

Charged Higgs boson decays with NLO corrections in the NMSSM

[Thi Nhung Dao, Margarete Mühlleitner, Shruti Patel, KS, Eur.Phys.J.C 81 (2021) 4, 340]

Kodai Sakurai (Tohoku U.)

Collaborators:

Thi Nhung Dao ^A, Margarete Mühlleitner ^B, Shruti Patel ^B

(^A: ICISE, Vietnam; ^B: KIT, Germany)

Introduction

- Supersymmetry: hierarchy problem, gauge coupling unification, dark matter, etc.
- The current LHC experiment data:
 - No SUSY particles have been found.
 - The discovered Higgs boson (125 GeV) behaves very SM-like.
 - The properties will be precisely measured in the future collider experiments such as HL-LHC, ILC, CEPC, FCC-ee.

Introduction

- Supersymmetry: hierarchy problem, gauge coupling unification, dark matter, etc.
 - The current LHC experiment data:
 - No SUSY particles have been found.
 - The discovered Higgs boson (125 GeV) behaves very SM-like.
 - The properties will be precisely measured in the future collider experiments such as HL-LHC, ILC, CEPC, FCC-ee.
 - This requires that theoretical predictions are evaluated accurately.
 - SUSY loop corrections to observables for the Higgs boson may be significant.
 - Not only the discovered Higgs boson but also extra Higgs boson should be precisely calculated.
 - Direct searches at LHC
 - Discrimination of NP models when extra Higgs is discovered.
- While there are a lot of works for the precise calculations in MSSM, we study next-to minimal supersymmetric model (NMSSM).

Precise calculations of Higgs bosons

Many studies of precise calculations of Higgs bosons have been performed in the NMSSM.

- Higgs boson masses
 - (Full 1-loop) [G. Degrandi, P. Slavich, NPB 825 (2010)], etc.
 - ($O(a_t a_s, a_t^2)$ 2-loop) [M. Mühlleitner, D. T. Nhung, H. Rzehak, K. Walz, JHEP05(2015)]
[T.N. Dao, R. Gröber, M. Krause, M. Mühlleitner, H. Rzehak], etc.
 - ($O((a_t + a_\lambda + a_t)^2)$ 2-loop) [T. N. Dao, M. Gabelmann, M. Mühlleitner, H. Rzehak, 2106.06990]
- Higgs boson decays
 - $H_i, A_i \rightarrow ff, VV$ (Full 1-loop) [F. Domingo, S. Heinemeyer, S. Paßehr, G. Weiglein, EPJC78(2018)]
 - $H_i \rightarrow H_j H_k$ (Full 1-loop) [D. T. Nhung, M. Mühlleitner, J. Streicher, K. Walz, JHEP11 (2013)]
($O(a_t a_s)$ 2-loop) [G. Belanger, V. Bizouard, F. Boudjema, G. Chalons, PRD96 (2017)]
[M. Mühlleitner, D. T. Nhung, H. Ziesche, JHEP 12 (2015)]
 - $A_i \rightarrow \tilde{t}\tilde{t}$ (Full 1-loop) [J. Baglio, C. Krauss, M. Mühlleitner, K. Walz, JHEP 10 (2015)]
 - All 2-body neutral Higgs decays, BRs [J. Baglio, T. N. Dao, M. Mühlleitner, 1907.12060]
 - $H_i^\pm \rightarrow WH_i$ (Full 1-loop) [T. N. Dao, L. Fritz, M. Krause, M. Mühlleitner, S. Patel, Eur.Phys.J. C80 (2020)]

→ Missing part for Higgs boson decays is charged Higgs boson decays.

- What is typical size of NLO corrections for each decay?

Charged Higgs decays in the NMSSM

Charged Higgs decays in the NMSSM can be categorized as follows:

1. Decays into SM particles

$$H^+ \rightarrow q\bar{q}', \quad (t\bar{b}, \quad t\bar{s}, \quad t\bar{d}, \dots)$$

$$H^+ \rightarrow \ell\nu, \quad (\tau^+\nu, \quad \mu^+\nu, \quad e^+\nu, \dots)$$

