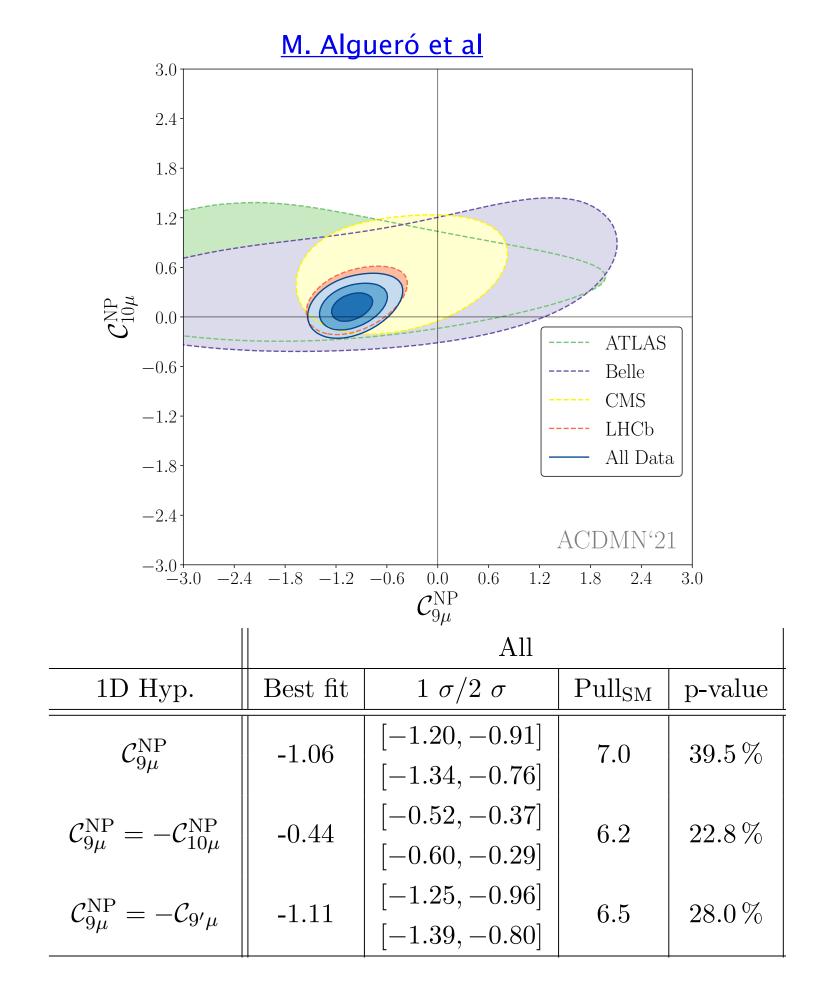
ich, together with the UKIVI matrix elements, describe the strength given operator enters the Hamiltonian.  $\mathcal{H}_{eff}$  is thus represented own as Operator Product Expansion (OPE), of effective vertices

own as Operator Product Expansion (OPE), of effective vertices ffective couplings to the leptons: Dominant SM semileptonic contribution.  $\text{ffective coupling to and } \text{Constants } \text{C}_i$ .

## Global $b \to s \ell^+ \ell^-$ fits

Global b—>sll fits show that all discrepancies are in consistent within the EFT approach, with significances
easily exceeding the conventional 5σ threshold.

