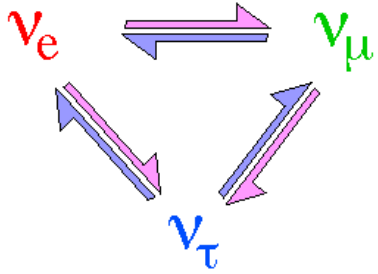


# LECTURE III & IV

- The standard  $3\nu$  scenario and its unknowns: status and prospects
- Neutrinos and beyond the Standard Model physics

# Standard 3ν scenario



$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} = U_{PMNS}(\theta_{12}, \theta_{23}, \theta_{13}, \delta, \dots) \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix}$$

$$\theta_{12} \sim 34^\circ$$

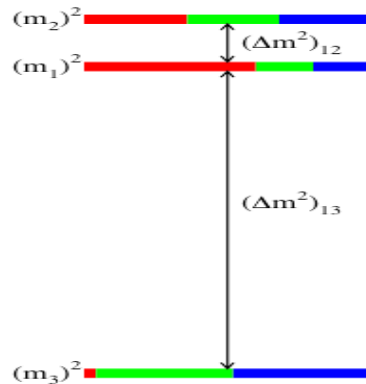
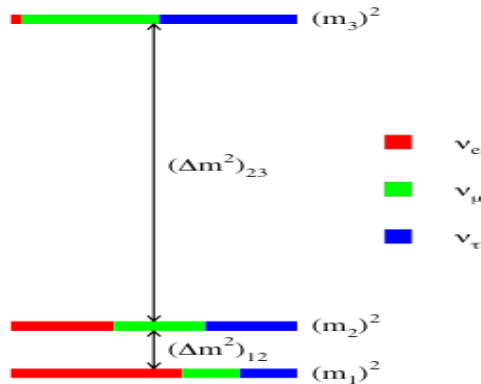
$$\theta_{23} \sim 42^\circ \text{ o } 48^\circ$$

$$\theta_{13} \sim 8.5^\circ$$

$$\delta \sim ?$$

normal hierarchy

inverted hierarchy



NO/NH

IO/IH

$$\updownarrow 7.5 \cdot 10^{-5} \text{eV}^2$$

$$2.5 \cdot 10^{-3} \text{eV}^2$$

**Caveat:** O(eV) neutrinos...reactor/short baseline anomalies still unresolved

# Outliers: LSND anomaly

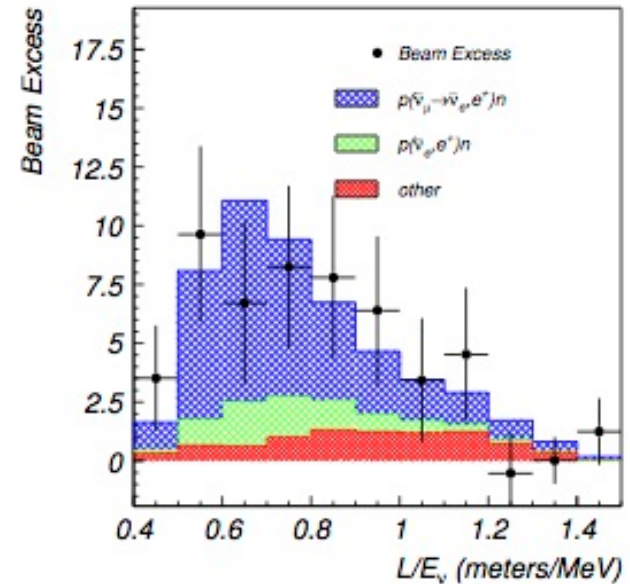
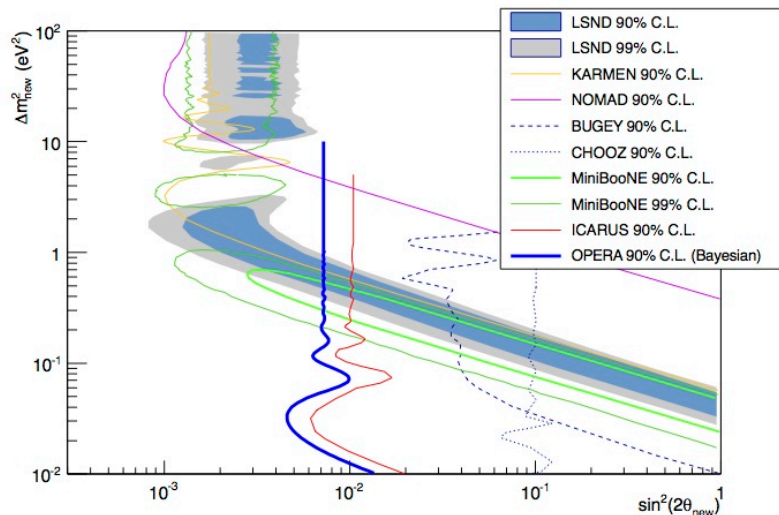
$$\pi^+ \rightarrow \mu^+ \nu_\mu$$

$$\nu_\mu \rightarrow \nu_e \text{ DIF } (28 \pm 6 / 10 \pm 2)$$

$$\mu^+ \rightarrow e^+ \nu_e \bar{\nu}_\mu$$

$$\bar{\nu}_\mu \rightarrow \bar{\nu}_e \text{ DAR } (64 \pm 18 / 12 \pm 3)$$

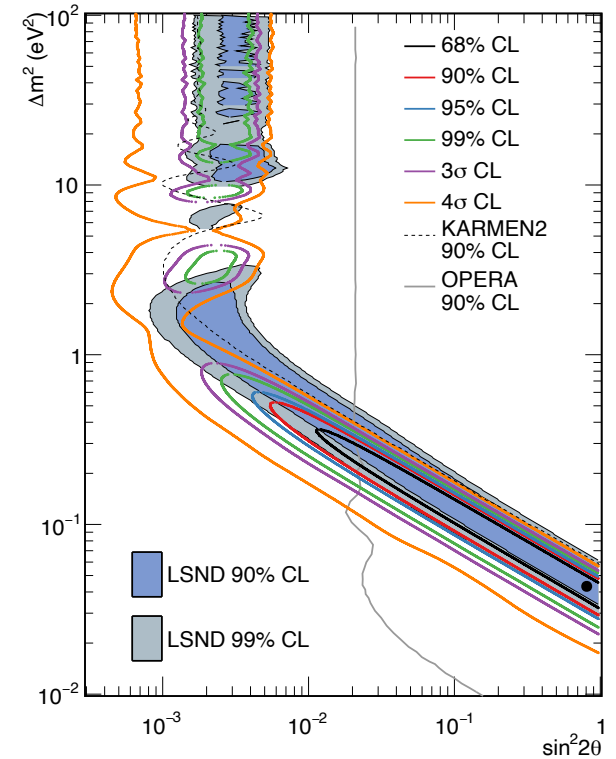
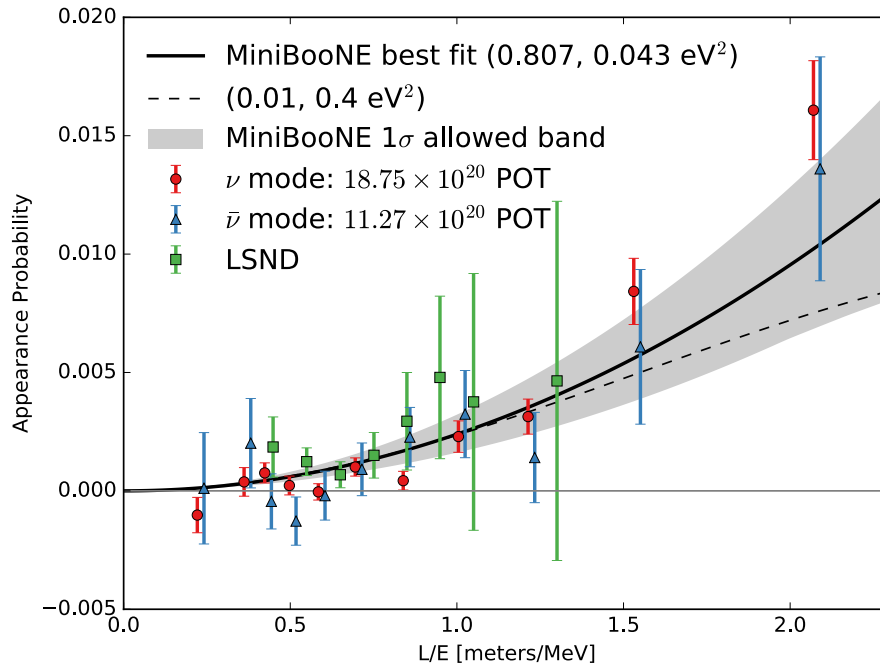
$$P(\nu_\mu \rightarrow \nu_e)$$



$$|\Delta m^2| \gg |\Delta m_{atm}^2|$$

# Outliers: LSND anomaly

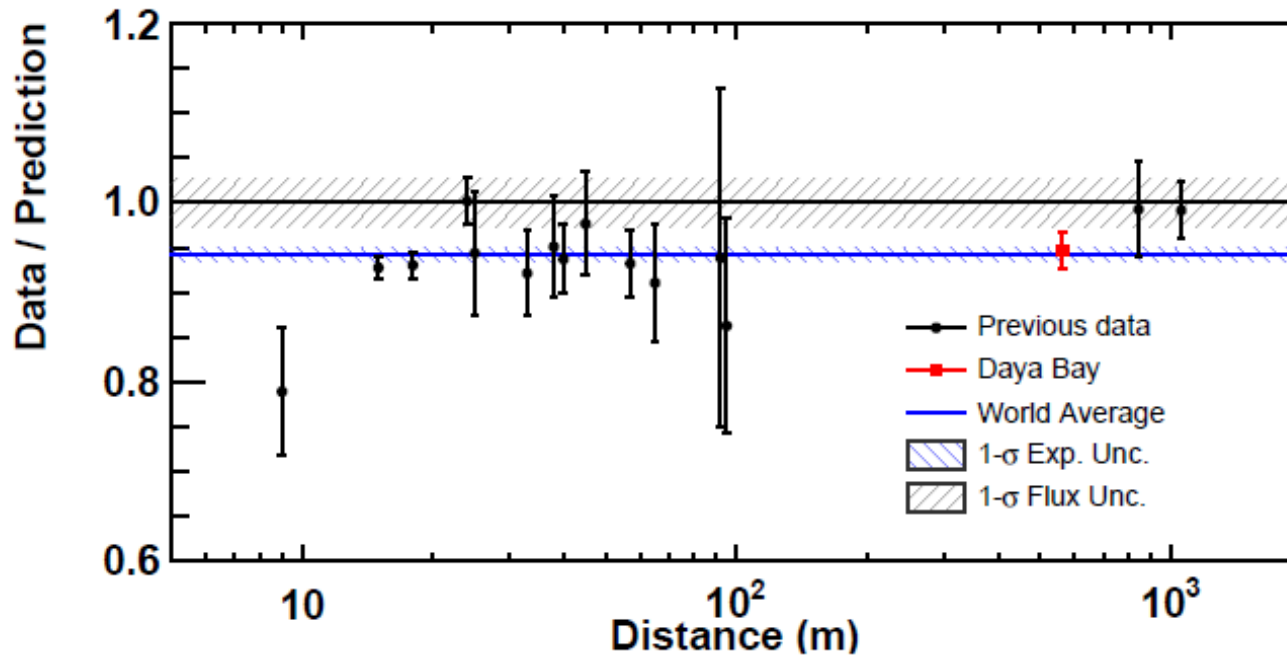
+ MiniBOONE



4.8 $\sigma$  discrepancy with SM !

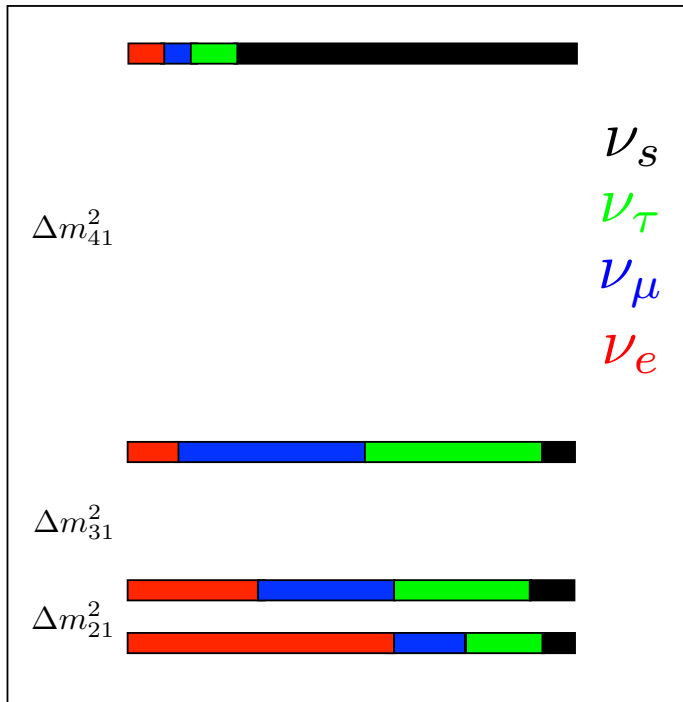
# Outliers: SBL reactor anomalies

Reactor  $P(\bar{\nu}_e \rightarrow \bar{\nu}_e)$



Re-evaluation of the predicted fluxes indicates an L-independent deficit (averaged oscillations ?)

