

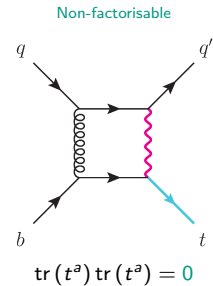
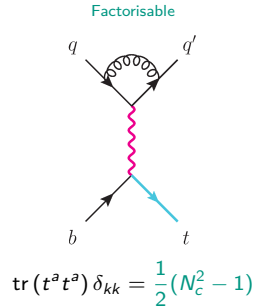
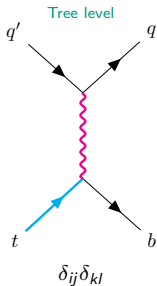
Non-factorisable two-loop contribution to t-channel single-top production

Based on [arXiv:2108.09222](https://arxiv.org/abs/2108.09222) with Kirill Melnikov, Jérémie Quarroz & Chen-Yu Wang

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Non-factorisable contribution to t-channel single-top production

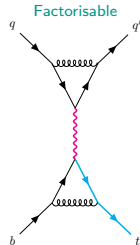
- Process interesting for studies of for example tbW vertex, V_{tb} constraints, Γ_t and m_t . See Monday's talks by *Victor Rodriguez Bouza and John Campbell*



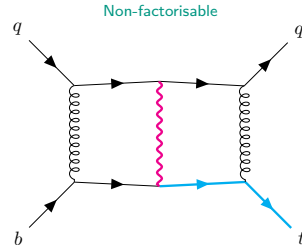
- Cross section at NLO known since a while *Harris et al. 2002; Campbell, Ellis, et al. 2004; Sullivan 2004; Cao and Yuan 2005; Sullivan 2005; Schwienhorst et al. 2011*
- No non-factorisable contribution to $\mathcal{A}_{\text{LO}} \otimes \mathcal{A}_{\text{NLO}}$

Non-factorisable contribution to t-channel single-top production

- For the NNLO correction to the cross section, we consider the interferences $\mathcal{A}_{\text{LO}} \otimes \mathcal{A}_{\text{NNLO}}$ and $\mathcal{A}_{\text{NLO}} \otimes \mathcal{A}_{\text{NLO}}$.
- Non-factorisable diagrams are colour-suppressed:



$$\text{tr}(t^a t^a) \text{tr}(t^b t^b) = \frac{1}{4}(N_c^2 - 1)^2$$

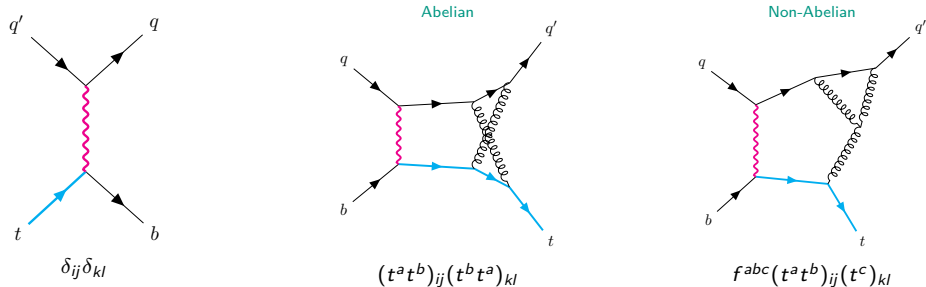


$$\text{tr}(t^a t^b) \text{tr}(t^a t^b) = \frac{1}{4}(N_c^2 - 1)$$

- Calculations of the factorisable contribution at NNLO exists. *Brucherseifer et al. 2014; Berger et al. 2016; Campbell, Neumann, et al. 2021*
- Non-factorisable contribution demonstrated to be enhanced for VBF Higgs production. *Liu, Melnikov, et al. 2019*

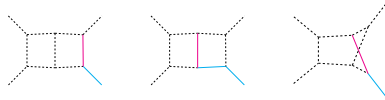
Non-factorisable diagrams at two loops

- For the NNLO correction to the cross section, we consider the interferences $\mathcal{A}_{LO} \otimes \mathcal{A}_{NNLO}$ and $\mathcal{A}_{NLO} \otimes \mathcal{A}_{NLO}$.



