

Quantum Sensors for Particle Physics

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Perspectives on Quantum Sensing and Computation for Particle Physics

Overview

Overview — Quantum Sensor Regimes

Through the dark matter lens

More generally

Quantum Calorimetry

Sub-meV Quantum Sensors: the truly quantum regime

Caveats:

DM motivates much of the quantum sensing universe

My own expertise is primarily in quantum calorimetry and dark matter

Much of the remainder drawn from DOE reports:

[Basic Research Needs for Dark Matter Small Projects New Initiatives \(2019\)](#)

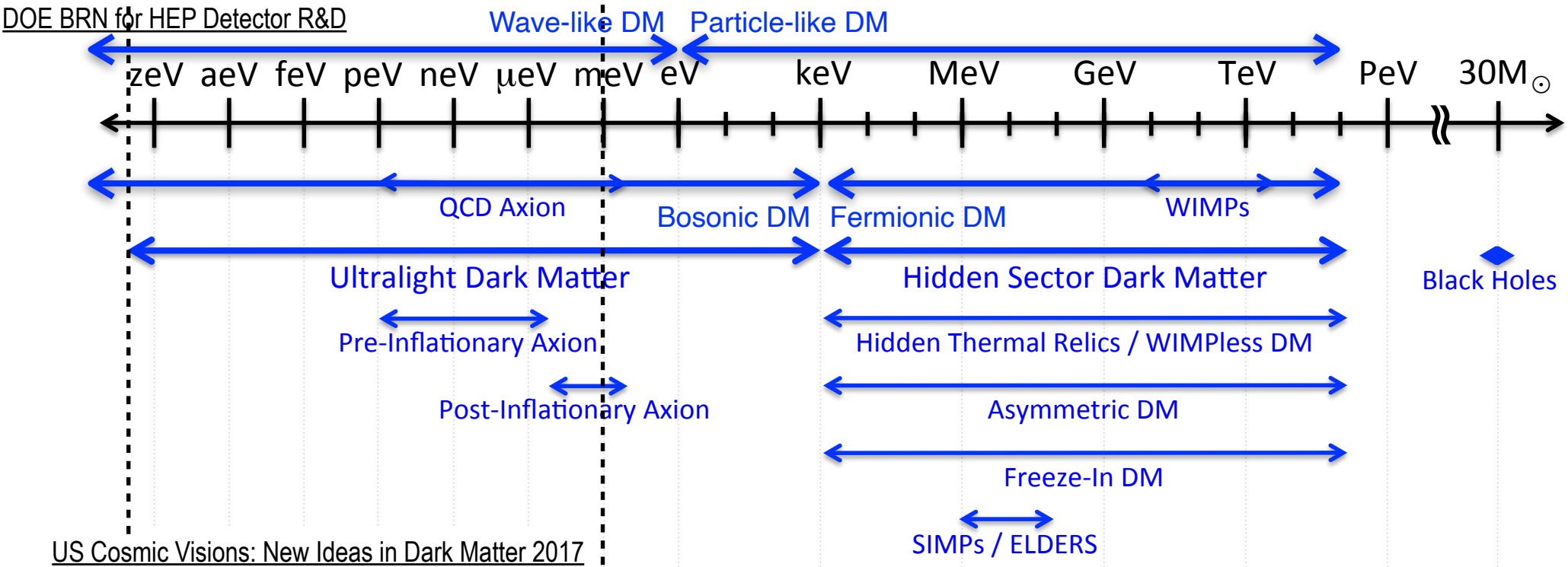
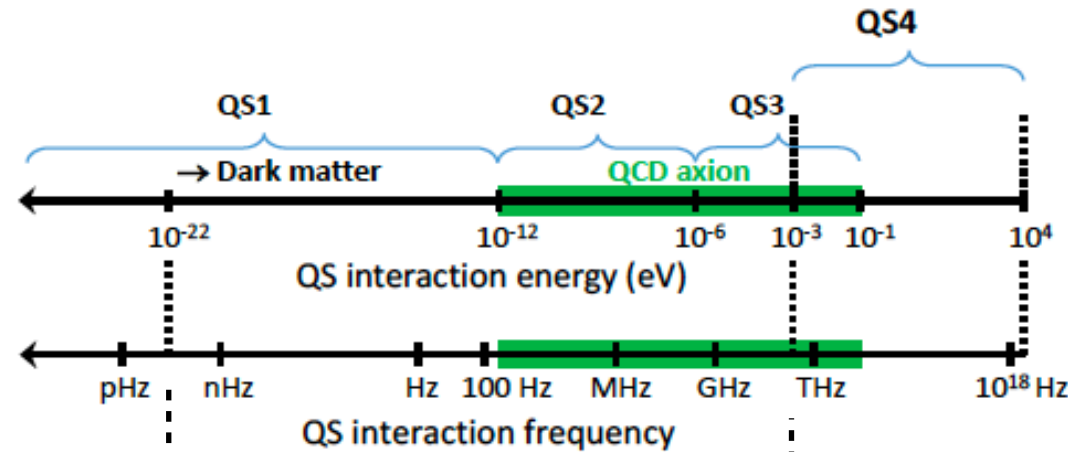
[Basic Research Needs for High Energy Physics Detector Research & Development \(2020\)](#)

Dark Matter → Quantum Sensors

QS for particle physics driven in large part by 2017 US Cosmic Visions *New Ideas in DM* white paper

Seeking to cover $E_{\text{dep}} \lesssim \text{keV}$ inaccessible to classic particle DM searches

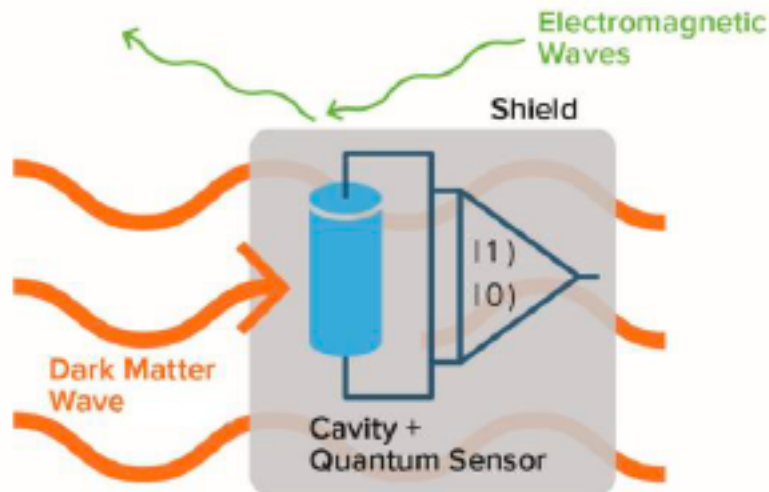
QS regime includes both wave-like (~classical field) and particle-like DM



US Cosmic Visions: New Ideas in Dark Matter 2017

Dark Matter → Quantum Sensors

Wave-like Dark Matter



Occupation number $\gg 1$: cannot identify individual DM particles, sense “Collective force from macroscopic numbers of particles”

New light scalar (axion-like particle (ALP) or vector (dark photon (DP))

Quantum zero-point fluctuations are the primary obstacle: a quantum problem

Fundamentally quantum approaches needed to circumvent

Particle-like Dark Matter



Occupation number $\ll 1$: individual particles can be seen*

Scattering of fermionic DM, absorption of heavier ALPs, DPs, and

“Quantum” effects:

Direct creation of single quantum (not via identifiable particle recoil)

Use of quantum sensors to see the very small energy depositions

$eV \rightarrow meV \rightarrow \mu eV?$

*Occupation number > 1 for bosonic DM does not imply particle-like absorption is not possible, but it makes classical field detection possible.