# SBL anomalies: 4<sup>th</sup> neutrino ?



$$P(\nu_\mu \rightarrow \nu_e) = O(|U_{e4}|^2 |U_{\mu4}|^2)$$

$$P(\nu_e \rightarrow \nu_e) = O(|U_{e4}|^2)$$

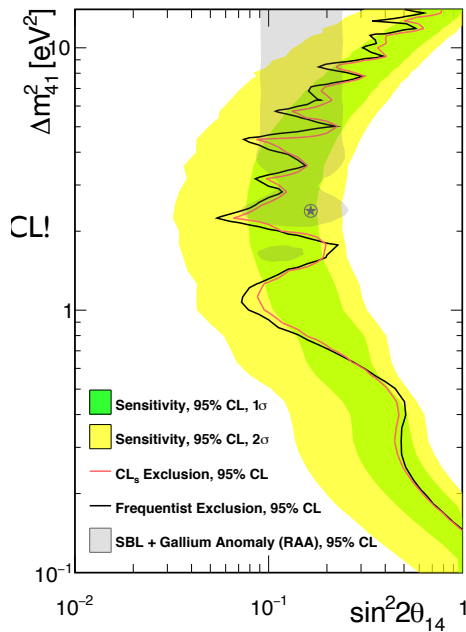
$$P(\nu_\mu \rightarrow \nu_\mu) = O(|U_{\mu4}|^2)$$

Oscillations at @meters for MeV neutrinos: **short baseline reactor experiment**

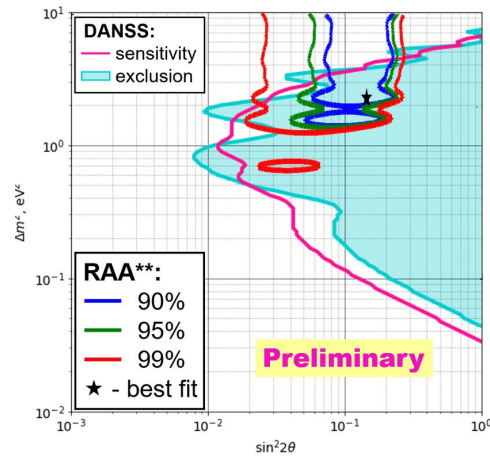
# SBL reactor anomaly Views

## New SBL reactor strategies: L-dep of signal

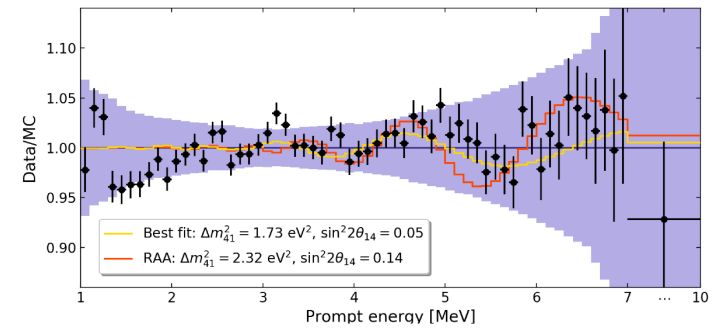
Prospect



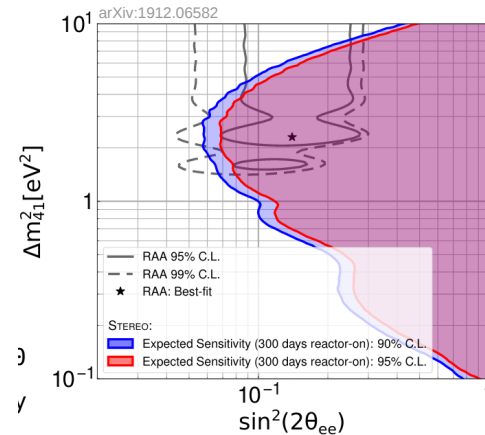
DANSS



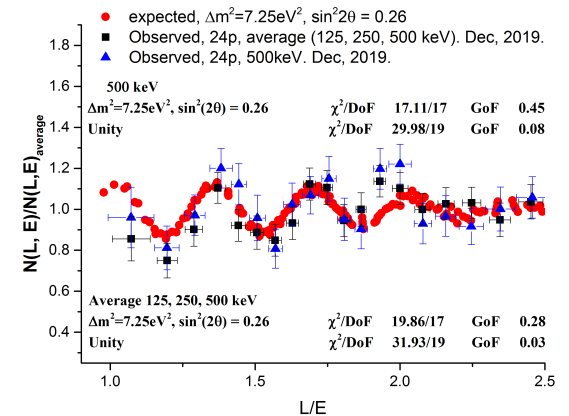
NEOS



Stereo



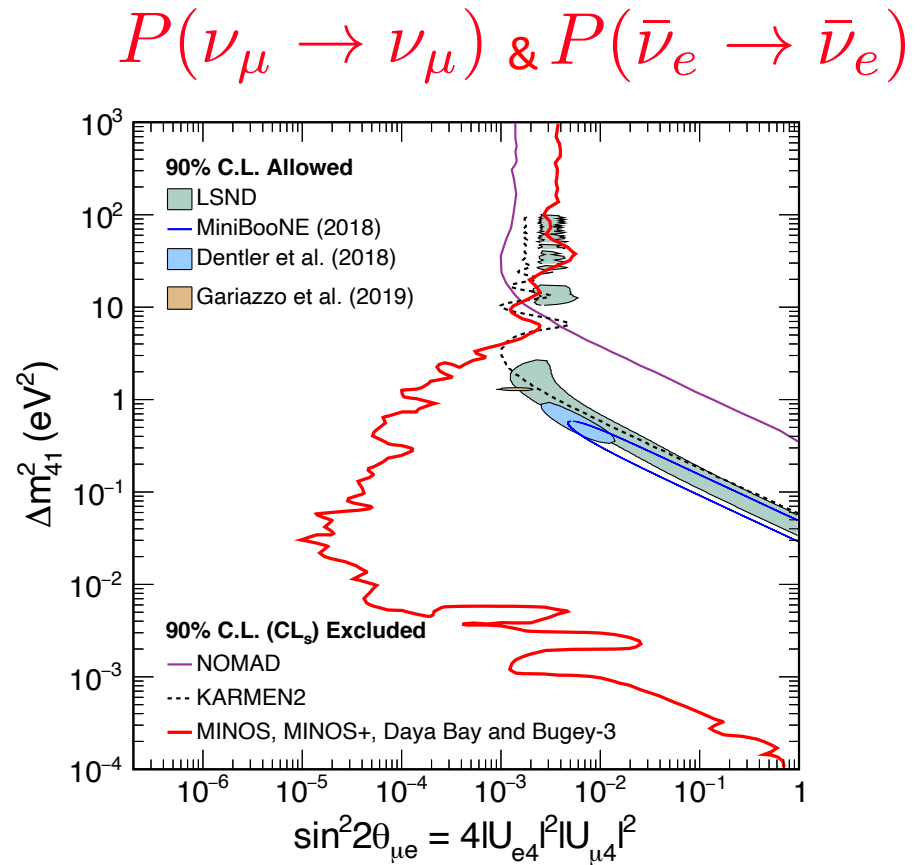
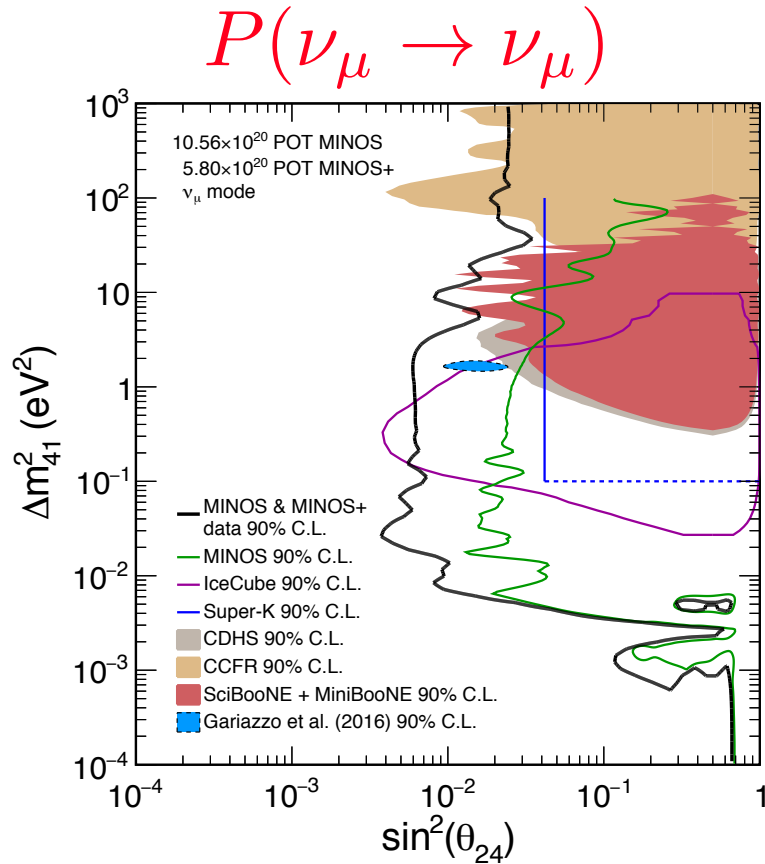
NEUTRINO-4



# O(eV) sterile neutrinos ?

Neutrino muons must disappear also but they don't

Minos, Minos+



O(eV) 4<sup>th</sup> neutrino is not a good fit (all things considered...)



**Exercise:** what about MSW resonances in the 4ν model ?  
Can the sterile oscillation be resonantly enhanced ? Estimate the resonance energy for Earth density and think where to look for this effect.

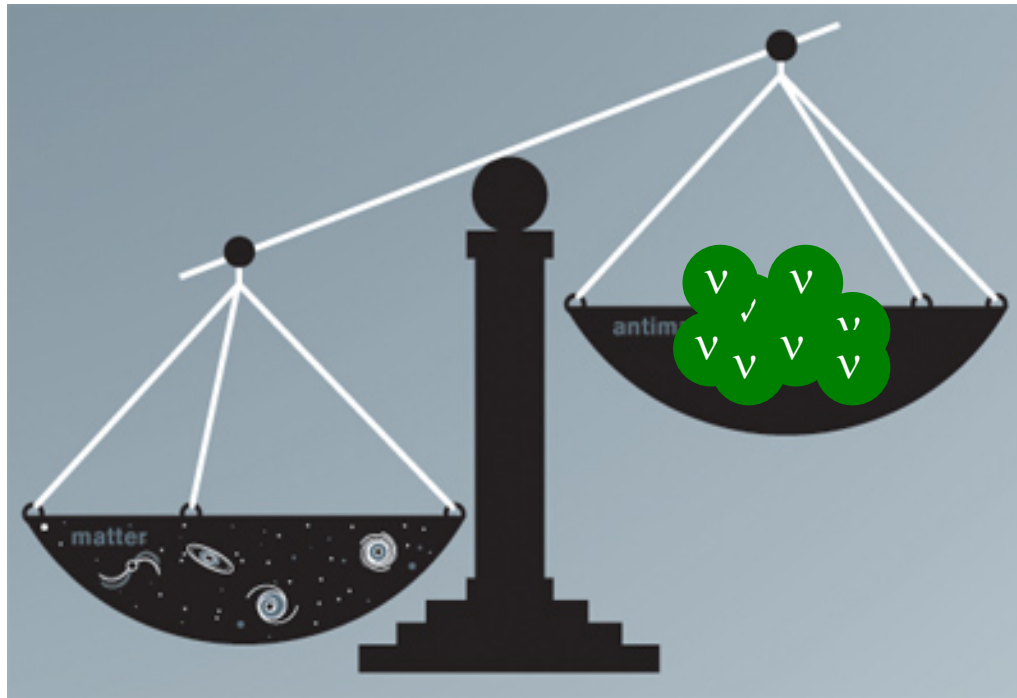
# The other big open questions

**Absolute mass scale:** minimum  $m_\nu$

Are neutrinos **Majorana** and if so, what **new physics** lies behind this fact ?

# Absolute $\nu$ mass scale

Best constraints at present from cosmology

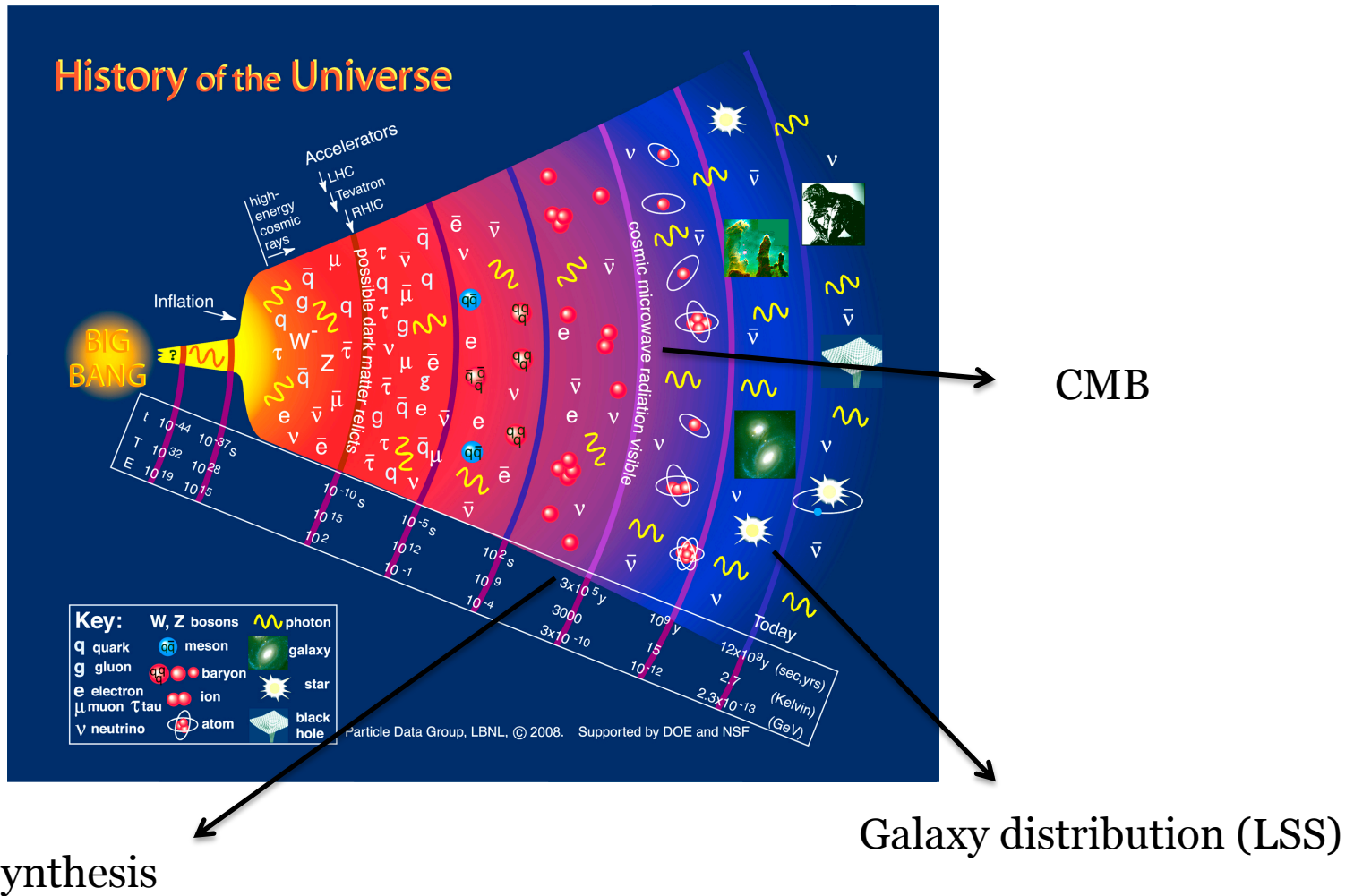


Planck '18

$$\sum m_\nu < 0.12 \text{ eV} \quad (95\%, \text{ Planck TT, TE, EE+lowE} \\ \text{+lensing+BAO}).$$

# Cosmological neutrinos

Neutrinos have left many traces in the history of the Universe



# Neutrinos @ nucleosynthesis (BBN)

Before LEP, the best constraint on  $N_\nu$  came from Big Bang nucleosynthesis!