- Only **18 non-vanishing diagrams**, generated in QGRAF *Nogueira 1993* and processed in FORM *Vermaseren 2000; Kuipers et al. 2015; Ruijl et al. 2017*.
- Integral reduction performed **analytically** with KIRA 2.0 *Klappert, Lange, et al. 2020* and FireFly *Klappert and Lange 2020; Klappert, Klein, et al. 2021*
- We find **428** masters in 18 different families

Master integral evaluation



- Based on the **auxiliary mass flow method** *Liu, Ma, and Wang 2018; Liu, Ma, Tao, et al. 2020; Liu and Ma 2021*

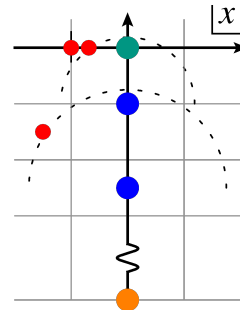
$$m_W^2 \rightarrow m_W^2 - i\eta.$$

- Solve differential equations at each kinematic point

$$\partial_x \mathbf{I} = \mathbf{M} \mathbf{I}, \quad x \propto -i\eta.$$

with boundary condition $x \rightarrow -i\infty$.

$$I = \sum_j^M \epsilon^j \sum_k^N \sum_l^I c_{jkl} x^k \ln^l x + \dots$$



Stepping from the boundary at $x \rightarrow -i\infty$, via **regular** points, to the **physical** mass. Step size is limited by **singularities** of the equation.

Results

- All **428** master integrals evaluated in ~ 30 minutes on a single core
- Comparison of poles at a typical phase space point

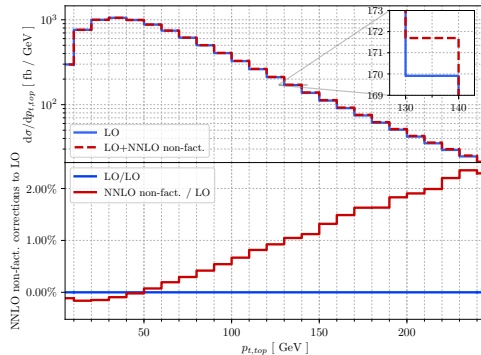
	ϵ^{-2}	ϵ^{-1}
$\langle \mathcal{A}^{(0)} \mathcal{A}_{\text{nf}}^{(2)} \rangle$	$-229.0940408654660 - 8.978163333241640i$	$-301.1802988944764 - 264.1773596529505i$
IR poles	$-229.0940408654665 - 8.978163333241973i$	$-301.1802988944791 - 264.1773596529535i$

- Double-virtual cross-section calculation from fixed grid of 100k points

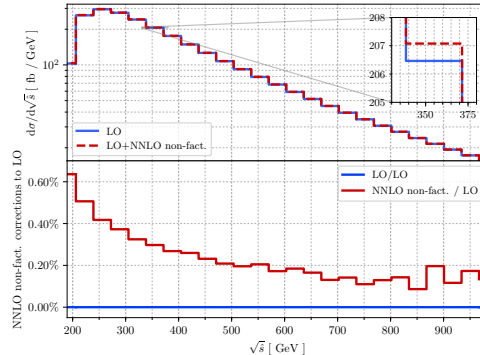
$$\sigma_{pp \rightarrow dt}^{ub} = \left(90.3 + 0.3 \left(\frac{\alpha_s(\mu_{\text{nf}})}{0.108} \right)^2 \right) \text{ pb}$$

- Correction of about 0.3% for $\mu_{\text{nf}} = 173$ GeV
- Real corrections are to be included and good choice of μ_{nf} is unclear

Results



(a) The top quark transverse momentum distribution.



(b) Differential cross section with respect to the partonic centre-of-mass energy $\sqrt{\hat{s}}$.

- Constant relative importance between factorisable and non-factorisable contributions throughout phase space