# Intermittency Analysis of charged particles generated in Xe-Xe collisions at $\sqrt{s}_{NN} = 5.44$ TeV using the AMPT Model



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### Outline

- Physics Motivation
- Methodology and Observables
- Observations
- Summary



### **Physics Motivation**

- As strongly coupled QGP expands initial state conditions of the collisions transfer into the final-state collective behaviour
- Large density fluctuations (QCD)

Fluctuations in the geometrical configurations (Spatial patterns)

#### One of the proposed measures

- to characterize QGP
- and hence to add to information about the QCD phase diagram/particle production mechanismm

To study Fluctuations in spatial patterns

#### To study Fluctuations in spatial patterns

#### scaling properties of multiplicity fluctuations over wide range of bin sizes using

#### **Normalized Factorial Moments (NFM)**

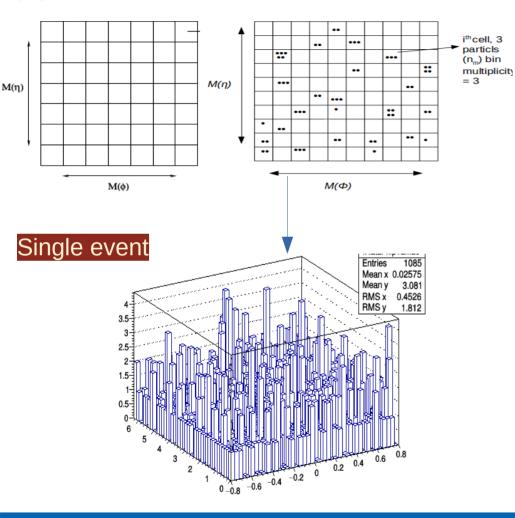
#### Intermittency analysis

#### **Primary Sources:**

- R.C. Hwa and C. B. Yang, Acta Physica Polonica B . Vol. 48 Issue 1 (2017).
- R.C. Hwa & C.B. Yang, PRC 85, 044914 (2012), nucl-ex:1411.6083
- R.C. Hwa and M.T. Nazirov, Phys. Lett. 69, 741 (1992).

#### Methodology & Observables

### Eta-phi phase space of an event is binned into M×M bins



#### Phase space $(\eta, \Phi)$ is divided into a square lattice

- M(η) and M(Φ) : Number of bins along eta and phi axis respectively.
- Charged particles are mapped into this 2D matrix
- Number of particles that go in each cell defines the bin multiplicity (n<sub>ie</sub>)

$$f_q(n_{ie}) = n_{ie}(n_{ie}-1)(n_{ie}-2)....(n_{ie}-q+1)$$

- *q* is the order of the moment (positive integer)
   *q* ≥ 2
- $n_{ie}$  is bin multiplicity,  $n_{ie} \ge q$ , *M* is number of bins.

### **Observables**

Normalized factorial Moments

$$F_{q}(M) = \frac{\frac{1}{N} \sum_{e=1}^{N} \frac{1}{M} \sum_{i=1}^{M} f_{q}(n_{ie})}{\left(\frac{1}{N} \sum_{e=1}^{N} \frac{1}{M} \sum_{i=1}^{M} f_{1}(n_{ie})\right)^{q}}$$

where M is number of bins, N is the number of events

M-Scaling 
$$\longrightarrow F_q(M) \propto M^{\varphi_q}$$

- → φ<sub>q</sub> is the intermittency index and can be acquired as a quantity characterizing fractality
- $F_{a}$  filters out statistical fluctuations
- Robust against the uniform efficiency losses
- A. Bialas and R. Peschanski, Nucl. Phys. B, 273 (1986).R.C. Hwa and C. B. Yang, Acta Physica Polonica B . Vol. 48 Issue 1 (2017).

Other than M-scaling, there exist another scaling which is

$$F_q(M) \propto F_2(M)^{\beta_q} \longrightarrow F$$
-Scaling

$$\beta_q \propto (q-1)^{\nu}$$

v: is scaling exponent (dimensionless exponent) gives economical summary of the system under studyAim

- Scaling @ LHC energies?
- v @ LHC energies?