Neutrinos decouple from the plasma @  $T_\nu \sim 1 \text{ MeV}$

$$\text{Rate} \left( \begin{array}{l} \nu_e n \leftrightarrow p e^- \\ e^+ n \leftrightarrow p \bar{\nu}_e \end{array} \right) \simeq \text{Expansion rate}(g_*(T))$$

$g_*(T) \equiv \#$  of relativistic degrees of freedom

Each neutrino species counts like one relativistic d.o.f.:

$g^*$  depends on  $N_\nu \Rightarrow T_\nu(N_\nu)$

# Neutrinos @nucleosynthesis

At  $@T_\nu \sim 1\text{MeV}$  the ratio neutrons/protons freezes and light elements start to form:

$$\frac{N_n}{N_p} = \exp\left(\frac{m_p - m_n}{T_\nu(N_\nu)}\right)$$

The abundance of light nuclei depends strongly on the ratio of n/p

$$Y_{4\text{He}} = \frac{\text{Mass of } ^4\text{He}}{\text{Total Mass}} = \frac{2N_n}{N_p + N_n}$$

# Neutrinos as DM

Neutrino distribution gets frozen at BBN when they are still relativistic

$$N_\nu \simeq N_{\bar{\nu}} \simeq \frac{4}{11} T_\gamma^3$$

Later on they become non-relativistic, but there are many of them

$$\Omega_\nu = \frac{\sum_i m_i}{93.5 \text{ eV}} h^{-2} < \Omega_m \rightarrow \sum_i m_i \leq 11.2 \text{ eV}$$

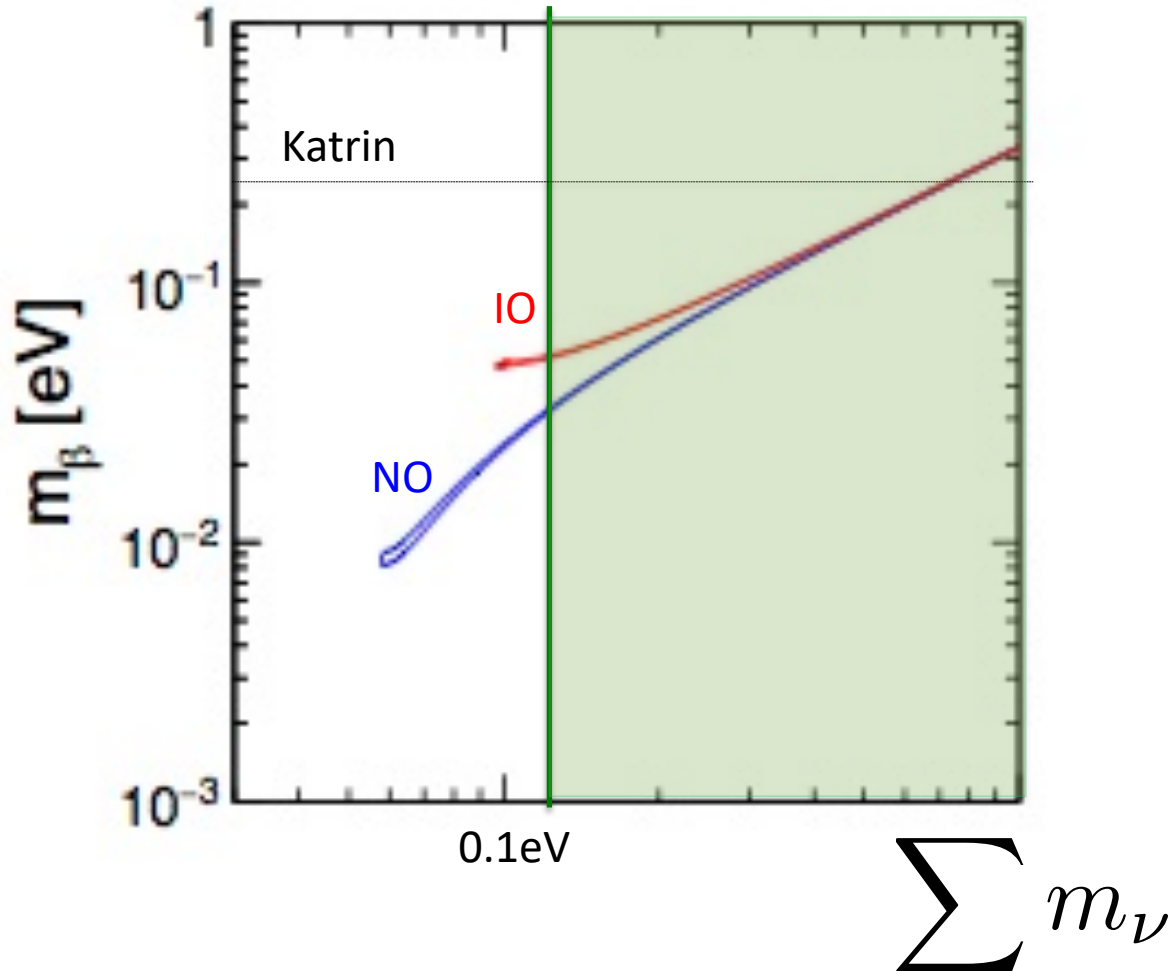
Gershtein, Zeldovich

Massive neutrinos O(eV) contribute significantly to  $\Omega_m$

They tend to produce a Universe with too little structure at small scales:  
**hot DM**

# Absolute $\nu$ mass scale

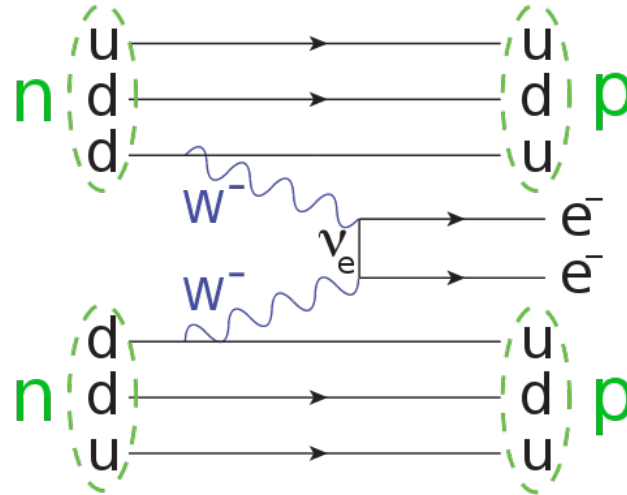
Neutrinos as light as 0.1-1eV modify the large scale structure and CMB





# Majorana nature: $\beta\beta 0\nu$

Plethora of experiments with different techniques/systematics: **EXO**, **KAMLAND-ZEN**, **GERDA**, **CUORE**, **NEXT** ...

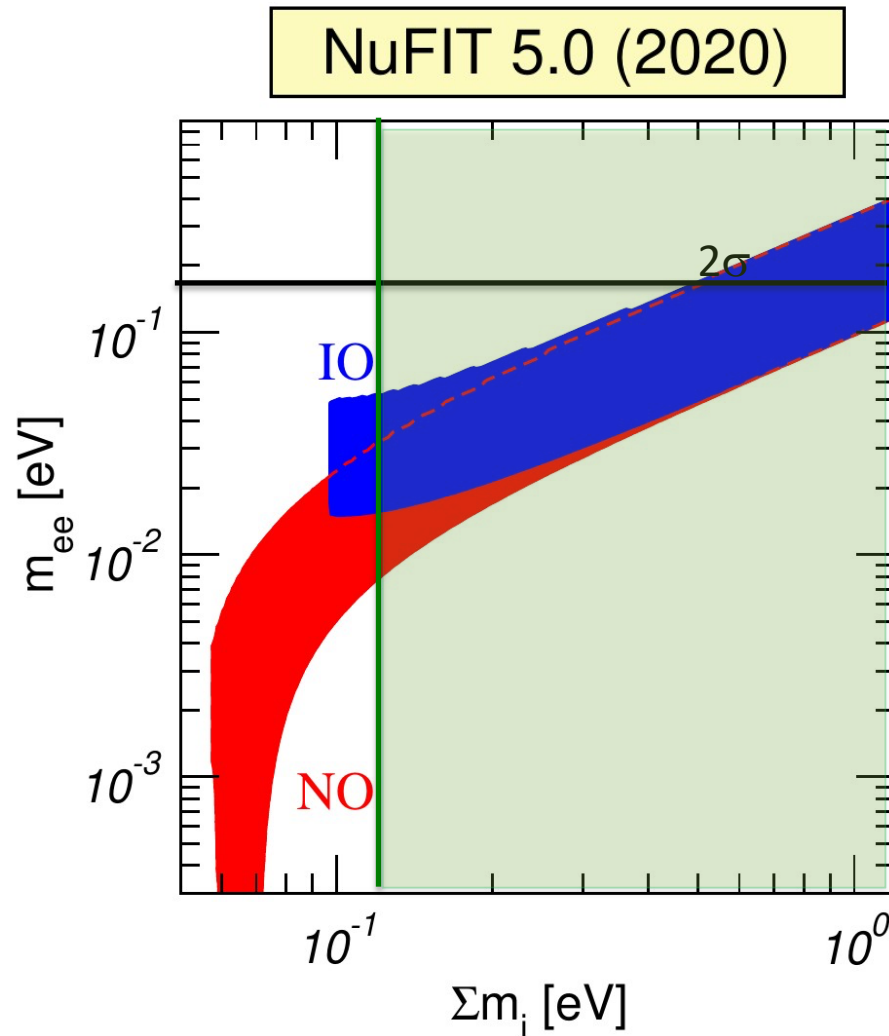


$$T_{2\beta 0\nu}^{-1} \simeq \underbrace{G^{0\nu}}_{\text{Phase}} \underbrace{|M^{0\nu}|^2}_{\text{Nuclear M.E.}} \underbrace{\left| \sum_i (V_{MNS}^{ei})^2 m_i \right|^2}_{|m_{ee}|^2}$$

$$|m_{ee}| = |c_{13}^2 (m_1 c_{12}^2 + m_2 e^{i\alpha} s_{12}^2) + m_3 e^{i\beta} s_{13}^2|$$

# Majorana nature: $\beta\beta 0\nu$

Plethora of experiments with different techniques/systematics: **EXO**, **KAMLAND-ZEN**, **GERDA**, **CUORE**, **NEXT** ...



Next generation of experiments @Ton scale to cover the IO region

# Massive neutrinos: a new flavour perspective

Why do they mix so differently ?

