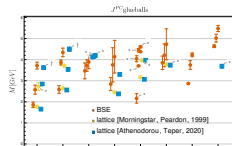
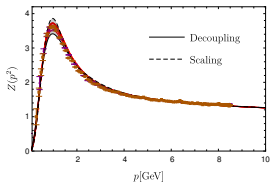
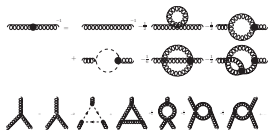


On the mass spectrum of glueballs with even charge parity



Markus Q. Huber

Institute of Theoretical Physics, Giessen
University

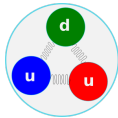
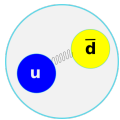
MQH, Phys.Rev.D 101, arXiv:2003.13703

MQH, Fischer, Sanchis-Alepuz, Eur.Phys.J.C 80,
arXiv:2004.00415

ACHT2021

April 21, 2021

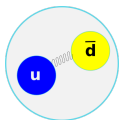
Bound states in QCD



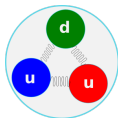
Mesons

Baryons

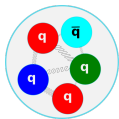
Bound states in QCD



Mesons



Baryons



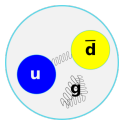
Pentaquarks

First observations 2015 (LHCb)



Tetraquarks

Increasing number of confirmed states



Hybrids



Glueballs

States of pure 'radiation'

Glueball observations

Experimental candidates, but situation not conclusive.

Scalar glueball: 0^{++} , mixing with scalar isoscalar mesons

Candidate reaction: $J/\psi \rightarrow \gamma + 2g$

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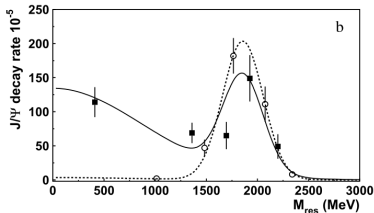
Scalar glueball: 0^{++} , mixing with scalar isoscalar mesons

Candidate reaction: $J/\psi \rightarrow \gamma + 2g$

Recent analysis of BESIII data [Sarantsev, Denisenko, Thoma, Klempt '21]:

$$M = 1865 \pm 25_{-30}^{+10} \text{ MeV},$$

$$\Gamma = 370 \pm 50_{-20}^{+30} \text{ MeV}$$



Future experiments: e.g., PANDA, GlueX

Glueball calculations

Yang-Mills theory

- “Isolated” problem: only gluons
- Clean picture: well-established lattice results

Glueball calculations

Yang-Mills theory

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Unquenching on the lattice [Gregory et al. '12]:

- Much higher statistics required (poor signal-to-noise ratio)
- Continuum extrapolation and inclusion of fermionic operators still to be done
- Mixing with $\bar{q}q$ challenging
- Tiny (e.g., 0^{++} , 2^{++}) to moderate unquenching effects (e.g., 0^{-+}) found

Glueball calculations

Yang-Mills theory

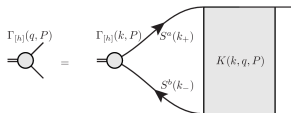
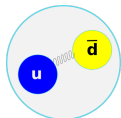
- “Isolated” problem: only gluons
- Clean picture: well-established lattice results
- **Functional methods: High quality input available for bound state equations**

Unquenching on the lattice [Gregory et al. '12]:

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Hadrons from bound state equations

Example: Meson

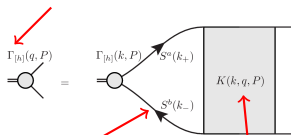
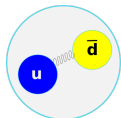


$$\text{Integral equation: } \Gamma(q, P) = \int dk \Gamma(k, P) S(k_+) S(k_-) K(k, q, P)$$

Hadrons from bound state equations

Bethe-Salpeter amplitude

Example: Meson



$$\text{Integral equation: } \Gamma(q, P) = \int dk \Gamma(k, P) S(k_+) S(k_-) K(k, q, P)$$

Ingredients:

- Quark propagator S

$$\begin{array}{c} \text{---} \circ \text{---} \\ S(p) \end{array}^{-1} = \begin{array}{c} \text{---} \\ S_0(p) \end{array}^{-1} + \begin{array}{c} \text{---} \circ \text{---} \\ S(q) \end{array} \begin{array}{c} \text{---} \circ \text{---} \\ S(p, q) \end{array} \begin{array}{c} \text{---} \circ \text{---} \\ S(q) \end{array}$$

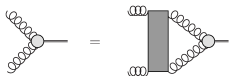
$D_{\mu\nu}(p-q)$

γ_μ γ_ν

- Interaction kernel K
- Constrained by symmetries

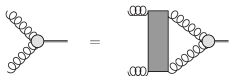
Nonperturbative diagram: full
momentum dependent dressings
→ numerical solution

Glueball BSE



Need loop and rectangle , solve for vertex . \rightarrow Mass

Glueball BSE



Need loop and rectangle , solve for vertex . \rightarrow Mass
 Not quite...

Glueball BSE



Gluons couple to ghosts \rightarrow Include 'ghostball'-part.

(First step: no quarks
 \rightarrow Yang-Mills theory)

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Construction of kernel

Consistency with input: Apply same construction principle.

Glueball BSE



Glueballs couple to ghosts \rightarrow Include 'ghostball'-part. (First step: no quarks
 \rightarrow Yang-Mills theory)

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Construction of kernel

Consistency with input: Apply same construction principle.

Previous BSE calculations for glueballs:

- ▶ [Meyers, Swanson '13]
- ▶ [Sanchis-Alepuz, Fischer, Kellermann, von Smekal '15]
- ▶ [Souza et al. '20]
- ▶ [Kaptari, Kämpfer '20]

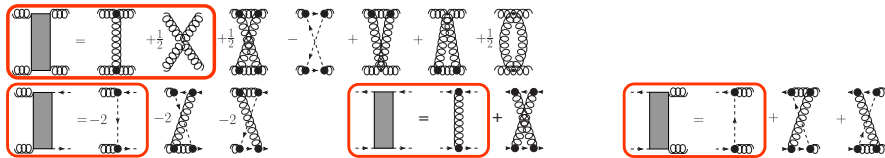
\Rightarrow **Input is important** for quantitative predictive power!


[MQH, Fischer, Sanchis-Alepuz '20]

Kernel construction

From 3PI effective action truncated to three-loops:

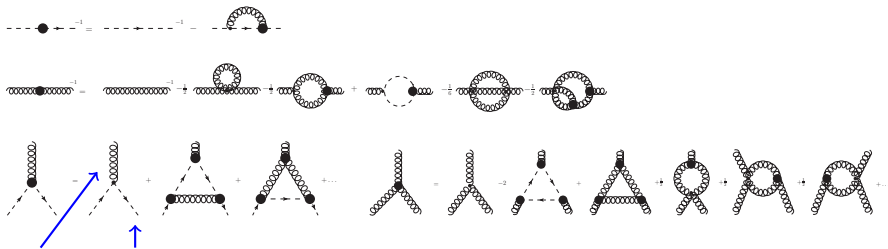
[Fukuda '87; McKay, Munczek '89; Sanchis-Alepuz, Williams '15; MQH, Fischer, Sanchis-Alepuz '20]



→ Need , , , .

- Some diagrams vanish for certain quantum numbers.
- Full QCD: Same for quarks → Mixing with mesons.

Equations of motion from 3-loop 3PI effective action



Gluon and ghost fields: Elementary fields of Yang-Mills theory in the **Landau gauge**

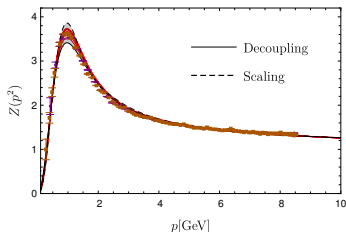
Self-contained system of equations with the scale as the only input.

Truncation \rightarrow 3-loop expansion of 3PI effective action [Berges '04]

- 4 coupled integral equations with full kinematic dependence.
- Sufficient numerical accuracy required for renormalization.
- One- and two-loop diagrams [Meyers, Swanson '14; MQH '17].