$$H^+ \rightarrow W^+V \quad (V = Z, \gamma)$$

2. Decays into extra Higgs bosons + W bosons

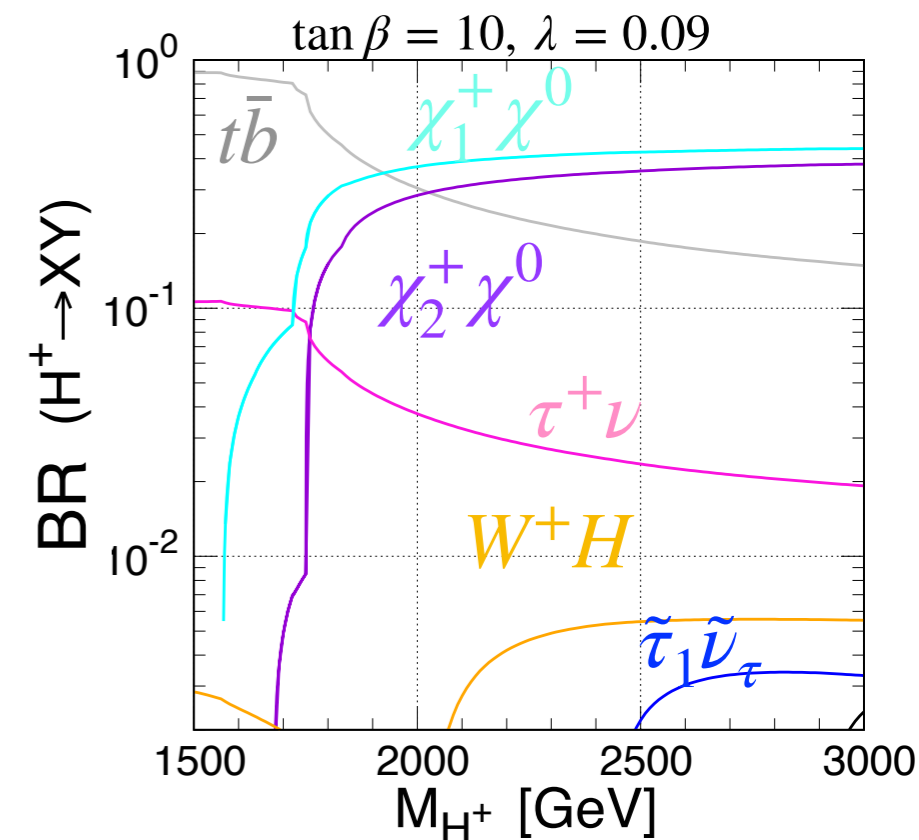
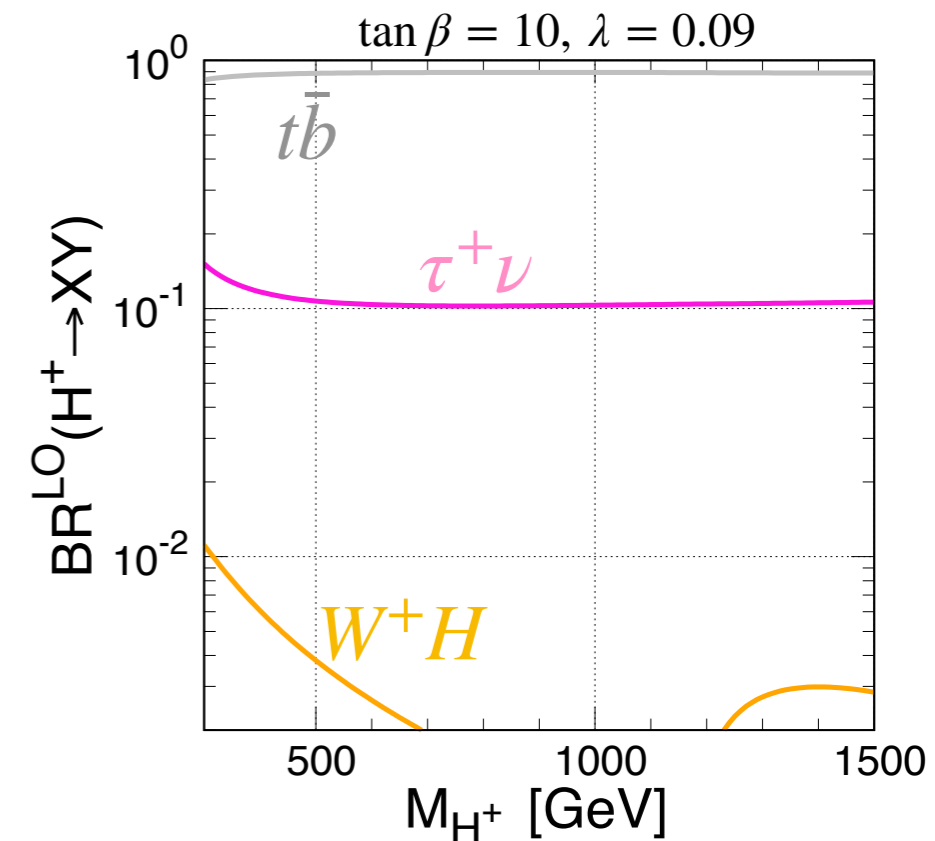
$$H^+ \rightarrow W^+H_i \quad (H_i : \text{Neutral Higgs bosons})$$

3. Decays into SUSY particles

$$H^+ \rightarrow \tilde{\chi}_i^0 \tilde{\chi}_j^+, \quad (i = 1, 2, j = 1, \dots, 5)$$

$$H^+ \rightarrow \tilde{q}\tilde{q}', \quad (\tilde{t}\tilde{b}, \dots)$$

$$H^+ \rightarrow \tilde{\ell}\tilde{\nu}, \quad (\tilde{\tau}^+\tilde{\nu}, \dots)$$



