

Bosons

Higgs boson

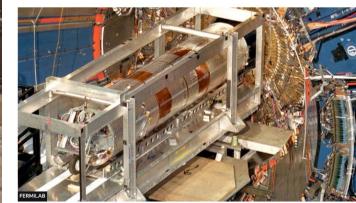
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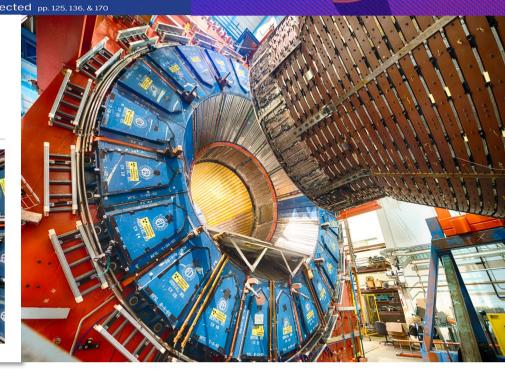
Shock result in particle experiment could spark physics revolution By Pallab Ghosh ^③ 7 April

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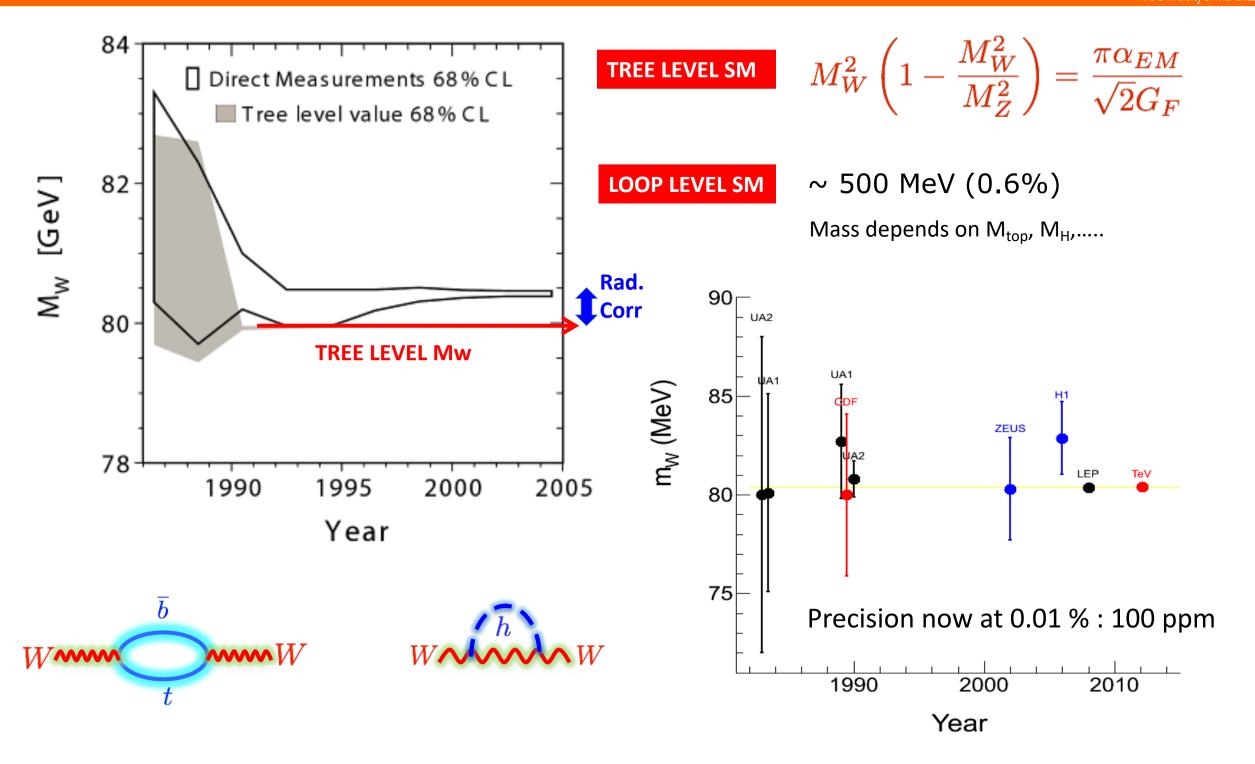
The Fermilab Collider Detector obtained a result that could transform the current theory of physics

