

Some special EM physics topics:

II. Some notes on charged particle stepping

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- Some special EM physics modelling topic to be discussed:
 - I. using secondary production thresholds (last time)
 - II. the corresponding stepping and its parameters (today)
 - III. multiple Coulomb scattering (next weeks)
- There are some ongoing optimisations in which these might help
- We also have some pending issues (regarding the "cuts", MSC)
- After some discussions with Marilena, it seemed to be a good idea to start working on these together by arranging these meetings
- These informal discussions might also help to tighten our collaboration



- II. Some notes on charged particle stepping (today)
 - Continuous step limit due to energy loss? Why?
 - How it is incorporated in **Geant4**. Why? Its parameters.
 - Relation to tracking cut or G4UserLimits.
 - How energy loss affects the **discrete** part of **the step limit**.
 - Some notes on the G4GammaGeneralProcess.





CONTINUOUS STEP LIMIT DUE TO ENERGY LOSS



- Introduce secondary photon production threshold:
 - secondary photons, with initial energy below a gamma production threshold(k<E_γ^{cut}), are not generated
 - E_γ^{cut} is the energy of the photon that has its absorption length equal to the cut (r, i.e. the radius of the sphere)
 - the corresponding energy (that would have been taken away from the primary, but remained inside the sphere) is accounted as CONTINUOUS energy loss of the primary particle along its trajectory
 - described by the radiative contribution of the (restricted) stopping power (dE/dx): mean energy loss due to subthreshold photon emissions in unit (path) length
 - i.e. when an electron makes a step with a given length *L*, one can compute the mean energy loss (due to sub-threshold photon emissions) along the step as *L* x *dE/dx* (would be true only if *E* = *const* along the step)
 - secondary photons, with initial energy above a gamma production threshold(k>E_γ^{cut}), are generated (since this energy is deposited outside of the sphere) (DISCRETE)
 - the emission rate is determined by the corresponding (restricted) cross section(σ)



$$\frac{\mathrm{d}E}{\mathrm{d}x}(E, E_{\gamma}^{\mathrm{cut}}, Z) = \mathcal{N} \int_{0}^{E_{\gamma}^{\mathrm{cut}}} k \frac{\mathrm{d}\sigma}{\mathrm{d}k}(E, Z) \mathrm{d}k$$

$$\sigma(E, E_{\gamma}^{\text{cut}}, Z) = \int_{E_{\gamma}^{\text{cut}}}^{E} \frac{\mathrm{d}\sigma}{\mathrm{d}k}(E, Z) \mathrm{d}k$$





 $\frac{\mathrm{d}E}{\mathrm{d}x}(E, E_{\gamma}^{\mathrm{cut}}, Z) = \mathcal{N} \int_{0}^{E_{\gamma}} k \frac{\mathrm{d}\sigma}{\mathrm{d}k}(E, Z) \mathrm{d}k$ $\sigma(E, E_{\gamma}^{\text{cut}}, Z) = \int_{E_{\gamma}^{\text{cut}}}^{L} \frac{\mathrm{d}\sigma}{\mathrm{d}k}(E, Z) \mathrm{d}k$ • Note - all quantities are "*restricted*" (covers only the sub/super secondary production threshold part of the interaction) - in case of ionisation, the minimum primary e- energy to be able to produce secondary e⁻ with initial energy > $E_{e^{-}}$ ^{cut} is $2x E_{e}$ -cut (due to the indistinguishable two e⁻ at the final state; the one with the lower energy is considered to be the secondary e-)

Continuous step limit due to energy loss







- Continuous step limit due to energy loss:
 - when using **condensed history** simulation, **continuous energy loss**es **impose a limit on the** charged particle **step** (beyond the discrete part)
 - an obvious choice of this limit **could be the** (pre-step, restricted) **range of the particle** (sum from collision and radiative, i.e. ionisation and bremsstrahlung)
 - this **would prevent** the particle **to go longer than the** (mean) **path length** at which **its energy** would **become zero due to** the (sub-threshold) **continuous energy loss**es
 - we **need to be even more strict** than this in order **to guarantee** the **stability** of the charged particle stepping: limit the allowed energy loss at most 20-25 % of the pre-step point energy of the particle (need to be relaxed at lower energies)
 - the *loss function* below ensures, that the charged particle is stopped at the appropriate position while its energy is deposited in the correct volume(s) !

$$\Delta S_{\text{eloss}}(E_0) = \begin{cases} R & \text{if } R < \rho_R \\ \alpha_R R + \rho_R (1 - \alpha_R)(2 - \rho_R/R) & \text{otherwise} \end{cases}$$

where $R := R(E_0)$ is the restricted range at energy E_0 , ρ_R final range and α_R rover range are two parameters.