## CKM

$$V_{\text{CKM}} = \begin{pmatrix} 0.97446 \pm 0.00010 & 0.22452 \pm 0.00044 & 0.00365 \pm 0.00012 \\ 0.22438 \pm 0.00044 & 0.97359^{+0.00010}_{-0.00011} & 0.04214 \pm 0.00076 \\ 0.00896^{+0.00024}_{-0.00023} & 0.04133 \pm 0.00074 & 0.999105 \pm 0.000032 \end{pmatrix}$$

$$J = (3.18 \pm 0.15) \times 10^{-5}$$

## PMNS

NuFIT 5.0 (2020)

$$|U|_{3\sigma}^{\text{w/o SK-atm}} = \begin{pmatrix} 0.801 \rightarrow 0.845 & 0.513 \rightarrow 0.579 & 0.143 \rightarrow 0.156 \\ 0.233 \rightarrow 0.507 & 0.461 \rightarrow 0.694 & 0.631 \rightarrow 0.778 \\ 0.261 \rightarrow 0.526 & 0.471 \rightarrow 0.701 & 0.611 \rightarrow 0.761 \end{pmatrix}$$

$$J \simeq 0.033 \sin \delta$$

# Why so different mixing ?

CKM

$$V_{CKM} \simeq \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

PMNS

$$|V_{PMNS}| \simeq \begin{pmatrix} \sqrt{\frac{2}{3}} & \sqrt{\frac{1}{3}} & 0 \\ \sqrt{\frac{1}{6}} & \sqrt{\frac{1}{3}} & \sqrt{\frac{1}{2}} \\ \sqrt{\frac{1}{6}} & \sqrt{\frac{1}{3}} & \sqrt{\frac{1}{2}} \end{pmatrix}$$

# Where the large mixing comes from ?



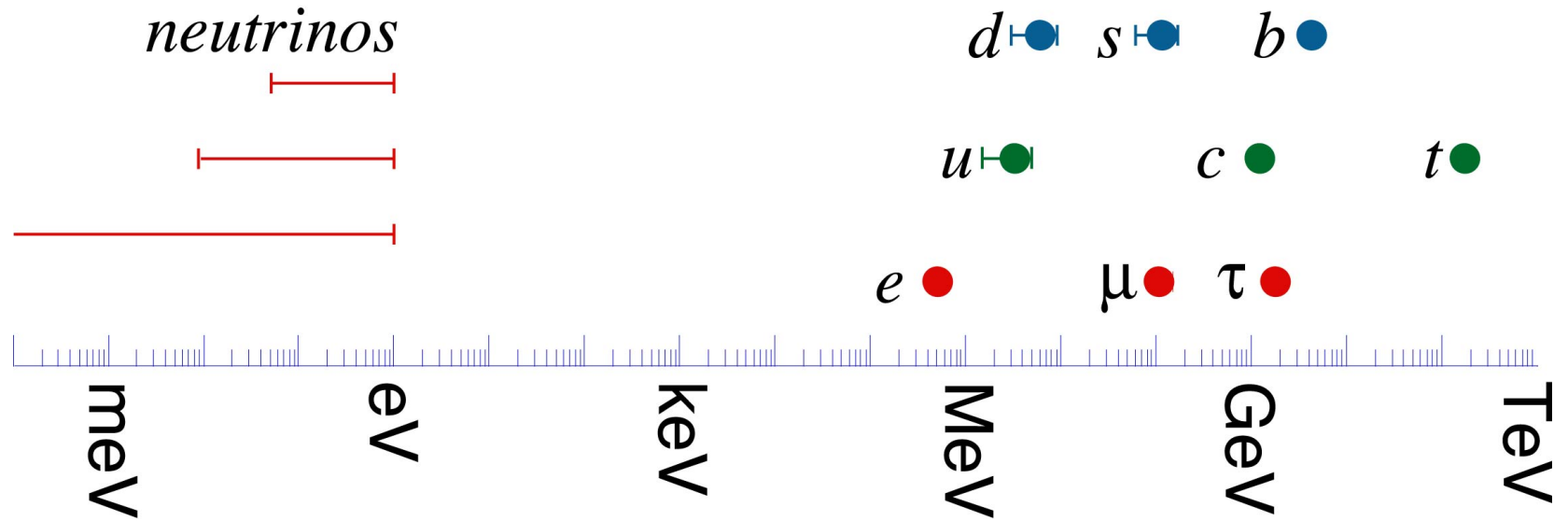
Anarchy for leptons

Discrete or continuous symmetries

Lepton-quark flavour connection in GUTs ?

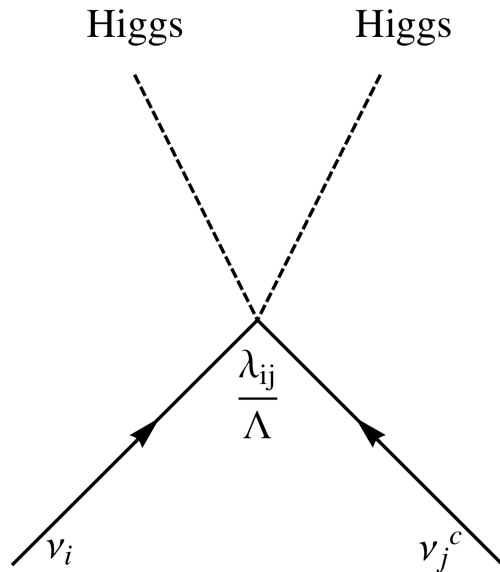
# Massive neutrinos: a new flavour perspective

Why are neutrinos so much lighter ?



They get their masses differently!

Neutrinos have tiny masses -> a new physics scale, what ?



$\nu$ SM ?

$$m_\nu = \lambda \frac{v^2}{\Lambda}$$



Scale at which new particles will show up

# What originates the neutrino mass ?

Could be  $\Lambda \gg v \dots$  the standard lore (theoretical prejudice ?)

$$\left. \begin{array}{l} \Lambda = M_{\text{GUT}} \\ \lambda \sim \mathcal{O}(1) \end{array} \right\} m_\nu \checkmark$$

Hierarchy problem

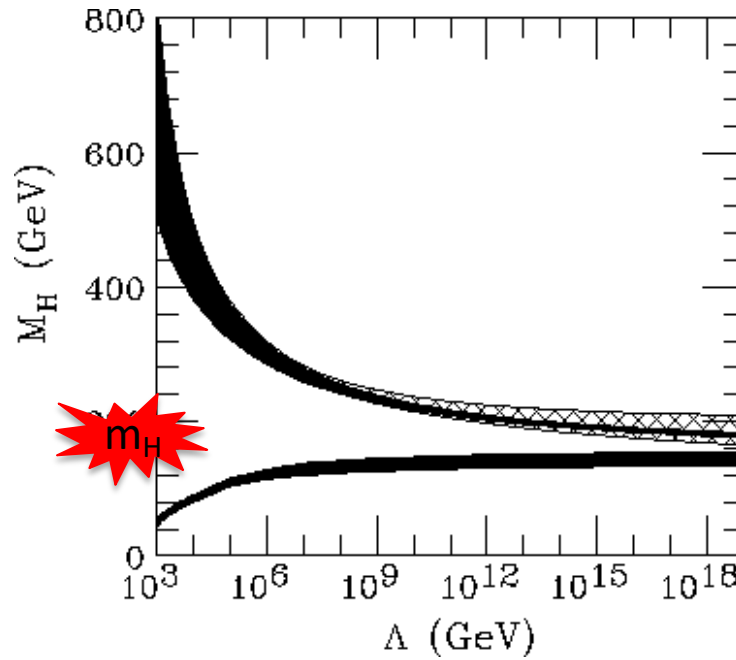
$$m_H^2 \propto \Lambda^2$$

Vissani

not natural in the absence of SUSY/other solution to the hierarchy problem



The Standard Model is healthy as far as we can see...



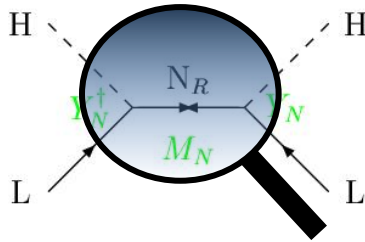
Could be naturally  $\Lambda \sim v$  ?

Yes !  $\lambda$  in front of neutrino mass operator must be small...

# Resolving the neutrino mass operator at tree level

E. Ma

Type I see-saw:  
a heavy singlet scalar

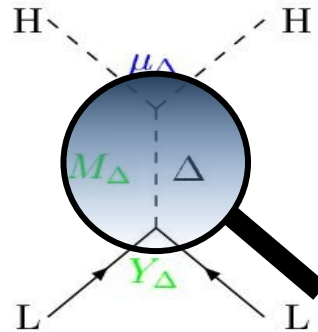


$$m_\nu = \frac{\lambda v^2}{\Lambda} \equiv Y_N^T \frac{v^2}{M_N} Y_N$$

Minkowski;  
Yanagida; Glashow;  
Gell-Mann, Ramond Slansky;  
Mohapatra, Senjanovic...

$$\lambda \sim O(Y^2)$$

Type II see-saw:  
a heavy triplet scalar

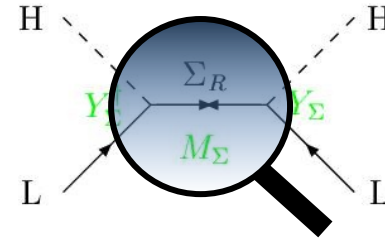


$$m_\nu = \frac{\lambda v^2}{\Lambda} \equiv Y_\Delta \frac{\mu_\Delta}{M_\Delta^2} v^2$$

Konetschny, Kummer;  
Cheng, Li;  
Lazarides, Shafi, Wetterich ...

$$\lambda \sim O(Y \mu/M_\Delta)$$

Type III see-saw:  
a heavy triplet fermion

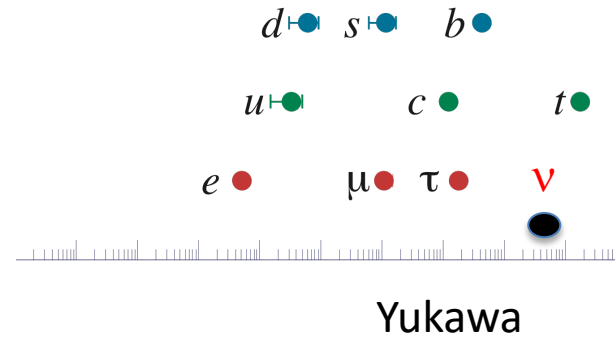


$$m_\nu = \frac{\lambda v^2}{\Lambda} \equiv Y_\Sigma^T \frac{v^2}{M_\Sigma} Y_\Sigma$$

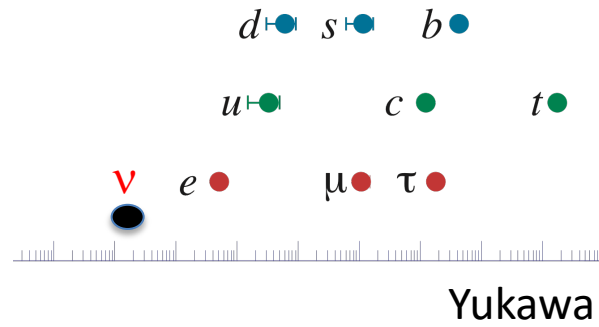
Foot et al; Ma;  
Bajc, Senjanovic...

$$\lambda \sim O(Y^2)$$

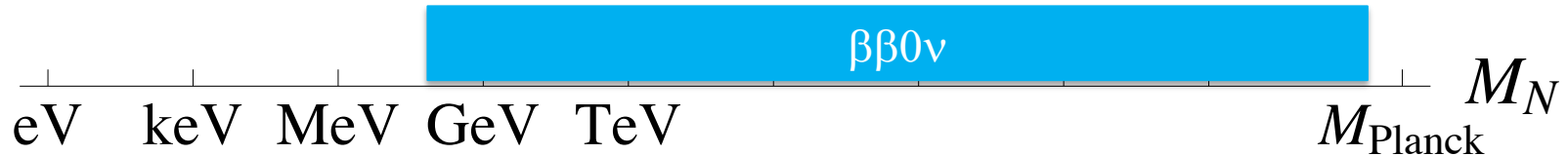
$$M_N \sim \text{GUT}$$



$$M_N \sim \nu$$



# Where is the new scale ?



## Generic predictions

- there is **neutrinoless double beta** decay at some level ( $\Lambda > 100\text{MeV}$ )

model independent contribution from the neutrino mass



# Where is the new scale ?

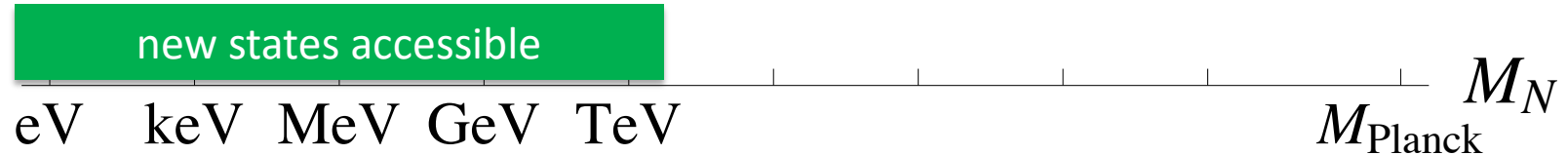


Generic predictions:

- a **matter-antimatter asymmetry** if there is **CP violation** in the lepton sector via **leptogenesis**

model dependent... 

# Where is the new scale ?



Generic predictions:

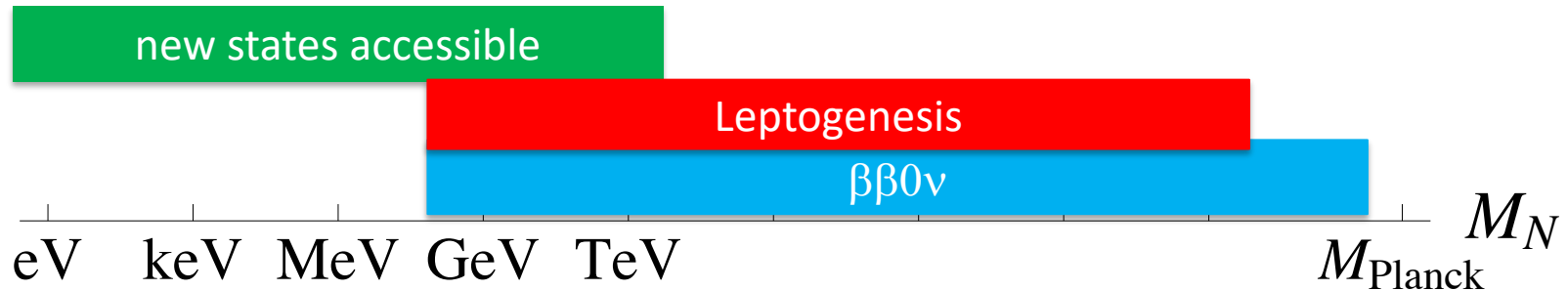
- there are other states out there at scale  $\Lambda$ : **new physics beyond neutrino masses**

potential impact in cosmology, EW precision tests, collider, rare searches,  $\beta\beta 0\nu$ , ...

model dependent...