Renormalization

- Higgs potential: mixed OS- $\overline{\text{DR}}$ scheme. [\[K. Ender, T. Graf, M. Muhlleitner, H. Rzehak, PRD85\(2012\)\] etc.](#)

$$V_H = |\lambda|^2 |S|^2 \left(H_u^\dagger H_u + H_d^\dagger H_d \right) + \left| \lambda (H_u i\sigma_2 H_d) + \kappa S^2 \right|^2 + \frac{1}{2} g_2^2 |H_u^\dagger H_d|^2 + \frac{1}{8} (g_1^2 + g_2^2) \left(H_u^\dagger H_u - H_d^\dagger H_d \right)^2$$

$$+ m_{H_u}^2 H_u^\dagger H_u + m_{H_d}^2 H_d^\dagger H_d + m_S^2 |S|^2 + \left[A_\lambda \lambda (H_u i\sigma_2 H_d) S + \frac{1}{3} \kappa A_\kappa S^3 + \text{h.c.} \right]$$

→ Input parameters(12): $T_{H_1}, T_{H_2}, T_{H_3}, M_W, M_Z, \alpha_{\text{em}}, \tan \beta, v_S, m_{H^\pm}, \lambda, \kappa, A_\kappa$

$T_{H_1}, T_{H_2}, T_{H_3}, M_{H^\pm}, M_Z, M_W, \alpha_{\text{em}}$ → On-shell $\hat{T}_{H_i} = 0, \hat{\Sigma}_{VV}(M_V^2) = 0, \hat{\Gamma}_{\bar{f}f}\gamma(0) = 0$

$\tan \beta, v_S, \lambda, \kappa, A_\lambda, A_\kappa,$

(Wave function renormalization → $\overline{\text{DR}}$
for H_d, H_u, S)

- Electroweakino sector: OS scheme or $\overline{\text{DR}}$ scheme [\[J. Baglio, T. N. Dao, M. Muhlleitner, EPJC 80 \(2020\) 10, 960\]](#)

M_1, M_2 : We have 7 OS conditions, i.e., $\sum_{ii} \chi_i^{\pm,0} (m_{\chi_i^{\pm,0}}^2) = 0$

→ OS1 : $\chi_i^+ \sim \tilde{W}^+ : \text{OS}, \chi_k^0 \sim \tilde{B} : \text{OS},$

OS2 : $\chi_i^0 \sim \tilde{W}^3 : \text{OS}, \chi_k^0 \sim \tilde{B} : \text{OS},$

Decay rates of charged Higgs bosons at NLO

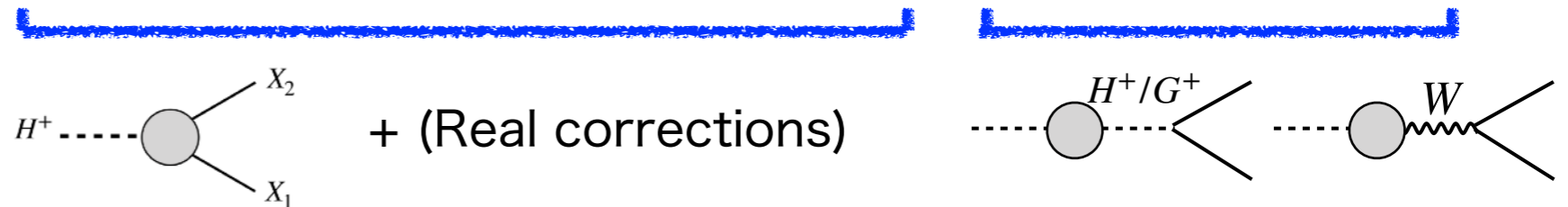
- We evaluated NLO EW (+SUSY EW) and NLO SUSY QCD corrections to the following processes:

$$H^\pm \rightarrow tb, \quad H^\pm \rightarrow \tau\nu, \quad H^\pm \rightarrow \chi_i^+ \chi_j^0, \quad H^+ \rightarrow \tilde{t}\tilde{b}, \quad H^+ \rightarrow \tilde{t}\tilde{\nu}, \quad H^+ \rightarrow WH_i$$

(i=1,2; j=1,...,5)

- Schematic formula for NLO decay rates

$$\Gamma(H^\pm \rightarrow X_1 X_2) = (\text{Resummed factors}) \times \Gamma_{\text{LO}}(H^\pm \rightarrow X_1 X_2) \times \left[1 + \Delta^{\text{QCD}} + \Delta^{\text{SUSYQCD}} + \Delta^{(\text{SUSY+}) \text{EW}} + \Delta^{\text{H}^+ \text{H}^-} + \Delta^{\text{H}^+ \text{G}^- / \text{W}^-} \right]$$



(Resummed factors): Δ_b corrections ($H^+ \rightarrow tb$), Z factor ($H^+ \rightarrow W H_i$)

[M. Carena, D. Garcia, U. Nierste, C. E.M. Wagner, NPB 577(2000)], etc. [K. E. Williams, G. Weiglein, Phys. Lett. B660 (2008)]

- resummed factors, Δ_{QCD} are calculated by NMSSMCALC(EW).

[J. Baglio, T. N. Dao, M. Muhlleitner, EPJC 80 (2020) 10, 960]

For $H^\pm \rightarrow tb$, it is known that SUSY particles corrections to the bottom Yukawa coupling are significant in large $\tan\beta$ regime.