Beyond DM: Quantum Sensor Regimes

ULDM = ultralight dark matter

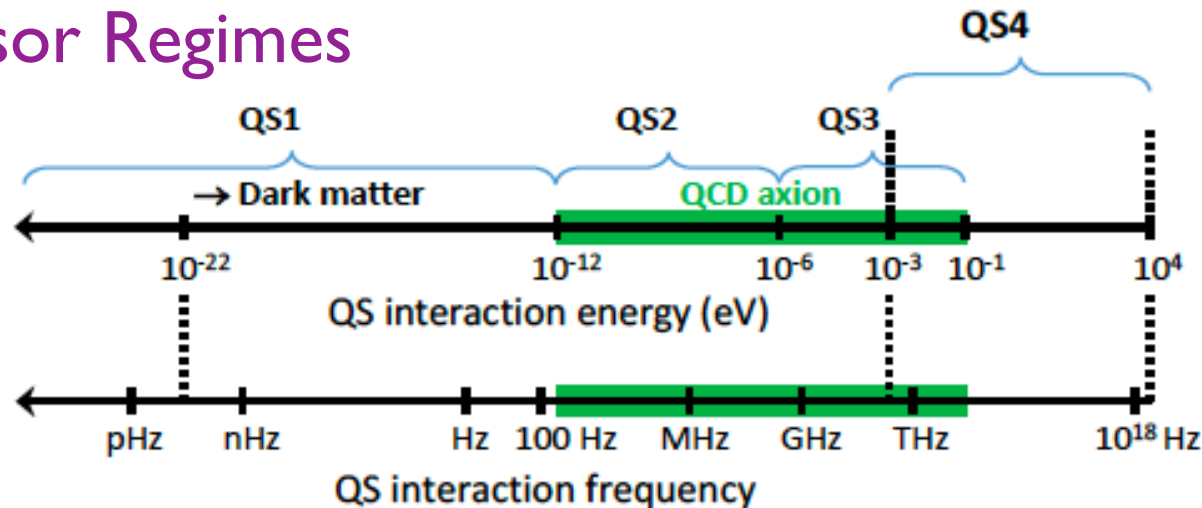
EDM = electric dipole moment

GW = gravitational waves

DE = dark energy

FSV = fund. symmetry violations

QND = quantum non-demolition



Regime/ Mass	Science Target	Quantum Sensor Technology	Quantum Protocol
QS1 < 10 ⁻¹² eV	ULDM, EDMs, GW, DE, FSV	Atomic and molecular spectroscopy, atom interferometers, mechanical sensors, clocks, atomic magnetometers, nuclear, electronic, and other spins, quantum defects in solids	Superposition, entanglement, squeezing, coherence
QS2 10 ⁻¹² to 10 ⁻⁶ eV	ULDM incl. axions, FSV, new forces/ particles	Nuclear, electronic, and other spins, electromagnetic quantum sensors, optical cavities, quantum defects in solids	Superposition, entanglement, backaction evasion, squeezing, coherence
QS3 10 ⁻⁶ to 10 ⁻¹ eV		Superconducting and other qubits, nuclear, electronic, and other spins, Rydberg atoms, quantum defects in solids	Parametric amplifiers, superposition, entanglement, squeezing, coherence, QND photon counting
QS4 10 ⁻³ to 10 ⁴ eV	ULDM scattering/ absorption, FSV, new forces/particles	Single-photon counters, low-threshold phonon and charge detectors, quantum defects in solids	Non-QND photon counting, high spatial resolution measurements of particle tracks

DOE BRN for HEP Detector R&D

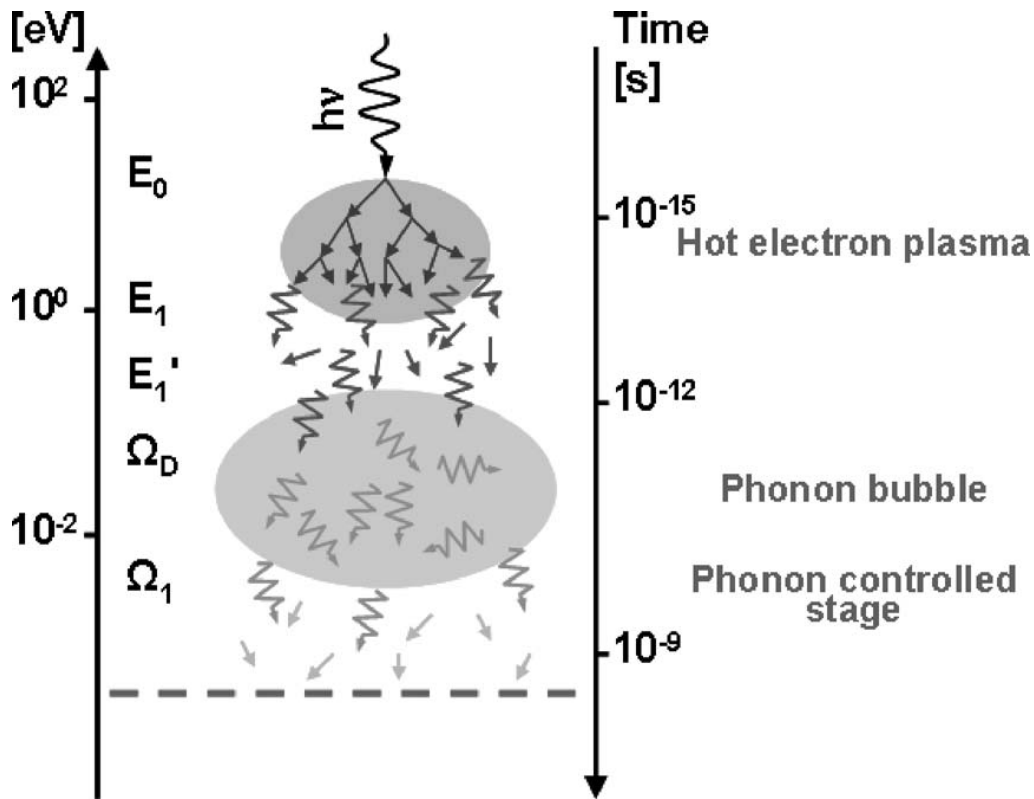
Quantum Calorimeters

“Logical conclusion” of existing solid-state modalities for sensing phonons, ionization, and scintillation

Semi-classical picture

Enough quanta produced to treat as continuous variable

Quantization not apparent in E_{observed} .



Kozorezov et al PRB: 75, 094513 (2007)

Quantum picture

Production of single charges or scintillation photons (eV quantization)

Arguable whether sensing single charges or photons is really “quantum” (e.g., PMTs)

Innovations that lead to “quantum” regime:

SNSPDs that detect single photons with near 100% efficiency

Reaching the regime of detection of
1-few athermal phonons (meV) or
1-few thermal phonons (μeV)

Energy depositions so small that “quantum”
techniques must be used to detect

Quantum Calorimeters: Kinematics and Scattering Modes

Energy-momentum transfer of galactic DM scattering kinematically limited by max v_{DM}

Nuclear recoils follow lines

Truncated where scattering with single nucleus invalid, $q < (\text{lattice constant})^{-1}$

Direct single phonon creation via coherent multi-site interaction valid for lower q

Also: coherent photon scattering w/nucleus

Electron recoils

FEG = free electron gas, valid for $E >$ binding energy

superconductor gap, semiconductor/insulator bandgap

Band gap (e.g. Si) shows energies accessible by e-h pair creation in semiconductors/insulators

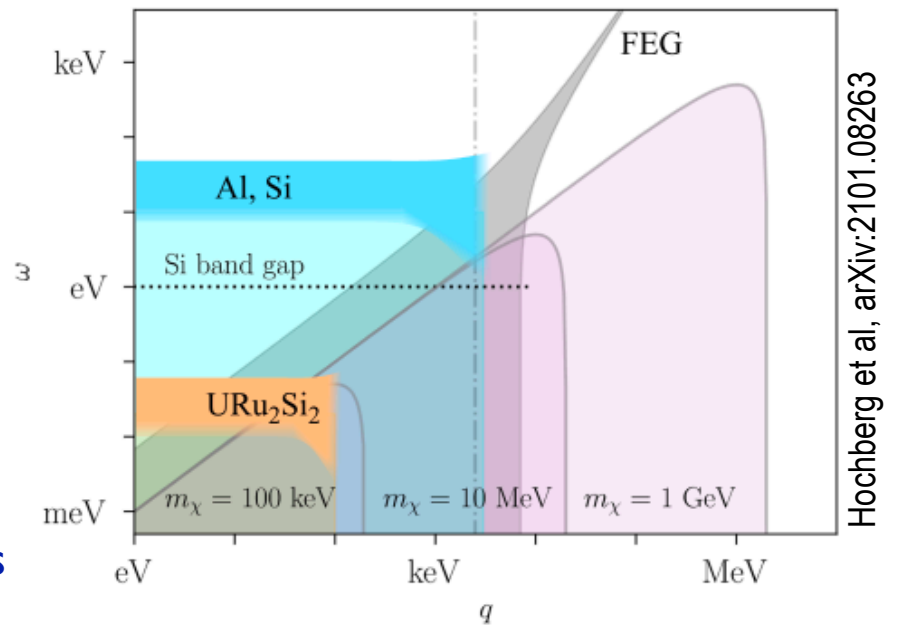
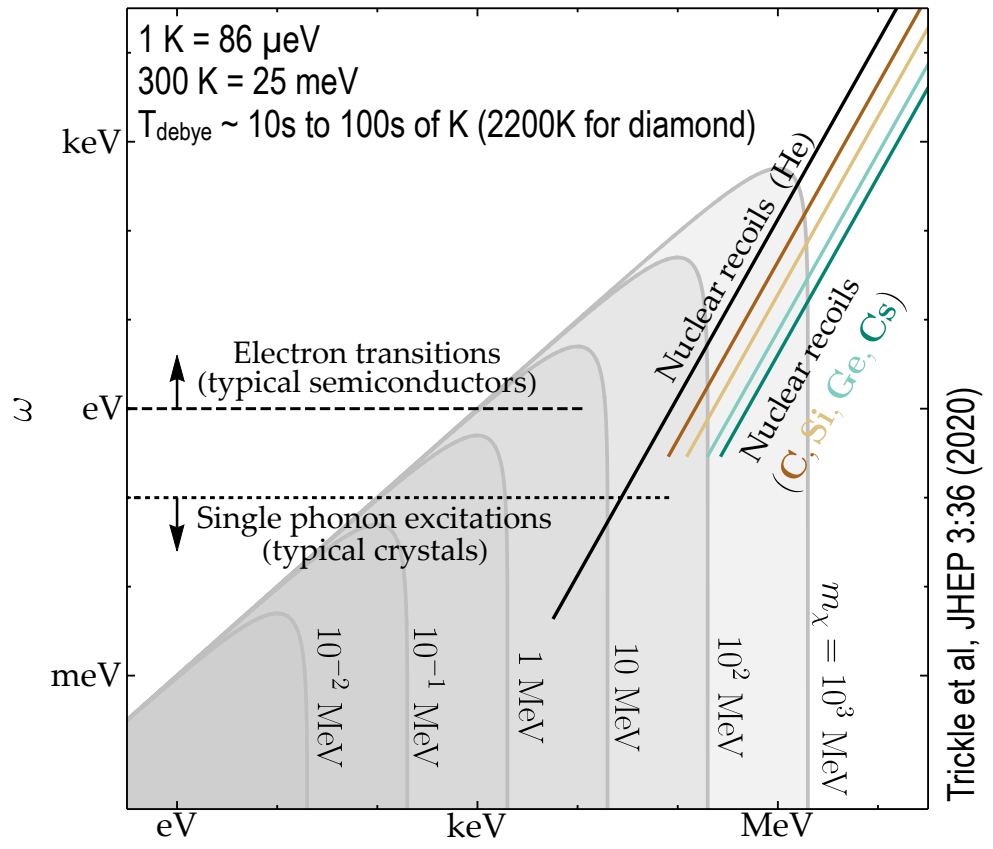
Suppressed for high q because wavefunction too large