- UI command to set the parameters:

/process/eLoss/StepFunction <roverRange-value> <finalRange-value> <unit>
/process/eLoss/StepFunctionMuHad" <roverRange-value> <finalRange-value> <unit>
- the lower the finalRange value the "later" (in range) the particle will be ranged-out (default 1 [mm])
- the lower the roverRange value the smaller the allowed steps will be (default 0.2, i.e. 20 % of R(E₀))

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$R:=R(E_0)$: restricted range including both collision and radiative contributions

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TRACKING CUTS, USERLIMITS V.S. THE LOSS FUNCTION



- Secondary production cuts vs tracking cut:
- Geant4 do not require any tracking cuts: charged particles (like e-/e+) are "ranged-out" thanks to the loss function discussed before (photons are absorbed)
- "ranged-out": **appropriate final position** (resulting appropriate location of the energy deposit)
- however, the user can easily introduce any limits on tracking by G4UserLimits
- a kinetic energy limit (tracking cut in energy) has also been introduced recently:
 - only for computing performance reasons: to kill low energy "loopers" (low density material and field)
 - this kinetic energy limit can be set to any (even to zero) energy values
 - ♦ different values for e⁻,e⁺ and for hadrons, muons
 - ♦ UI commands:

```
/process/em/lowestElectronEnergy 100 eV
```

```
/process/em/lowestMuHadEnergy 10 keV
```

 particles are killed when their kinetic energy drops below the limit and their energy is deposited (at the given point!)







- Tracking cut (e.g. G4UserLimits):
 - special user (limit) process, can impose tracking cut in kinetic energy, range, etc.
 - checked at the step-limit phase of the step, i.e. at the pre-step point: if the given condition evaluates to be true, this special process is selected, the particle is stopped immediately and its kinetic energy is deposited at the given point
 - + many different type of limits even per detector region; one extra step; inappropriate final position: the corresponding energy deposit (or a fraction of it) might be assigned to the wrong volume
- Tracking cut (e.g. kinetic energy limit):
- checked within the step (after the continuous energy loss and the discrete interaction) and the particle is stopped immediately when its kinetic energy drops below the (global) kinetic energy limit and its kinetic energy is deposited at the given point
- + no any extra steps; global, inappropriate final position: the corresponding energy deposit (or a fraction of it) might be assigned to the wrong volume
- (note: since the corresponding kinetic energy limit is usually small (default 1 [keV]) it do not cause any problems)



No tracking cut (e.g. relying on the loss function):



- whenever the **particle rage drops below** the **final range** parameter value, an extra step, with a continuous step limit equal to the particle range, is proposed
- + only appropriate final position and energy deposit location;
 requires an extra (last) step



- No racking cut (e.g. relying on the loss function):
- the only way to **ensure appropriate final position of the charged particle** (and energy deposit location; see below)
- whenever the **particle rage drops below** the *final range* parameter value, an extra step, with a continuous step limit equal to the particle range, is proposed
- + only appropriate final position and energy deposit location;
- + /process/eLoss/useCutAsFinalRange true Can be used to set *final range* = Cut
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GFANT4

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HOW ENERGY LOSS AFFECTS THE DISCRETE PART OF THE STEP LIMIT



Effects of the energy loss to the discrete part of the step limit

 assuming a single interaction, e.g. bremsstrahlung with a constant restricted macroscopic cross section Σ along the step:

according to the *Beer-Lambert law*, the p.d.f. of the corresponding **interaction length** s i.e. the probability that the interaction will happen at s, s + ds is p(s)ds

$$p(s) = \Sigma \exp[-s\Sigma] \to s \in \mathbb{E}xp\{\Sigma\}$$

with a mean or expected value of

$$\mathbb{E}(s) = \frac{1}{\Sigma} \equiv \lambda \text{ mean free path}$$

however, the restricted macroscopic cross section is not constant along the step: depends on the particle energy, that (in case of charged particles like e⁻⁾ is constantly changing along the s step!



