


Gintropy and the LGGR model for income and particle spectra

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What is it?



What is gintropy

a textual definition

- Formulas resembling entropy terms
- by calculating the Gini-index *def. due to Corrado Gini*
- (Gross Inequality Natural Index)
- **Its properties are entropy-like, but it is not an entropy.**

Definitions

Gini index due to Corrado Gini

Let $P(x)$ be a normalized PDF. The GINI is defined as:

$$G \equiv \frac{\langle |x - y| \rangle}{\langle |x + y| \rangle} = \frac{1}{2 \langle x \rangle} \int_0^{\infty} dx \int_0^{\infty} dy |x - y| P(x) P(y). \quad (1)$$

Cumulant population fraction (from the rich end):

$$\bar{C}(0) = 1.$$

$$\bar{C}(x) \equiv \int_x^{\infty} dy P(y). \quad (2)$$

Cumulant wealth per average (from the rich end):

$$\bar{F}(0) = 1.$$

$$\bar{F}(x) \equiv \frac{1}{\langle x \rangle} \int_x^{\infty} dy y P(y). \quad (3)$$

Pareto point

80/20: 80% of the wealth is possessed by 20% of the folk.

$$\text{Total population: } N_{\text{tot}} = \int_0^{\infty} dx N(x), \text{ total wealth: } X_{\text{tot}} = \int_0^{\infty} dx x N(x).$$

By this $P(x) = N(x)/N_{\text{tot}}$ and $\langle x \rangle = X_{\text{tot}}/N_{\text{tot}}$.



At any p -Pareto point: $\bar{C}(x_p) = p$; and $\bar{F}(x_p) = (1 - p)$

$(1 - p)X_{\text{tot}}$ is owned by pN_{tot} .

The original definition (1) can be expressed by using the cumulatives as

$$G = \int_0^{\infty} dx P(x) \int_x^{\infty} dy \frac{y-x}{\langle x \rangle} P(y) = \int_0^{\infty} dx P(x) \left[\bar{F}(x) - \frac{x}{\langle x \rangle} \bar{C}(x) \right] \quad (4)$$

Proof uses the $x \leftrightarrow y$ symmetry under the integral.

GINI

alternative expressions $\int_0^{\infty} dx P(x) \dots = \int_0^1 d\bar{C} \dots$ $\int_0^{\infty} dx \frac{x}{\langle x \rangle} P(x) \dots = \int_0^1 d\bar{F} \dots$

The cumulative of the cumulative:

$$\bar{h}(x) \equiv \int_x^{\infty} dy \bar{C}(y) = \int_x^{\infty} dy \int_y^{\infty} dz P(z) = \int_x^{\infty} dz \int_x^z dy P(z) = \int_x^{\infty} dz (z - x) P(z). \quad (5)$$

Indeed, $\bar{h}(x) = \langle x \rangle \bar{F}(x) - x \bar{C}(x)$ and $\bar{h}(0) = \langle x \rangle$.

We have the derivatives: $P(x) = -d\bar{C}/dx$, $xP(x) = -\langle x \rangle d\bar{F}/dx$ and therefore $x/\langle x \rangle = d\bar{F}/d\bar{C}$. Also $\bar{C} = -d\bar{h}/dx$.



$$G = \int_0^{\infty} dx \bar{F}(x) P(x) - \int_0^{\infty} dx \frac{x}{\langle x \rangle} \bar{C}(x) P(x) = \int_0^1 \bar{F} d\bar{C} - \int_0^1 \bar{C} d\bar{F}. \quad (6)$$

GINI

expressed via cumulative population

Using that $P(x) = \frac{d^2}{dx^2} \bar{h}$ and $\bar{C}(x) = -d\bar{h}/dx$, we integrate by parts

$$\langle x \rangle G = \int_0^{\infty} dx \bar{h} \frac{d^2 \bar{h}}{dx^2} = \bar{h}(0) \bar{C}(0) - \int_0^{\infty} dx \bar{C}^2(x). \quad (7)$$