(Ex.) MSSM

Effective lagrangian for bottom Yukawa in gauge basis:

$$\mathcal{L}_{eff} = y_b h_1^0 b \bar{b} + \Delta y_b h_2^0 b \bar{b}.$$

After h_1^0 and h_2^0 acquire VEV,

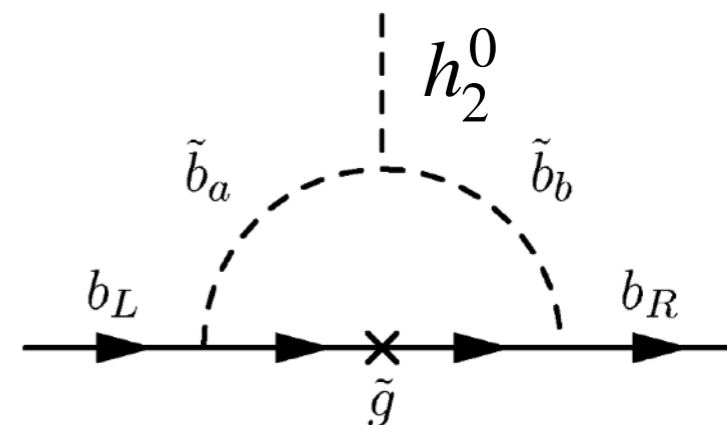
$$\rightarrow m_b = \frac{v_1}{\sqrt{2}}(y_b + \Delta y_b \tan\beta) \equiv \frac{y_b}{\sqrt{2}} v_1 (1 + \Delta m_b)$$