# Where is the new scale ?

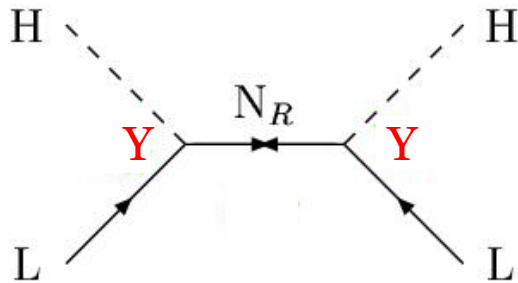


The EW scale is an interesting region: **new physics underlying the matter-antimatter asymmetry could be predicted & tested !**

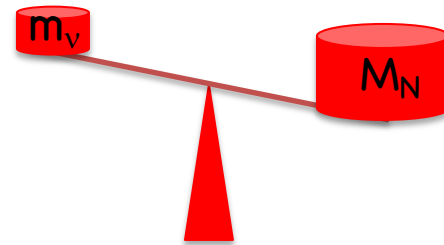
# Minimal model of neutrino masses:

Type I seesaw: SM+right-handed neutrinos

$$\mathcal{L}_\nu = -\bar{l}Y\tilde{\Phi}N_R - \frac{1}{2}\bar{N}_R M N_R + h.c.$$



$$n_R \geq 2$$



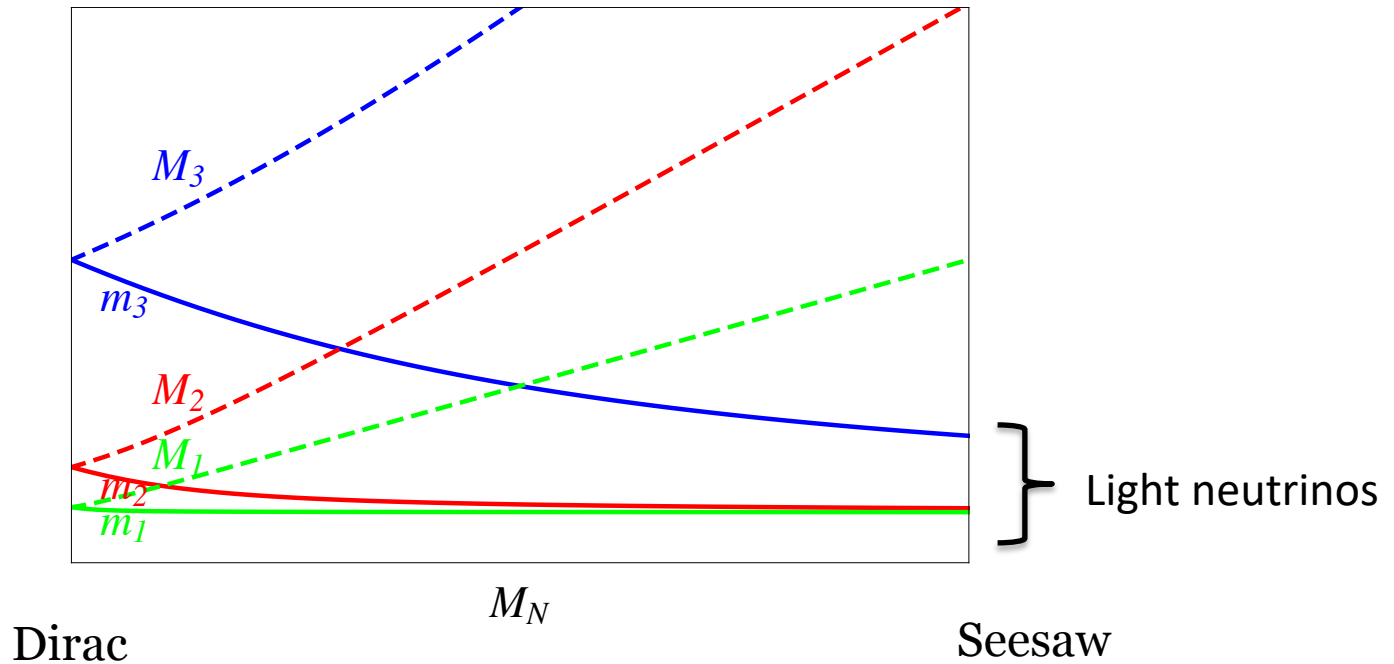
$$m_\nu = \lambda \frac{v^2}{\Lambda} \equiv Y^T \frac{v^2}{M} Y$$

Minkowski; Yanagida; Glashow; Gell-Mann, Ramond Slansky; Mohapatra, Senjanovic...



# Type I seesaw models

$n_R = 3$  : 18 free parameters (6 masses+6 angles+6 phases)  
out of which we have measured 2 masses and 3 angles...

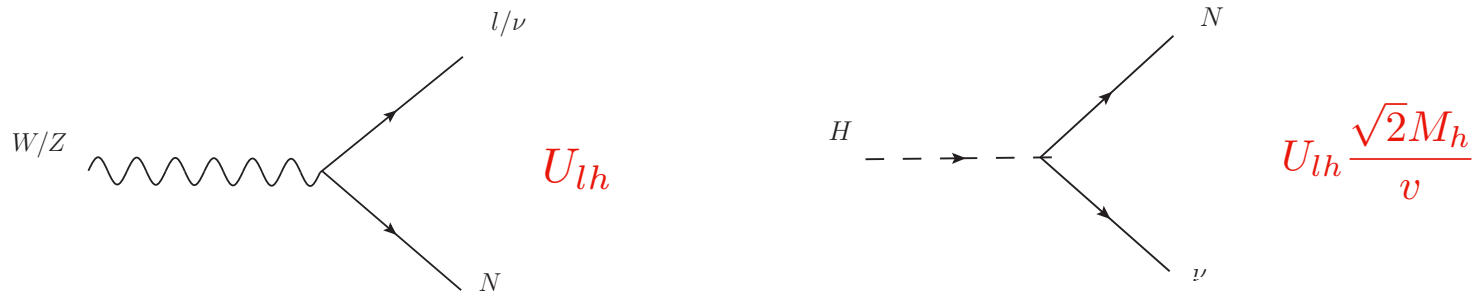


# Type I seesaw models

Phenomenology (beyond neutrino masses) of these models depends on the heavy spectrum and the size of active-heavy mixing:

$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} = U_{ll} \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix} + U_{lh} \begin{pmatrix} N_1 \\ N_2 \\ N_3 \end{pmatrix}$$

# Type I seesaw models



$$U_{lh} \simeq \underbrace{iU_{\text{PMNS}}\sqrt{m_l}}_{\text{light param}} \underbrace{R \frac{1}{\sqrt{M_h}}}_{\text{heavy param}}$$

Casas-Ibarra

**R**: general orthogonal complex matrix (contains all the parameters we cannot measure in neutrino experiments)

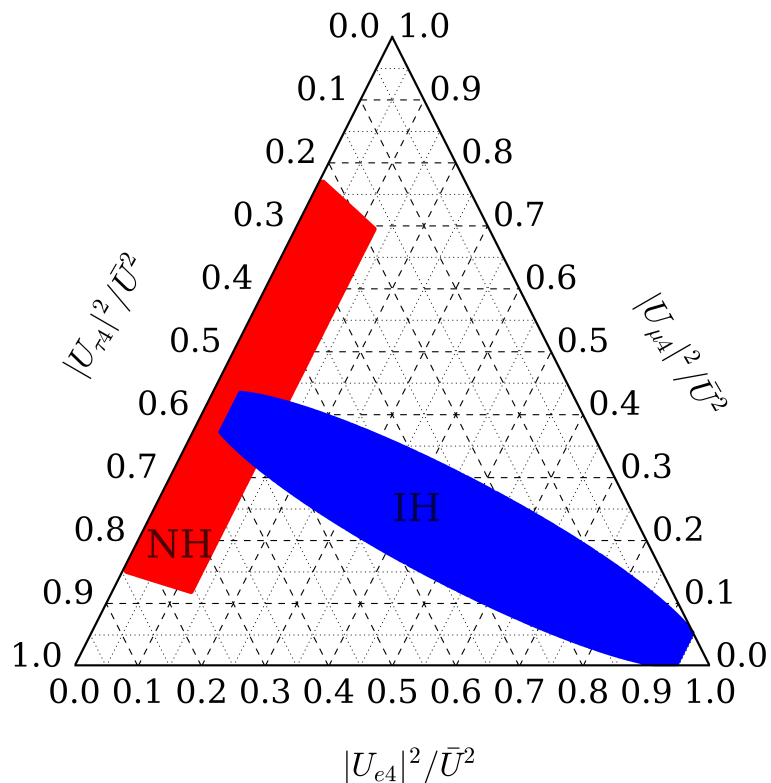
Strong correlation between active-heavy mixing and neutrino masses:

$$|U_{lh}|^2 \sim \frac{m_l}{M_N} \quad (\text{but naive scaling too naive for } n_R > 1 \dots)$$

# Seesaw correlations:

flavour ratios of heavy lepton mixings strongly correlated with ordering,  $U_{\text{PMNS}}$  matrix:  $\delta, \phi_1$

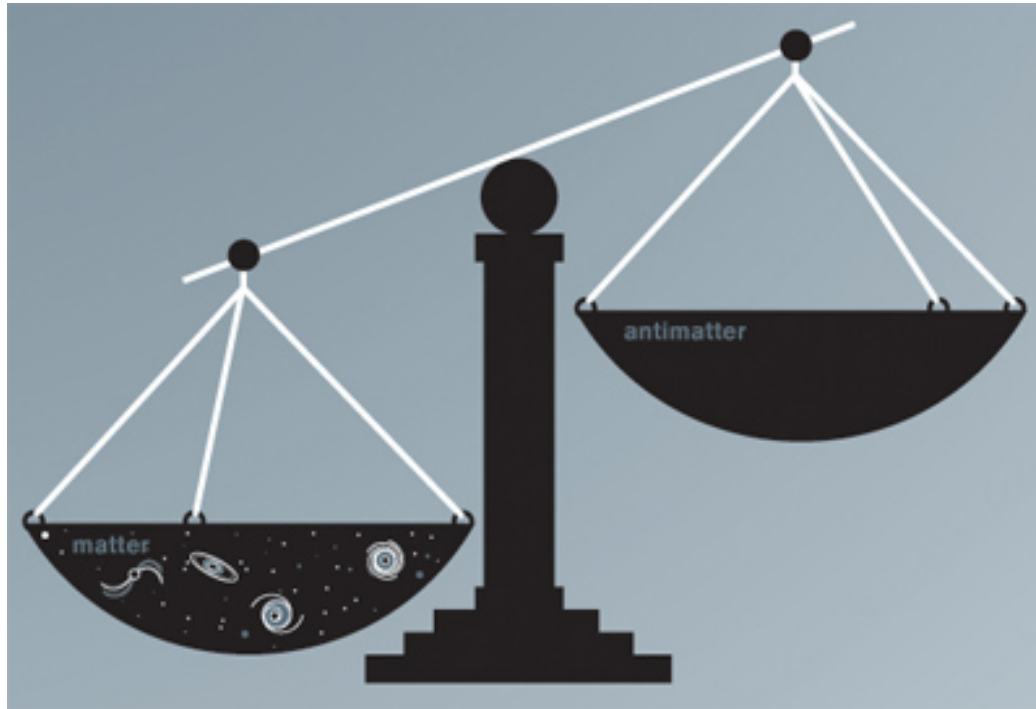
$n_R=2$ :



Caputo, PH, Lopez-Pavon, Salvado arxiv:1704.08721

# Baryon asymmetry

The Universe seems to be made of matter



$$\eta \equiv \frac{n_B - n_{\bar{B}}}{n_\gamma} = 6.21(16) \times 10^{-10}$$

# Baryon asymmetry

Can it arise from a symmetric initial condition with same matter & antimatter ?

## Sakharov's necessary conditions for baryogenesis

- ✓ Baryon number violation (B+L violated in the Standard Model)
- ✓ C and CP violation (both violated in the SM)
- ✓ Deviation from thermal equilibrium (at least once: electroweak phase transition)

It does not seem to work in the SM with massless neutrinos ...

CP violation in quark sector far too small, EW phase transition too weak...

# Leptogenesis

Models with massive neutrinos generically lead to generation of lepton and therefore baryon asymmetries



I Standard leptogenesis in out-of-equilibrium decay  $M_N > 10^7 \text{ GeV}$  Fukuyita, Yanagida

II Resonant leptogenesis  $M_N > 100 \text{ GeV}$  Pilaftsis...

III Leptogenesis from neutrino oscillations

$0.1 \text{ GeV} < M_N < 100 \text{ GeV}$

Akhmedov, Rubakov, Smirnov;  
Asaka, Shaposhnikov,...