Blue shaded region: plasmons

Mostly outside allowed region except for heavy f materials (e.g. URu₂Si₂)

Light/Dark Photon absorption

Coupling to unit cell dipole moment in polar materials



Quantum Calorimeters: Interaction vs. Creation vs. Sensing

Interaction modalities

DM

Nucleus scattering

single-site

multi-site

Electron scattering

single-site

multi-site

(dielectric absorption function)

Neutrinos

Nucleus scattering

single-site

multi-site

Photons

Electron scattering

single-site

multi-site

(dielectric absorption function)

Cavity or antenna mode creation

Creation modalities

Single e-h pair creation

Single scintillation photon
creation

Single optical phonon creation

Single acoustic phonon
creation

Single Cooper pair breaking

Single plasmon creation

Exotic quanta

Dirac materials, magnons, ...

Single conversion photon
creation

Sensing Modalities: mostly cryogenic

Transition-Edge Sensors
(TESs)

Kinetic Inductance Detectors
(KIDs)

Skipper CCDs

Superconducting Nanowire
Single Photon Detectors
(SNSPDs)

Parametric amplifiers
for squeezing

Quantum non-demolition
(QND) photon detectors

Quantum Calorimeters Today

Calorimeters based on Transition-Edge Sensors

TESs provide very sharp resistance vs. T curve

Electothermal feedback can be used to stabilize
Electrical signal measures received energy

Coupled to calorimetric substrate via athermal phonon collectors

Now approaching the quantum regime

Large bias voltage provides amplification via Neganov-Trofimov-Luke phonons: enables single e-h pair detection

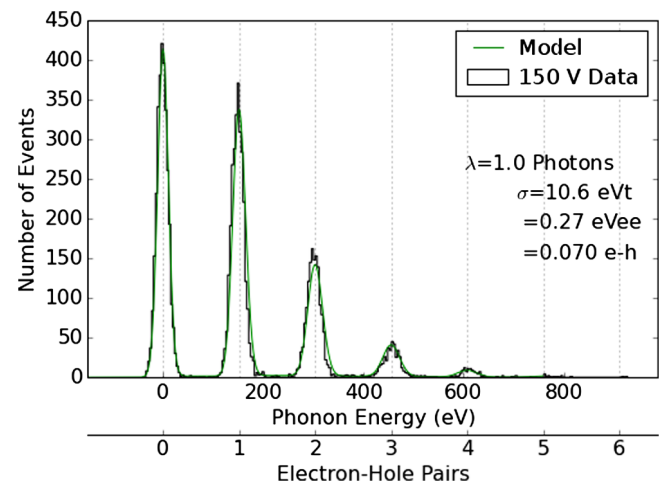
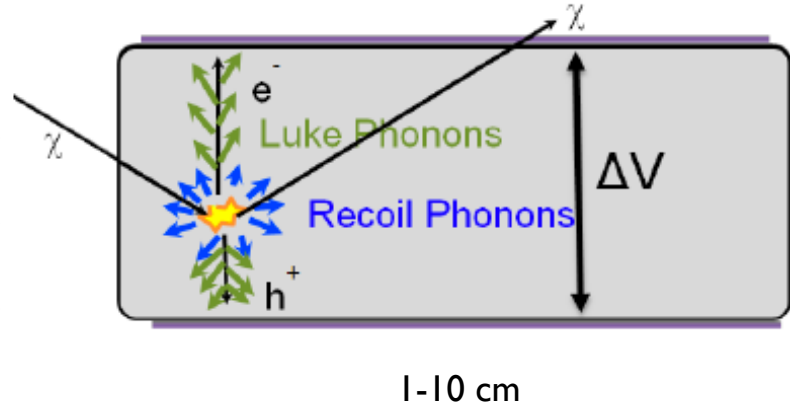
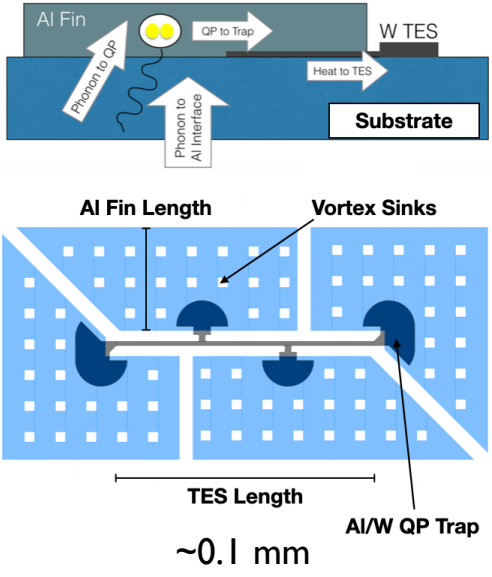
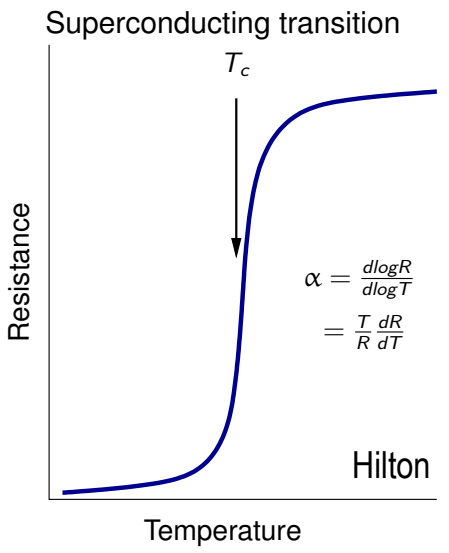
But subject to significant leakage currents: single eh pair detection unreliable

Native resolution ~ 3 eV achieved

Reduced T_c , improved design expected to yield sub-eV resolutions, eventually approaching 10 meV

Single eh pair detection without leakage

Single phonon detection



PRL 121: 0513401 (2018)

Quantum Calorimeters Today

Kinetic Inductance Detectors (KIDs)

Superconductors have an AC inductance due to inertia of Cooper pairs

KID = superconducting film incorporated into LC resonator to sense change in L

Energy resolution:
sub-eV \rightarrow meV
thresholds w/o HV

Direct sensitivity to
pair-breaking phonons

Large resonators obviate
phonon collectors

Gapped density of states

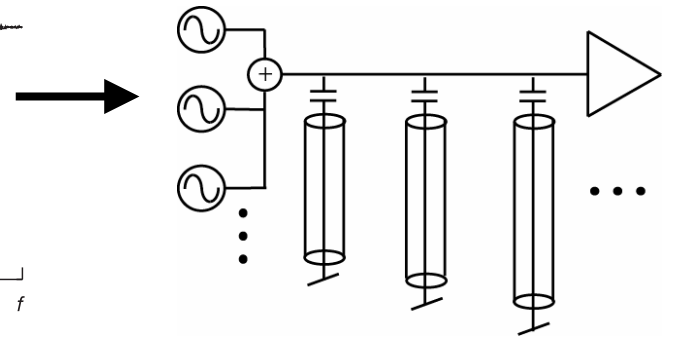
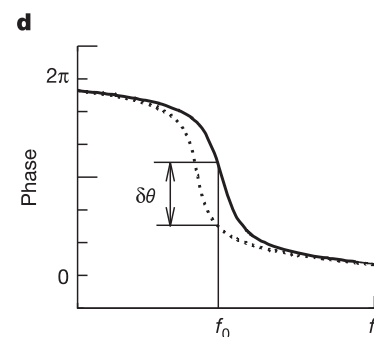
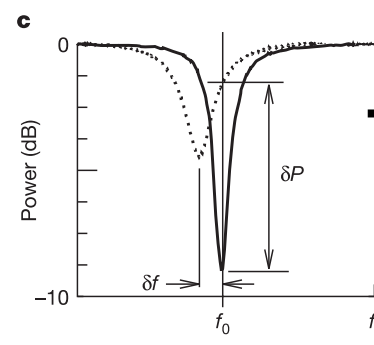
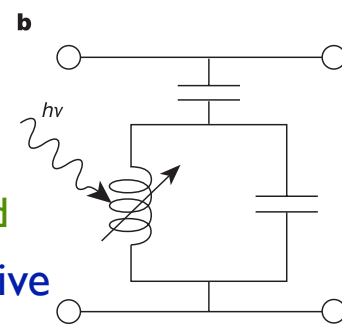
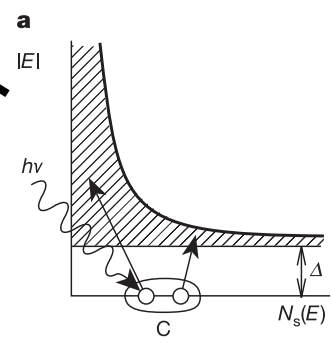
Thermal quasiparticles
exponentially suppressed