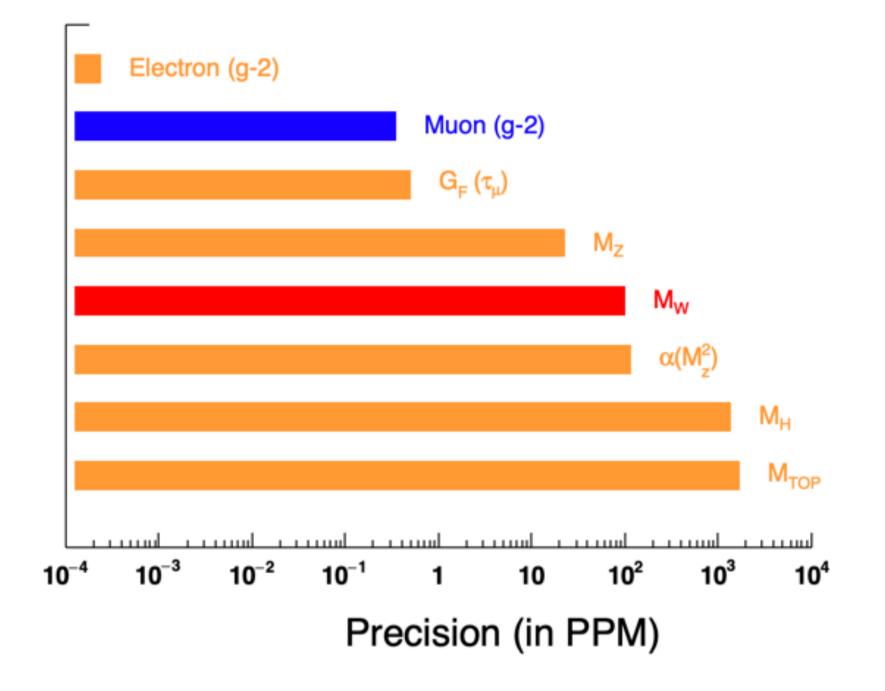


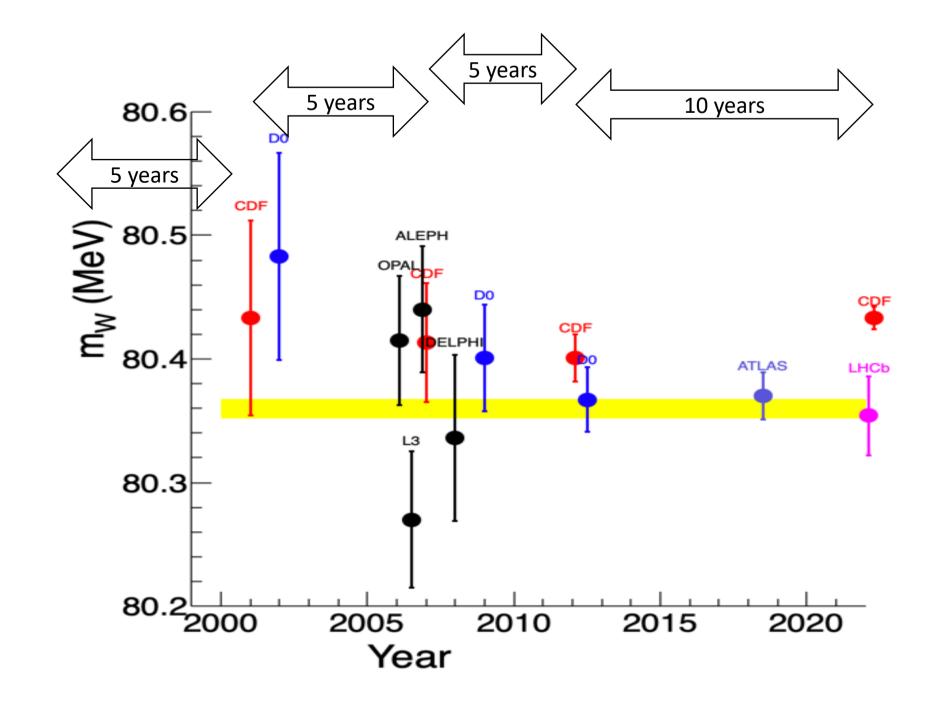
(CDF) W Mass : Experimental Review

Mark Lancaster : University of Manchester

W Mass in SM

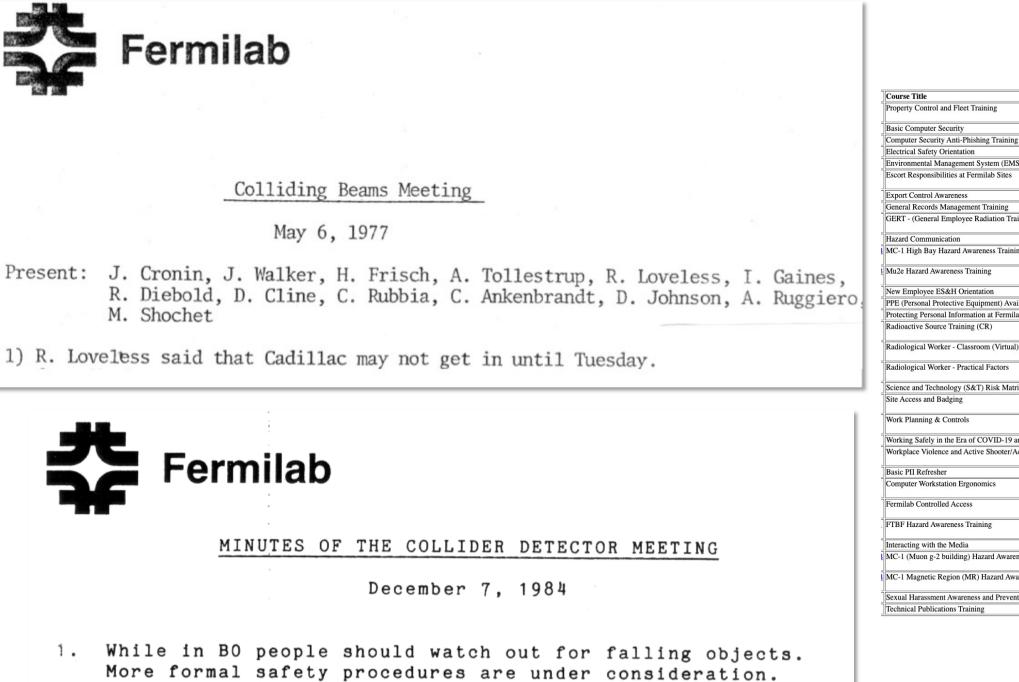




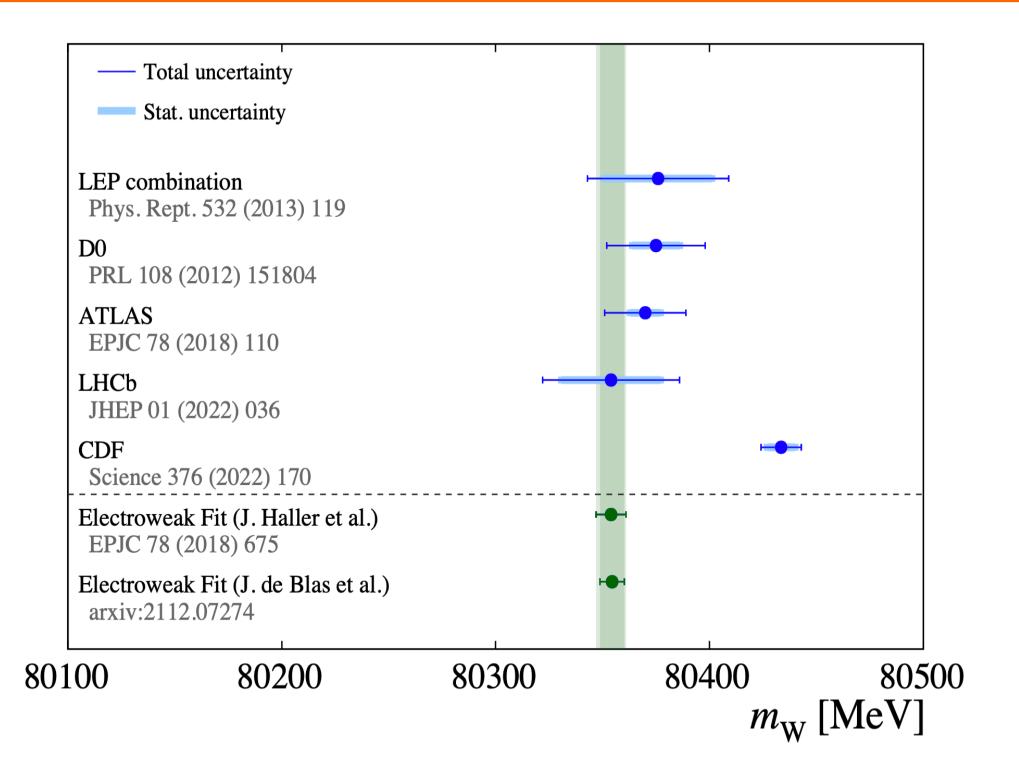


Complete Date

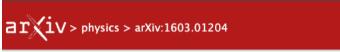
02/10/2019



02/10/2019 08/28/2006 Environmental Management System (EMS) 08/28/2006 01/03/2022 04/10/2022 04/10/2022 GERT - (General Employee Radiation Training) 01/03/2022 08/28/2006 MC-1 High Bay Hazard Awareness Training 01/03/2022 01/03/2022 08/28/2006 PPE (Personal Protective Equipment) Availability And Use 08/28/2006 Protecting Personal Information at Fermilab 05/30/2010 01/03/2022 Radiological Worker - Classroom (Virtual) 01/03/2022 01/03/2022 Science and Technology (S&T) Risk Matrix Lab-wide Training 01/03/2022 10/11/2020 01/03/2022 Working Safely in the Era of COVID-19 and the Return to On-site Work 01/03/2022 Workplace Violence and Active Shooter/Active Threat Awareness Training 07/19/2022 06/04/2019 01/03/2022 01/03/2022 01/03/2022 01/03/2022 MC-1 (Muon g-2 building) Hazard Awareness Training 01/03/2022 MC-1 Magnetic Region (MR) Hazard Awareness Training 01/03/2022 Sexual Harassment Awareness and Prevention for Fermilab Users, Visitors and Contract Employees 06/27/2016 03/24/2015



Ambulances



Physics > Physics and Society

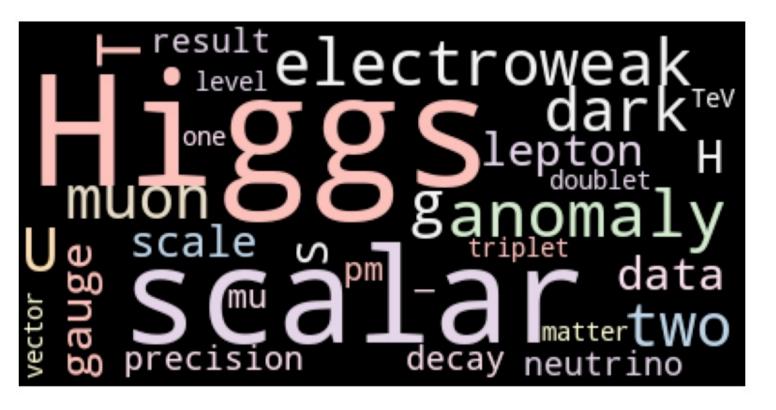
[Submitted on 3 Mar 2016]

A Theory of Ambulance Chasing

Mihailo Backović

Ambulance chasing is a common socio-scientific phenomenon in particle physics. I argue that despite the seeming complexity, it is possible to gain insight into both the qualitative and quantitative features of ambulance chasing dynamics. Compound-Poisson statistics suffices to accommodate the time evolution of the cumulative number of papers on a topic, where basic assumptions that the interest in the topic as well as the number of available ideas decrease with time appear to drive the time evolution. It follows that if the interest scales as an inverse power law in time, the cumulative number of papers on a topic is well described by a di-gamma function, with a distinct logarithmic behavior at large times. In cases where the interest decreases exponentially with time, the model predicts that the total number of papers on the topic will converge to a fixed value as time goes to infinity. I demonstrate that the two models are able to





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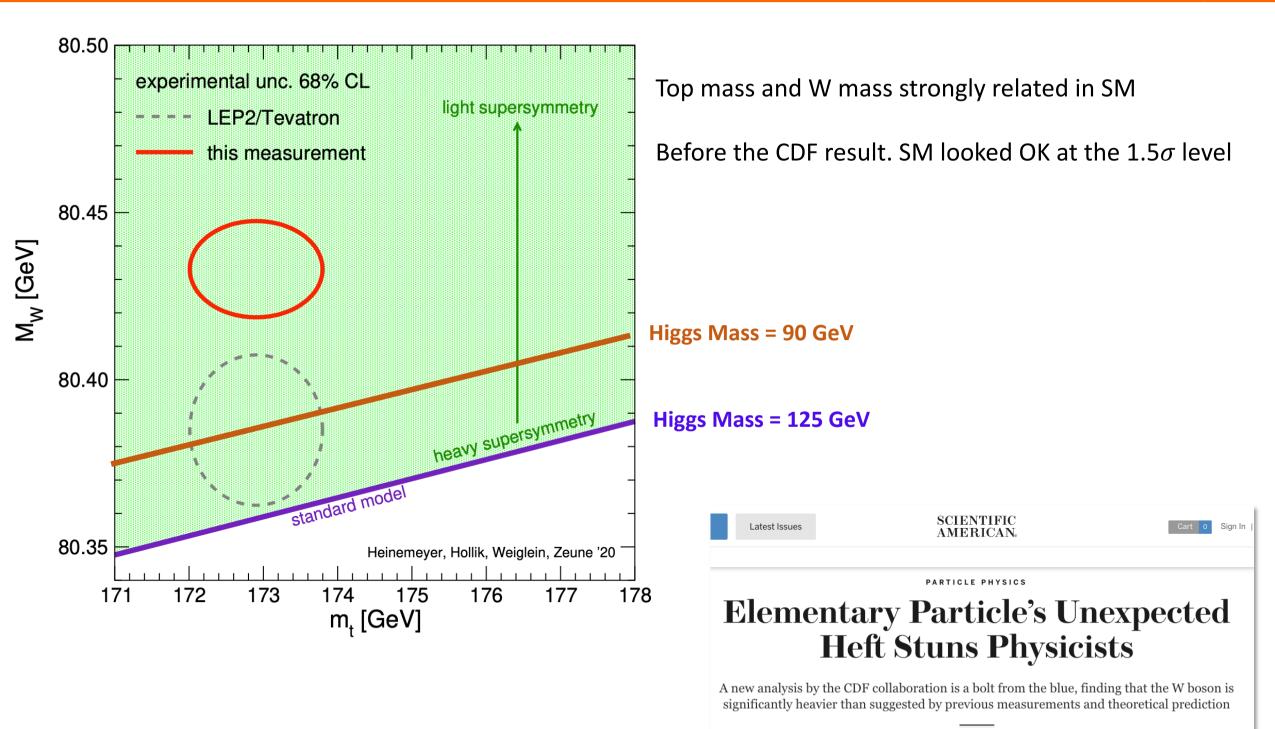
~ 100 papers since April with theories of why Mw should be above the SM.

About 20% of these also try and explain the high (g-2) at the same time

Many of these BSM ideas date back 20 years : e.g. Higgs triplets

Why the excitement ?

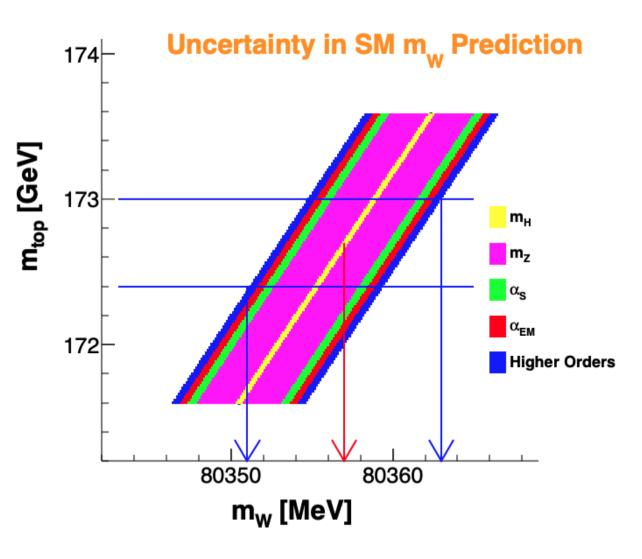




By Daniel Garisto on April 7, 2022

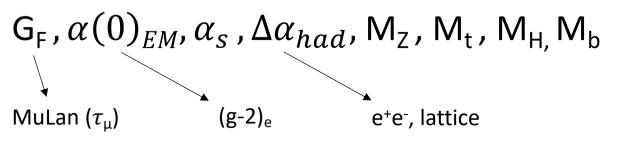


Predicted Mw has uncertainty 6-8 MeV : half from theory, half from exp. measurements Top mass and Z mass uncertainty are now the dominant drivers in SM Mw uncertainty Our friend : $\Delta \alpha_{had}$ is also important. This is basically (g-2) HVP