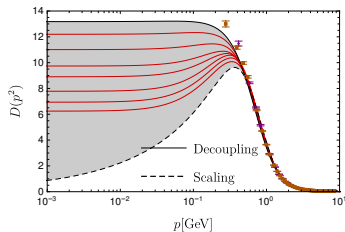
Landau gauge propagators

Gluon dressing function:

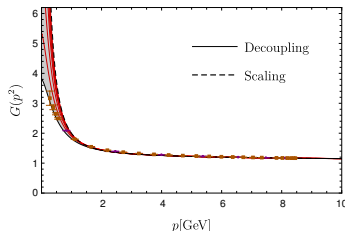


- Family of solutions:
Nonperturbative completions of Landau gauge [Maas '10]?
- Realized by condition on $G(0)$
[Fischer, Maas, Pawłowski '08; Alkofer, MQH, Schwenzer '08]

Gluon propagator:



Ghost dressing function:

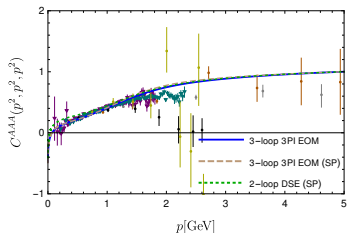


[Sternbeck '06; MQH '20]

Concurrence of functional methods

Exemplified with three-gluon vertex.

3PI vs. 2-loop DSE:

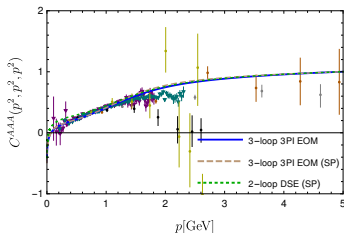


[Cucchieri, Maas, Mendes '08; Sternbeck et al. '17;
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Concurrence of functional methods

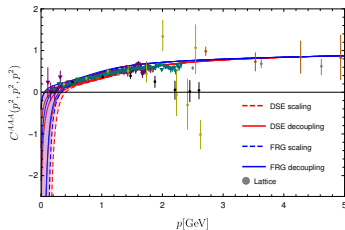
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DSE vs. FRG:

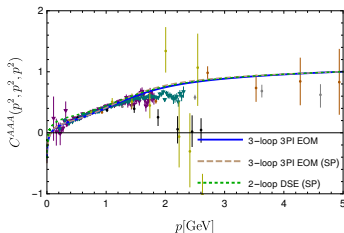


[Cucchieri, Maas, Mendes '08; Sternbeck et al. '17; Cyrol et al. '16; MQH '20]

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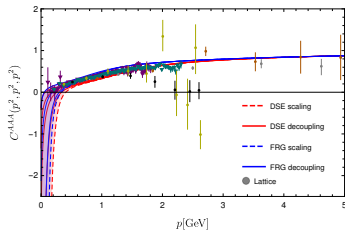
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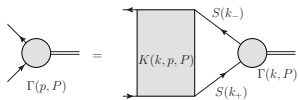


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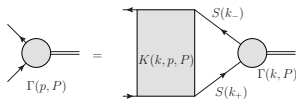
Beyond this truncation

- Further dressings of three-gluon vertex [Eichmann, Williams, Alkofer, Vujanovic '14]
- Effects of four-point functions [MQH '16, MQH '17, Corell et al. '18, MQH '18]

Solving a BSE



Solving a BSE

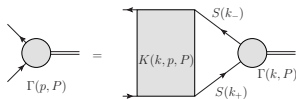


Consider the eigenvalue problem (Γ is the BSE amplitude)

$$\mathcal{K} \cdot \Gamma(P) = \lambda(P) \Gamma(P).$$

$\lambda(P^2) = 1$ is a solution to the BSE \Rightarrow Glueball mass $P^2 = -M^2$

Solving a BSE



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Calculation requires quantities for

$$k_{\pm}^2 = P^2 + k^2 \pm 2\sqrt{P^2 k^2} \cos \theta = -M^2 + k^2 \pm 2iM\sqrt{k^2} \cos \theta.$$

\Rightarrow Complex momentum arguments.