### 

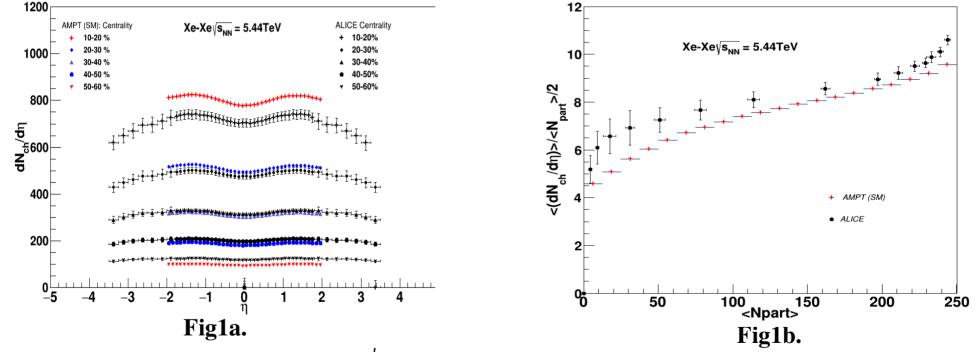
2. R.C. Hwa and C.B. Yang, PRC 85, 044914 (2012).

### **Observations: AMPT**

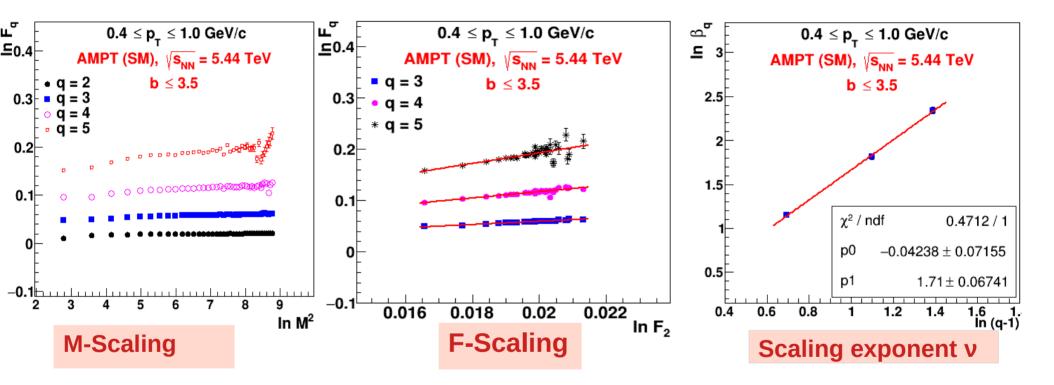
• Is a hybrid transport model and is designed to model the heavy-ion collisions available at

relativistic energies

• AMPT with string melting version

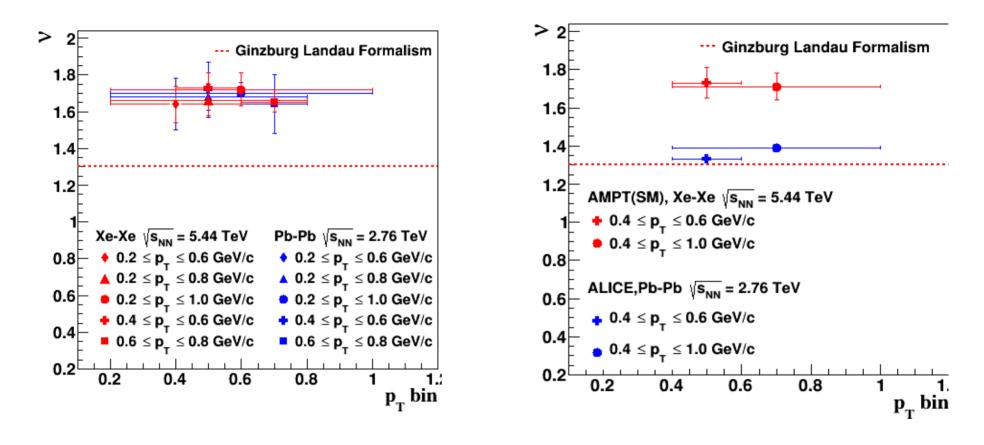


- 500K minimum biased Xe-Xe events at  $\sqrt{s_{NN}} = 5.44$  TeV generated, impact parameter  $0 \le b \le 3.5$  fm.
- Pseudorapidity density distribution (Fig1a), particle density vs number of participants (Fig1b)



• The M-scaling (left) and F-scaling (middle) behaviour of the NFM  $(F_q)$  are studied for  $p_T$  bins [0.4-1.0] and the scaling index (right) value are calculated for the same.