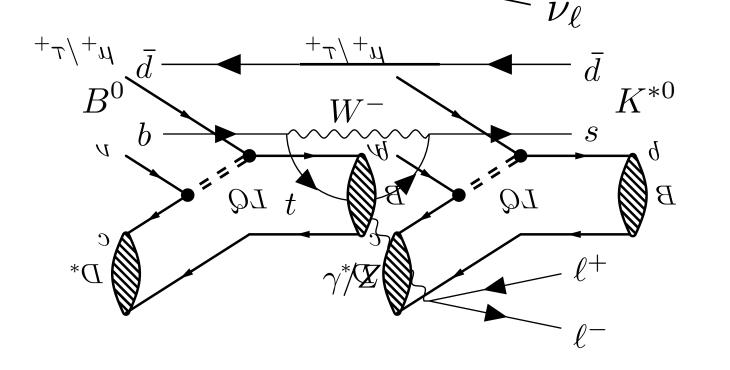


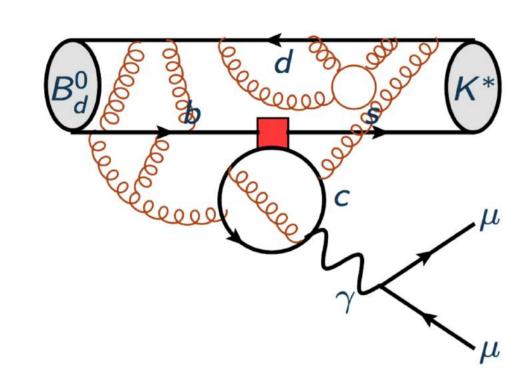


No discovery claim from LHCb at this point. Why not?

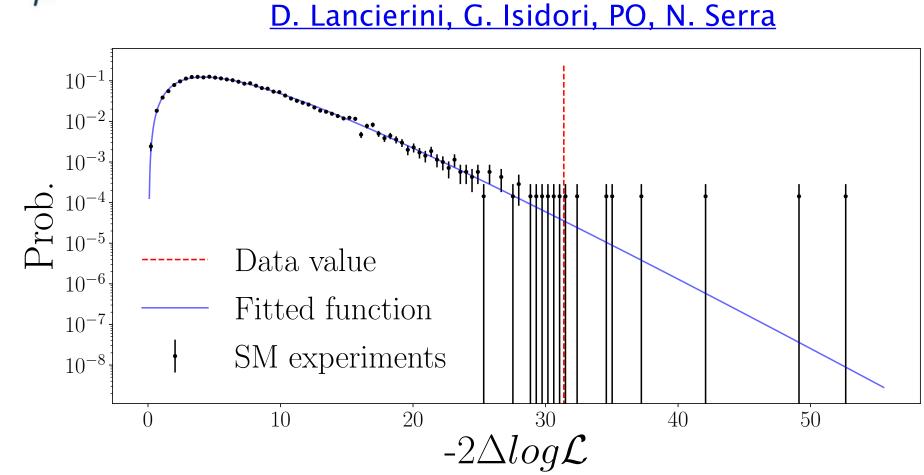
# ASM complication

- Also a SM contribution which can negatively interfere with the semileptonic amplitude.
  - Does not affect LFU atios, but could cause other anomalies we see.