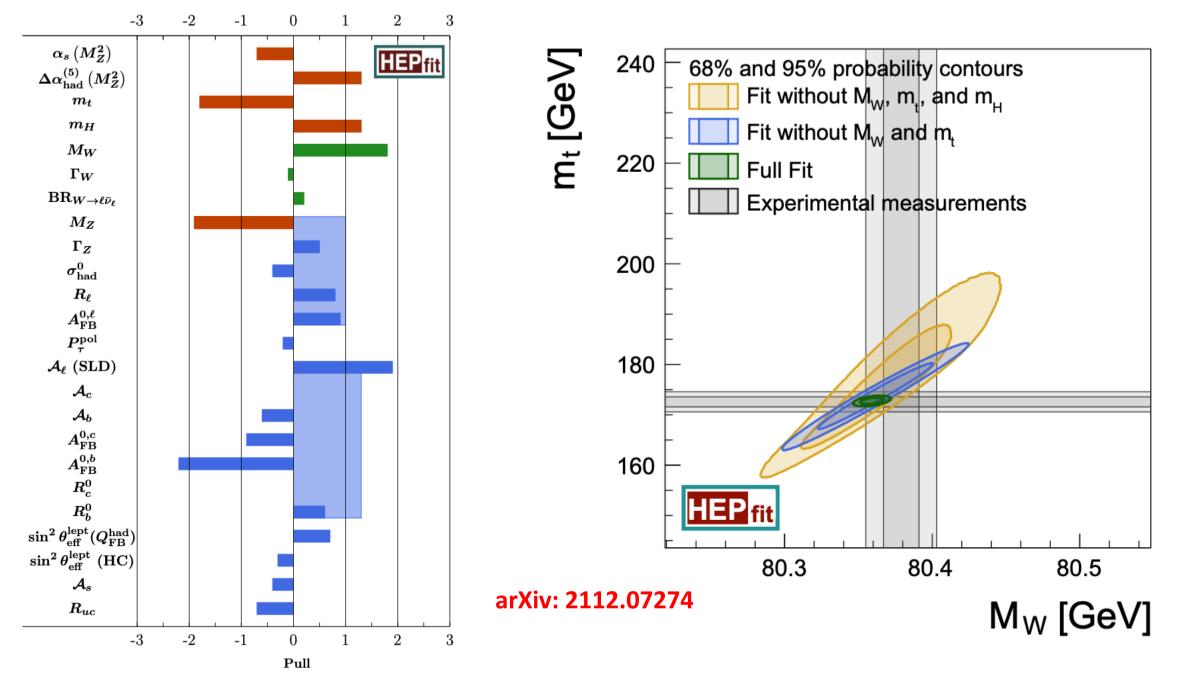
Could this be wrong ?

The inputs are:



 $\Delta \alpha_{had}$ comes from analysis of same e⁺e⁻ data (or lattice calculation) as that determining (g-2) HVP.

SM Electroweak Fits / Constraints

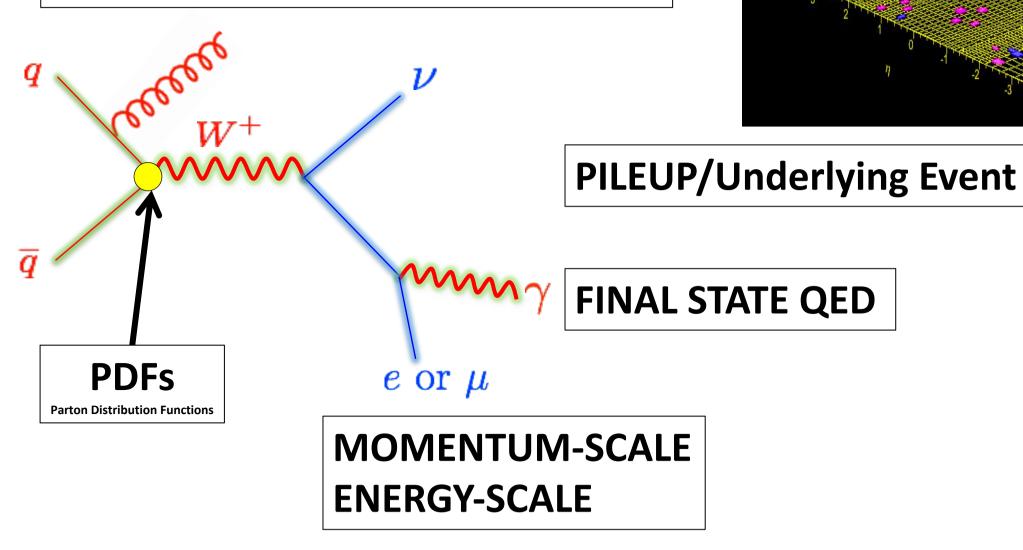


N.B. peculiarly in break with tradition this uses BMW lattice QCD result to constrain: $\Delta \alpha_{had}$ and not the e⁺e⁻ cross section data

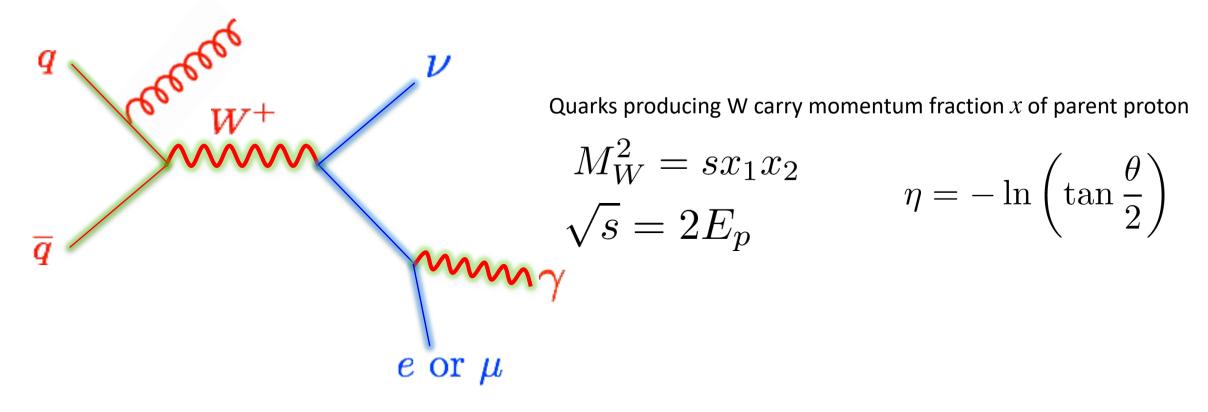
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INITIAL STATE RADIATION (aka RECOIL)

- BOTH QCD AND QED
- TO NON PERTURBATIVE REGION IE $PT(W) \rightarrow 0$







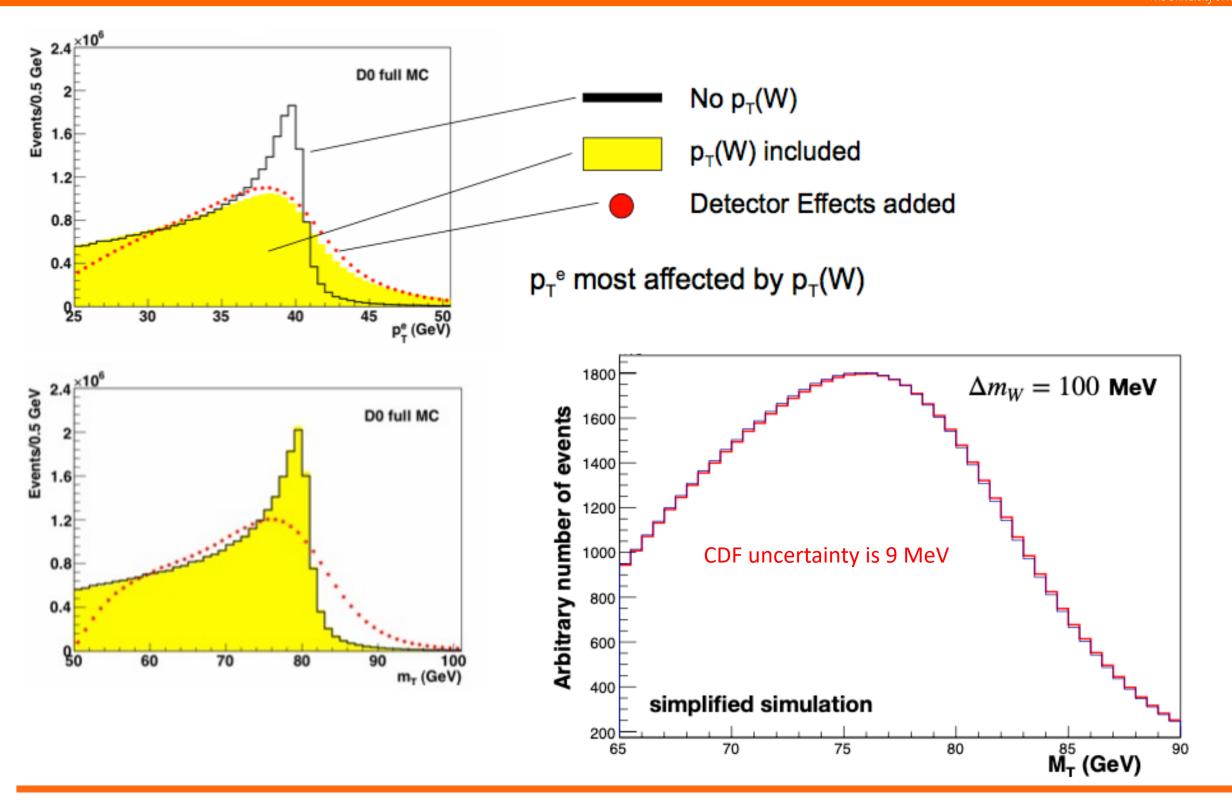
x values at LHC are much lower and so uncertainties (and shifts) from parton distribution functions are different Only measure 2-vectors in plane transverse to beam since there is no measured momentum constraint along the colliding beam direction.