This modifies tree level $\bar{t}bH^+$ vertex relations:

$$\bar{t}P_R b H^+ : \frac{m_b}{v} \tan\beta \rightarrow \frac{m_b}{v} \tan\beta \frac{1}{1 + \Delta m_b}$$

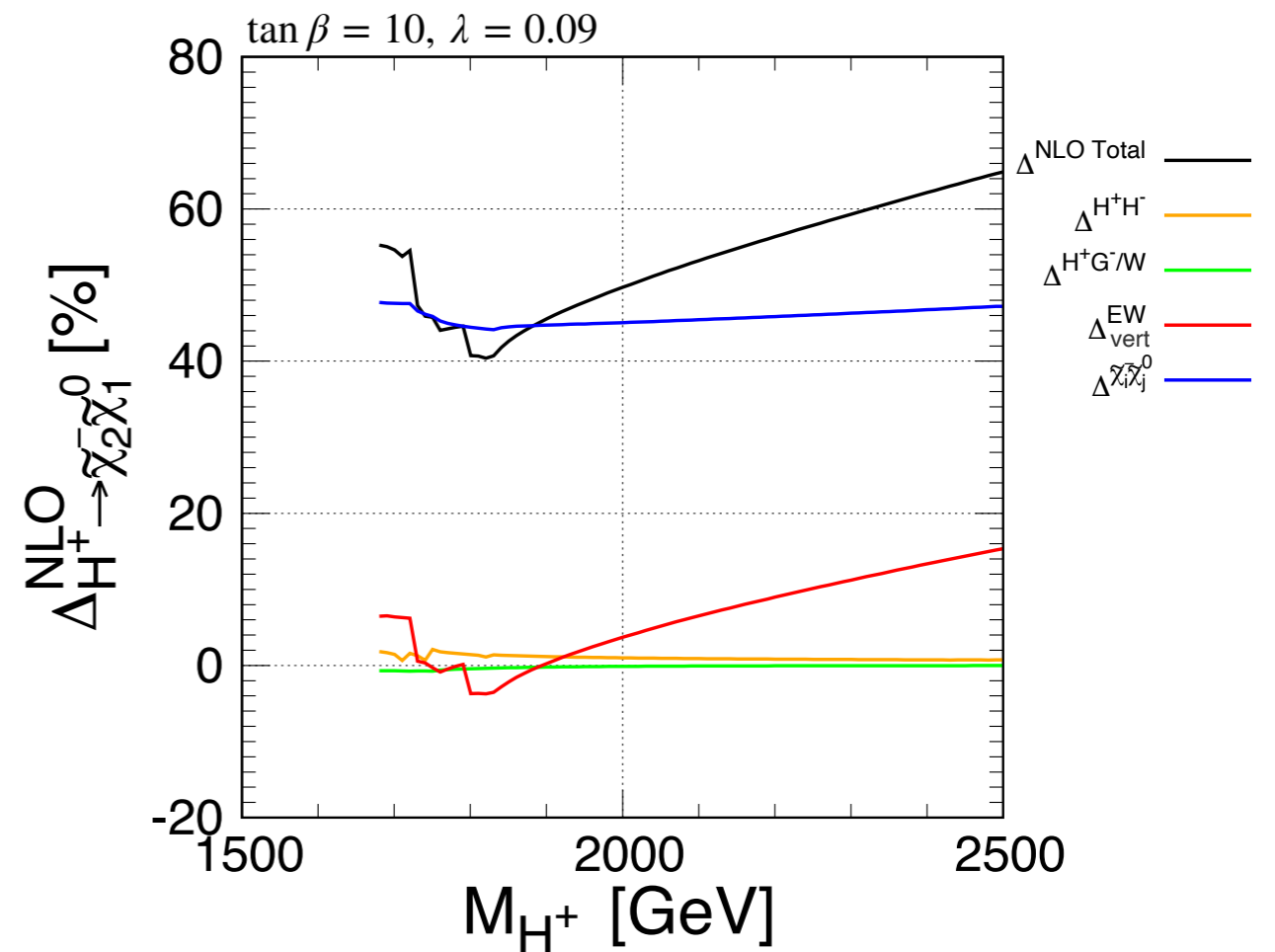
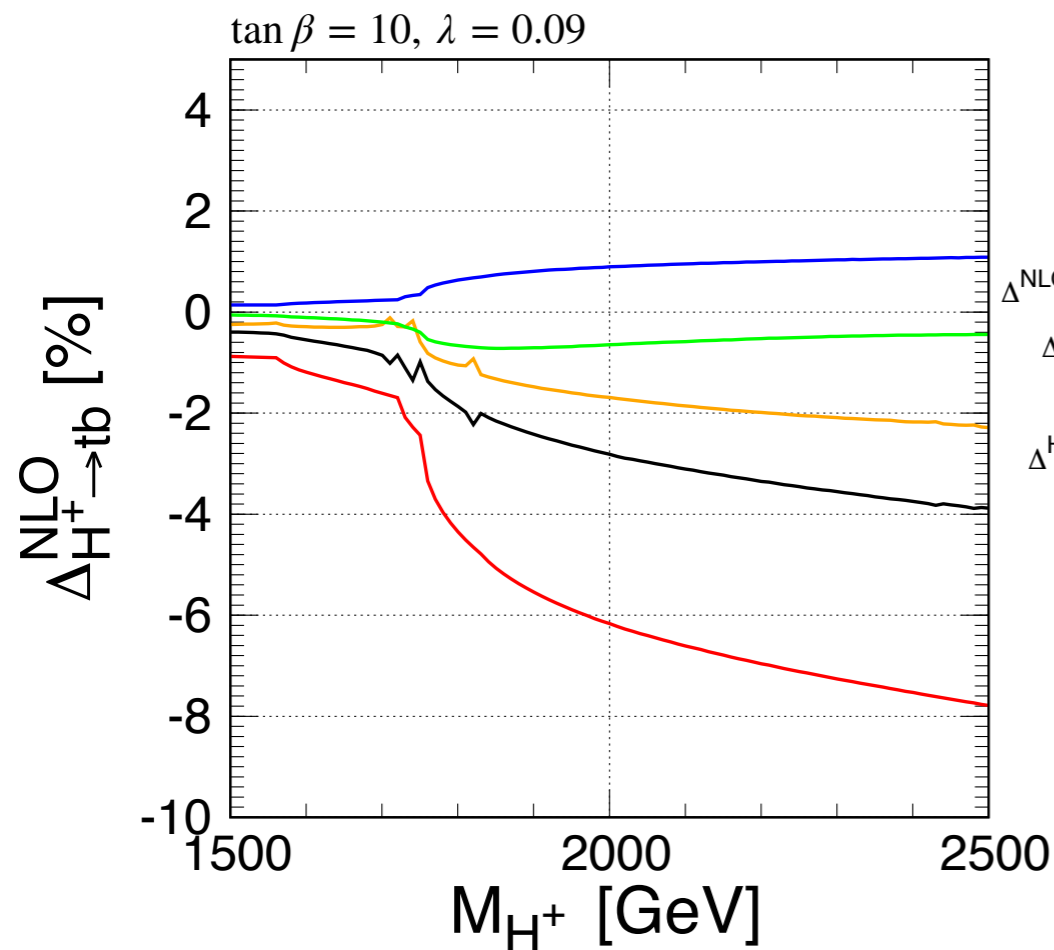
Δm_b is proportional to $\tan\beta$.