### **Observations**



ALICE Results: Poster Session (today) Pythia+Angantyr: Flash Talk (2-8-22)

### Summary

•Scaling properties of the charged particles generated in the mid rapidity region in Xe-Xe collisions at

 $\sqrt{s_{_{NN}}}$  = 5.44 TeV have been studied in the framework of intermittency analysis.

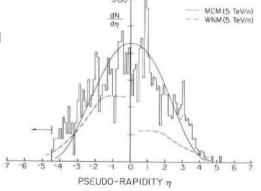
- A power-law growth of NFM (ln F<sub>q</sub>) with ln M is observed to be absent at high M values in all the p<sub>T</sub> bins.
- However F-scaling is observed.
- With no phase transition physics implemented in (SM) AMPT model
  - Scaling behaviour in line with intermittency is absent.
  - Scaling exponent value (v) is greater than the value predicted by theory for second order phase transition.

Thanks ....

## Back up

## **Motivation: Fluctuations and Intermittency**

- Large local density fluctuations exist in the process of space-time evolution in high-energy collisions
- To decide whether these fluctuations are dynamical, i.e. larger than expected from Poisson noises,
  - Bialas and Peschanski have suggested the use of normalized factorial moments (NFM)
- These dynamical fluctuations in high-energy collisions can be manifested as an abnormal scaling
  - when the corresponding collision system is a fractral



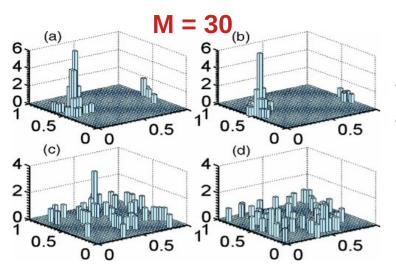
A. Bialas and R. Peschanski, Nucl. Phys. 273, 703 (1986).; 308, 857 (1988)

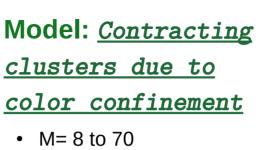
R. C. Hwa et al, PRL 69, 741 (1992); R. C. Hwa et al, PRC 85, 044914 (2012).

Scaling properties of multiplicity fluctuations over wide range of bin sizes (Intermittency Analysis)

#### **Model Predictions**

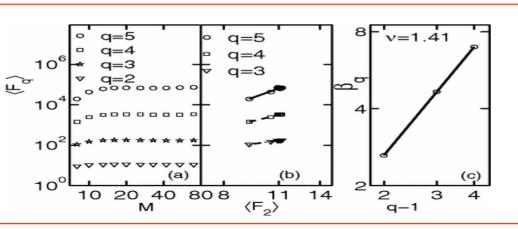
#### R. C. Hwa et al, PRC 85, 044914 (2012)





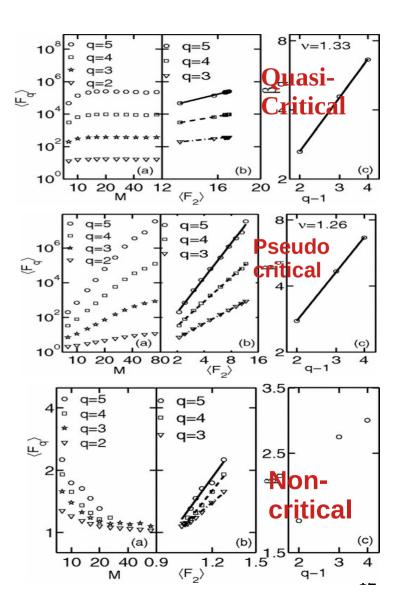
- $\Delta p_{T} = 0.04, 0.07, 0.1$ GeV/c around  $p_{T} = 1$ GeV/c
- Quar-hadron phase transition