Initial beams have no transverse momentum so \mathbf{p}_{T} is conserved i.e. $\vec{p}_{T}^{\ l} + \vec{p}_{T}^{\ \nu} + \vec{p}_{T}^{\ ISR} + \vec{p}_{T}^{\ FSR} = 0$ $\vec{p}_{T}^{\ \nu} = -\left(\vec{p}_{T}^{\ l} + \vec{U}\right)$

We get the W mass by comparing the transverse quantities $\, p_T^{\,\,
u}, p_T^{\,\, l}, m_T \,$ with simulation

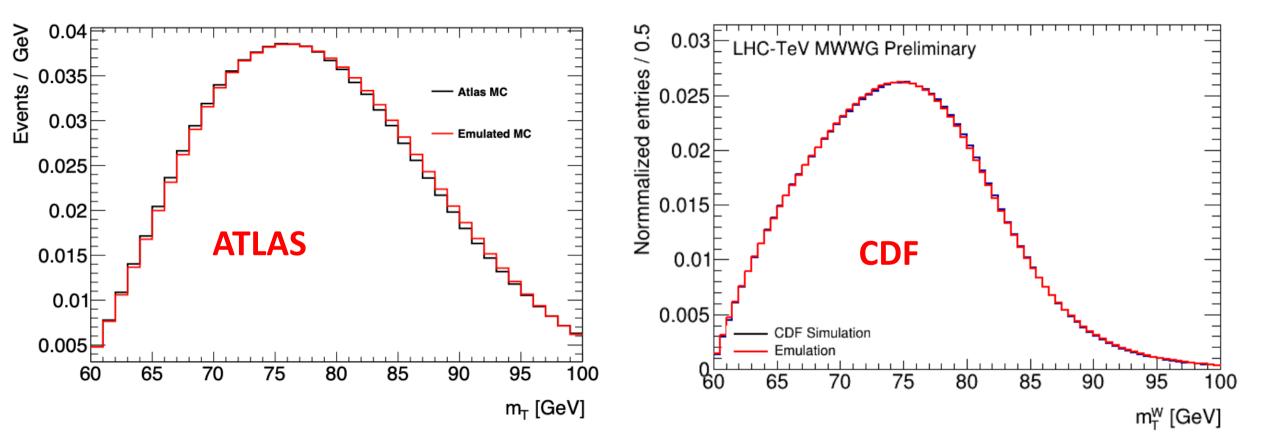
 $m_T = \sqrt{2p_T^{\ \nu} p_T^{\ l} (1 - \cos \Delta \phi_{l\nu})}$ Transverse mass

Transverse Quantities





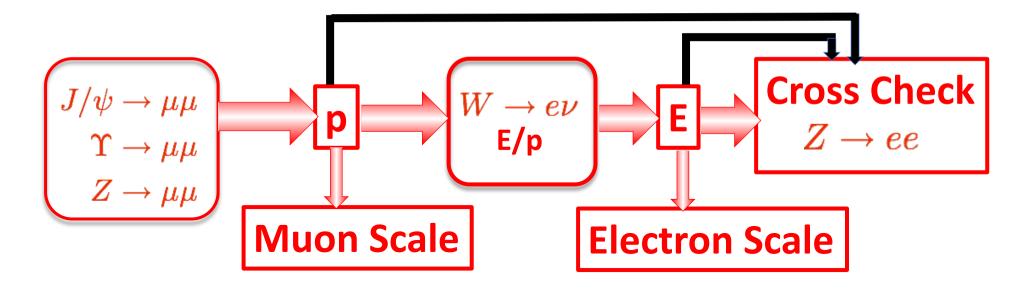




To achieve the same statistical precision as CDF/D0 requires x10 the data at LHC.

Setting the energy and momentum scale





ATLAS is the same except it does not use J/ Ψ or Υ data at low momentum to constrain non-linearity in the momentum scale instead it uses the vast Z statistics to look at Z mass in different momentum bins. W data is ~ 5 GeV lower in momentum than Z data.

LHCb only uses muons and does it the same as CDF.

D0 only uses electrons and sets scale only with $Z \rightarrow ee$ (no E/p cross-check, no J/ Ψ or Υ).

ATLAS and LHCb have better detectors than CDF: better resolution, less leakage.

Arguably CDF has the most internal consistency checks in the scale determination

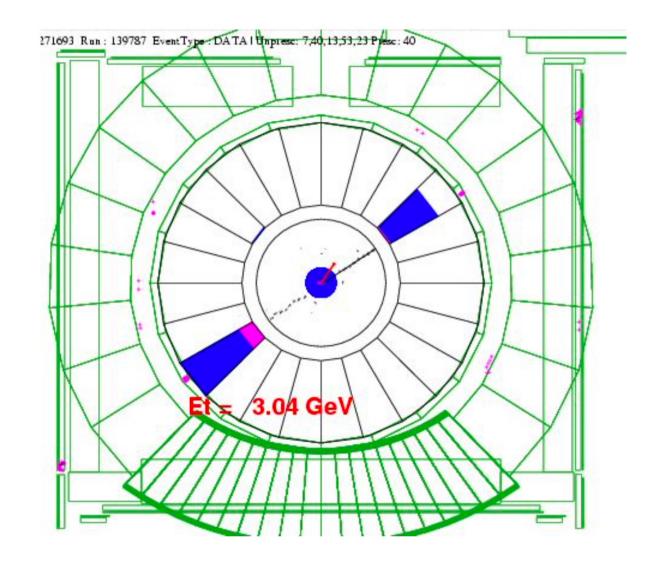
CDF: Momentum Scale





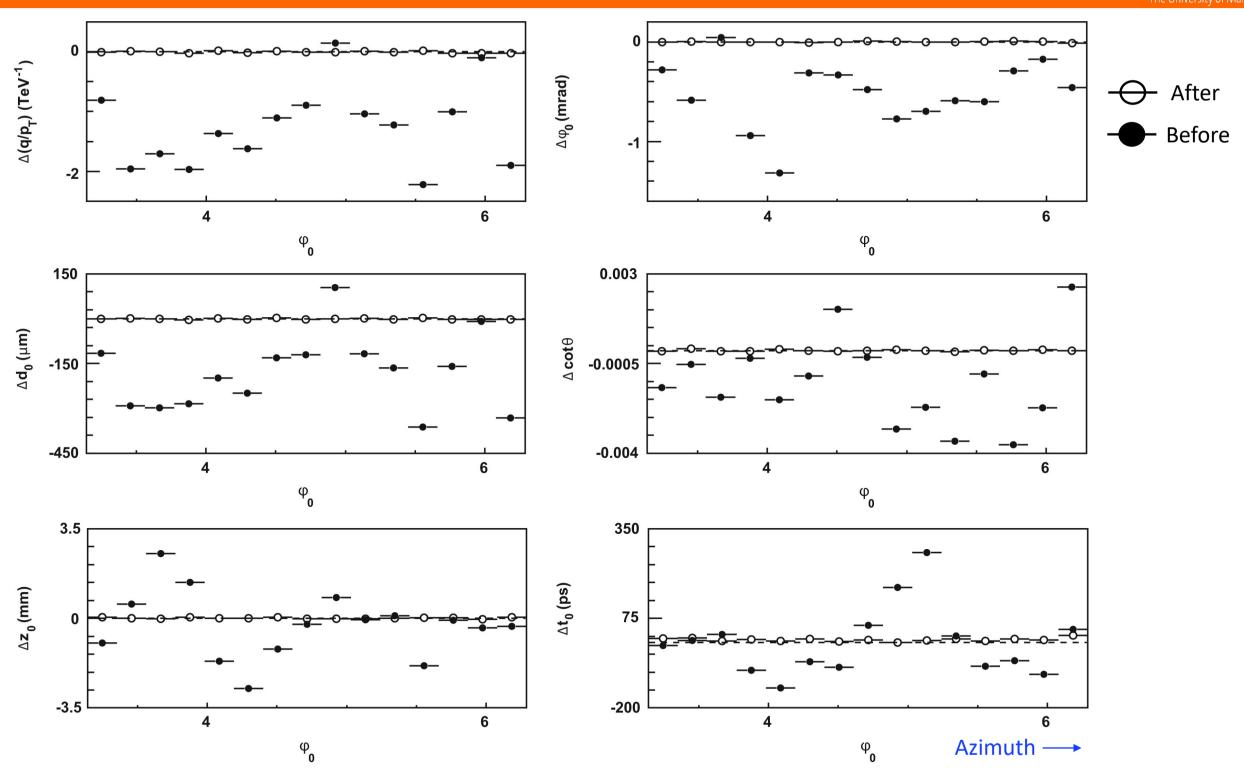


Years of work to calibrate and align the drift chamber and remove biases due to gravity, twists, bends

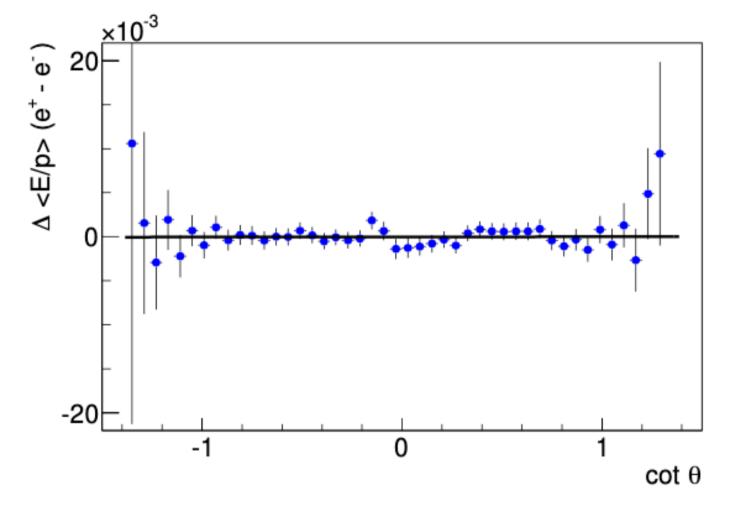


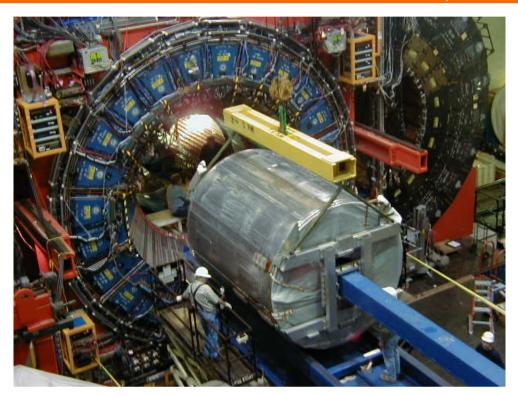
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CDF: Momentum Scale



After corrections no evidence of "false curvatures"