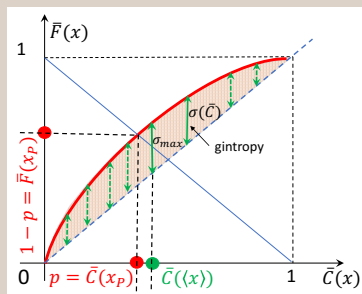
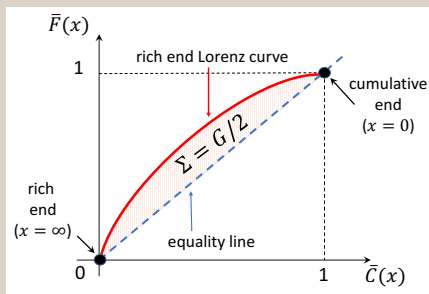
Replacing the boundary conditions we arrive at

$$G = 1 - \frac{1}{\langle x \rangle} \int_0^{\infty} dx \bar{C}^2(x) = \frac{1}{\langle x \rangle} \int_0^{\infty} dx \bar{C}(1 - \bar{C}). \quad (8)$$

For scaling $P(x) = \frac{1}{\langle x \rangle} f\left(\frac{x}{\langle x \rangle}\right)$, \bar{C} and G do not depend directly on $\langle x \rangle$.

Lorenz curve and gintropy

$$G = \int_0^1 (\bar{F}d\bar{C} - \bar{C}d\bar{F}) \quad \int_0^1 \bar{F}d\bar{C} = 1/2 + \Sigma, \quad \int_0^1 \bar{C}d\bar{F} = 1/2 - \Sigma.$$



GINI

expressed via gintropy: $G = 2 \int_0^1 \sigma(\bar{C}) d\bar{C}$, $\Sigma = \int_0^1 \bar{F} d\bar{C} - 1/2 = \int_0^1 (\bar{F} - \bar{C}) d\bar{C}$.

It can be shown that the half-moon area,

$$\Sigma \equiv \int_0^1 \sigma(\bar{C}) d\bar{C} \equiv \int_0^1 (\bar{F}(x) - \bar{C}(x)) d\bar{C}, \quad (9)$$

is exactly $G/2$. The integrand is alike an entropy-density we call it **gintropy**.

From the Lorentz curve geometry:

$$\int_0^1 \sigma d\bar{C} = \int_0^1 \sigma d\bar{F} = \frac{1}{2}G. \quad (10)$$

Properties of gintropy

general

- 1 The gintropy is never negative: $\sigma = \bar{F} - \bar{C} = C - F \geq 0$ inspecting the integral

$$\sigma = \int_x^\infty \left(\frac{y}{\langle x \rangle} - 1 \right) P(y) dy = \int_0^x \left(1 - \frac{y}{\langle x \rangle} \right) P(y) dy \geq 0$$

take first form for $y \geq x \geq \langle x \rangle$, the second form for the opposite case.

- 2 $\sigma(x)$ is maximal at $x = \langle x \rangle$: $d\sigma/dx = (1 - x/\langle x \rangle) P(x)$ changes sign there.

- 3 Convexity: $\frac{d\sigma}{d\bar{C}} = \frac{x}{\langle x \rangle} - 1$;

$$\frac{d^2\sigma}{d\bar{C}^2} = \frac{1}{\langle x \rangle} \frac{dx}{d\bar{C}} = -\frac{1}{\langle x \rangle P(x)} < 0.$$

- 4 For some PDF it looks like entropy density (see examples later).

GINI examples

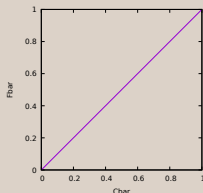
Communism: all incomes are equal

$P(x) = \delta(x - a)$ delivers $\langle x \rangle = a$, $\bar{C}(x) = \Theta(a - x)$ and $\bar{h}(x) = (a - x)\Theta(a - x)$.

This leads to $\bar{F} = (\bar{h} + x\bar{C}) / \langle x \rangle = \Theta(a - x)$ and by that

$$\sigma(x) = \bar{F}(x) - \bar{C}(x) = 0. \quad (11)$$

As a consequence also $\mathbf{G} = \mathbf{0}$.



GINI examples

Communism++ : some of them are more equal.