# Sakharov conditions

- ✓ CP violation (up to 6 new CP phases in the lepton sector)

$$Y = U_{\text{PMNS}}^* \sqrt{m_\nu} R \sqrt{M_h} \frac{\sqrt{2}}{v}$$

(**R**: 3 complex angles + **U<sub>PMNS</sub>**: 3 phases)

- ✓ B+L violation from **sphalerons**  $T > T_{\text{EW}}$

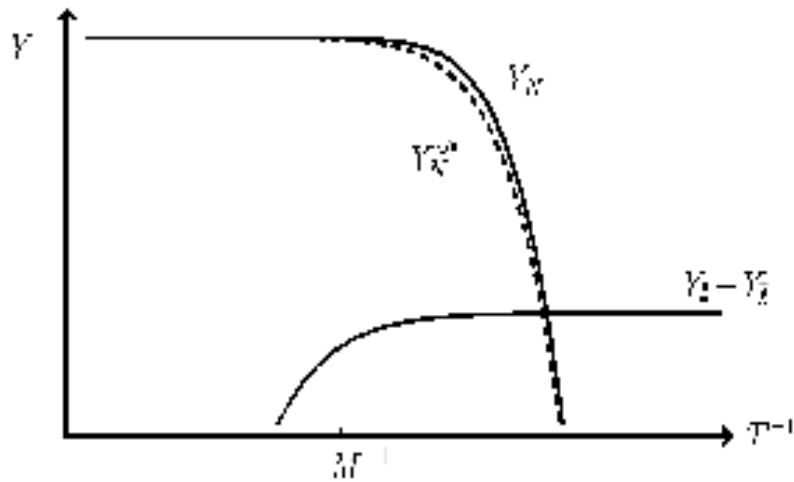
+ L (high-scales)

+  $L\alpha$  (high and low scales)

- ✓ Out of equilibrium condition: different for low and high scales



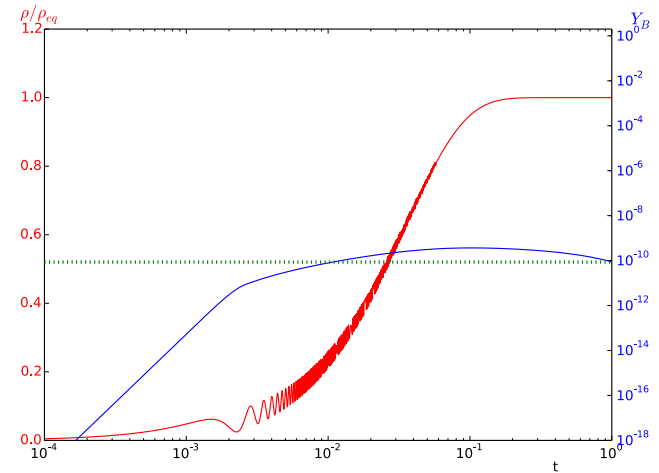
## High-scale leptogenesis (larger Y)



$$\Gamma_N \leq H(M_N)$$

(decay rate < hubble expansion)

## Low-scale leptogenesis (smaller Y)



$T_{EW}$

$$\Gamma_s(T_{EW}) \leq H(T_{EW})$$

(scattering rate < hubble expansion)

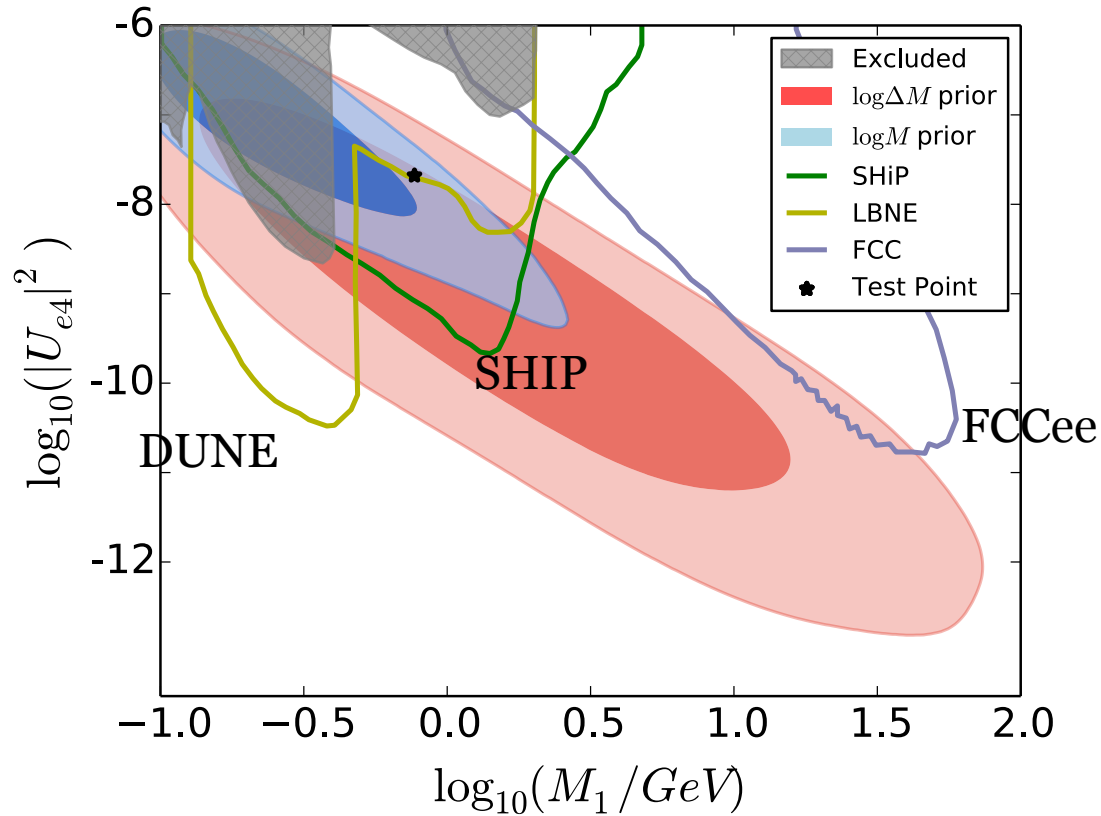
# Testability/predictivity ?

- $Y_B$  cannot be determined from neutrino masses and mixings only
- More information from the heavy sector is needed:

High-scale scenarios: very difficult for  $M_N > 10^7$  GeV

Low-scale scenarios:  $N$ 's can be produced in the lab  
and could be in principle detectable !

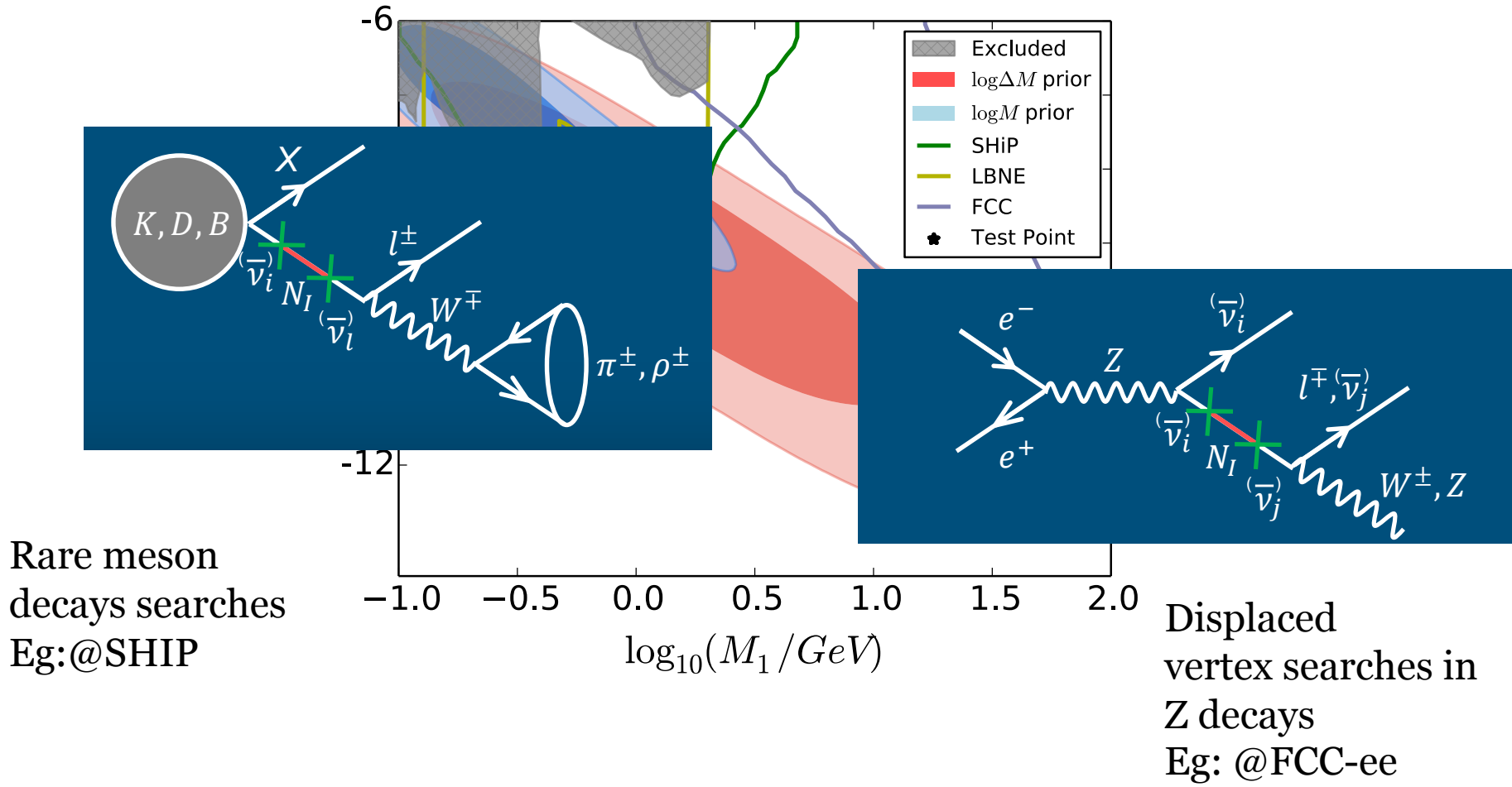
In the minimal model with just  $n_R=2$  neutrinos (IH)



PH, Kekic, Lopez-Pavon, Racker, Salvado

Colored regions: posterior probabilities of successful  $Y_B$  (for not too degenerate states)

# In the minimal model with just $n_R=2$ neutrinos (IH)



# Predicting $Y_B$ in the minimal model $n_R=2$ ?

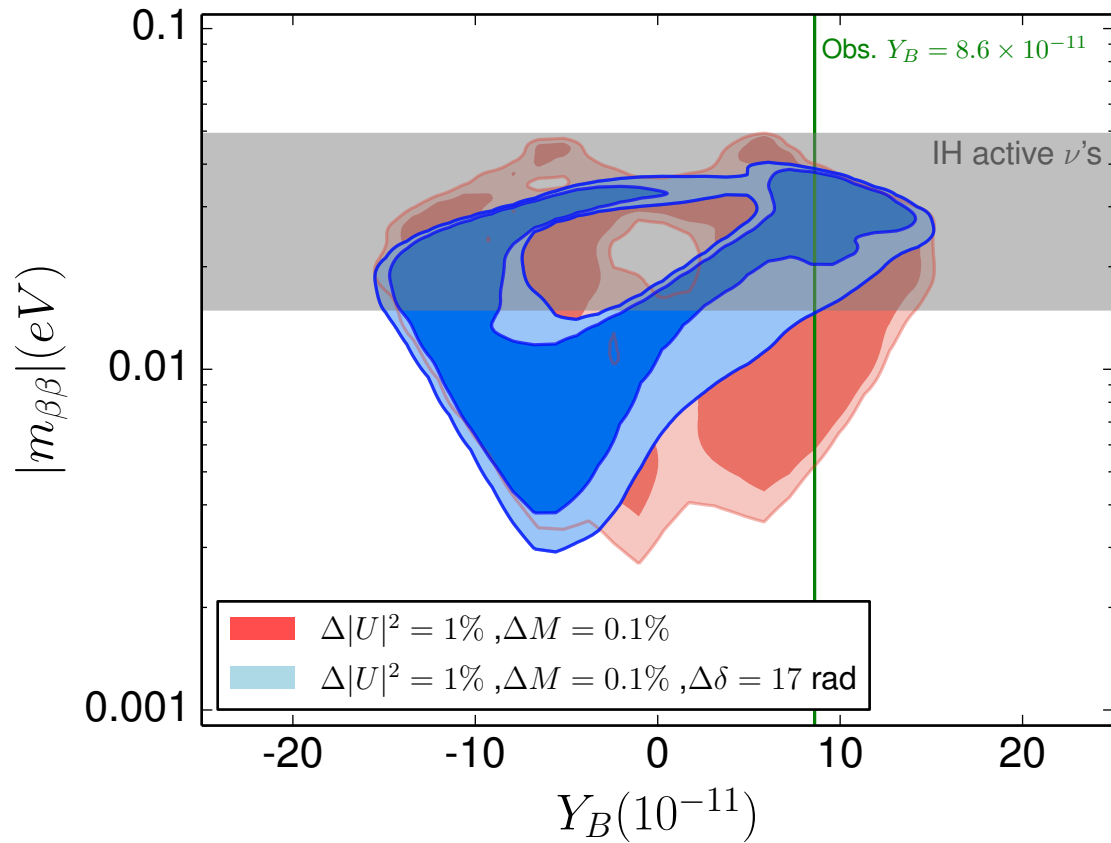
Assume a point within SHIP reach that gives the right baryon asymmetry

- SHIP measurement could provide (if states not too degenerate)

$$M_1, M_2, |U_{e1}|^2, |U_{\mu1}|^2, |U_{e2}|^2, |U_{\mu2}|^2$$

- Future neutrino oscillations:  $\delta$  phase in the  $U_{PMNS}$

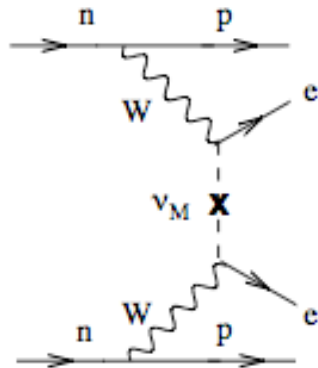
# Predicting $Y_B$ in the minimal model $n_R=2$ (IH)



PH, Kekic, Lopez-Pavon, Racker, Salvado

# Predicting $Y_B$ in the minimal model $n_R=2$

Heavy states also contribute to the  $\beta\beta_{0\nu}$  amplitude...



$$m_{\beta\beta} = \underbrace{\sum_{i=1}^3 [(U_{PMNS})_{ei}]^2 m_i}_{\text{Light states}} + \underbrace{\sum_{i=j}^3 U_{ej}^2 M_j \frac{\mathcal{M}^{0\nu\beta\beta}(M_j)}{\mathcal{M}^{0\nu\beta\beta}(0)}}_{\text{Heavy states}}$$