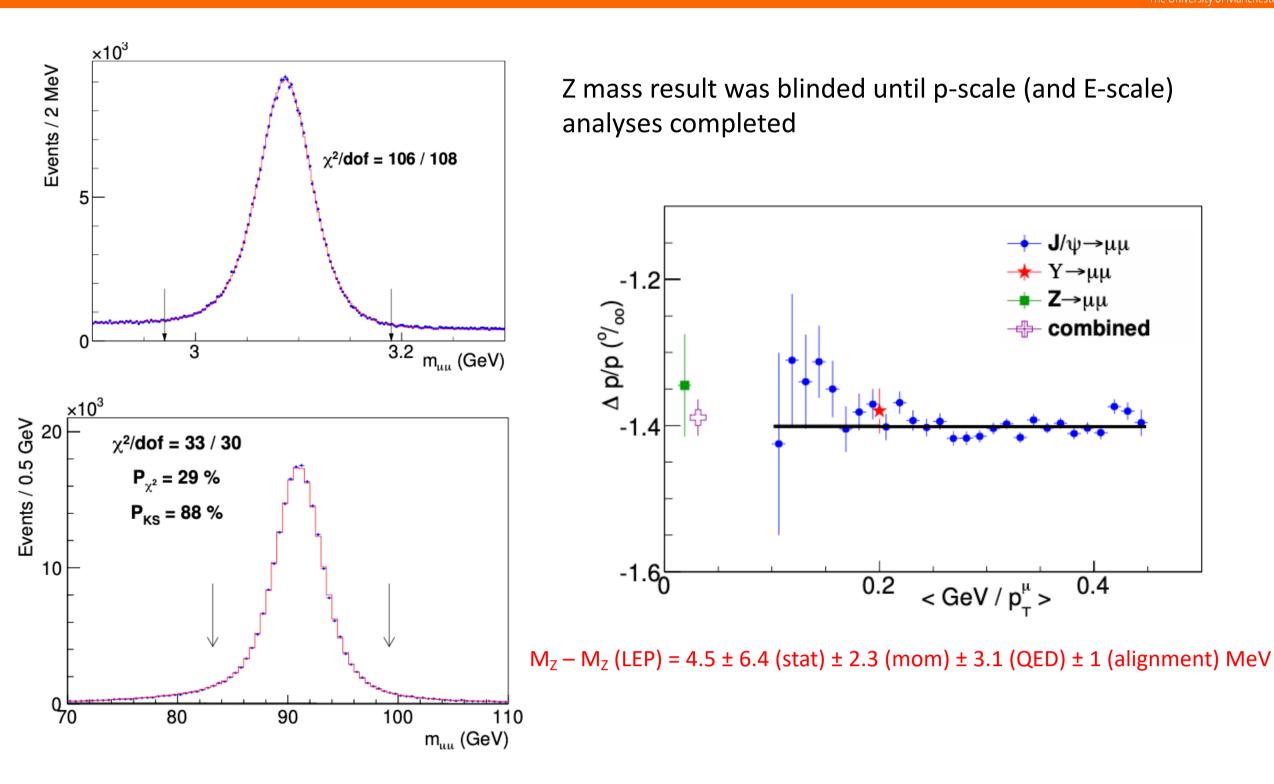




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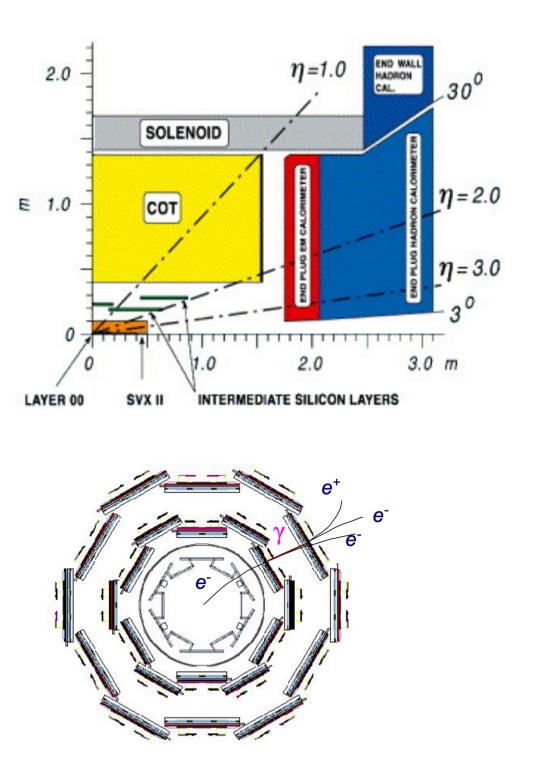
CDF: Momentum Scale



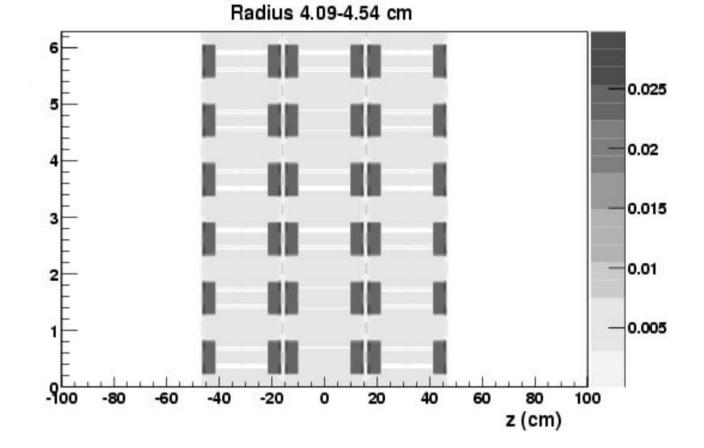


CDF: Energy Scale



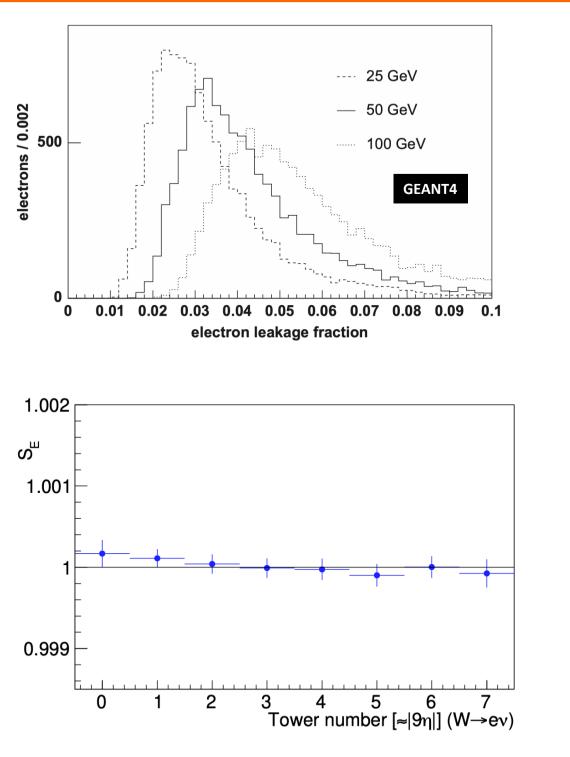


Done via E/p which requires detailed model of mass and Z of all passive material prior to calorimeter