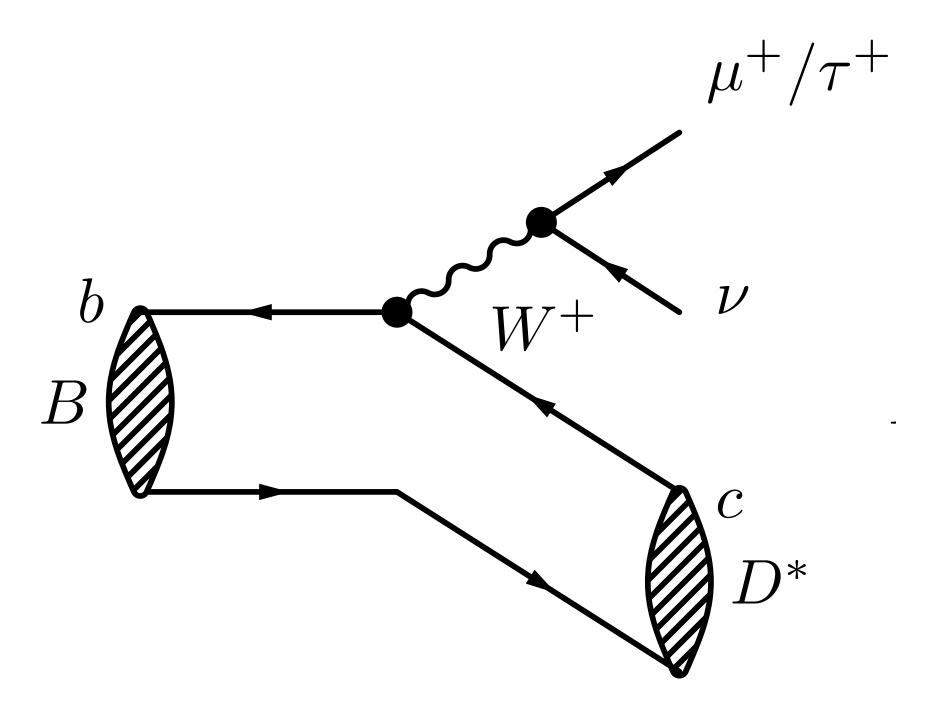


- Try to construct as conservative approach as possible.
  - Hyper conservative hadronic treatment (absorb into SM definition).
  - Test against general NP as possible to take into account the so-called Look-Elsewhere-Effect.



• Global significance is  $\sim$ 4 $\sigma$ , which is not at the discovery threshold but still high.

## Tree-level $b \to c \ell \nu$ transitions

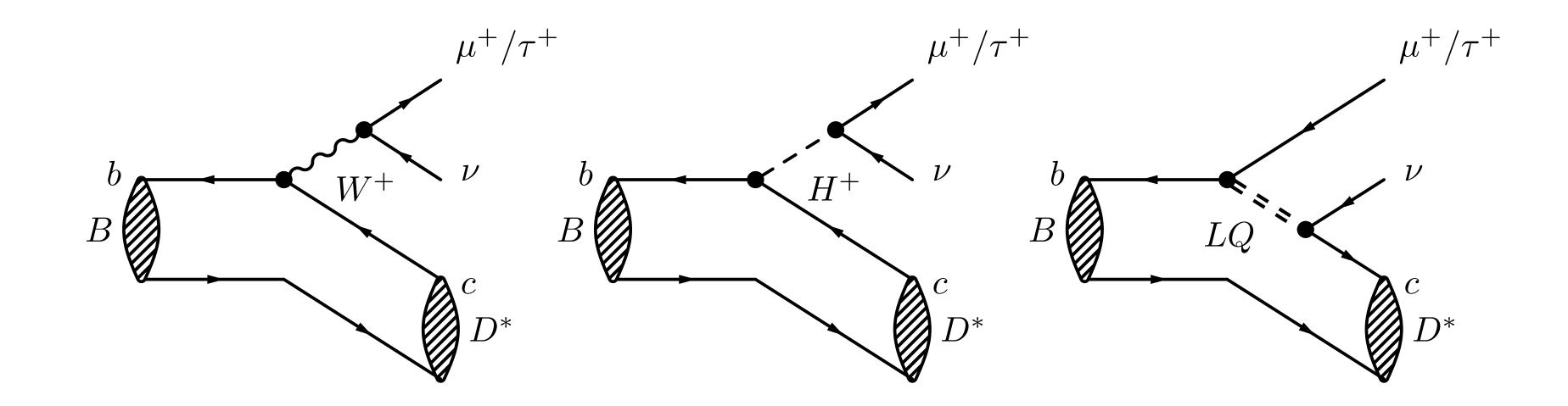


## Lepton universality ratio R(D\*)

Now compare tau-leptons with muons/electrons.