Extrapolation of $\lambda(P^2)$

Extrapolation method

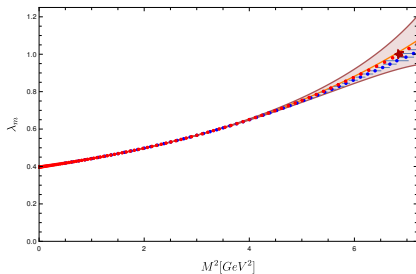
- Extrapolation to time-like P^2 using Schlessinger's continued fraction method (proven superior to default Padé approximants) [Schlessinger '68]
- Average over extrapolations using subsets of points for error estimate

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Test extrapolation for solvable system: Heavy meson



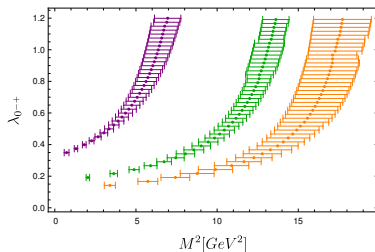
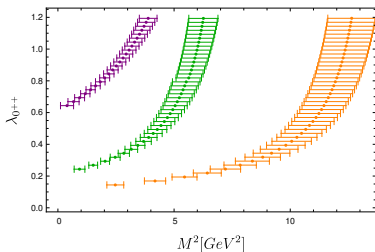
[MQH, Sanchis-Alepuz, Fischer '20]

Extrapolation of $\lambda(P^2)$ for glueballs

Higher eigenvalues: Excited states.

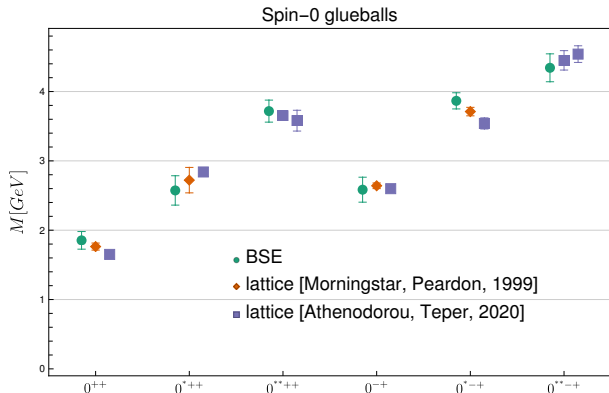
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Higher eigenvalues: Excited states.



Physical solutions for $\lambda(P^2) = 1$.

Glueballs masses for $0^{\pm+}$

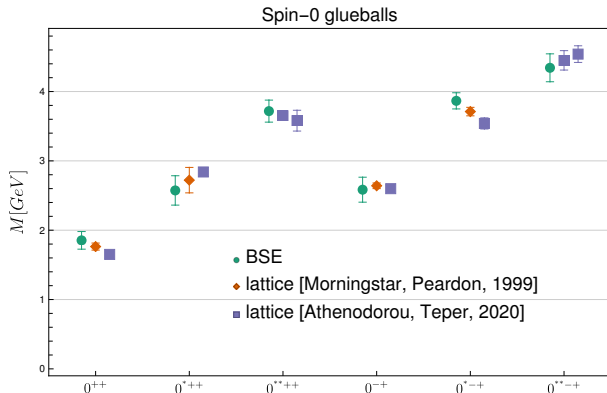


Lattice 0^{**++} :
 Conjectured based on
 irred. rep. of octahedral
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All results for $r_0 = 1/418(5)$ MeV.

[MQH, Fischer, Sanchis-Alepuz '20]

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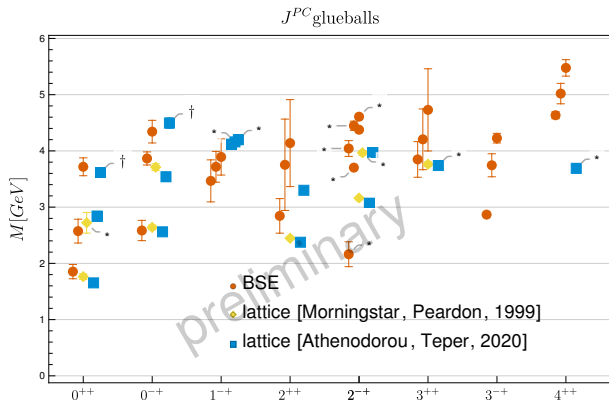
All results for $r_0 = 1/418(5)$ MeV.

[MQH, Fischer, Sanchis-Alepuz '20]

Under conjecture that choice of solution is a gauge choice: **Explicit test of gauge independence!**

Tested that results are independent of family of solutions.