50/50 – 100/0

Two-peak-PDF, $P(x) = p\delta(x - a) + (1 - p)\delta(x - b)$, delivers $\langle x \rangle = pa + (1 - p)b$.

$$\bar{C}(x) = p\Theta(a - x) + (1 - p)\Theta(b - x) \quad (12)$$

Having the value 1 for $x \leq a$ and $(1 - p)$ for $x \in [a, b]$, otherwise 0.

$$G = \frac{1}{\langle x \rangle} \int_a^b p(1 - p) dx = \frac{(b - a)p(1 - p)}{\langle x \rangle} = \frac{(\langle x \rangle - a)(b - \langle x \rangle)}{(b - a)\langle x \rangle}. \quad (13)$$

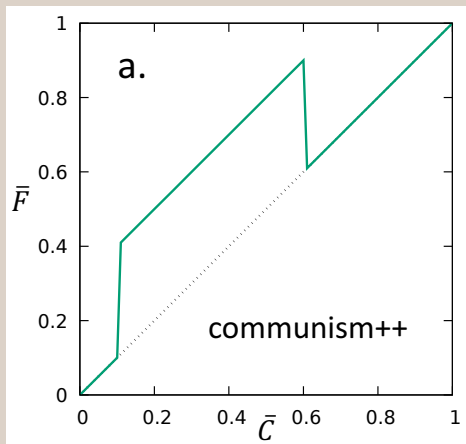
Gintropy is a box: $\sigma(\bar{C}) = G\Theta(\bar{C}(b) \leq \bar{C} \leq \bar{C}(a))$.

$$\bar{C}(b) = \frac{1-p}{2}, \quad \bar{C}(a) = 1 - \frac{p}{2}.$$

Pareto fraction: $\bar{C}_P = \frac{1-G}{2}, \quad \bar{F}_P = \frac{1+G}{2}, \quad G_{\max}(p) = \frac{\sqrt{b} - \sqrt{a}}{\sqrt{b} + \sqrt{a}}.$

Lorenz curve examples

Communism++ : some of them are more equal.



GINI examples

Eco-window:

50/50 – 62/38

The PDF has the form: $P(x) = \frac{1}{b-a} \Theta(a \leq x \leq b)$. Then the cumulative

$$\bar{C}(x) = \begin{cases} 1 & (x < a) \\ \frac{b-x}{b-a} & \text{otherwise} \\ 0 & (b < x) \end{cases} \quad (14)$$

Obviously $\langle x \rangle = (a + b)/2$ and the GINI

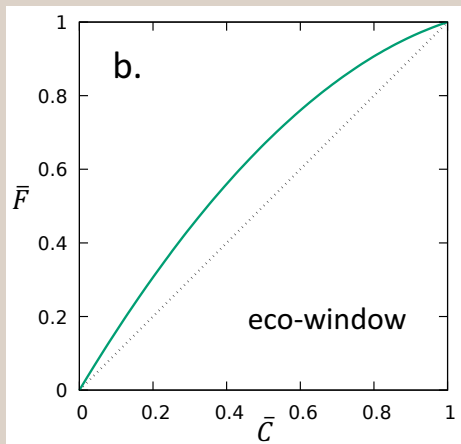
$$G = 1 - \frac{1}{\langle x \rangle} \left[a + \int_a^b \frac{(b-x)^2}{(b-a)^2} dx \right] = \frac{1}{3} \frac{b-a}{b+a}. \quad (15)$$

gintropy: $\sigma(\bar{C}) = \frac{b-a}{b+a} \bar{C}(1 - \bar{C})$.

For $a = 0$: $G = 1/3$, $\sigma = \bar{C}(1 - \bar{C})$ and the Pareto ratio: $\bar{F}_P/\bar{C}_P = 62/38$

Lorenz curve examples

Eco-window : equal chance between min and max.



GINI examples

Natural

68/32

The PDF scales: $P(x) = \frac{1}{\langle x \rangle} e^{-x/\langle x \rangle}$.

The corresponding tail-cumulative probability: $\bar{C}(x) = e^{-x/\langle x \rangle}$. The GINI becomes

$$G = 1 - \frac{1}{\langle x \rangle} \int_0^{\infty} e^{-2x/\langle x \rangle} dx = \frac{1}{2}. \quad (16)$$

Entropy formula is constructed as follows:

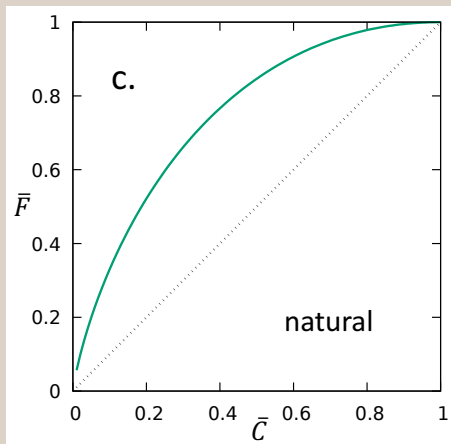
$$\bar{h} = \int_x^{\infty} e^{-y/\langle x \rangle} dy = \langle x \rangle e^{-x/\langle x \rangle}. \quad (17)$$

$$\langle x \rangle \bar{F} = \bar{h} + x \bar{C} = (x + \langle x \rangle) e^{-x/\langle x \rangle} \text{ and } \sigma = \frac{x}{\langle x \rangle} e^{-x/\langle x \rangle}$$

$$\sigma = -\bar{C} \ln \bar{C} \quad !$$

Lorenz curve examples

Natural : exponential PDF.



GINI examples

Capitalism: the rich gets richer

Start: $\bar{C}(x) = (1 + Ax)^{-B-1}$ leads to $\bar{h}(x) = \frac{1}{AB}(1 + Ax)^{-B}$. From this $\langle x \rangle = \bar{h}(0) = 1/AB$.

GINI becomes

$$G = 1 - AB \int_0^{\infty} (1 + Ax)^{-2B-2} dx = \frac{B+1}{2B+1}. \quad (18)$$

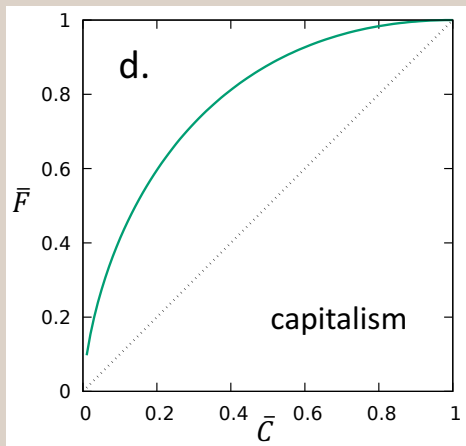
The gintropy

$$\sigma = Ax(B+1)(1 + Ax)^{-B-1} = (B+1) \left(\bar{C}^{\frac{B}{B+1}} - \bar{C} \right). \quad (19)$$

With $q = \frac{B}{B+1}$ we have $\sigma(\bar{C}) = \frac{1}{1-q}(\bar{C}^q - \bar{C})$ and $G = \frac{1}{q+1}$.

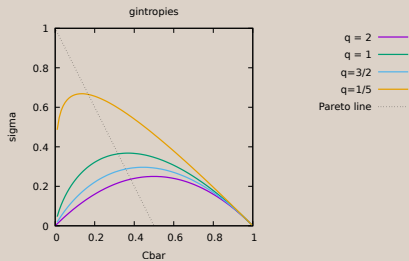
Lorenz curve examples

Capitalism : linear preference for the rich.



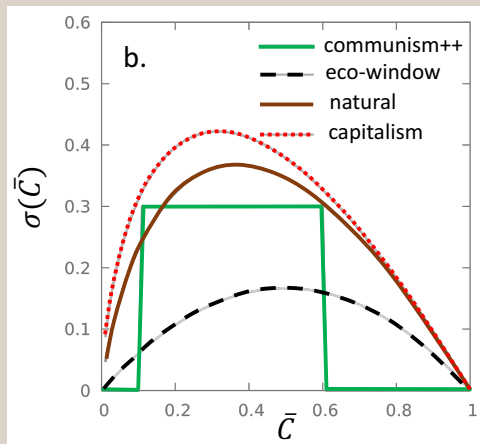
Gintropy

demonstrations



Gintropy

Comparison of examples



”An $e^{-\varepsilon/T}$ fit to ε -data justifies a thermal model”

is a FALLACY

Non-equilibrium stationary PDF, superstatistics

Non-thermal exponential PDF

with $T = E / \langle n \rangle$.

$$\rho(\varepsilon) = \left\langle \frac{\Omega_n(E - \varepsilon)}{\Omega_n(E)} \right\rangle = \left\langle \left(1 - \frac{\varepsilon}{E}\right)^n \right\rangle. \quad (20)$$

Statistics over several events:

$$\rho(\varepsilon) = \sum_{n=0}^{\infty} \left(1 - \frac{\varepsilon}{E}\right)^n P_n(E). \quad (21)$$

delivers the characteristic function of the PDF.