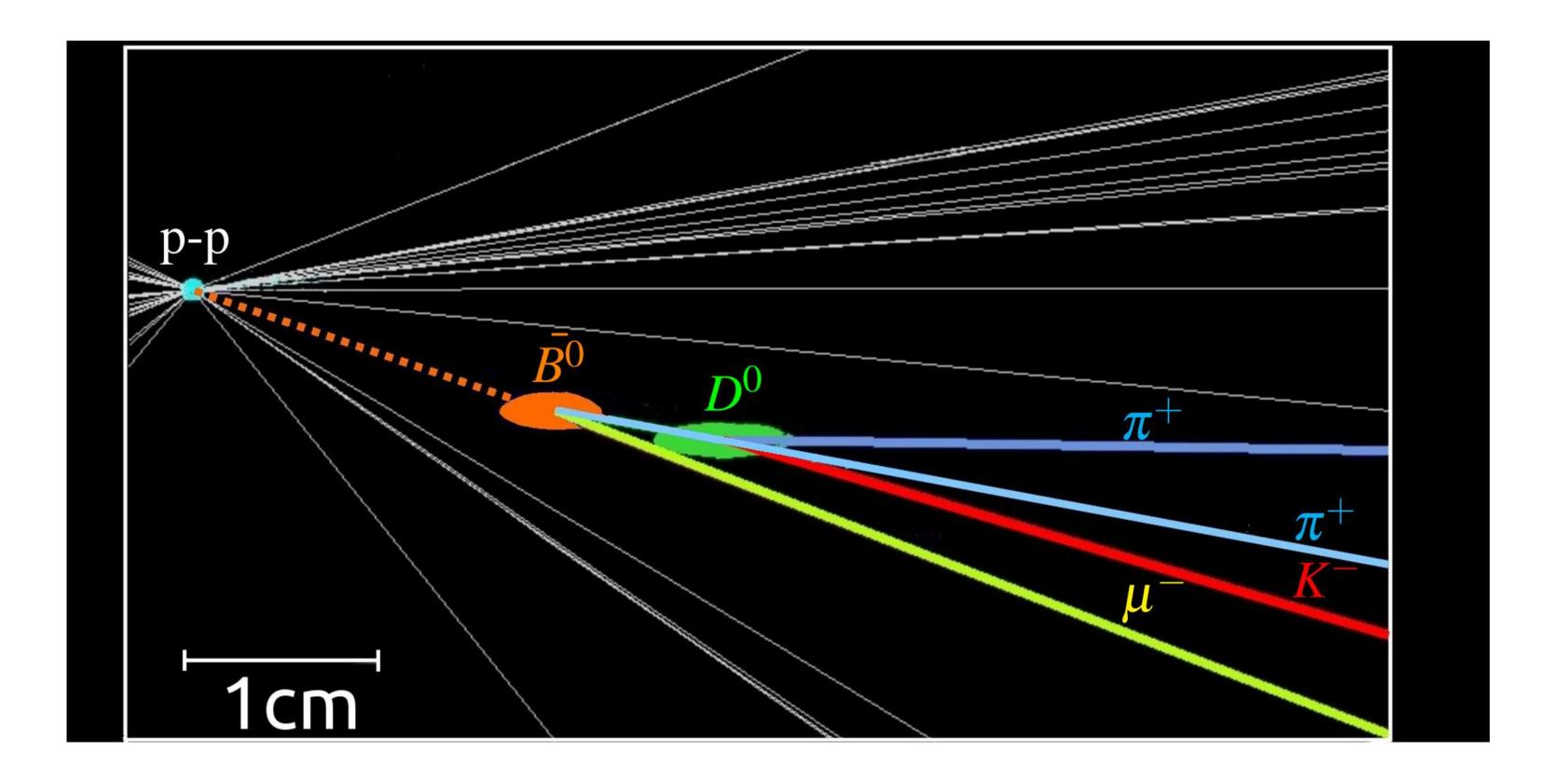
$$R(D^{(*)}) = \frac{\mathcal{B}(B \to D^{(*)}\tau\nu)}{\mathcal{B}(B \to D^{(*)}\ell\nu)}$$

· R(D\*) sensitive to any physics model favouring 3rd generation leptons (e.g. charged Higgs).