$\rightarrow \Delta m_b$ is significant if $\tan\beta \gg 1$ and this should be resummed.



Pure NLO corrections: $BR(H^+ \rightarrow tb)$, $BR(H^+ \rightarrow \tilde{W}^+ B)$

$$\Gamma(H^\pm \rightarrow X_1 X_2) = (\text{Resummed factors}) \times \Gamma_{\text{LO}}(H^\pm \rightarrow X_1 X_2) \times \left[1 + \Delta^{\text{QCD}} + \Delta_{\text{vert}}^{\text{SQCD}} + \Delta_{\text{vert}}^{\text{EW}} + \Delta^{\text{H}^+ \text{H}^-} + \Delta^{\text{H}^+ \text{G}^- / \text{W}^-} \right]$$



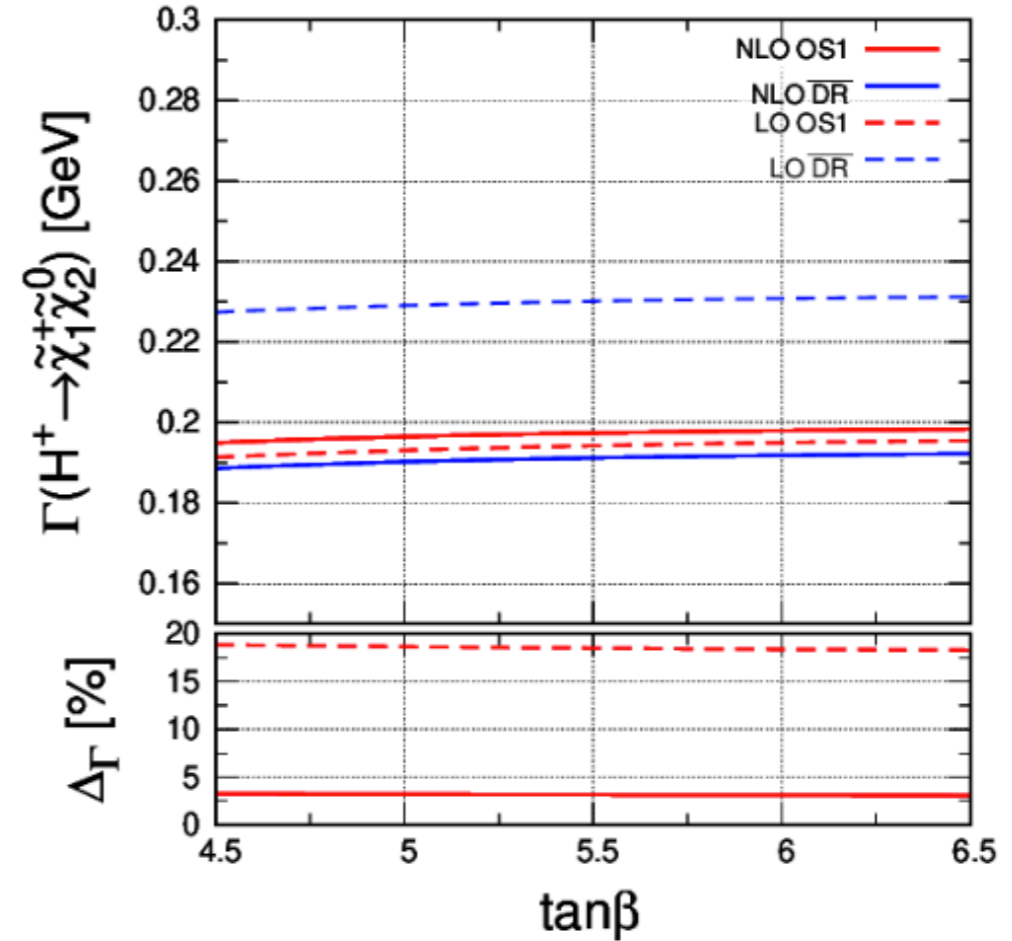
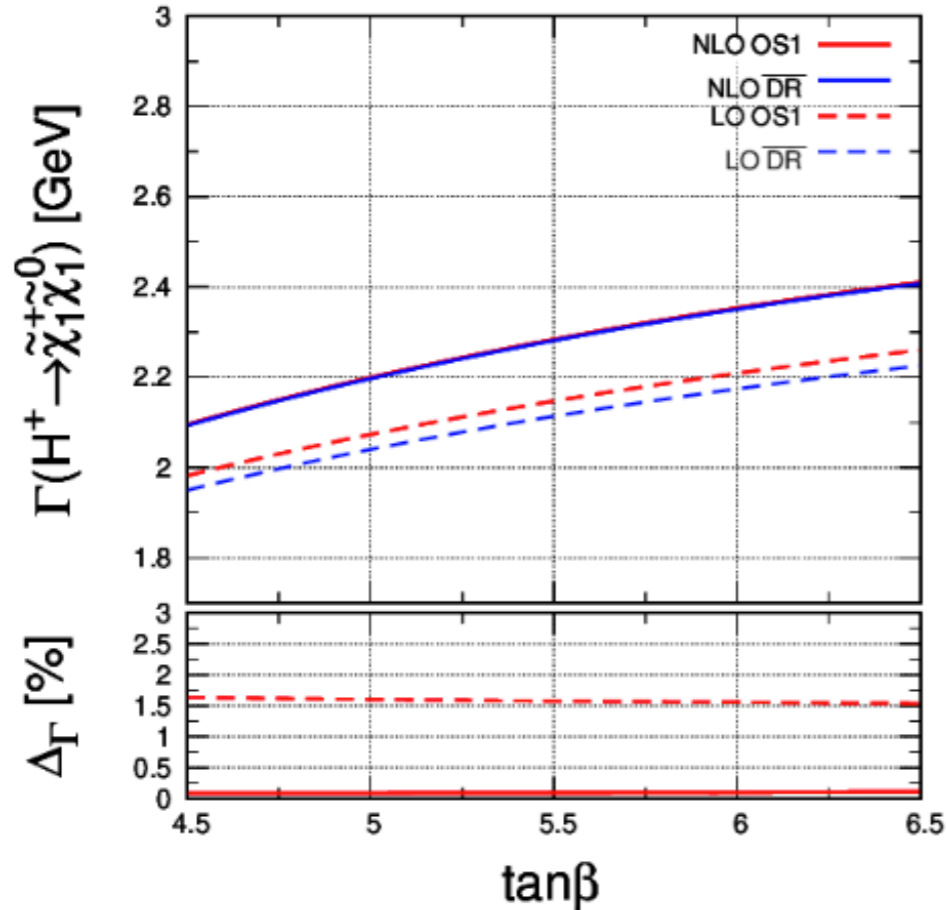
- For $H^+ \rightarrow tb$, EW corrections and SQCD corrections are destructive.