Fundamentally non-dissipative

Amenable to QIS techniques
(e.g. squeezing, QND)

Noise is limited by

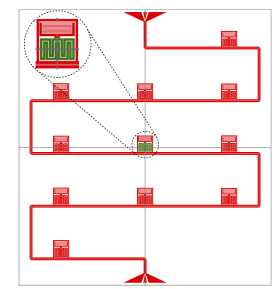
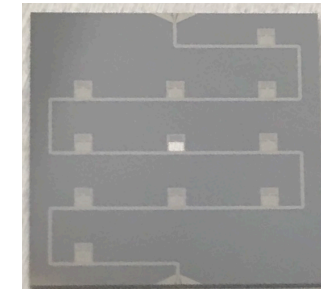
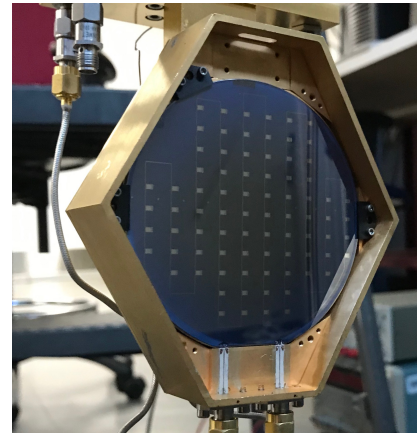
quasiparticle population
fluctuations
amplifier noise



Multiplexing:

- KIDs are $Q > 10^5$ resonators
- \rightarrow Readout many with one cryo line/amplifier; most electronics at 300K
- \rightarrow Highly position-resolved phonon detection

75 mm
x 1 mm



22 mm
x 1 mm

Quantum Calorimeters Today

Skipper CCDs

CCD with two readout innovations

- High-frequency differencing to reduce impact of $1/f$ amplifier noise
- Non-destructive multiple read cycles to reduce electronics noise by \sqrt{N}

Provides similar single-eh pair sensitivity
 Currently being applied for DM searches, low-light-level astronomy

Superconducting Nanowire Single Photon Detectors (SNSPDs)

- Threshold detector for single photons
- Very narrow (~ 100 nm) superconducting meander biased close to transition
- Absorption of photon drives normal

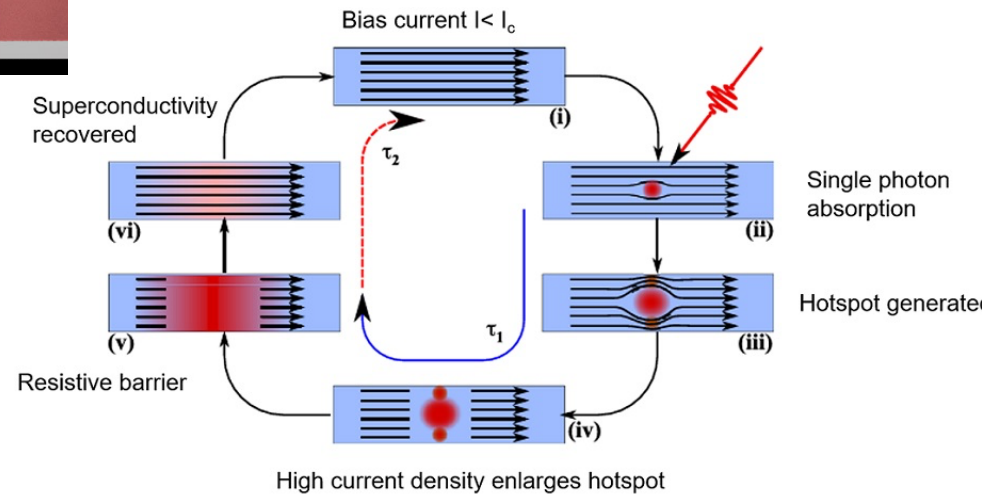
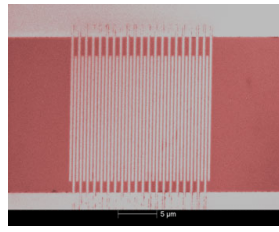
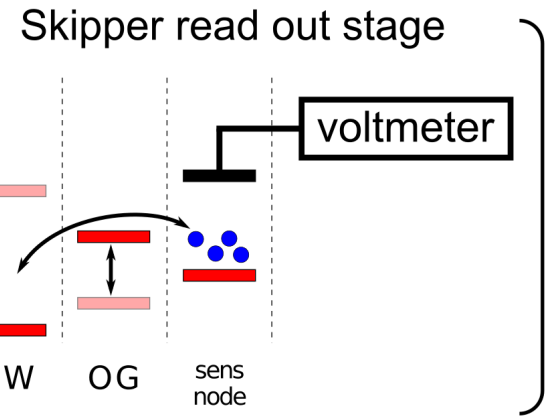
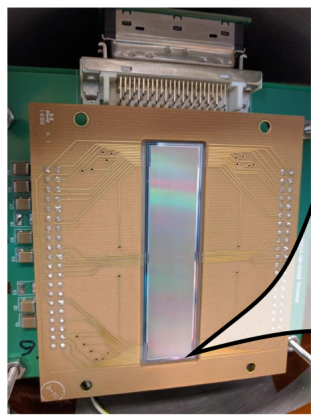
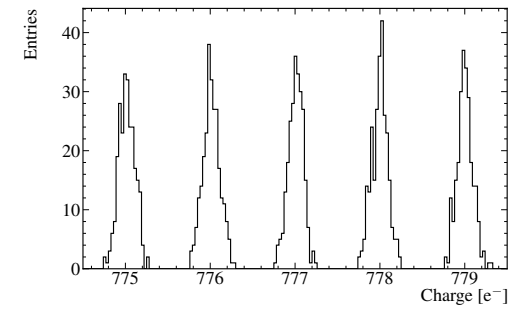
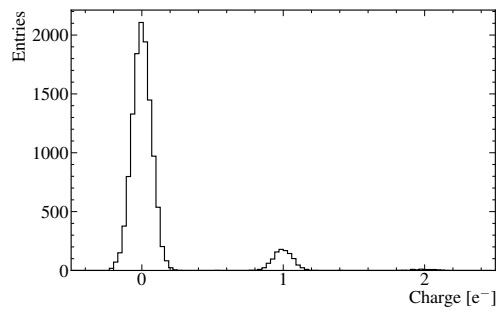
ps timing resolution

Provides high-efficiency, high-fidelity photon counter for QIS applications

WSi demonstrated with 100 meV threshold

Very low dark count rate demonstrated, applicable for DM searches

But very small volume



<https://sensei-skipper.github.io/#SkipperCCD>

<https://singlequantum.com/technology/snspd/>

Quantum Calorimeters in the Future

Single optical phonon creation/detection

Polar materials:

- > 1 atom/unit cell → optical phonons (10s of meV)
- polar → unit cell EDM can couple to dark photons

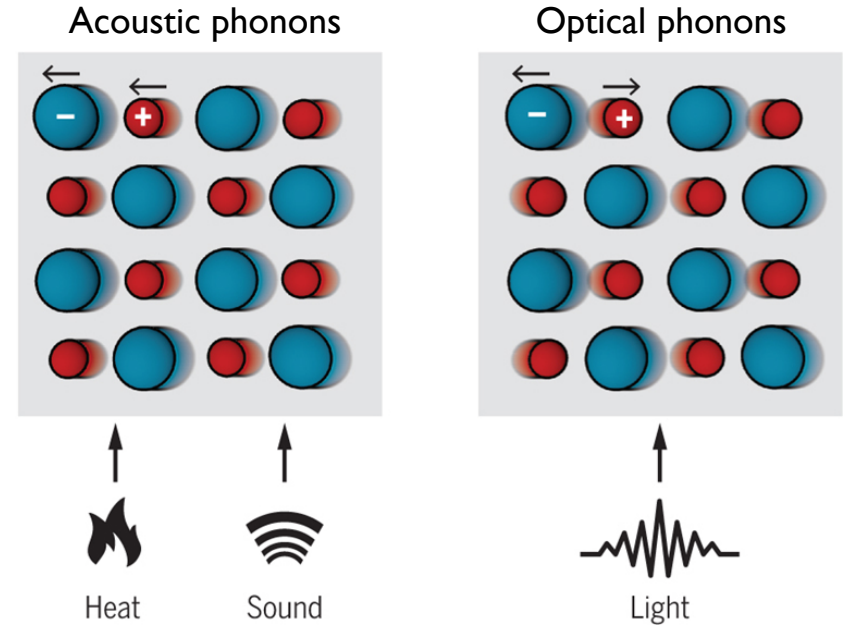
Optical phonon creation provides access to electron scattering down to few keV m_{DM}