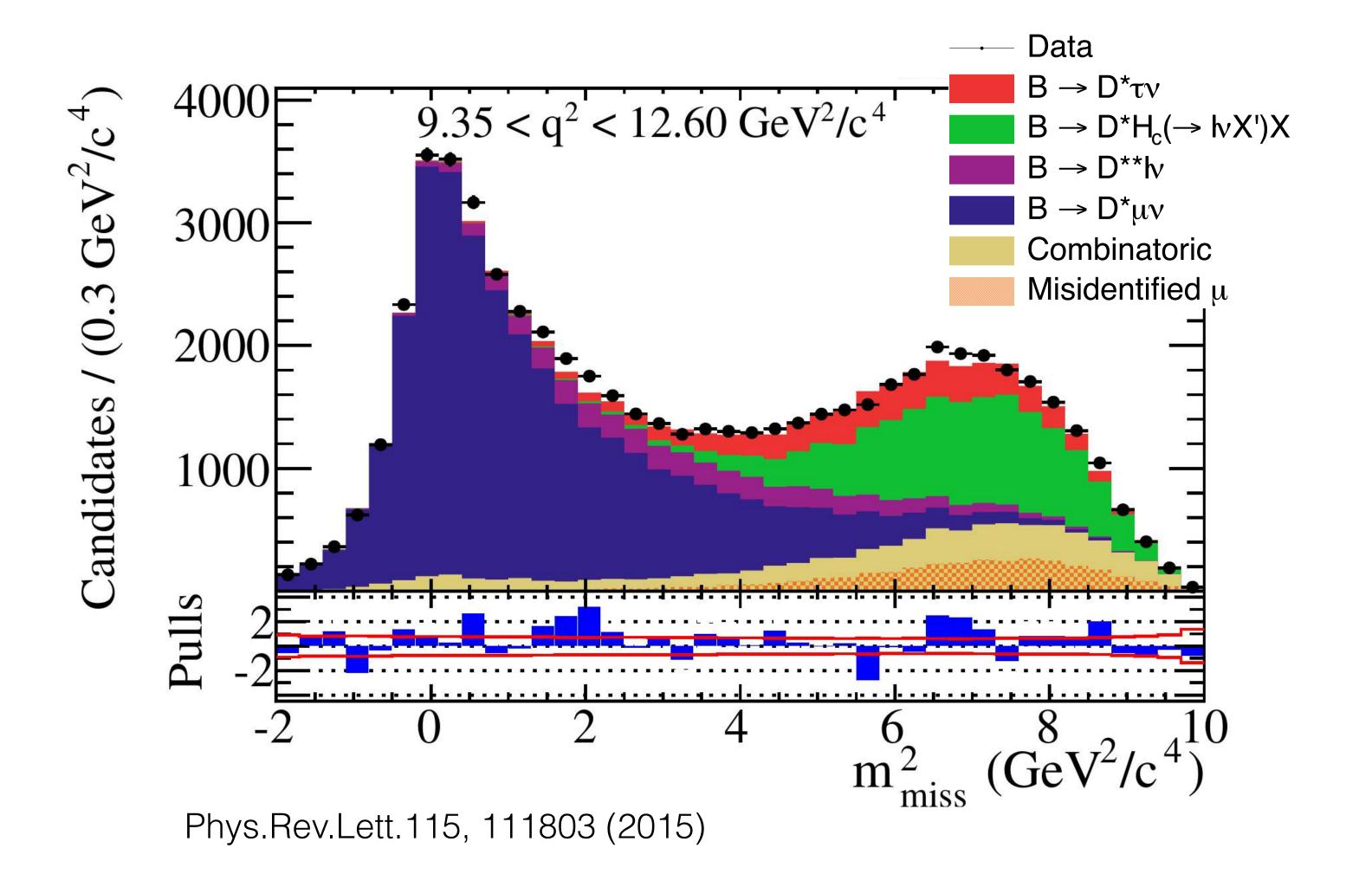
## The problem with neutrinos at a hadron collider

- At least two neutrinos in the final state.
  - No sharp peak to fit in any distribution.