Effects of the energy loss to the discrete part of the step limit



 having an energy E₀ at the pre-step point and E' = E₀ - ΔE_{loss} at the post-step point





- a fictive δ interaction (no effect) is introduced
- its macroscopic cross section $\Sigma^{\delta}(E)$ is such that const. = $\Sigma = \Sigma^{\delta}(E) + \Sigma^{real}(E)$
- Σ is already const. along the step and the interaction is either the *real* or the δ
- at the post-step point, the probability of the *real* interaction is $p(real) = \Sigma^{real}(E')/\Sigma$
- the only restriction is that $\Sigma \geq \Sigma^{\text{real}}(E)$ along the step (otherwise $p(\delta) < 0$)



- energy E₀ at the pre-step point and E' = E₀ ΔE_{loss} at the post-step point: instead of sampling the interaction till the next discrete *real* (e.g. brem.) event:
 - a fictive δ interaction (no effect) is introduced
 - its macroscopic cross section $\Sigma^{\delta}(E)$ is such that const. = $\Sigma = \Sigma^{\delta}(E) + \Sigma^{real}(E)$
 - Σ is already const. along the step and the interaction is either the *real* or the δ
 - at the post-step point, the probability of the real interaction is $p(\text{real}) \neq \Sigma^{\text{real}}(E')/\Sigma$
 - the only restriction is that $\Sigma \geq \Sigma^{\text{real}}(E)$ along the step (otherwise $p(\delta) < 0$)





Good, but can results in a high fraction of "delta" steps (when far from the max.)!

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Using the "maximum allowed energy loss" along the step (from the step function) !

- energy E_0 at the pre-step point and $E' = E_0 \Delta E_{loss}$ at the post-step point: instead of sampling the interaction till the next discrete *real* (e.g. brem.) event:
 - a fictive δ interaction (no effect) is introduced
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SOME NOTES ON THE G4GAMMAGENERALPROCESS



S GEANT4

Why we could use the "sum" of the two ("*real*" and " δ ") processes above?

if $s_i = \{s_1, \ldots, s_M\}$ independent stochastic variables, such that $s_i \in \mathbb{E}xp\{1/\lambda_i\}, \forall i = 1, \ldots, M$, then $\eta \equiv \min\{s_1, \ldots, s_M\} \in \mathbb{E}xp\left\{\sum_{i=1}^M 1/\lambda_i\right\}$

Proof: by definition, the cumulative distribution function of η at any x is equal to the probability that

$$P(\eta < x) = 1 - P(\eta \ge x) = 1 - P(s_1 \ge x, \dots, s_M \ge x) = 1 - \prod_{i=1}^M P(s_i \ge x)$$

$$= 1 - \prod_{i=1}^{M} [1 - P(s_i < x)] = 1 - \prod_{i=1}^{M} [1 - (1 - \exp(-x/\lambda_i))]$$
$$= 1 - \exp\left[-x \sum_{i=1}^{M} 1/\lambda_i\right] \to \eta \in \mathbb{E} xp\left\{\sum_{i=1}^{M} 1/\lambda_i\right\}$$

• This can be applied in other cases!



Some notes on the G4GammaGeneralProcess(UI cmd:/process/em/UseGeneralProces true)

- This can be applied in other cases: <u>above a given gamma energy</u>, the sum of the 4 (macroscopic) cross sections is smooth:
 - instead of evaluating the 4 macroscopic cross sections, sampling the corresponding 4 interaction lengths (s₁, s₂, s₃, s₄) and taking their minimum s as the current physics step length at each step
 - one might use the single sum of macroscopic cross section value to sample directly the minimum ${\bf s}$ length
 - then evaluate the 4 interactions only if needed, i.e when eventually physics limits the step
 - this can save up several run time interpolations (especially in highly granular detectors) without altering the results





THAT'S IT FOR TODAY