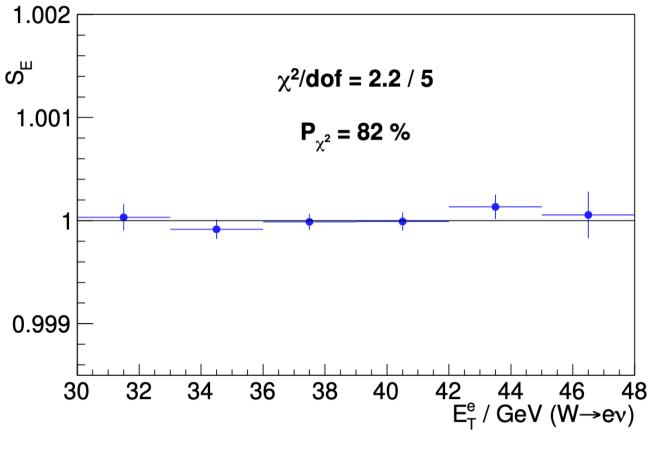
CDF: Energy Scale



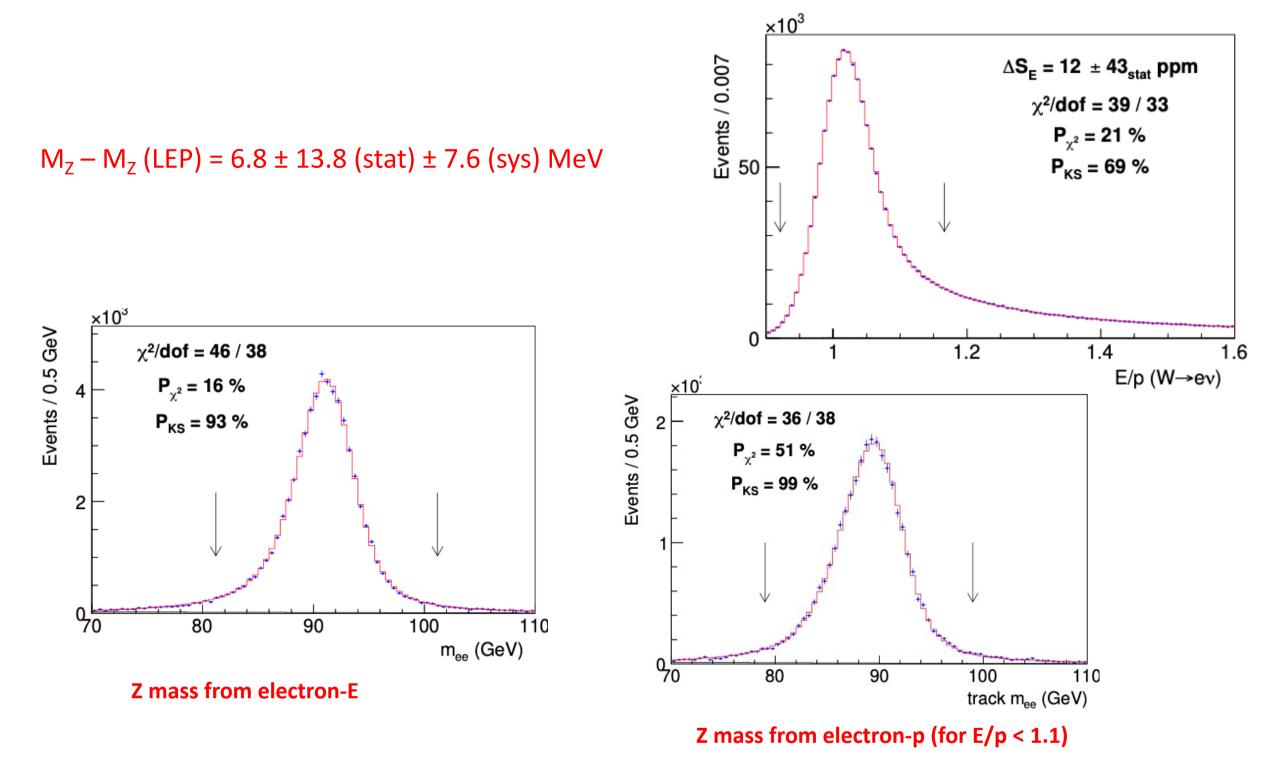


Also need to model

- shower leakage
- QED effects (incl LPM suppression)
- detector response

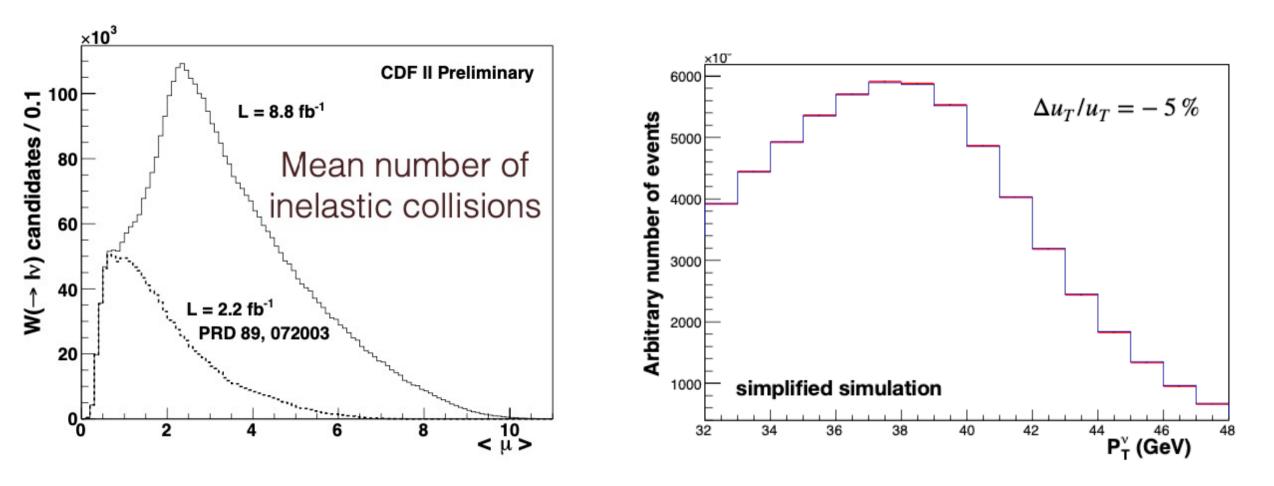


CDF: Energy Scale



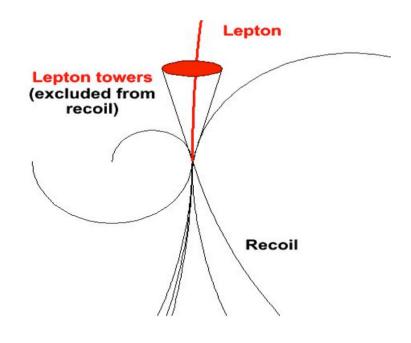


Unlike momentum and energy scale calibration that needs to be good to 0.01% the calibration of the non electron/muon part of the event i.e. initial state QCD radiation, underlying event energy (from additional interactions) only needs to be good to 0.5% or so

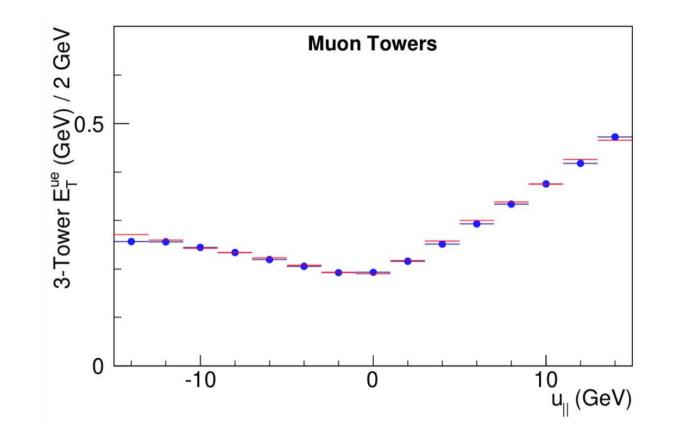


At ATLAS/CMS the mean number of inelastic collisions was ~ x5 that of CDF/D0





Calorimeter towers with lepton are removed from recoil sum and correction for underlying energy removed is made.

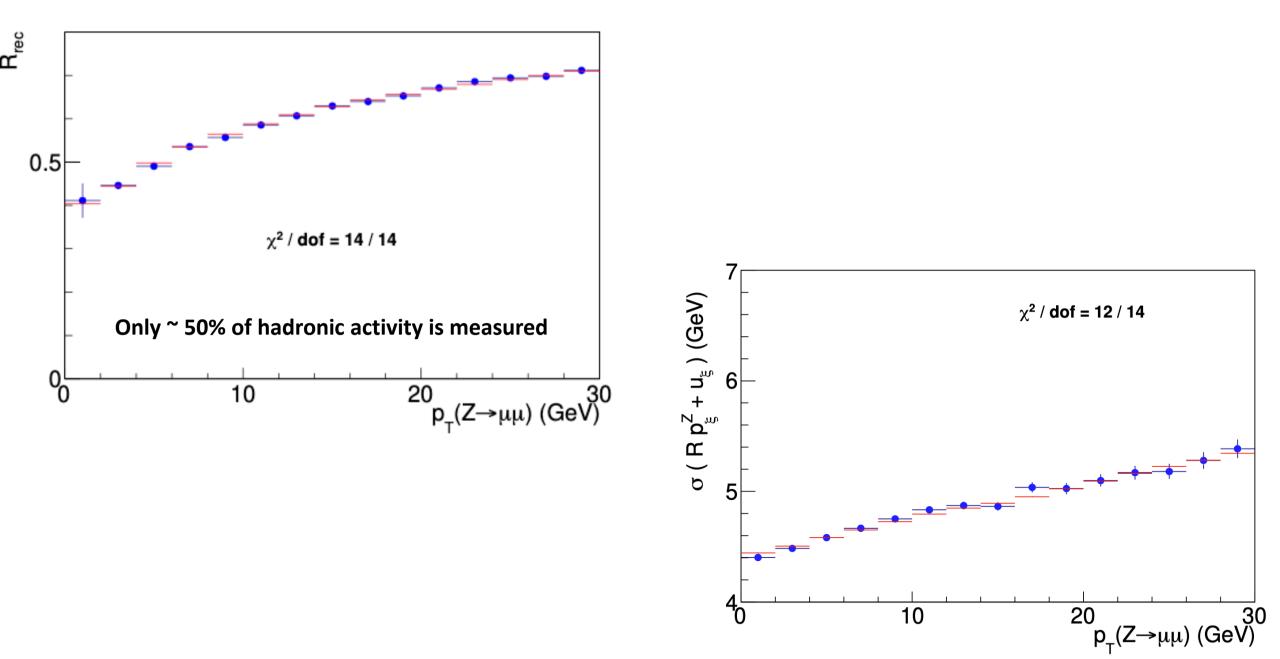


Calorimeter alignment is important

QCD ISR, underlying event is calibrated on Z events (as a function of p_T of the Z)

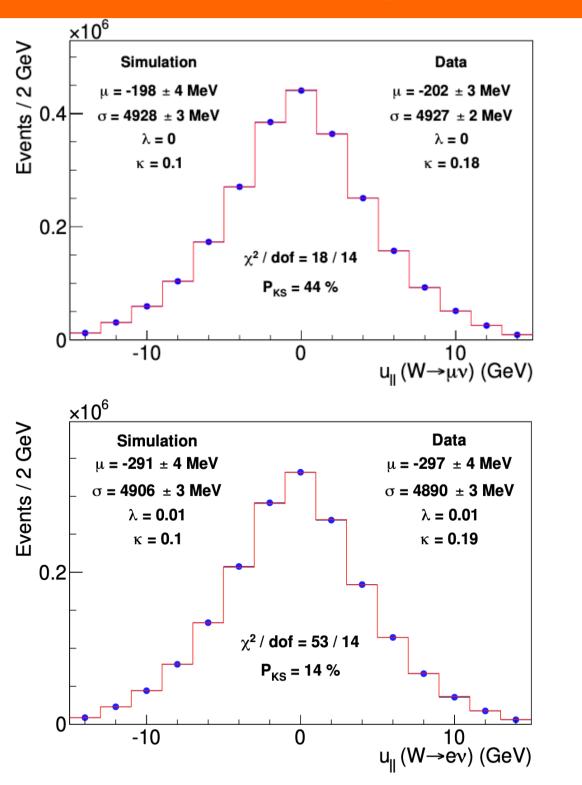
The model to do this ad-hoc: D0, ATLAS, LHCb do the same.