- For $H^\pm \rightarrow \tilde{\chi}_2^\pm \tilde{\chi}_1^0$, WFRs for electroweakinos are significant.

$$\Delta^{\tilde{\chi}^+ \tilde{\chi}^0} \sim \dots + \dots$$

Impact for scheme difference for charged Higgs decays

[T. N. Dao, M. Mühlleitner, S. Patel, KS]



- Scheme difference is defined by $\Delta_\Gamma \equiv \left| \frac{\Gamma^{\text{OS1}} - \Gamma^{\text{DR}}}{\Gamma^{\text{OS1}}} \right|$.

$$\Delta_\Gamma^{1\text{-loop}}(H^+ \rightarrow \tilde{\chi}_1^+ \tilde{\chi}_1^0) \sim 0\%, \quad \Delta_\Gamma^{1\text{-loop}}(H^+ \rightarrow \tilde{\chi}_1^+ \tilde{\chi}_2^0) \sim 2.5\%,$$

→ We can see that theoretical error becomes small by including 1-loop corrections.

Scan analysis for NLO corrections

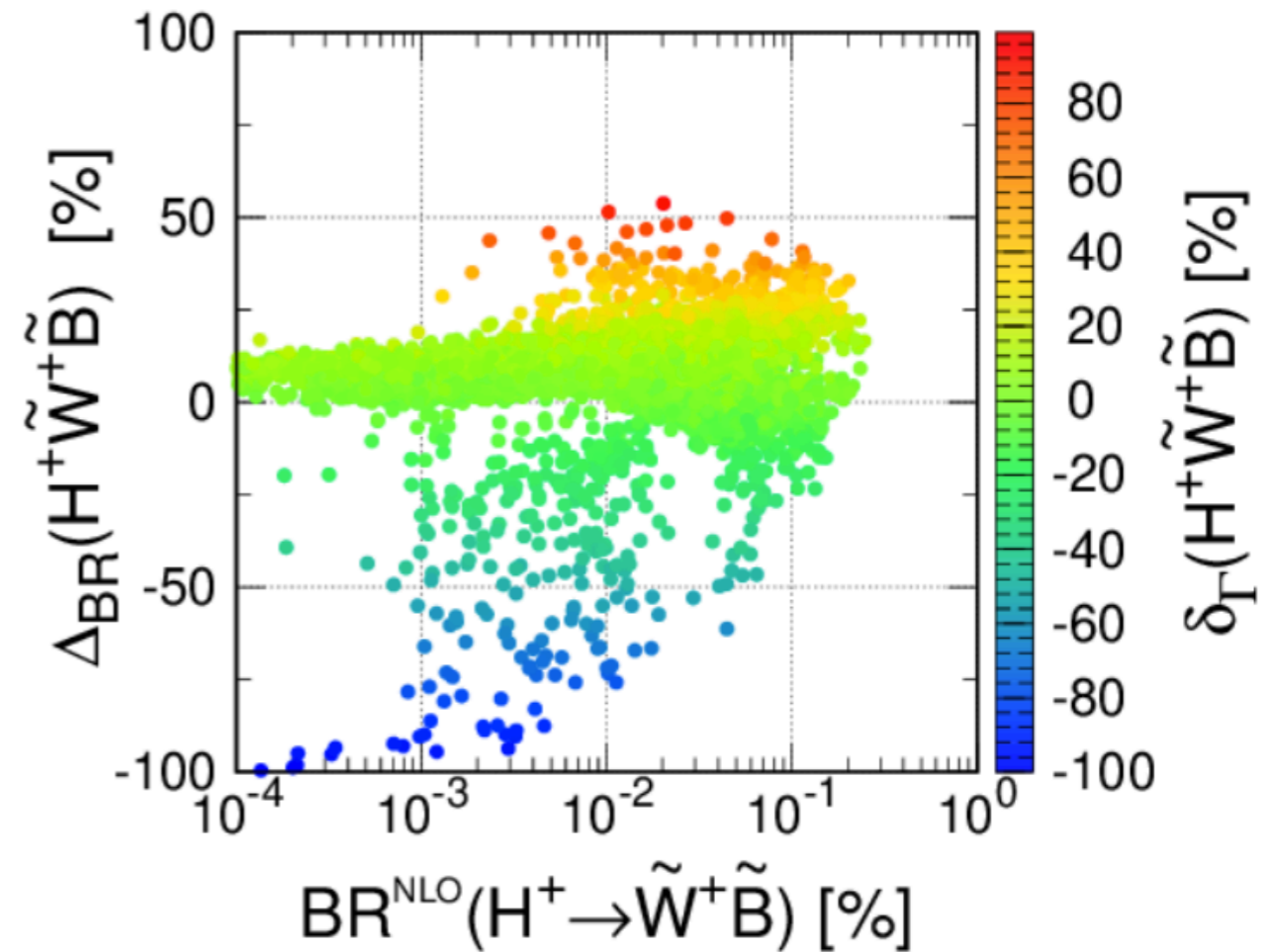
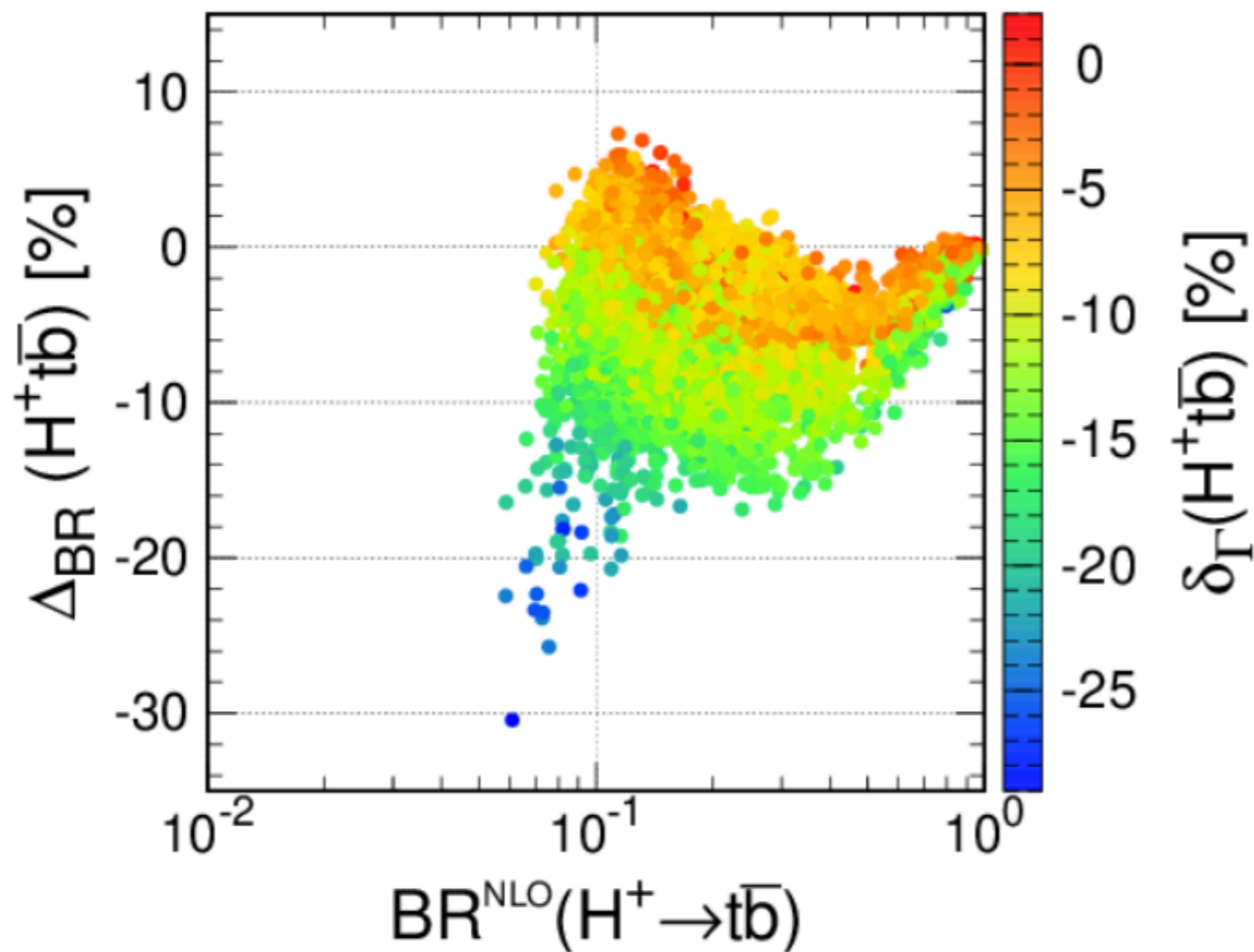
Scan range:

$$0.5 \text{ TeV} < m_{H^+} < 3 \text{ TeV}, \quad 1 < t_\beta < 20, \quad 0 < \lambda, |\kappa| < 0.7,$$

also, soft breaking parameters, A terms are scanned.

$$\Delta_{\text{BR}} = \frac{\text{BR}^{\text{NLO}} - \text{BR}^{\text{LO}}}{\max(\text{BR}^{\text{NLO}}, \text{BR}^{\text{LO}})}, \quad \delta_\Gamma = \frac{\Gamma^{\text{NLO}}}{\Gamma^{\text{LO}}} - 1$$

[T. N. Dao, M. Mühlleitner, S. Patel, KS]



- For $H^+ \rightarrow t\bar{b}$, maximum size of NLO corrections is $\sim -30\%$.
- For $H^+ \rightarrow \tilde{W}^+ \tilde{B}$, large corrections, $|\Delta_{\text{BR}}| \sim 100\%$, can appear.

Summary

- The Higgs boson will be precisely measured at the future collider experiments, such as the HL-LHC and the ILC.
 - This means that the theoretical predictions should also accurately evaluated not only for the discovered Higgs bosons but also extra Higgs bosons.
- We study NLO (SUSY +)EW and SUSY QCD corrections for various charged Higgs bosons decays in the complex NMSSM.
 - $H^+ \rightarrow tb$: maximally $\sim|30\%$ NLO corrections are obtained.
 - $H^+ \rightarrow \tilde{W}^+ \tilde{B}$: large corrections due to mixing of electroweakinos can be found in $BR < 1\%$.