Single acoustic phonon creation/detection

DM interacting with many nuclei coherently

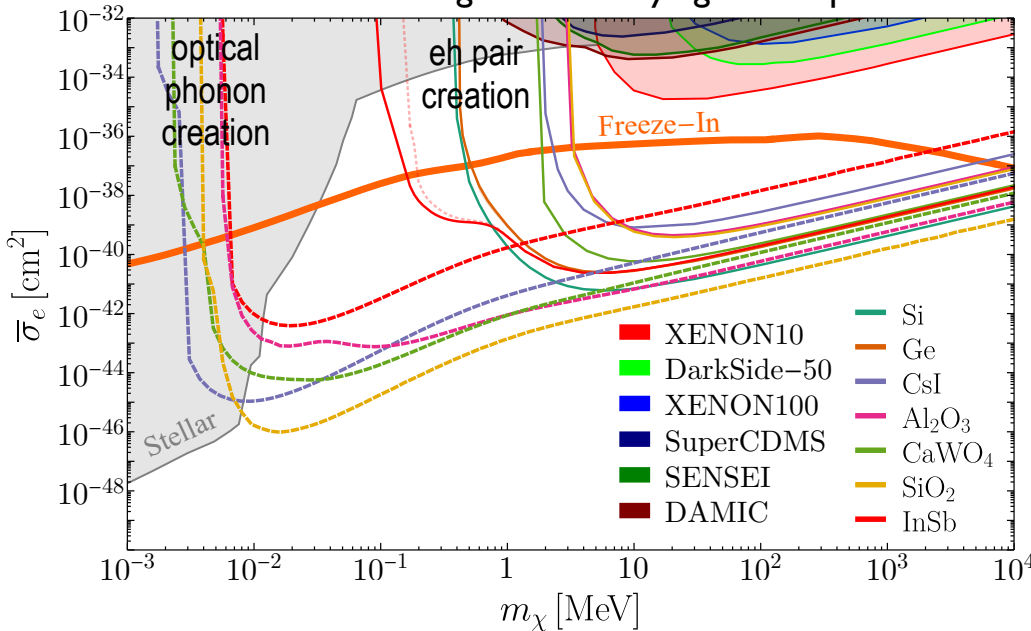
Requires meV-100 meV thresholds

Eventual reach down to few keV m_{DM}

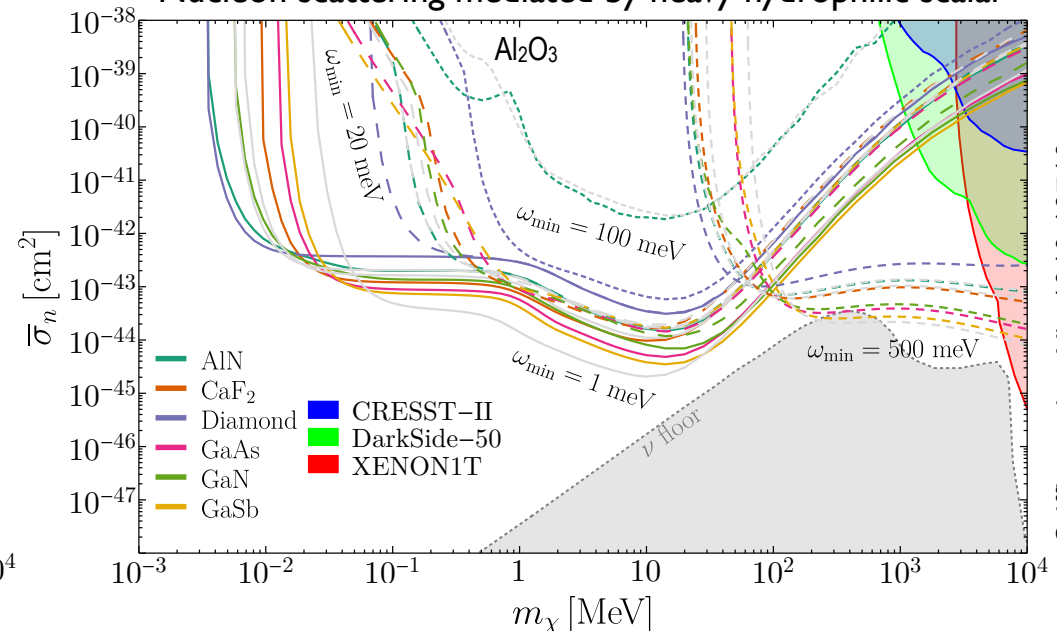


<https://images.app.goo.gl/vvBWnerR6noCnseCXA>

Electron scattering mediated by light dark photon

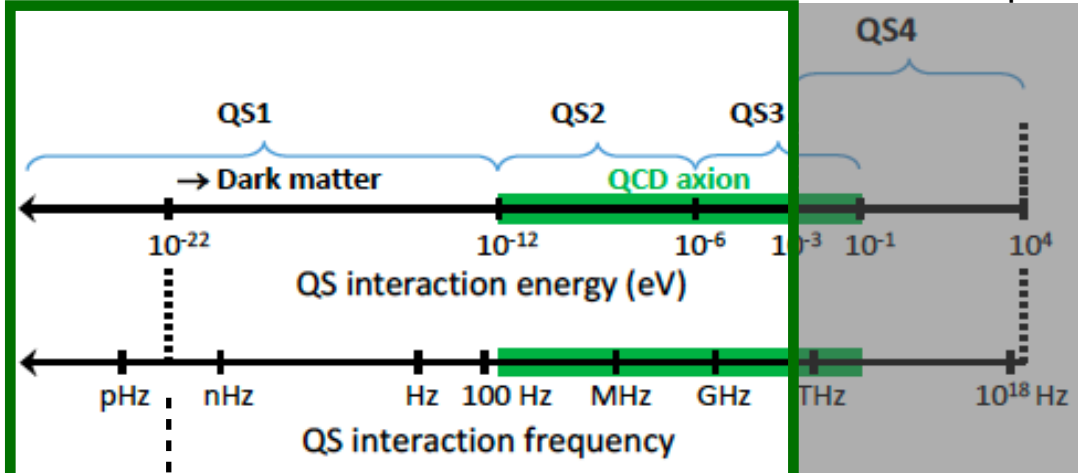


Nucleon scattering mediated by heavy hydrophilic scalar

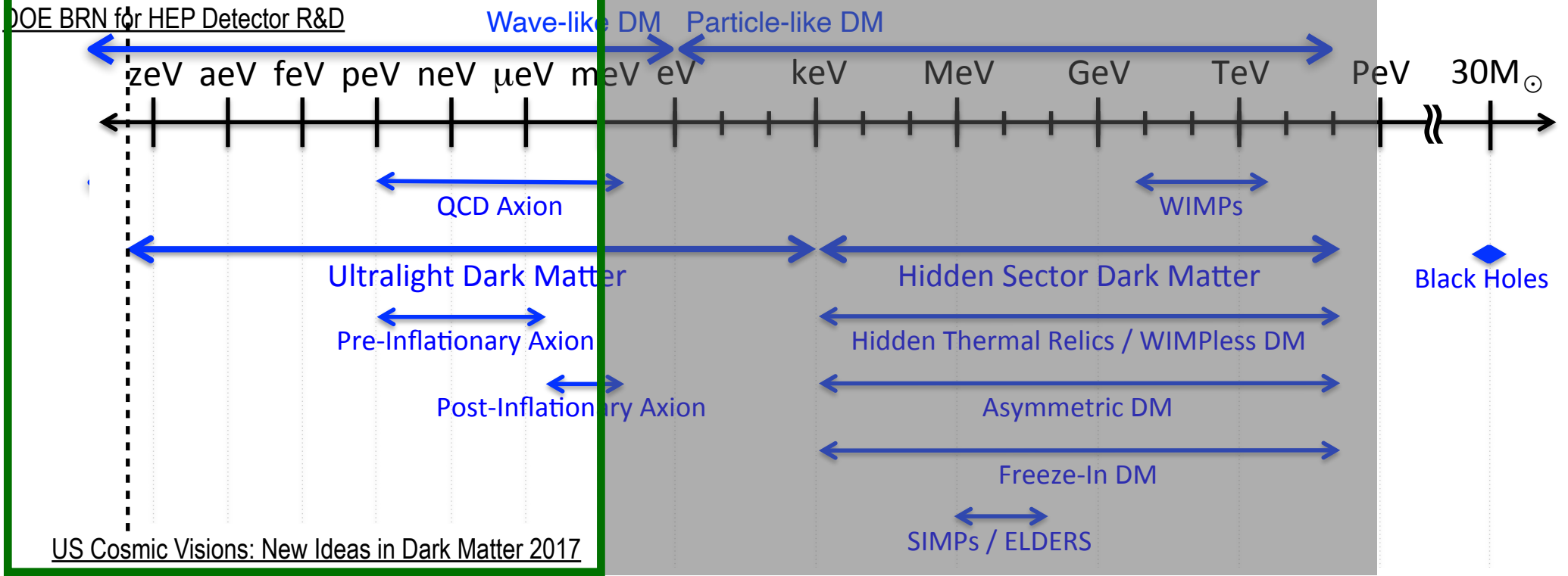


Sub-meV Quantum Sensors (→ Wave-like DM)

Natural limit for quantum calorimeters set by smallest available gapped excitation: superconductor pair-breaking energy ~ 0.1 meV.