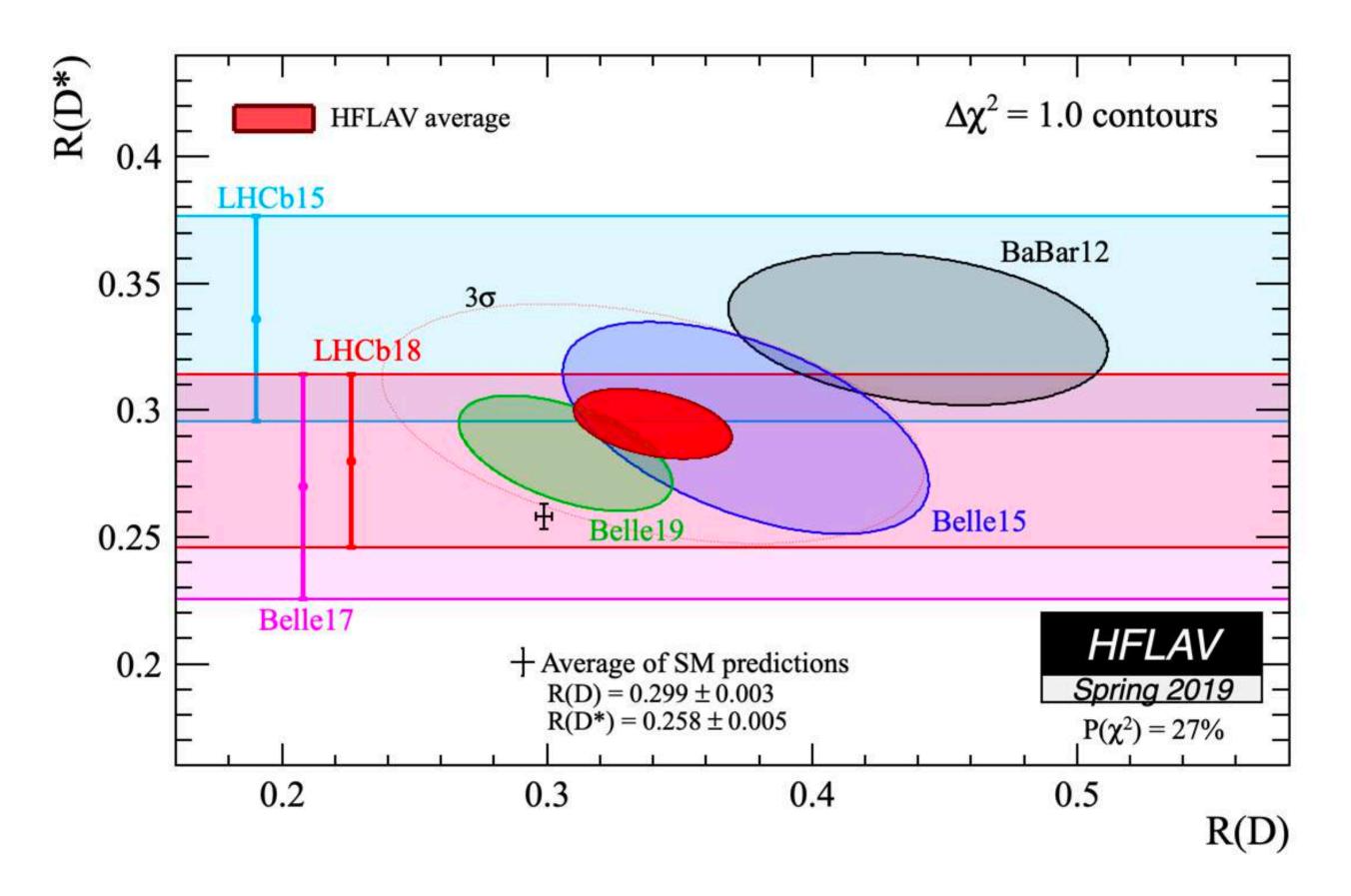
## Signal extraction

Fit the missing mass, discriminating between muon and tauonic mode.