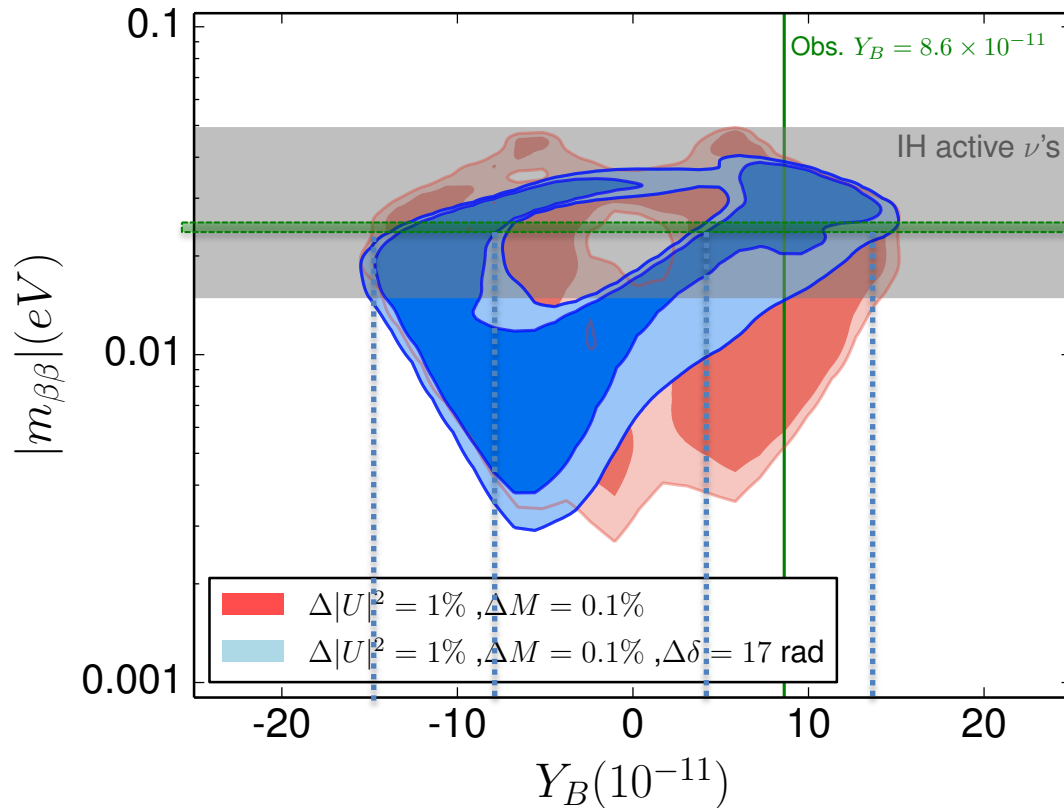
$$M_j \rightarrow \infty \quad \frac{\mathcal{M}^{0\nu\beta\beta}(M_j)}{\mathcal{M}^{0\nu\beta\beta}(0)} \propto \left( \frac{100 \text{ MeV}}{M_j} \right)^2$$

the heavy contribution is sizeable for  $M_i$  of O(GeV)

Blennow, Fernandez-Martinez, Lopez-Pavon, Menendez;  
Lopez-Pavon, Pascoli, Wong; Lopez-Pavon, Molinaro, Petcov

The non standard contributions bring essential information of some CP phases and other unknown parameters

# Predicting $Y_B$ in the minimal seesaw model $M \sim \text{GeV}$

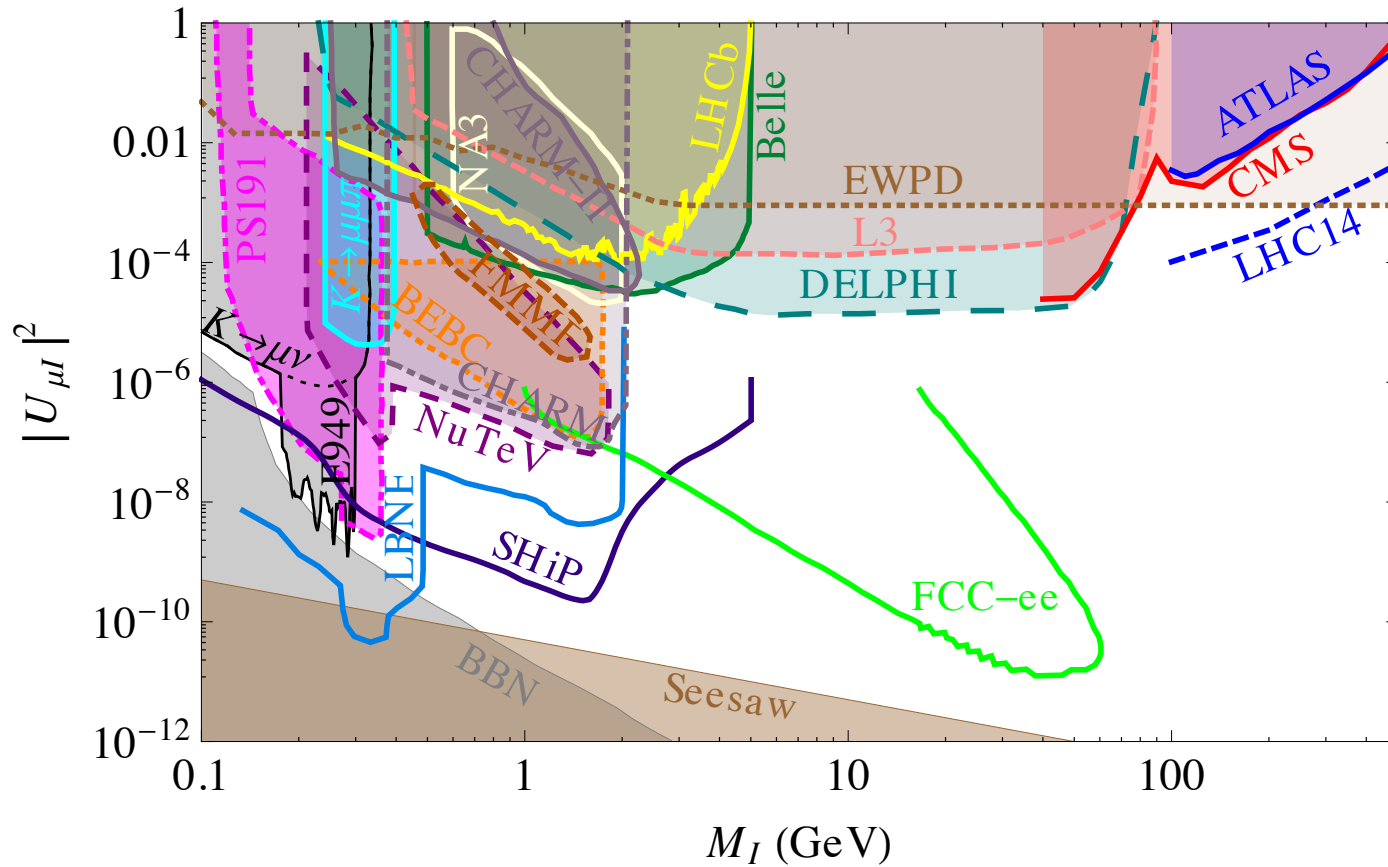


PH, Kekic, Lopez-Pavon, Racker, Salvado  
arxiv:1606.06719

**The GeV-miracle:** the measurement of the mixing to  $e/\mu$  of the sterile states, neutrinoless double-beta decay and  $\delta$  in neutrino oscillations have a chance to give a prediction for  $Y_B$



# Looking for neutrino mass mediators in the EW region



Reviews Atre, Han, Pascoli, Zhang; Gorbunov, Shaposhnikov; Ruchayskiy, Ivashko; Deppisch, Dev, Pilaftsis,...

Bounds in the EW only interesting if  $|U_{\alpha i}|^2 \gg \frac{m_\nu}{M_i} \Leftrightarrow R \gg 1$

- In some cases **unnatural**:

eg: **cancellation between tree level and 1 loop contribution to neutrino masses**

Lopez-Pavon, Pascoli, Wang

- But also technically natural textures:

**protected by an approximate global  $U(1)_L$**

Example  $n_R=2$ :  $L(N_1)=+1, L(N_2)=-1$

$$\begin{pmatrix} 0 & Yv & 0 \\ Yv & 0 & M_N \\ 0 & M_N & 0 \end{pmatrix}$$

$$-\mathcal{L}_\nu \supset \bar{N}_1 M N_2^c + Y \bar{L} \tilde{\Phi} N_1 + h.c.$$

Does not induce neutrino masses: Y unbounded by them

# Seesaw models + approx Lepton number

Wyler, Wolfenstein; Mohapatra, Valle; Branco, Grimus, Lavoura, Malinsky, Romao; Kersten, Smirnov; Abada et al; Gavela et al; Dev, Pilaftsis....many others

$$\begin{pmatrix} 0 & Y_1 v & \epsilon Y_2 v \\ Y_1 v & \mu' & M_N \\ \epsilon Y_2 v & M_N & \mu \end{pmatrix}$$

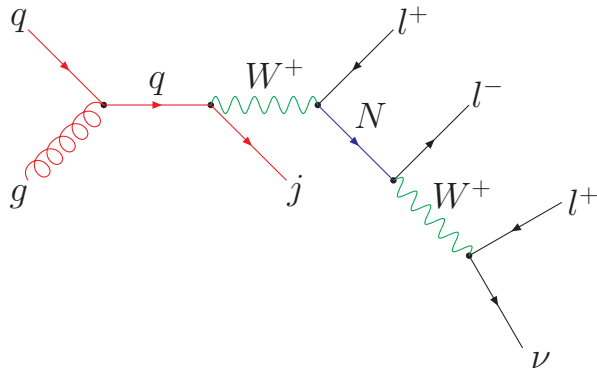
direct seesaw  
inverse seesaw  
extended seesaw

They are all a subclass of type I seesaw models with the generic features:

- quasi-Dirac heavy states
- LNV (neutrino masses, same-sign W decays, etc)  $\sim O(\mu, \mu', \epsilon)$
- Yukawa hierarchies

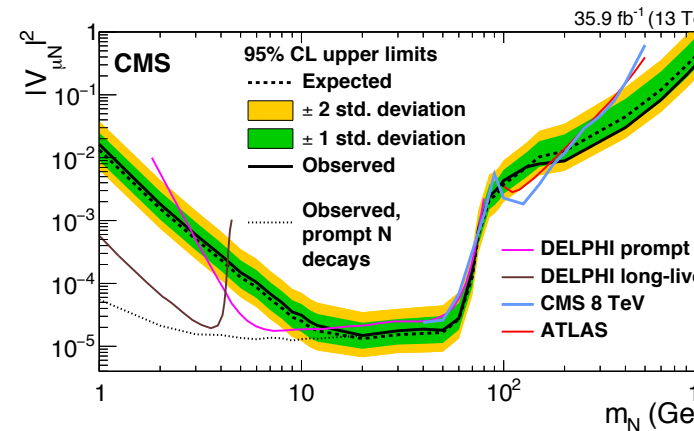
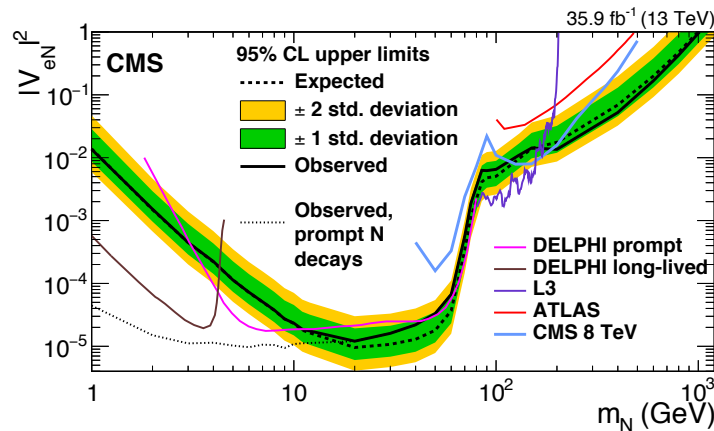
Look for LNC processes ! Can we test their Majorana nature ?