CDF: Recoil / Underlying Event Scale

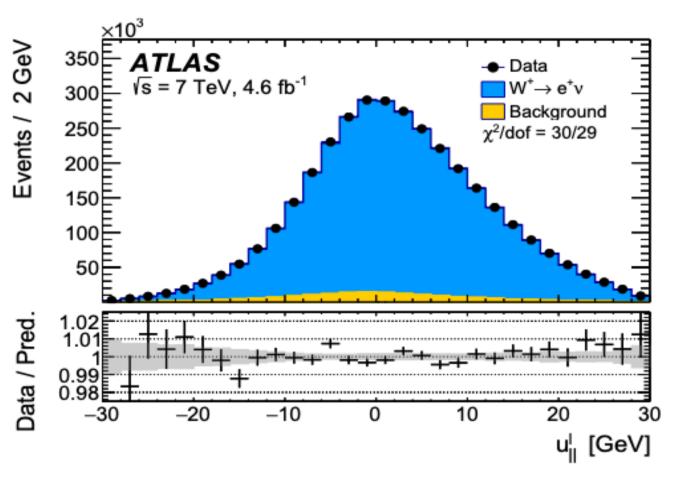


CDF: Recoil / Underlying Event Scale



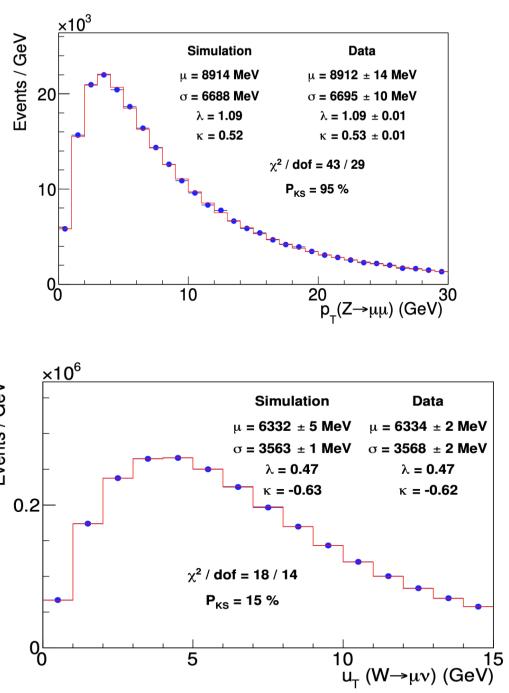


 $m_T \approx 2p_T^l + u_{\parallel}$



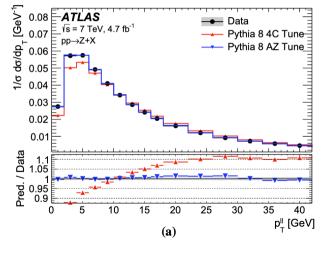
CDF: Recoil / Underlying Event Scale

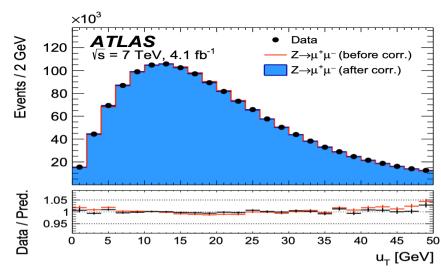


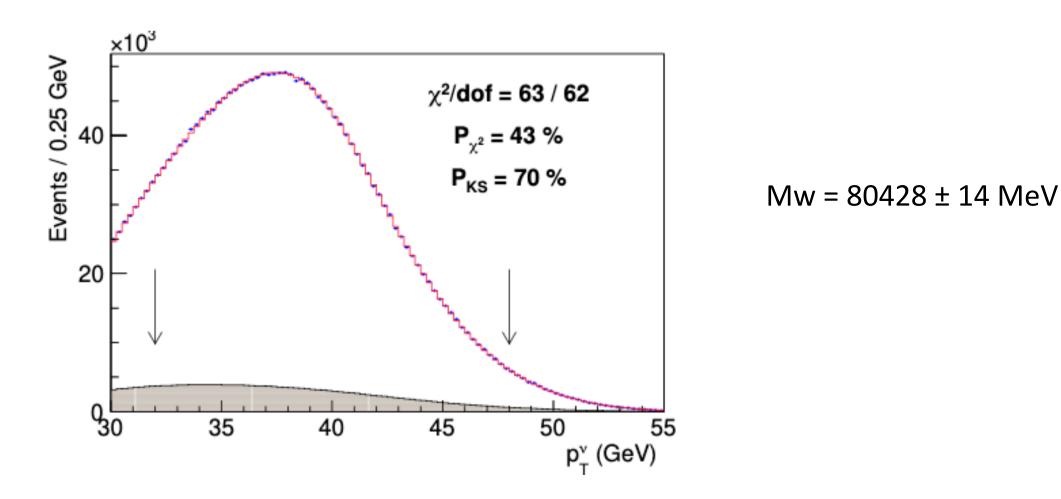


 p_T of the W (and Z) boson has significant contribution from low-momentum radiation only modelled via an ad-hoc nonperturbative contribution and a QCD NNLL resummation.

CDF/D0: PYTHIA/RESBOS, ATLAS: PYTHIA/DYNNLO/POWHEG/MINNLO_{PS} LHCb: PYTHIA/POWHEG/DYTurbo



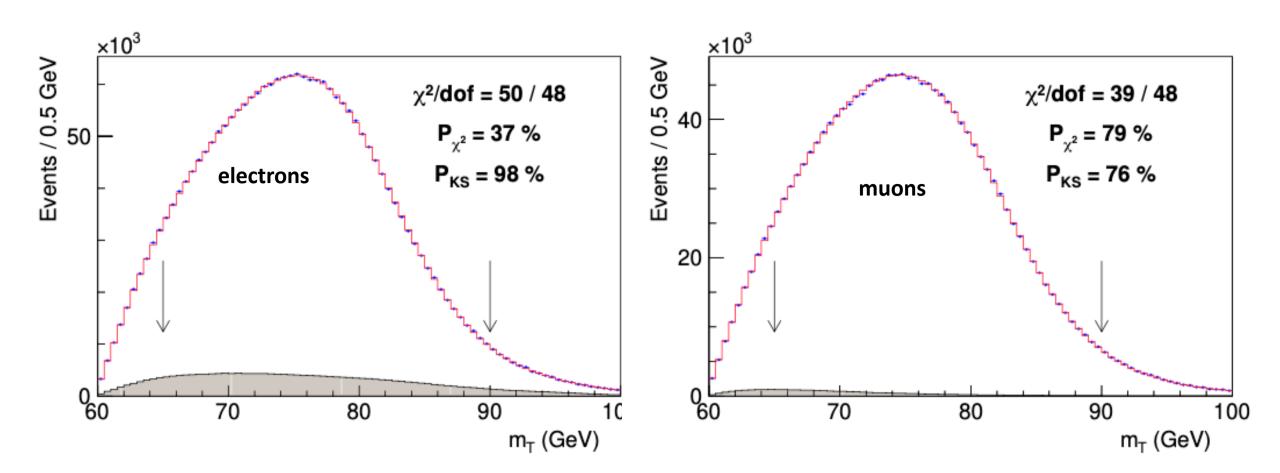




Determination from transverse mass and lepton $p_{\rm T}$ is more precise : 10 MeV and 12 MeV and very correlated

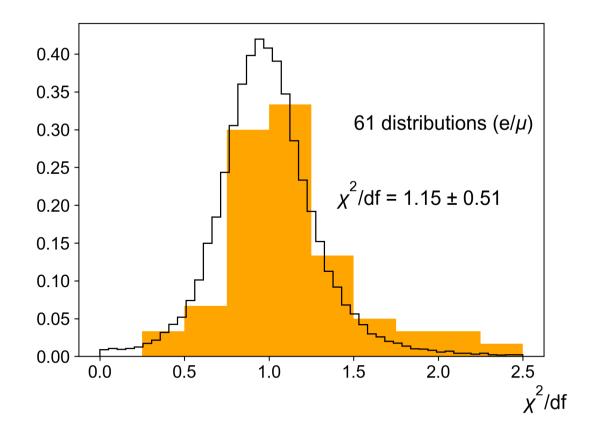


These fits were done with a blinded Mw offset.



Difference in Mw between muons and electrons : $13.3 \pm 15.1 \text{ MeV}$

	Muons	Electrons	
Positive-Negative Lepton	-7.8 ± 22.4 MeV	14.7 ± 22.6 MeV	
Second 4/fb - First 4/fb of data	5.2 ± 22.4 MeV	63.2 ± 31.0 MeV	
8/fb – published 2/fb of data	50.4 ± 24.6 MeV	5.1 ± 28.3 MeV	



CDF Systematics



Source	Uncertainty	(MeV)
Lepton energy scale	3.0	
Lepton energy resolution	1.2	
Recoil energy scale	1.2	Stati
Recoil energy resolution	1.8	Seve
Lepton efficiency	0.4	
Lepton removal	1.2	Meth
Backgrounds	3.3	& rec
p_T^Z model	1.8	
p_T^W/p_T^Z model	1.3	
Parton distributions	3.9	
QED radiation	2.7	
W boson statistics	6.4	
Total	9.4	

	Uncertainty (MeV)
	3.0	
tion	1.2	
	1.2	Statistical and systematic uncertainty ~ same.
ion	1.8	Several of systematics are driven by available Z stats
	0.4	
	1.2	Methodology particularly in setting the energy, momentum
	3.3	& recoil scales is basically the same for CDF, DO, ATLAS, LHCb
	1.8	
	1.3	
	3.9	
	2.7	
	6.4	
	9.4	

What is mildly different is choice of parton distributions

and the modelling of the transverse momentum of the W and Z and W decay (polarization) : QCD.

This changes the central value and means the values from different experiments can't simply be compared (combined) but must be corrected to the same underlying PDF and QCD model.

CERN-LPCC-2022-06

Towards a combination of LHC and Tevatron W-boson mass measurements

The LHC-Tevatron *W*-boson mass combination working group

In this note methodological and modelling considerations towards a combination of the ATLAS, CDF and D0 measurements of the *W*-boson mass are discussed. As they were performed at different moments in time, each measurement employed different assumptions for the modelling of *W*-boson production and decay, and different fits of the parton distribution functions of the proton. Methods are presented to accurately evaluate the effect of PDFs and other modelling variations on existing measurements, allowing to extrapolate them to any PDF set and to evaluate the corresponding uncertainties. Based on this approach, the measurements can be corrected to a common modelling reference and to the same PDFs, and subsequently combined accounting for PDF correlations in a rigourous way.

ATLAS is 15 ± 19 MeV above the SM LHCb is 0 ± 32 MeV above the SM D0 is 22 ± 24 MeV above the SM CDF is 79 ± 11 MeV above the SM

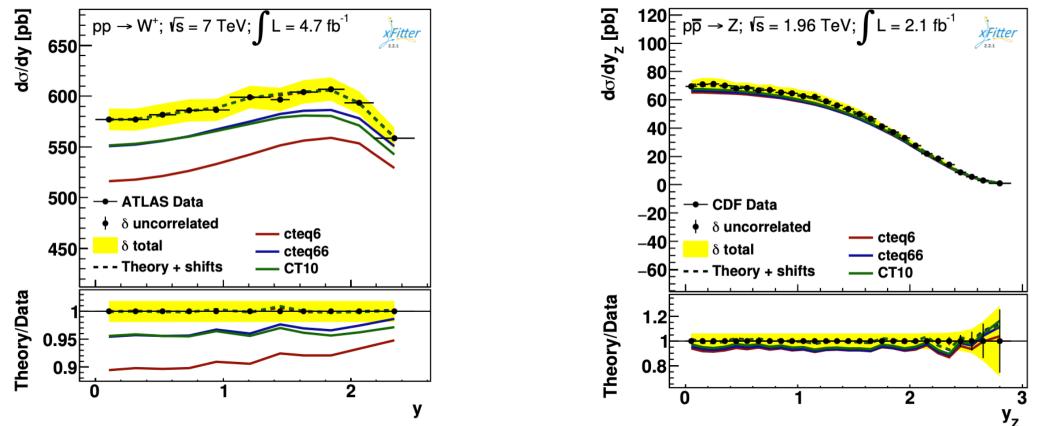
Preliminary findings were presented at ICHEP.