## Hints of an excess

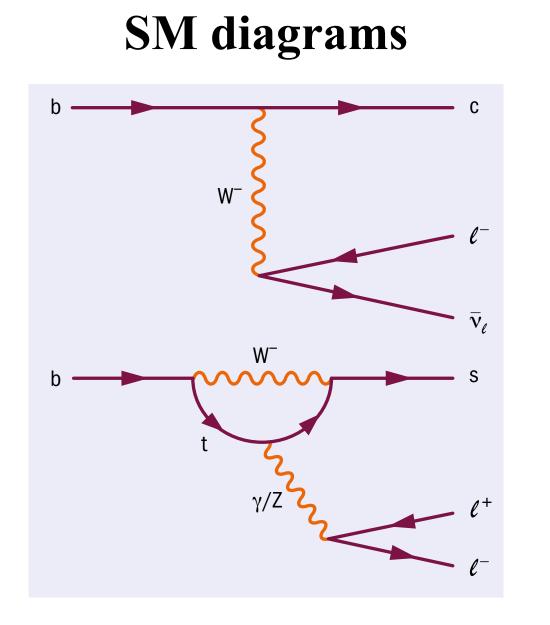
• Three different experiments see an excess in the number of semitauonic candidates.

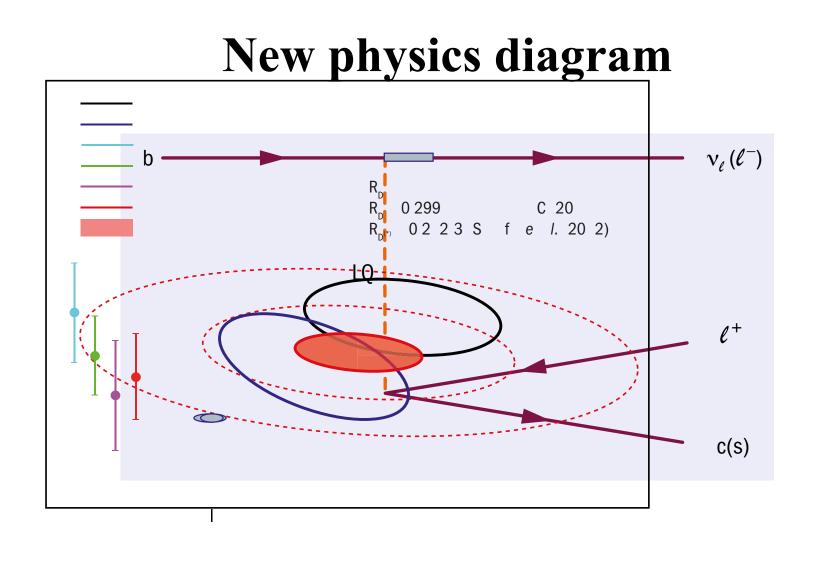


• Combined deviation around 3  $\sigma$ , strongly motivates new measurements in different decay channels.

### What could this all mean?

- We have two sets of anomalies in charged and neatral current semileptonic B decays.
- They both point towards a violation of lepton universality.
- Possible to explain both anomalies with a single new particle (leptoquark) of around 2TeV mass.





- If these anomalies are confirmed it would mean a new force of Nature.
  - Extraordinary claims need extraordinary evidence...

#### What's next?

- We have not yet fully exhausted the current dataset at LHCb. Upcoming measurements:
  - Measurement of R<sub>K(\*)</sub> in the high q<sup>2</sup> region. Talk by Vadym Denysenko (UZH)
  - Measurement of new R ratios with different hadron species (R<sub>Kπ</sub>, R<sub>Kππ</sub>, R<sub>Φ</sub>, R<sub>D</sub>, R<sub>D</sub>, R<sub>J/ψ</sub>)
  - Angular analysis of B<sup>0</sup>—>K<sup>\*0</sup>e+e- decays. Discussed by Zhenzi Wang (UZH) and by Michele Atzeni (UZH)
  - Lepton universality tests with baryons (Λ<sub>b</sub> baryons).
  - Searches for  $b \to s \tau^+ \tau^-$  Talk by Martin Andersson (UZH)
  - + many more ...
- Exciting times to be on LHCb!