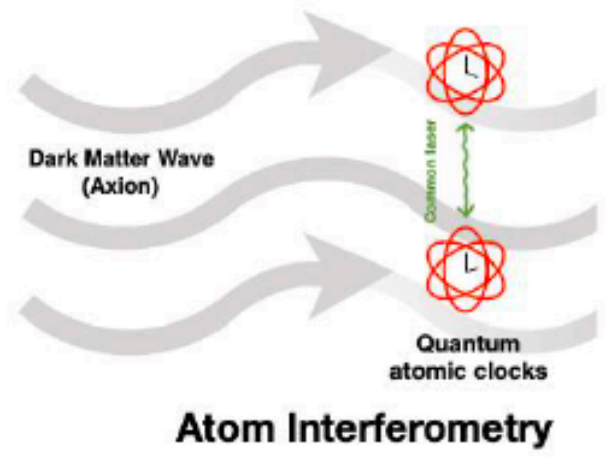
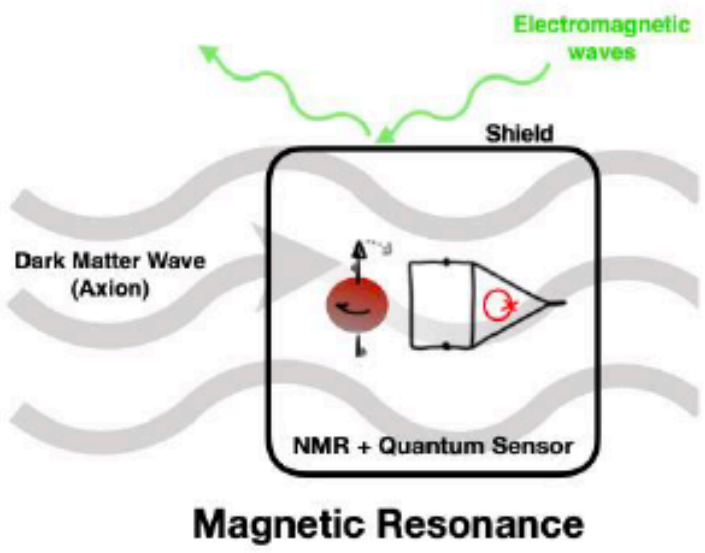
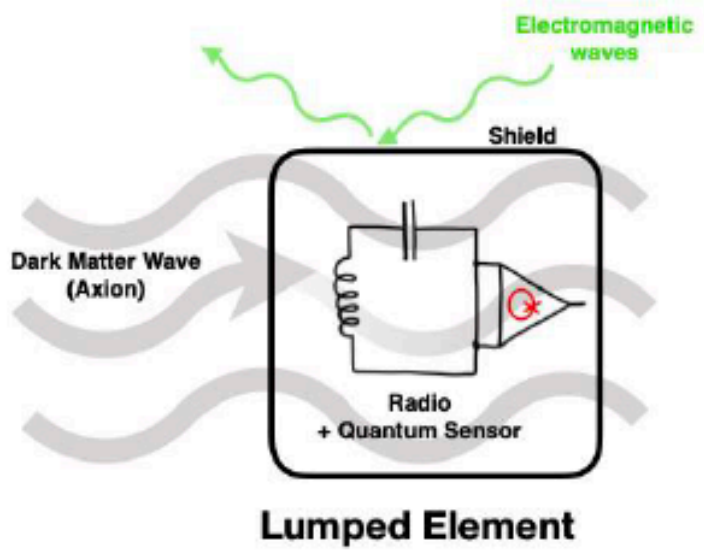
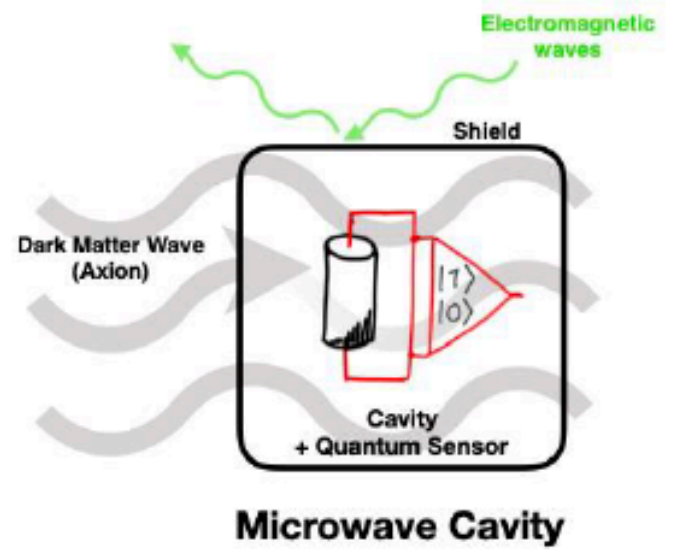
Smaller gap materials could push quantum calorimetry 1-2 more decades to thermal limit set by dilution refrigerators, 5-10 mK ~ 1 μ eV. Perhaps metamaterials can have spectrum with smaller gap.



US Cosmic Visions: New Ideas in Dark Matter 2017

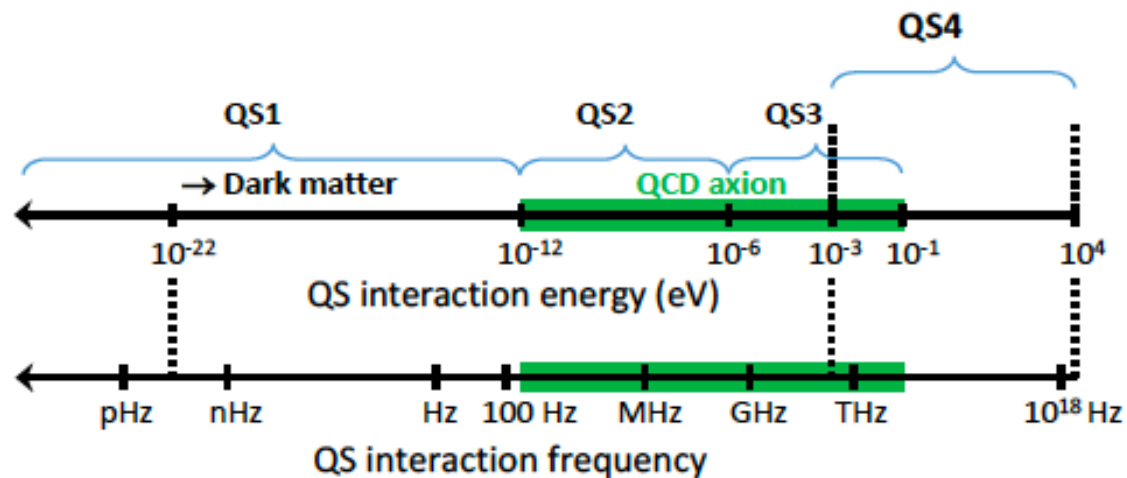
Sub-meV Quantum Sensors (→ Wave-like DM)

Wide range of techniques depending on coupling and mass: EM, QCD, spin, other



Same approaches amenable to fifth-force/light-shining-through-walls experiments

Sub-meV Quantum Sensors: ULDM via EM Coupling



DOE BRN for HEP Detector R&D

> 1 μeV (240 Hz) mass range

< 240 Hz impractical for photon conversion given EM-saturated environment

Classic quantum-limited photon detection problem

practical limit of commercial dilution refrigerators \approx 5-10 mK

→ Smallest thermal energy per mode accessible \sim 1 $\mu\text{eV} \sim$ 240 MHz

→ Two regimes:

1 μeV to 100 meV (240 MHz to 24 THz, QS3)

Detection modality can be placed in quantum ground state

e.g., resonant cavity with no thermally generated photons, only 1/2 photon vacuum fluctuation

1 peV to 1 μeV (240 Hz to 240 MHz, QS2)

Thermal state for detection modality unavoidable

Upconvert to GHz for readout to make use of established techniques for sub-SQL amplification

Sub-meV Quantum Sensors: ULDM via EM Coupling

1 μeV to 100 meV (240 MHz to 24 THz, QS3) example: HAYSTAC axion search ($g_{a\gamma\gamma}$)

Classical amplifier approach:

Standard quantum limit (SQL): a phase-insensitive amplifier adds 1/2 photon noise in each quadrature due to internal modes that enable gain in both quadratures

Quantum sensing approach:

Phase-sensitive amplification (Caves, 1982) circumvents this

Attenuate one quadrature while amplifying the other, preserving uncertainty area in quadrature phase space.