Glueball masses for $J^{\pm+}$



Lattice:

*: identification with some uncertainty

†: conjecture based on irred. rep of octahedral group

[MQH, Fischer, Sanchis-Alepuz, in preparation]

Summary

Parameter-free determination of glueball masses from functional methods.

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Thank you for your attention.

Glueballs as bound states

Hadron masses from correlation functions of **color singlet operators**.

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Example: For $J^{PC} = 0^{++}$ glueball take $O(x) = F_{\mu\nu}(x)F^{\mu\nu}(x)$:

$$D(x - y) = \langle O(x)O(y) \rangle$$

- → Lattice: Mass from this correlator by exponential Euclidean time decay.
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Hadron masses from correlation functions of **color singlet operators**.

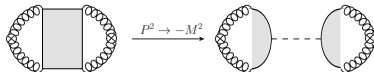
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- Complicated object in a diagrammatic language: 2-, 3- and 4-gluon contributions

Put total momentum on-shell and consider individual 2-, 3- and 4-gluon contributions. → Each can have a pole at the glueball mass.

A^4 -part of $D(x - y)$, total momentum on-shell:



Charge parity

Transformation of gluon field under charge conjugation:

$$A_{\mu}^a \rightarrow -\eta(a)A_{\mu}^a$$

where

$$\eta(a) = \begin{cases} +1 & a = 1, 3, 4, 6, 8 \\ -1 & a = 2, 5, 7 \end{cases}$$

Color neutral operator with two gluon fields:

$$A_{\mu}^a A_{\nu}^a \rightarrow \eta(a)^2 A_{\mu}^a A_{\nu}^a = A_{\mu}^a A_{\nu}^a.$$

$$\Rightarrow C = +1$$

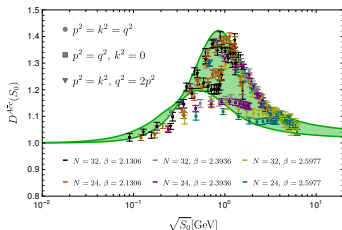
Negative charge parity, e.g.:

$$\begin{aligned} d^{abc} A_{\mu}^a A_{\nu}^b A_{\rho}^c &\rightarrow -d^{abc} \eta(a)\eta(b)\eta(c) A_{\mu}^a A_{\nu}^b A_{\rho}^c = \\ &= -d^{abc} A_{\mu}^a A_{\nu}^b A_{\rho}^c. \end{aligned}$$

Only nonvanishing elements of the symmetric structure constant d^{abc} : zero or two indices equal to 2, 5 or 7.

Landau gauge vertices

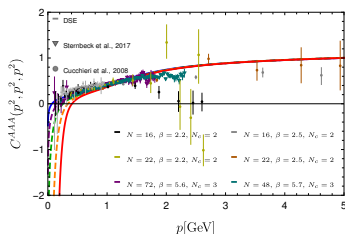
Ghost-gluon vertex:



[Maas '19; MQH '20]

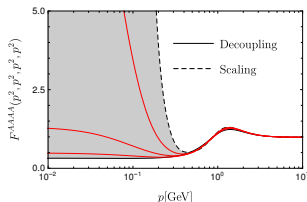
- Nontrivial kinematic dependence of ghost-gluon vertex
- Simple kinematic dependence of three-gluon vertex
- Four-gluon vertex from solution

Three-gluon vertex:



[Cucchieri, Maas, Mendes '08; Sternbeck et al. '17; MQH '20]

Four-gluon vertex:

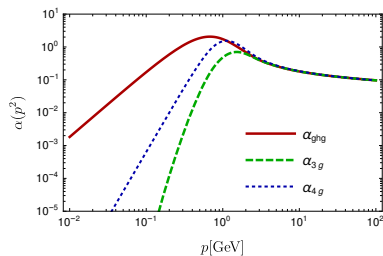


[MQH '20]

Some properties of the Landau gauge solution

[MQH '20]