Two issues:

- each experiment has used a different parton distribution function
- no single parton distribution function describes all data used in the PDF fits

"The ATLAS 7 TeV precision W/Z data are not included in CT18, due to their tension with other data sets in the global fit"



Tevatron: only 5% of events involve c/s quarks and same number of W⁺ and W⁻ LHC: 25% of events involves c/s quarks and x1.4 W⁺ vs W⁻

CDF data much better described by current PDFs than ATLAS. Shifts when change PDFs up to 10 MeV

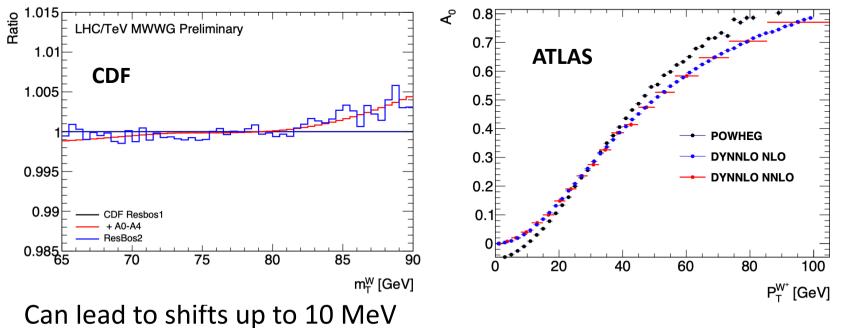
Higher Order QCD Effects on Lepton Angular Distribution

MANCHESTER 1824

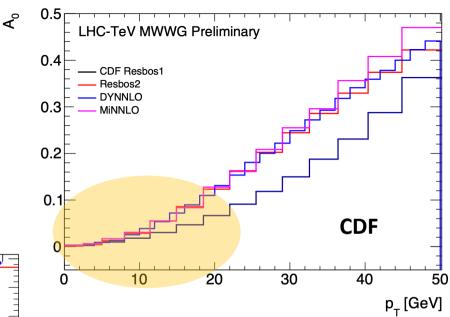
$$\frac{\mathrm{d}\sigma}{\mathrm{d}p_{\mathrm{T}}^{Z}\,\mathrm{d}y^{Z}\,\mathrm{d}m^{Z}\,\mathrm{d}\cos\theta\,\mathrm{d}\phi} = \frac{3}{16\pi} \frac{\mathrm{d}\sigma^{U+L}}{\mathrm{d}p_{\mathrm{T}}^{Z}\,\mathrm{d}y^{Z}\,\mathrm{d}m^{Z}} \\ \left\{ (1+\cos^{2}\theta) + \frac{1}{2}\,A_{0}(1-3\cos^{2}\theta) + A_{1}\,\sin2\theta\,\cos\phi \right. \\ \left. + \frac{1}{2}\,A_{2}\,\sin^{2}\theta\,\cos2\phi + A_{3}\,\sin\theta\,\cos\phi + A_{4}\,\cos\theta \right. \\ \left. + A_{5}\,\sin^{2}\theta\,\sin2\phi + A_{6}\,\sin2\theta\,\sin\phi + A_{7}\,\sin\theta\,\sin\phi \right\} \right\}.$$

CDF/D0 used NLO predictions for these A parameters.

There are now NNLO (and resumed) predictions for these which are being studied.



At leading order all but A4(=2) is zero. A_{5,6,7} are zero until NNLO and are negligible afterwards.





High Energy Physics – Phenomenology

[Submitted on 5 May 2022]

ResBos2 and the CDF W Mass Measurement

Joshua Isaacson, Yao Fu, C.-P. Yuan

The recent CDF W mass measurement of 80,433 \pm 9 MeV is the most precise direct measurement. However, this result deviates from the Standard Model predicted mass of 80,359.1 \pm 5.2 MeV by 7σ . The CDF experiment used an older version of the ResBos code that was only accurate at NNLL+NLO, while the state-of-the-art ResBos2 code is able to make predictions at N³ LL+NNLO accuracy. We determine that the data-driven techniques used by CDF capture most of the higher order corrections, and using higher order corrections would result in a decrease in the value reported by CDF by at most 10 MeV.

In conclusion, two of the major criticisms leveled against the theory calculations involved in the ResBos program cannot explain the deviation from the SM that is reported by CDF. We found that the data-driven techniques used by the CDF experiment help to reduce the effects of higher order corrections. The estimated shift due to including these corrections is at most 10 MeV, and may reduce the disagreement from 7o to 6o. The PDF uncertainty is found to be consistent with the numbers quoted by CDF. We have addressed the most important questions related to the theory calculations.....

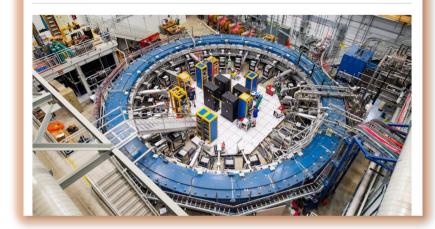
Muon g-2 and Mw



Muons: 'Strong' evidence found for a new force of nature

By Pallab Ghosh Science corresponden





The muon g-2 and W-mass anomalies explained and the electroweak vacuum stabilised by extending the minimal Type-II seesaw http://arxiv.org/abs/2206.11771v1

Leptoquark-vectorlike quark model for mm (CDF), (g-2)_µ, R_K anomalies and neutrino mass: http://arxiv.org/abs/2205.03917v2

A leptoquark and vector-like quark extended model for the simultaneous explanation of the W boson mass and muon g-2 anomalies: http://arxiv.org/abs/2205.02088v1

CDF W boson mass and muon g-2 in type-X two-Higgs-doublet model with a Higgs-phobic light pseudoscalar: http://arxiv.org/abs/2205.01701v2

Compatibility of muon g-2, W mass anomaly in type-X 2HDM : <u>http://arxiv.org/abs/2205.01437v1</u>

Muon and electron g-2 anomalies in a flavor conserving 2HDM with an oblique view on the CDF M_W value : http://arxiv.org/abs/2205.01115v1

 $\text{The CDF W-mass, muon g-2, and dark matter in a U(1)_{L_\mu-L_\tau} \text{Model with vector-like leptons: } \frac{http://arxiv.org/abs/2204.13027v1}{http://arxiv.org/abs/2204.13027v1} \text{Model with vector-like leptons: }$

Implications of CDF W-mass and (g-2)_µ on U(1)_{L_µ-L_T} model : <u>http://arxiv.org/abs/2204.09585v1</u>

Combined explanation of W-mass, muon g-2, R_K* and R_D* anomalies in a singlet-triplet scalar leptoquark model: http://arxiv.org/abs/2204.09031v1

The 2HD+a model for a combined explanation of the possible excesses in the CDF Mw measurement and (g-2) with Dark Matter : http://arxiv.org/abs/2204.08406v1

A model explaining the new CDF II W boson mass linking to muon g-2 and dark matter: http://arxiv.org/abs/2204.07411v1

W boson mass and muon (g-2) in a lepton portal dark matter model: http://arxiv.org/abs/2204.07022v2

A joint explanation of W-mass and muon g-2 in 2HDM: http://arxiv.org/abs/2204.06505v2

Correlating W-Boson Mass Shift with Muon {g-2} in the 2HDM: http://arxiv.org/abs/2204.05303v2

NMSSM neutralino dark matter for W-boson mass and muon g-2 and the promising prospect of direct detection : http://arxiv.org/abs/2204.04356v2

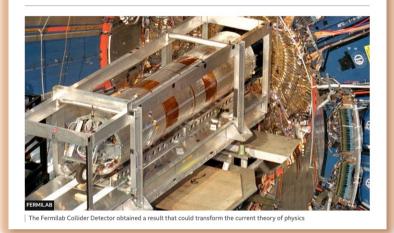
Explaining The Muon g-2 Anomaly and New CDF II W-Boson Mass in the Framework of (Extra)Ordinary Gauge Mediation http://arxiv.org/abs/2204.04286v3

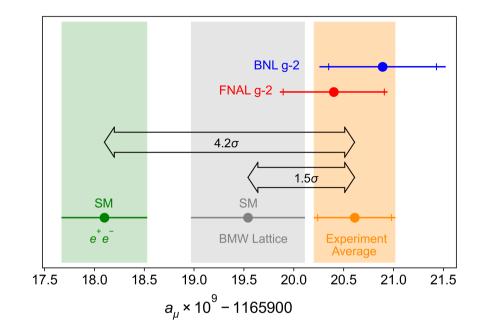
Low energy SUSY confronted with new measurements of W-boson mass and muon g-2: http://arxiv.org/abs/2204.04202v3

The W boson Mass and Muon g-2: Hadronic Uncertainties or New Physics? : http://arxiv.org/abs/2204.03996v2

Shock result in particle experiment could spark physics revolution







Mark Lancaster: W Mass



 $\Delta \alpha_{had}$ that goes in EWK fits is also what determines the SM prediction for (g-2)

$$a_l^{\text{had, LOVP}} = \frac{\alpha^2}{3\pi^2} \int_{s_{th}}^{\infty} \frac{\mathrm{d}s}{s} R(s) K_l(s) \,, \qquad \qquad \textbf{g-2}$$

$$\Delta \alpha_{\rm had}^{(5)}(M_Z^2) = -\frac{\alpha M_Z^2}{3\pi} \operatorname{P} \int_{s_{th}}^{\infty} \mathrm{d}s \frac{R(s)}{s(s-M_Z^2)} \,, \qquad \qquad \text{EWK FIT}$$

R(s) are measured $e^+e^- \rightarrow$ hadrons cross sections

If you change the SM prediction for (g-2) e.g. by using a lattice calculation and not measured e⁺e⁻ cross sections then you also change your EWK fit predictions e.g. of Mw.



SM predicted Mw = 80355 ± 6 MeV and Higgs predicted = 95 ± 20 GeV CDF measured Mw = 80433 ± 9 MeV and Higgs measured = 125.2 ± 0.1 GeV

You can remove the SM tension in (g-2) by changing the SM prediction ie blindly increasing the measured e⁺e⁻ cross sections (ie moving in direction of BMW lattice calculation).

If you do this then the SM predicted Mw reduces to about 80340 and so does the predicted Higgs to about 60 GeV. So you remove BSM from (g-2) and need BSM to explain why the measured W (even without CDF) and Higgs masses are so much higher than the SM predictions....

Conversely to predict a high Mw in the SM you need to reduce the measured e⁺e⁻ cross sections but then your SM (g-2) value goes lower and the (g-2) tension increases

Tricky to accommodate both a high measured g-2 and a high measured Mw **without SM inconsistencies** which can potentially be removed instead with BSM.



We need a precise measurement from CMS....

Experimental techniques used by CDF, D0, ATLAS, LHCb are all very similar

However, each experiment has a different model for PDFs, $p_T(W)$, angular coefficients and QED radiation.

These are all "theoretical" and so one can "easily" reweight a given measurement to another model.

This needs to be done before any meaningful comparison with SM can be made i.e. a common model has to be agreed by all experiments.....and then the published values shifted. This is not something the PDG can do.

There are effects at the 10 MeV level in these models.

However, assuming a combination with uncertainty 10 MeV then a combined Mw > 80400 MeV looks tricky for the SM and then we have Muon g-2



MINUTES OF THE COLLIDER DETECTOR MEETING

May 25, 1984

1. CDF has run out of money.





MINUTES OF THE COLLIDER DETECTOR MEETING

November 9, 1984

 There will be a workshop to discuss upgrades to the CDF detector in early January.