No fundamental limit to noise in amplified quadrature; practical limit set by parasitic losses, 0.01-0.1xSQL achievable

Phase-sensitive amplification applied to vacuum can attenuate vacuum fluctuations in one quadrature

SQ = vacuum squeezer, AMP = phase-sensitive amplifier

a) vacuum state before SQ

(b)-(d) w/SQ and AMP:

b) vacuum state after SQ: noise squeezed in X quadrature, amplified in Y

c) state after addition of signal from cavity

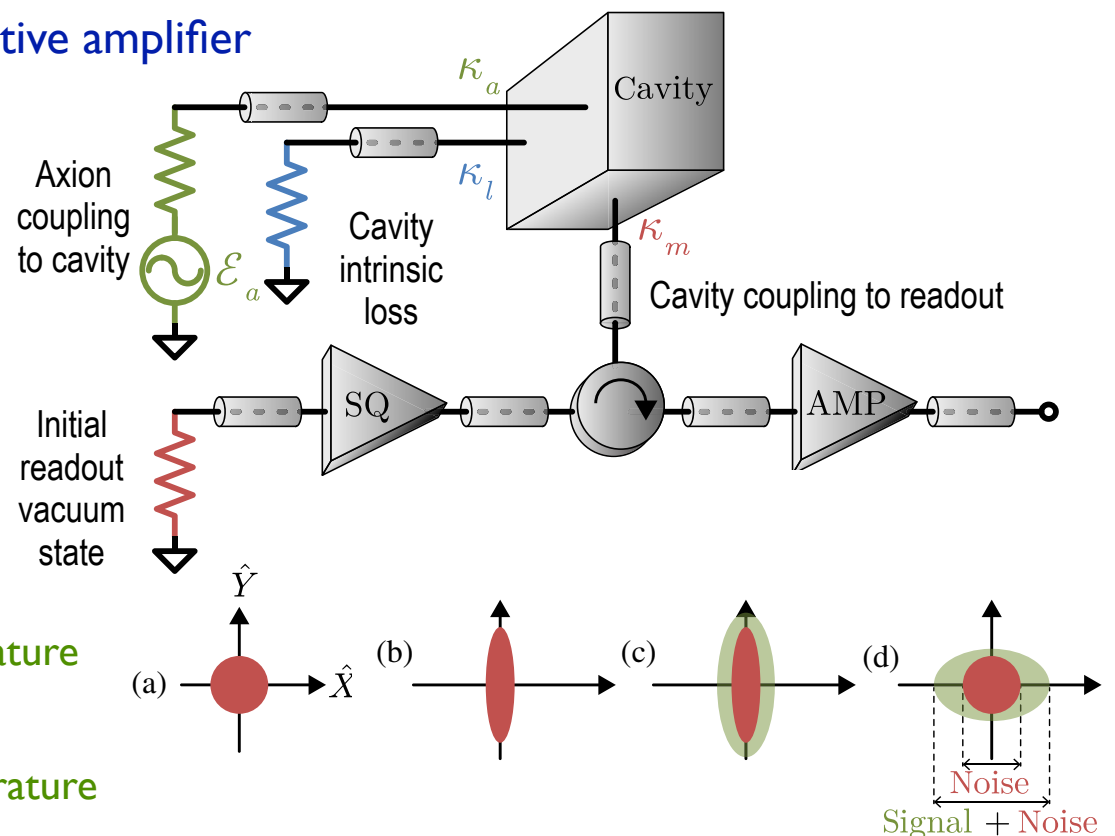
d) state after phase-sensitive amplifier: amplify in X, attenuate in Y, without adding significant noise

w/o SQ:

Add 1/2 photon vacuum noise in each quadrature

w/o AMP (i.e., a standard amplifier)

Add 1/2 photon amplifier noise in each quadrature



Sub-meV Quantum Sensors: ULDM via EM Coupling

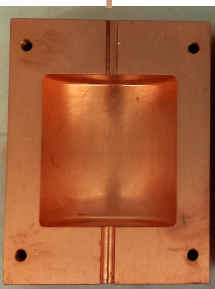
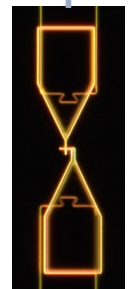
I μeV to 100 meV (240 MHz to 24 THz, QS3) example: QND photon counting ($g_{a\gamma\gamma}$)

Logical limit of squeezing:

- photon counting senses only the amplitude quadrature
- insensitive to vacuum fluctuations
- quantum non-demolition approach ensures high fidelity to eliminate false positives

e.g., Couple a transmon qubit to axion conversion cavity (Dixit et al 2021)

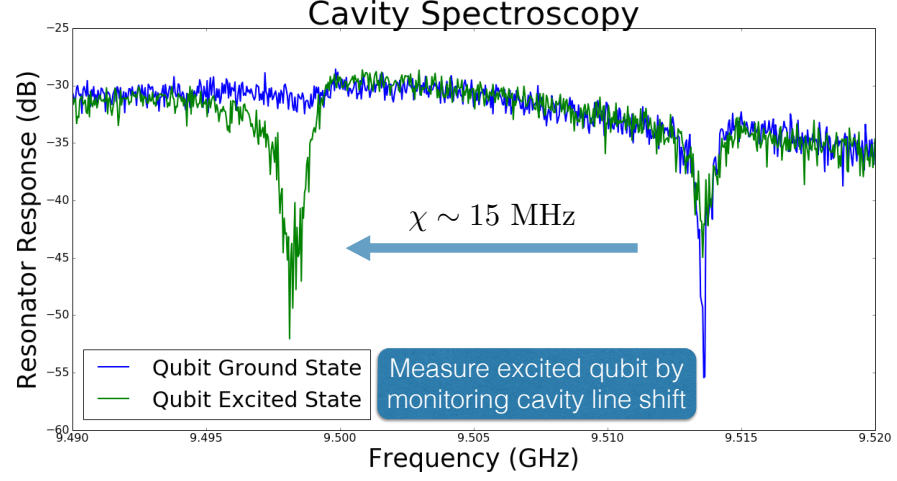
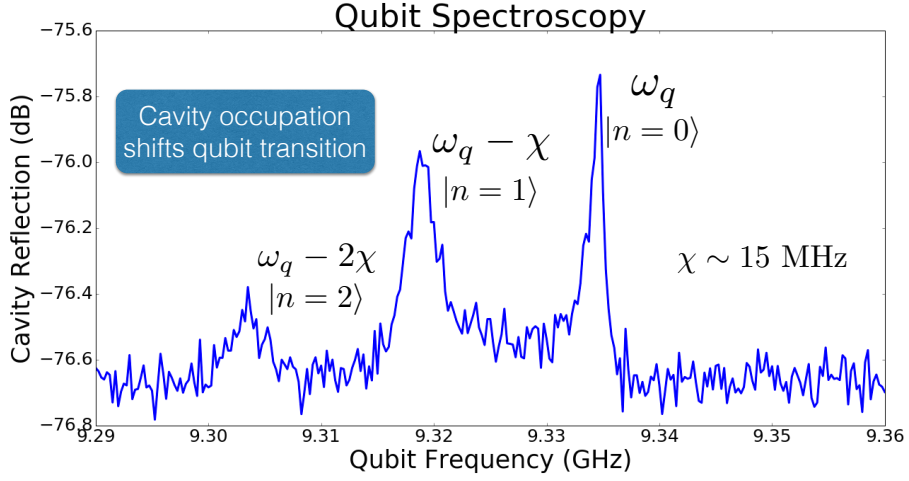
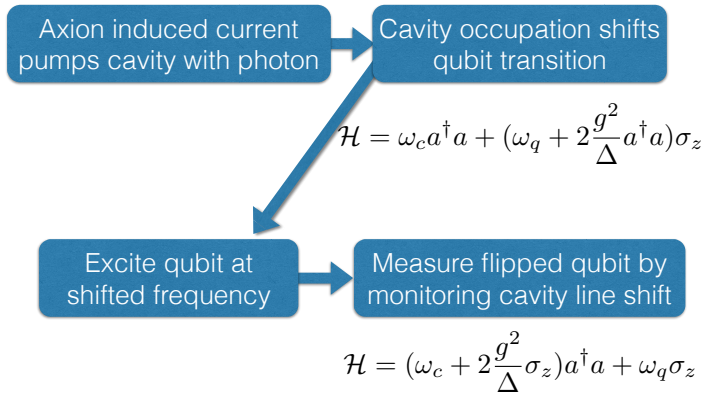
$$\mathcal{H} = \omega_c a^\dagger a + \omega_q \sigma_z + 2 \frac{g^2}{\Delta} a^\dagger a \sigma_z$$

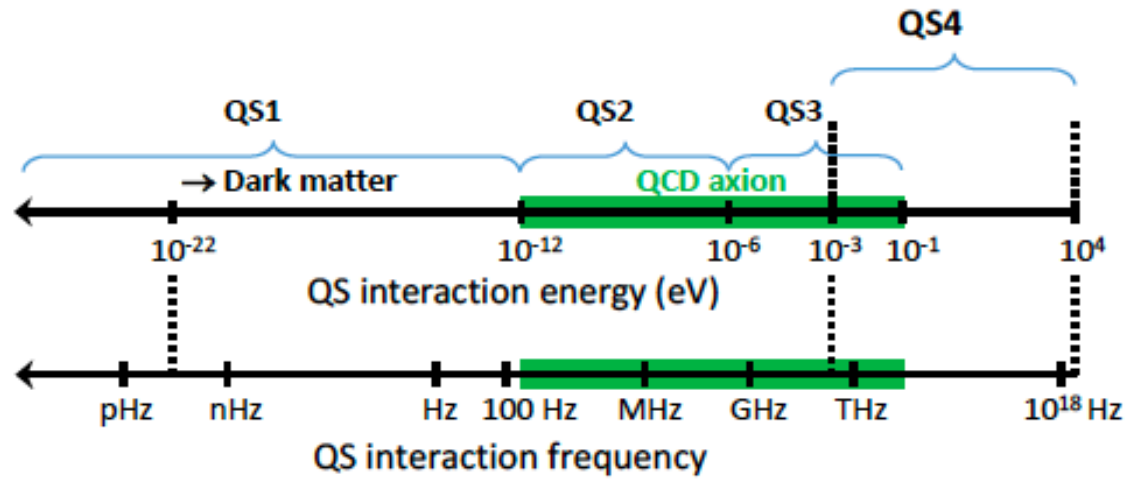
g : cavity-qubit coupling

$\Delta = \omega_q - \omega_c$ detuning

$\chi = \frac{g^2}{\Delta}$ Stark Shift



Sub-meV Quantum Sensors: ULDM via QCD & Spin Couplings



Exploit the weak coupling between ULDM and nuclear strong charge and/or spin

Axions:

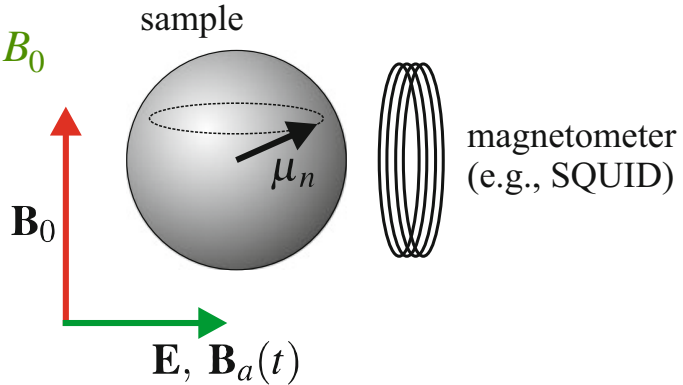
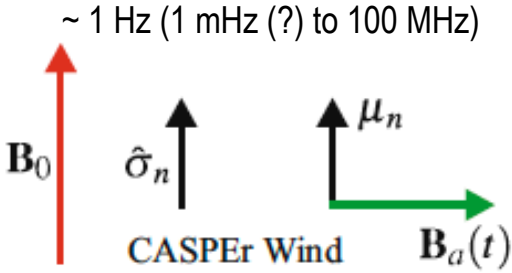
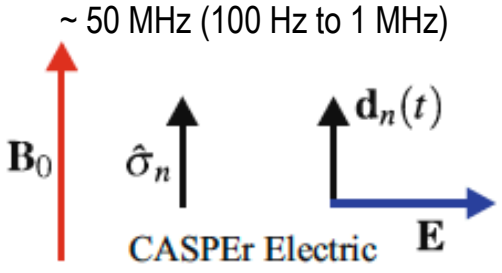
g_d axion-gluon coupling yields a time-varying nuclear EDM along spin axis

g_{aNN} axion-nuclear-spin coupling causes interaction of nuclear spin with gradient of axion field (axion wind)

In both cases, consider axion to be a classical time-varying field $a(\vec{r}, t) = \left(\frac{\hbar \sqrt{2 \rho_{DM}}}{m_a c} \right) \cos \omega_a t$

Examples:

CASPER: solid state NMR spin precession when $\omega_a = \Omega_{Larmor} = \gamma_n B_0$



Atomic magnetometers, spin-polarized torsion balances suitable for lower frequencies (10^{-8} to 100 Hz)

Sub-meV Quantum Sensors: Other Ideas

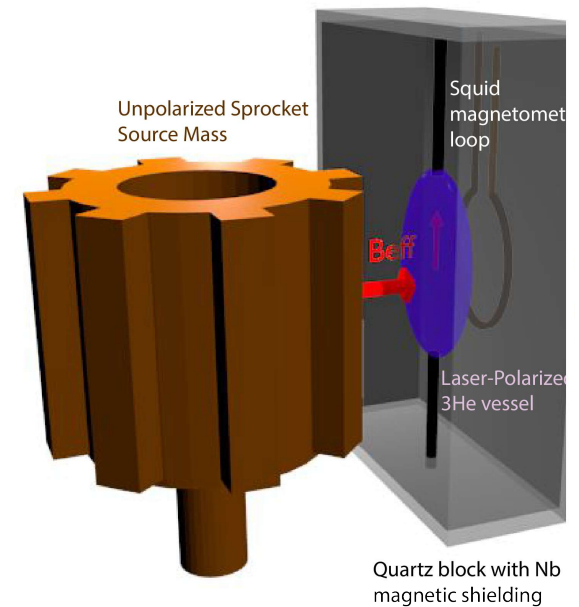
Fifth Forces: Replace time-varying cosmological DM axion-like particle or dark photon field with one created in the lab

e.g. ARIADNE

Apply time-varying ALP fifth force to laser-polarized ^3He cell via rotating sprocket source mass

NMR measurement of ^3He (a la CASPEr)

SQUID magnetometer pickup

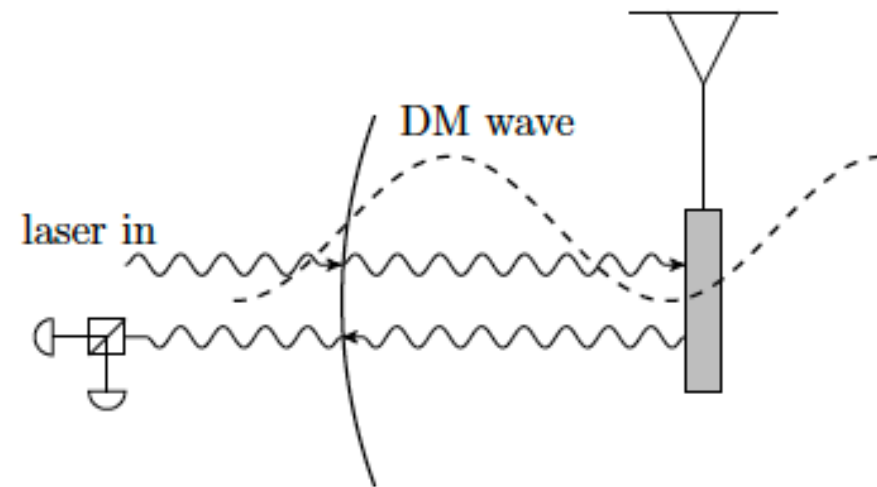


Opto-Mechanical Approaches for ULDM

Use laser-interferometric techniques to measure position difference between fixed and suspended masses composed of different materials

e.g., g_{aNN} axion-nuclear-spin coupling would result in different DM forces on materials of same mass

Can benefit from vacuum squeezing/squeeze amplification techniques as used in LIGO (and HAYSTAC)



Carney et al, arXiv:1908.04797

Sub-meV Quantum Sensors: Atomic Clocks, Atom Interferometers, Entanglement, etc.

Signatures of new physics:

DM and symmetry violation

time-dependent, equivalence-principle-violating accelerations

different DM couplings for different types of atoms
(e.g., g_{aNN})

symmetry violations yielding different forces on different kinds of atoms

time variation of effective fundamental constants

e.g., time-varying atomic transition frequencies

nuclear spin precession

Technical advances in atom interferometers

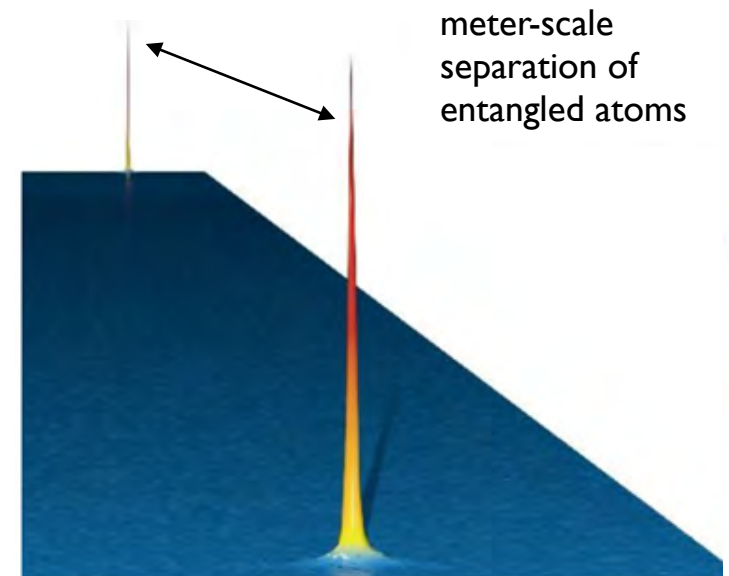
macroscopic delocalization (~ 1 m scale)

gradiometer configurations to cancel environmental systematics

single-photon atom interferometry

sub-nK cooling

dual species



Entanglement as a tool

Squeezing is one type of entanglement:
preparation of a specific type of quantum state

Generically, entanglement provides robustness against random noise processes, decoherence, etc.

Entanglement of distant sensors may provide ways to search for dark matter, better understand dark energy, and/or do new types of astronomical observations

Conclusions

The field of quantum sensors is very broad-ranging, employing a wide range of techniques from condensed matter physics, atomic/molecular/optical physics, and quantum information science

The interaction between particle physicists and these other fields has been intellectually exciting and very fruitful

Quantum Sensors are opening up significant new parameter space for precision measurement searches for new particle physics, resulting in the initiation of a wide range of new experiments