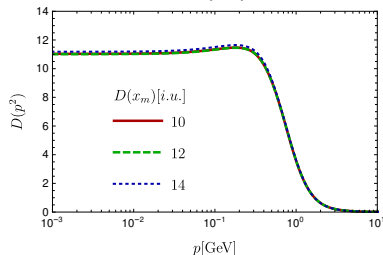
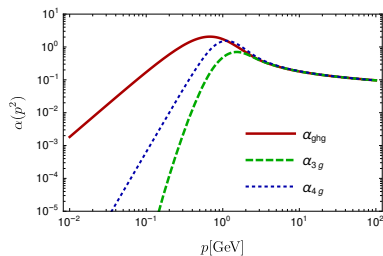
- Slavnov-Taylor identities (gauge invariance): Vertex couplings agree down to GeV regime



Some properties of the Landau gauge solution

[MQH '20]

- Slavnov-Taylor identities (gauge invariance): Vertex couplings agree down to GeV regime
- Renormalization: First parameter-free subtraction of quadratic divergences
 \Rightarrow **One unique free parameter** (family of solutions)



Landau gauge propagators in the complex plane

Propagators for complex momenta

- **Reconstruction** from Euclidean results: mathematically ill-defined, bias in solution
- **Direct calculation** from functional methods possible, e.g., contour deformation or spectral DSEs [Horak, Pawłowski, Wink '20]

Landau gauge propagators in the complex plane

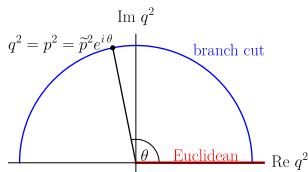
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Contour deformation: Special technique to respect analyticity (avoid branch cuts in the integrand)

- ▶ QED3 [Maris '95 (QED)]
- ▶ Quark propagator [Alkofer, Fischer, Detmold, Maris '04]
- ▶ Self-consistent solution: **Ray technique**, YM propagators [Strauss, Fischer, Kellermann '12; Fischer, MQH '20]
- ▶ Glueball correlators [Windisch, Alkofer, Haase, Liebmann '13; Windisch, MQH, Alkofer '13]
- ▶ Meson decays [Weil, Eichmann, Fischer, Williams '17; Williams '18]
- ▶ Spectral functions at $T > 0$ [Pawłowski, Strodthoff, Wink '18]
- ▶ Quark-photon vertex [Miramontes, Sanchis-Alepuz '19]
- ▶ Scalar scattering amplitude [Eichmann, Duarte, Pena, Stadler '19]

Landau gauge propagators in the complex plane



Appearance of branch cuts for complex momenta forbids integration directly to cutoff.

Deformation of integration contour necessary [Maris '95]. Recent resurgence: [Alkofer et al. '04; Windisch, MQH, Alkofer, '13; Williams '19; Miramontes, Sanchis-Alepuz '19; Eichmann et al. '19], ...

Ray technique for self-consistent solution of a DSE: [Strauss, Fischer, Kellermann; Fischer, MQH '20].

Landau gauge propagators in the complex plane

Technique to resp. analyticity (avoid branch cuts in integrand): **Contour deformation**

Simpler truncation:

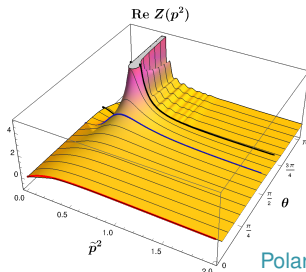
$$\text{wavy line with a black dot}^{-1} = \text{wavy line}^{-1} - \frac{1}{2} \text{wavy line} \text{ loop wavy line} + \text{wavy line} \text{ dashed loop wavy line}$$

Landau gauge propagators in the complex plane

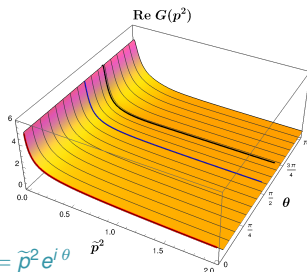
Technique to resp. analyticity (avoid branch cuts in integrand): **Contour deformation**

Simpler truncation:

$$\text{---}^{\circ} \text{---}^{-1} = \text{---}^{\circ} \text{---}^{-1} - \frac{1}{2} \text{---}^{\circ} \text{---}^{-1} + \text{---}^{\circ} \text{---}^{-1}$$



Polar coordinates: $p^2 = \tilde{p}^2 e^{i\theta}$



- Current truncation leads to a pole-like structure in the gluon propagator.
- Analyticity up to 'pole' confirmed by various tests (Cauchy-Riemann, Schlessinger, reconstruction)
- No proof of existence of complex conjugate poles due to simple truncation.