# LNC @LHC: trilepton + missing energy



Del Aguila, Aguilar-Saavedra; ...Chen,Dev;  
Izaaguirre, Shuve...many more

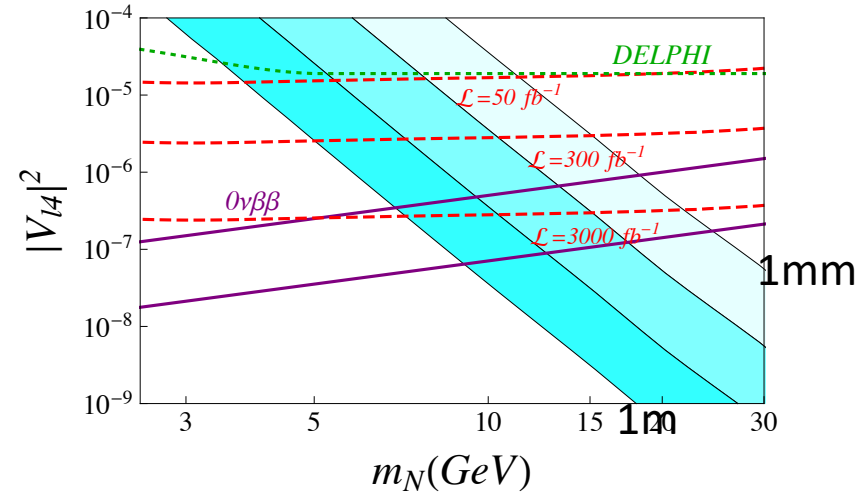
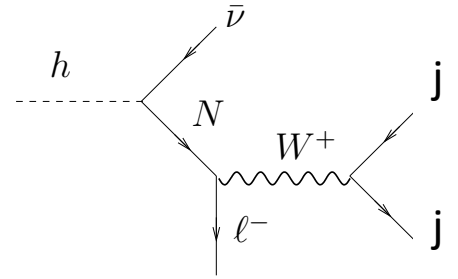
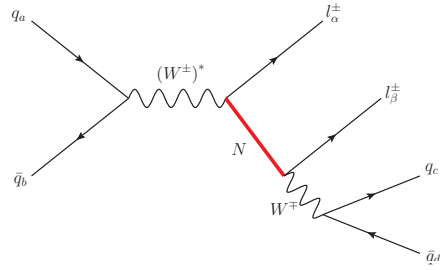
CMS '18



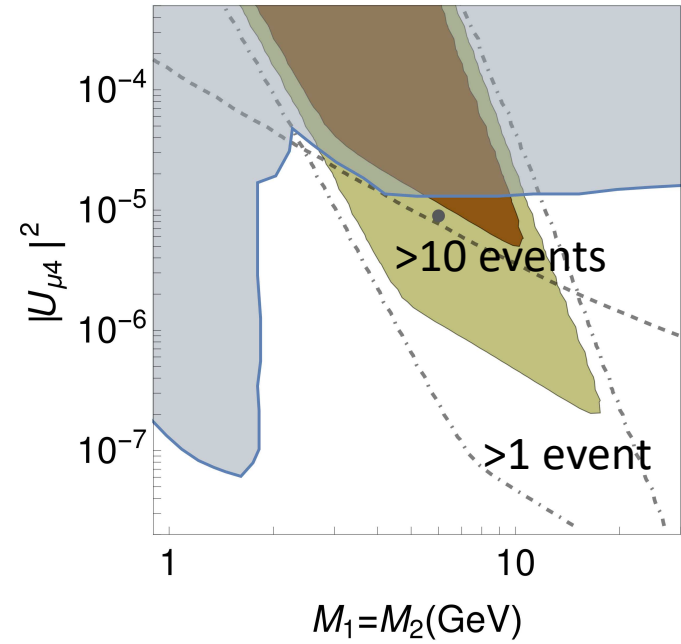
Reaching significantly lower mixings (& lower masses) via displaced decays

Helo, Kovalenko, Hirsch ; Gago, PH, Jonez-Perez, Losada, Moreno;  
Blondel, Graverini,Serra, Shaposhnikov;Antush, Cazzato, Ficher;...

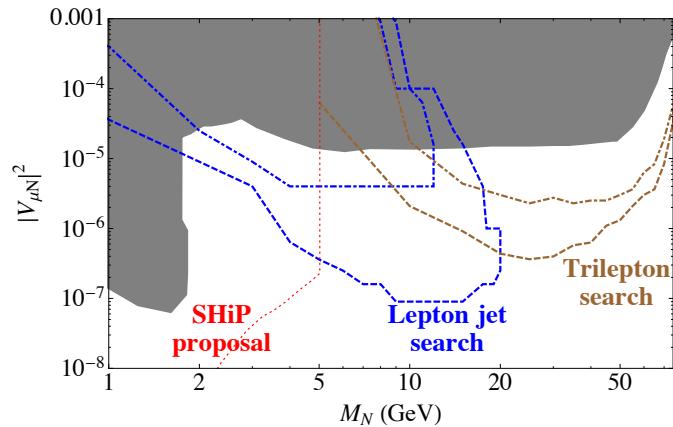
# Golden signal: Displaced Vertices



Helo, Kovalenko, Hirsch

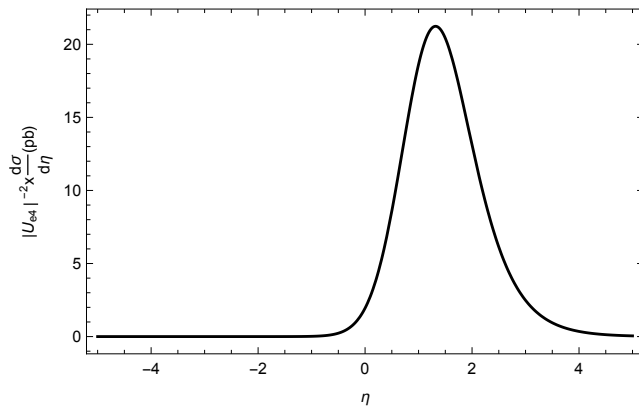
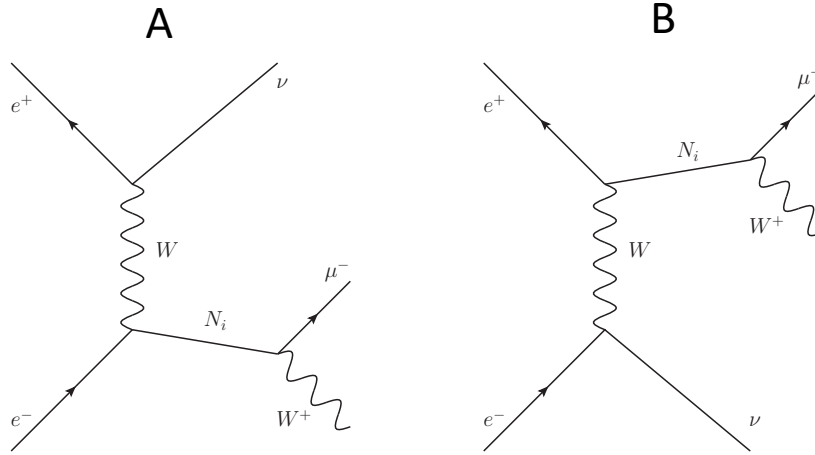


Gago, PH, Jones-Perez, M.Losada, A. Moreno

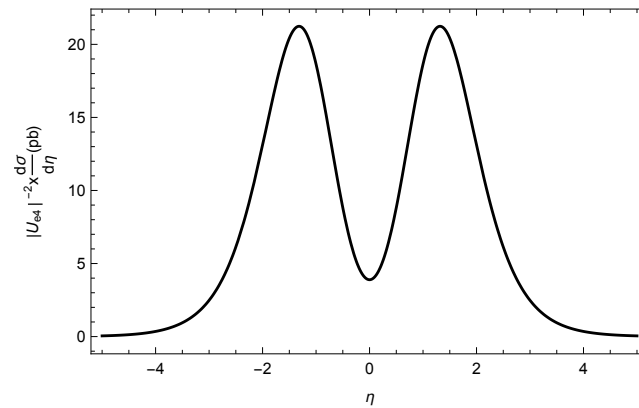


Izaguirre, Shuve

# Majorana vs pseudo-Dirac @ e+e-

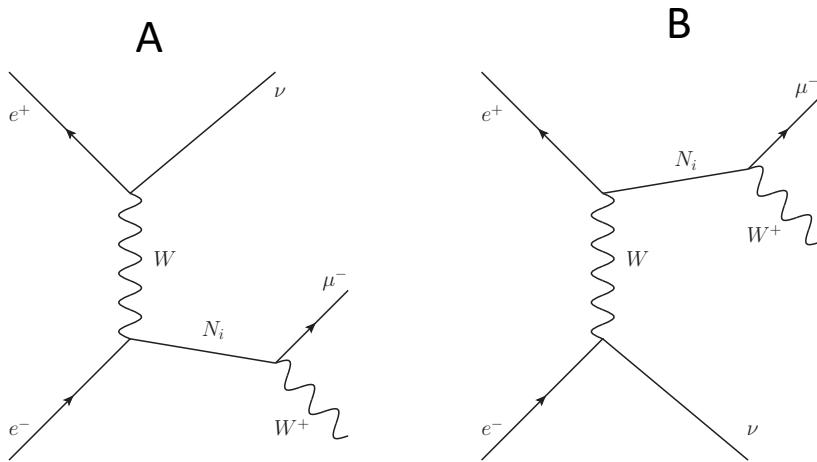


A  $\Leftrightarrow$  LNC

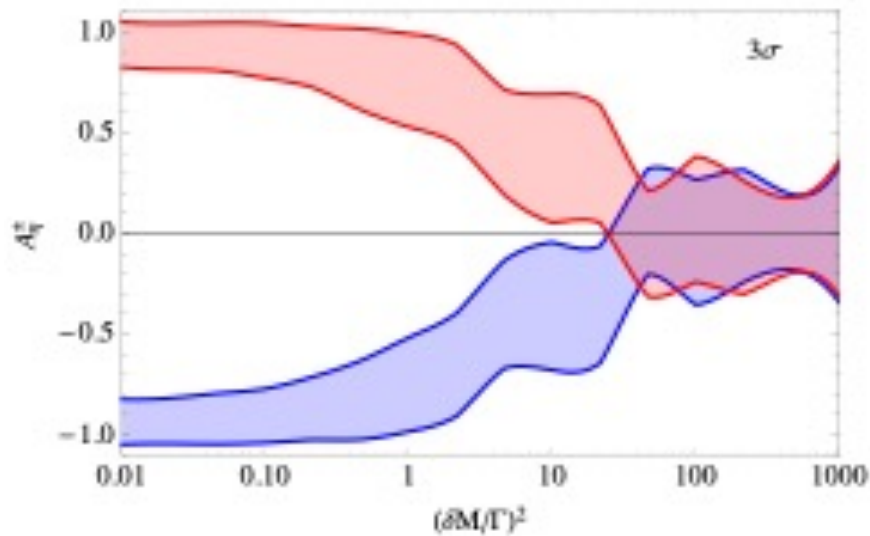


A+B  $\Leftrightarrow$  LNV

# Majorana vs pseudo-Dirac @ e+e-



$$A_\eta^\pm = \frac{N^\pm(\eta > 0) - N^\pm(\eta < 0)}{N_{\text{tot}}^\pm},$$



# Beyond the minimal model

Many possibilities:

Examples: type I + extra  $Z'$ ,  
type II, III  
left-right symmetric models  
GUTs, etc

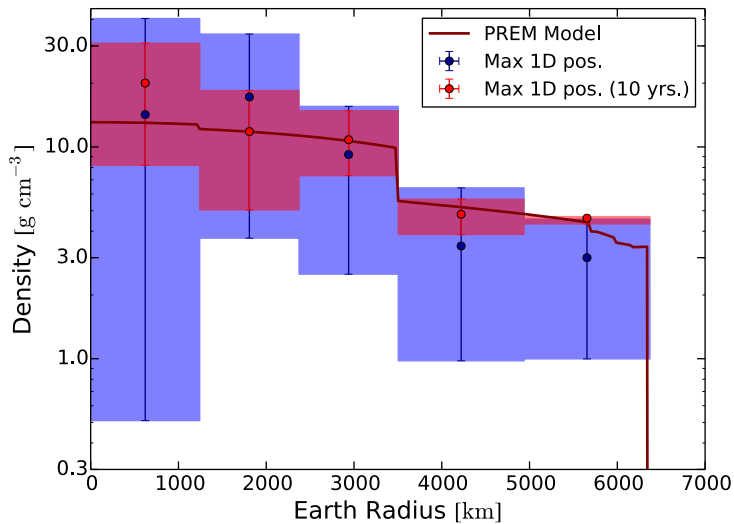
Keung, Senjanovic; Pati, Salam, Mohapatra, Pati; Mohapatra, Senjanovic;  
Ferrari et al + many recent refs...

- Generically new gauge interactions can enhance the production in colliders: richer phenomenology
- But also make leptogenesis more challenging (out-of-equilibrium condition harder to meet)



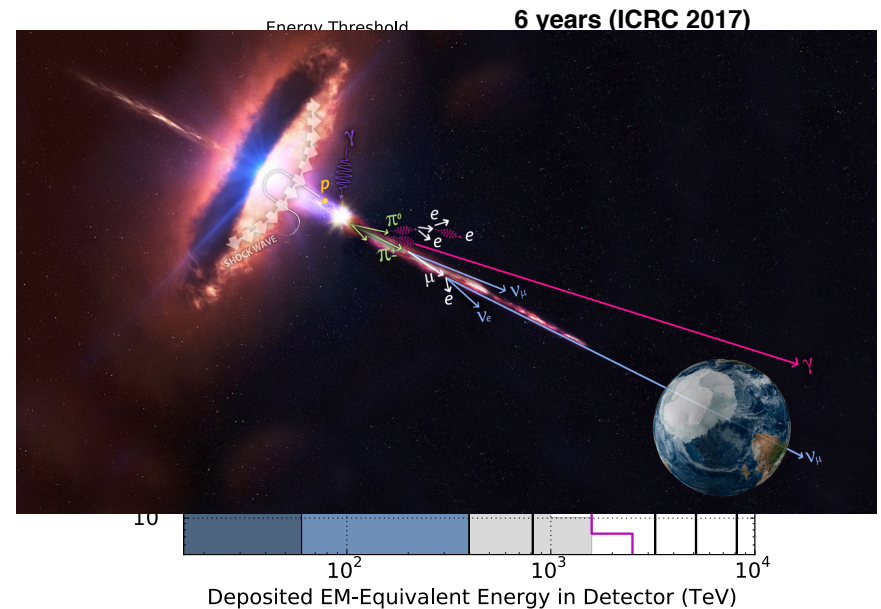
# New era of $\nu$ physics: neutrino astronomy, geology,...

## Understand the Earth



Donini, Palomares-Ruiz, Salvado, 1803.05901

## Understand Astrophysical sources



Icecube '17

Whole new lecture !

# Conclusions

- The results of many beautiful experiments have demonstrated that  $\nu$  are (for the time-being) the less standard of the SM particles
- Many fundamental questions remain to be answered however:  
Majorana nature of neutrinos and scale of new physics? CP violation in the lepton sector? Source of the matter-antimatter asymmetry ?  
Lepton vs quark flavour ?
- A new scale  $\Lambda$  could explain the smallness of neutrino and other mysteries such as the matter-antimatter asymmetry, DM, etc
- Complementarity of different experimental approaches:  $\beta\beta\nu$ , CP violation in neutrino oscillations, direct searches in meson decays, collider searches of displaced vertices, etc...holds in well motivated models with a low scale  $\Lambda$  (GeV scale very interesting)

These tiny pieces of reality have brought many surprises, maybe they will continue with their tradition...