**M-Scaling** 



**F-Scaling** 

**Critical** 

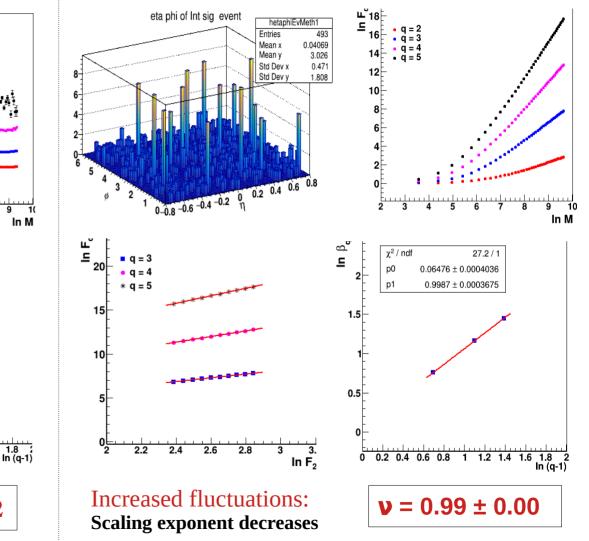


### **Toy MC: Baseline value of exponent**

**Toy MC** щ Ц Toy Model eta phi of an event hetaphiEv q = 3 493 Entries 0.01257  $\mathbf{q} = 4$ Mean x 0.08 - • q = 5 Mean y 3.127 Std Dev x 0.4695 Std Dev y 1.884 2.5 0.06 0.04 0.02 n 0 2 0.4 0.6 0-0.8 -0.6 -0.4 -0.2 -0.02 யீ ⊆0.1 Toy Model പ്പ  $\chi^2$  / ndf 2.149/1 2 p0  $-0.08611 \pm 0.01837$ q = 3 2.5 q = 4p1  $1.565 \pm 0.02204$ 0.08 \* q = 5 0.06 1.5 0.04 0.02 0.5 0.2 0.4 0.6 0.8 0.006 1.2 0.0045 0.005 0.0055 In F<sub>2</sub>  $v = 1.565 \pm 0.022$ No/least fluctuations

1.8

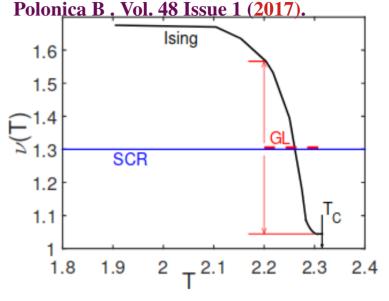
#### **Toy MC: 1% fluctuations in each events**



#### Scaling exponent (v): Data vs Theory and model predictions

ν	Different Models
1.304	Without kinetic term in GL potential
	<b>GL Formalism:</b> R.C. Hwa, M.T. Nazirov, Phys. Rev. Lett. 69, 741 ( <b>1992</b> )
1.316 ± 0.012	With kinetic term in GL potential PLB 297, 35-58 (1992)
1.41	<ul> <li>SCR Model : R. C. Hwa and C. B. Yang, Phys. Rev. C 85, 044914(2012)</li> <li>Ising Model and SCR Model : R. C. Hwa and C. B. Yang, Acta Physica Polonica B 48, (2016)</li> </ul>
1.55 ± 0.12	EMU01 and KLM Collaboration EMU01: M.I. Adamovic PRL, 65,412 (1990) KLM: PRL, 62, 733 (1889); PRC 40, 2449 (1989)
1.79±0.10	R. Sharma and R. Gupta, AHEP, Article ID 6283801 (2018)
	(small $p_T bin \sim 0.2$ , 0.4 to 0.6 GeV/c, M: 5 to 30) $ \eta  < 0.8$ , full azimuthal coverage
<b>RHIC energies</b> 1.75 ±0.12 EPOS3 $\nu = 1.743$ from Ur $\nu = 1.94$ from AM	R. Gupta and Salman Khurshid Malik AHEP. Article IDQMD model.S. Bhattacharyya, EPJP 136, 471 (2021).PT model.X. Y. Long et al, NPA 920, 33-34 (2013).
	$ \eta  < 0.8$ , full azimuthal coverage

R.C. Hwa and C. B. Yang, Acta Physica



The curve v(T) provides a model-dependent interpretation of the value of v in terms of

- temperature  $T_{c}$  (=2.315 J/k<sub>B</sub>) is the transition temperature in • Ising system
- Dashed red line is the GL value (average of the Ising values between T = 2.2 and  $T_{o}$ )
- However in real system transition occurs at lower temperature (=2.21 J/k<sub>B</sub>) 14