For $P_n(E) = a^n e^{-a} / n!$, Poissonian, $\langle n \rangle = a$ and

$$\rho(\varepsilon) = e^{-\frac{\langle n \rangle}{E} \varepsilon} \quad (22)$$

Non-thermal Tsallis–Pareto PDF

with $T = E / \langle n \rangle$ and $q = 1 + 1/(k + 1)$.

Statistics over several events:

$$\rho(\varepsilon) = \sum_{n=0}^{\infty} \left(1 - \frac{\varepsilon}{E}\right)^n P_n(E). \quad (23)$$

delivers the characteristic function of the PDF.

For NBD, $P_n(E) = \binom{k}{n} f^n (1+f)^{-n-k-1}$, one obtains

$$\rho(\varepsilon) = \left(1 + \frac{\langle n \rangle}{k+1} \frac{\varepsilon}{E}\right). \quad (24)$$

Phase Space and Entropy

Boltzmann – Gibbs – Planck – Einstein

Volume in n dimensions and linear size E : $\Omega_n(E) = e^{S_n(E)}$.

$$\begin{aligned} \rho(\varepsilon) &= \left\langle e^{S_n(E-\varepsilon) - S_n(E)} \right\rangle = \left\langle e^{-\varepsilon S'_n(E) + \frac{\varepsilon^2}{2} S''_n(E) + \dots} \right\rangle \\ &= \left\langle 1 - \varepsilon S'_n + \frac{\varepsilon^2}{2} \left((S'_n)^2 + S''_n \right) + \dots \right\rangle \quad (25) \end{aligned}$$

Interpretation of Tsallis parameters

Compare it with

$$\left(1 + (q-1)\frac{\varepsilon}{T}\right)^{-\frac{1}{q-1}} = 1 - \frac{\varepsilon}{T} + q\frac{\varepsilon^2}{2T^2} + \dots \quad (26)$$

interprets

$$\frac{1}{T} = \langle S'_n \rangle = \langle \beta_n \rangle, \quad (27)$$

and

$$q = T^2 \langle (S'_n)^2 + S''_n \rangle = T^2 \langle \beta_n^2 \rangle - \frac{dT}{dE}. \quad (28)$$

Finally

$$q = 1 + \frac{\Delta\beta_n^2}{\langle \beta_n \rangle^2} - \frac{1}{C}.$$

Construct another entropy

so that $S \rightarrow K(S)$

$$K(S_{\text{Renyi}}) = S_{\text{Tsallis}}$$

$$\rho_K(\varepsilon) = \left\langle e^{K(S(E-\varepsilon)) - K(S(E))} \right\rangle = 1 - \varepsilon K_1 + \frac{\varepsilon^2}{2} K_2 + \dots \quad (29)$$

Comparison with Tsallis–Pareto PDF delivers

$$\frac{1}{T_K} = \langle \beta \rangle K'; \quad \frac{q_K}{T_K^2} = (K'' + (K')^2) \langle \beta^2 \rangle - K' \frac{\langle \beta \rangle^2}{C}. \quad (30)$$

e.g. $K(S) = \frac{e^{(1-q)S} - 1}{1-q}$ leads to Tsallis and Renyi entropies ($\Delta\beta^2 = 0$).

Find $K(S)$ for $q_K = 1$, and then $K(S) = \sum_i p_i K(-\ln p_i)$.

LGGR model

local growth global reset

Transitions: from n to $n + 1$ (local) and from any n to 0 (global).

$$\frac{\partial}{\partial t} P = -\frac{\partial}{\partial x} (\mu P) - \gamma P. \quad (31)$$

Stationary PDF-s:

- γ and μ rates constant:

$$P(\infty, x) = Q(0) \exp -\frac{\gamma}{\mu} x$$

- γ constant $\mu(x) = \sigma(x + b)$ linear:

$$P(\infty, x) = Q(0) \left(1 + \frac{x}{b}\right)^{-1-\gamma/\sigma}$$

Summary

- There are exp-s which are **non-thermal**.
- T stems from $E / \langle n \rangle$, q from $\Delta n^2 / \langle x \rangle$.
- Entropy mapping $S \rightarrow K(S)$ to **ensure** $q_K = 1$.
- Several $K(S) = \sum_i p_i K(-\ln p_i)$ formulas are possible.
- Gini index \rightarrow **gintropy** reconstructs entropy formulas



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