Higgs and Standard Model Experimental Measurements at the LHC
Part 1

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The Standard Model at the LHC

June 2021

CMS Preliminary

Production Cross Section, $\sigma$ [pb]

- 7 TeV CMS measurement ($L \leq 5.0$ fb$^{-1}$)
- 8 TeV CMS measurement ($L \leq 19.6$ fb$^{-1}$)
- 13 TeV CMS measurement ($L \leq 137$ fb$^{-1}$)
- Theory prediction
- CMS 95%CL limits at 7, 8 and 13 TeV

All results at: http://cern.ch/go/pNj7

- Broad and spectacular confirmation of the Standard Model (and perturbative QCD/factorization)
The Standard Model at the LHC

- **LHC provides:**
  - Unprecedented production rates for $W$, $Z$, $t\bar{t}$, high energy photons/jets with respect to previous colliders
  - Access to the Higgs for the first time (and across a range of production modes)
  - A wide range of rare and complex SM processes/final states
Outline

Today:
- Z as a standard candle
- Some background on object reconstruction and identification
- Some background on detector simulation and Monte Carlo
- Tag and probe for efficiency measurements
- Precision electroweak measurements with $W$ and $Z$: $m_W$, $\sin^2 \theta_W$
- Precision measurements of $W$ and $Z$ cross sections and constraints on PDFs

Thursday and Friday:
- More on Jet/MET/tau reconstruction/identification
- Overview of Jet/multiboson/top measurements
- Higgs measurements
Single W and Z(\(\gamma^*\)) production

- W branching ratios
  - \(W \rightarrow \ell\nu\) (\(\ell = e, \mu, \tau\)): \(\sim 11\%\) per flavour
  - \(W \rightarrow \text{hadrons } (q\bar{q})\): \(\sim 67\%\)

- Z branching ratios
  - \(Z \rightarrow \ell^+\ell^-\) (\(\ell = e, \mu, \tau\)): \(\sim 3.4\%\) per flavour
  - \(Z \rightarrow \nu\nu\): \(\sim 20\%\)
  - \(Z \rightarrow \text{hadrons } (q\bar{q})\) : \(\sim 70\%\)

- Significant branching ratios with charged leptons in the final state
- Widths are non-negligible (\(\Gamma_W \sim 2.1\) GeV, \(\Gamma_Z \sim 2.5\) GeV)
Single W and Z($\gamma^*$) production

- **Why measure this process?**
  - Z especially is a “standard candle” processes which can be used to calibrate simulation and reconstruction, derive correction factors for charged lepton energy/momentum scale, efficiencies, etc
  - Large cross section allows continuous monitoring of detector/reconstruction performance
  - Inclusive and differential production cross sections are tests of perturbative QCD, and sensitive to parton distribution functions
  - Precision electroweak measurements: $m_W$, $\sin^2 \theta_W$

Very Early detector performance plots (CMS-DP-2010/016) and early xsec measurements

(10.1007/JHEP01(2011)080)
Digression: Object Reconstruction, Identification and Mis-identification

Main “high level” objects:
- Jets (±b or c tagging)
- Missing transverse momentum (aka Missing Energy aka MET), e.g. from neutrinos in final state
- (Isolated high $p_T$) photons
- (Isolated high $p_T$) electrons
- (Isolated high $p_T$) muons
- (Isolated high $p_T$) taus
Digression: Object Reconstruction, Identification and Mis-identification

- **What is actually measured in the detector: Stable* particles**
  *given relativistic boost and size of the detector
  - Charged hadrons
  - Stable neutral hadrons (e.g. neutral Kaons)
  - Photons
  - Electrons
  - Muons

- **Important special cases:**
  - $\pi^0$ is the lightest and most copiously produced neutral hadron, but promptly decays to $\gamma\gamma$ (99%) or $e^+e^-,\gamma$ (1%)
  - $\tau$ has a short but measurable lifetime (decay length 87 $\mu$m) → decays to slightly displaced electrons or muons + neutrinos ($\sim 18\%$ each) or hadrons + neutrino

- Jets are a collection of all of the above, but mostly charged hadrons, photons (mainly from $\pi^0$) and neutral hadrons in very roughly 60/30/10 proportions on average (but with large fluctuations from jet to jet)
Particle Identification in General Purpose Detectors

Key:
- **Muon**
- **Electron**
- **Charged Hadron (e.g. Pion)**
- **Neutral Hadron (e.g. Neutron)**
- **Photon**
**Prompt High $p_T$ Electrons:**

- Clusters of energy in electromagnetic calorimeter (grouped to recover bremsstrahlung/secondary conversions), matched to reconstructed track
- Hardware trigger from calorimeter
- Electron-like shower profile and track properties
- No large deposits in hadronic calorimeter behind
- Well-isolated (e.g. sum of transverse energy/momentum in a cone around the electron)
Main sources of Misidentified Prompt Electrons:

- Heavy flavour decays (e.g. $B$ and $D$ hadrons) producing displaced electrons
- Photon conversions (attempt to reconstruct or identify them, but difficult to do efficiently)
- Early showering of charged hadrons in EM calorimeter (e.g. via inelastic charge exchange $\pi^+ p \rightarrow \pi^0 n$)
Prompt High $p_T$ Photons:

- Clusters of energy in electromagnetic calorimeter (grouped to recover conversions/bremsstrahlung)
- Hardware trigger from calorimeter
- Photon-like shower profile
- No large deposits in hadronic calorimeter behind
- Well-isolated (e.g. sum of transverse energy/momentum in a cone around the photon)

n.b. shower profile can look significantly different depending on whether the photon converts to an electron positron pair before reaching the calorimeter
Main sources of Misidentified Prompt Photons:

- $\pi^0 \rightarrow \gamma\gamma$ (at high energies, the decay is collimated and tends to merge into a single shower)
- **Electrons** where primary track is not reconstructed, or misidentified as belonging to a conversion

Transverse shower width (parallel to B-field)
Muon Reconstruction/Identification

**Prompt High $p_T$ Muons:**
- Reconstructed track in inner tracker and muon chambers
- Hardware trigger from muon chambers
- No large deposits in calorimeter
- Well-isolated (e.g. sum of transverse energy/momentum in a cone around the muon)

**Main sources of Misidentified Prompt Muons:**
- **Heavy flavour** decays (e.g. $B$ and $D$ hadrons) producing displaced muons
- **Decay in flight** of charged hadrons (e.g. $\pi^+/K^+ \rightarrow \mu^+\nu$), can be suppressed with track quality, “kink-finding”
- **“Punch through”** of charged hadrons (negligible with enough hadronic interaction lengths upstream)
A “Monte Carlo Sample” as produced by the LHC experiments typically consists of a full chain of Monte Carlo Generator → Detector Simulation → “digitization” → reconstruction, to produce events which look as close as possible to data given the input physics assumptions.

**Monte Carlo Generator:** Simulate proton collisions up to stable particle level (PDFs, matrix element/hard interaction, hadronization, prompt decays, MPI/underlying event, etc)

More details on Monte Carlo generators themselves in Rikkert’s lectures.

**Detector simulation:** Simulate the interaction of the generated particles with the detector using **Geant4**

- Energy loss
- Multiple scattering
- Bremsstrahlung
- Photon conversions
- Nuclear interactions
- Electromagnetic and hadronic showers
- Many many other small details with input from many sources of experimental data on interactions of particles with matter
Aside from the details of the physics model (e.g. modelling details of showers in calorimeters can be challenging), simulation quality depends on accuracy of \textbf{geometry} and \textbf{material model} \rightarrow notoriously difficult.

One method of checking this is with reconstructed nuclear interactions (nuclear interaction probability depends on material density).
Monte Carlo + Detector Simulation as an Analysis Tool

- A “Monte Carlo Sample” as produced by the LHC experiments typically consists of a full chain of Monte Carlo Generator → Detector Simulation → “digitization” → reconstruction, to produce events which look as close as possible to data given the input physics assumptions.

- **Digitization:** Simulate the readout/electronics of the detector
  - Energy deposits in active detector elements converted to raw hits/ADC counts etc
  - Electronics noise, inefficiencies, dead channels, etc can be simulated at this stage
  - Pileup is typically also overlaid at this stage (from independently simulated Minimum Bias events)

- **Reconstruction:** Unpack the raw data and run the reconstruction chain up to high level objects (four-vectors, ID variables, etc)
  - Ideally this is algorithmically exactly the same between data and simulation
  - Typically depends on a large set of calibration and/or alignment constants depending on the detector
  - Try to reproduce in the MC known inaccuracies and precision limitations on the data calibration constants
In general: The Monte Carlo at the LHC experiments is good, but not perfect.

Accuracy and uncertainties associated with the generator part depends very much on the generator and process.

Detector simulation/response for well-reconstructed objects is not terrible → use the Monte Carlo as a starting point and derive (hopefully small) residual corrections from data which can be used in the analysis.

Residual systematic uncertainties may be limited by the degree of (in)accuracy of the simulation → particularly difficult/high precision cases may benefit from dedicated refinement efforts.

Mis-identified objects tend not to be well predicted by the Monte Carlo, depend on details of jet flavour composition in QCD multijet events, tails of jet fragmentation functions, probability of rare interactions in the detector, etc → strong preference for data-driven methods to predict the rate and kinematic distributions of these backgrounds, especially in precision measurements.
Z → ℓ^+ ℓ^- is a very special standard candle: the presence of two leptons in the event gives the possibility to select the event based on one lepton (the tag) in order to construct an unbiased sample from the second lepton in the event (the probe).

In particular this allows the efficiency of various reconstruction and selection steps to be directly measured in data.

Trigger efficiencies can also be measured in this way (typically using single lepton triggers for the tag).

Efficiencies are typically measured in bins of $p_T$ and $\eta$ of the probe.
Lepton Efficiencies: Tag and Probe

Concrete example: Electron identification efficiency

- **Tag**: reconstructed electrons passing all ID and isolation requirements
- **Probe**: reconstructed electrons with no ID requirements applied
- **Passing probe**: Probe passing the ID requirements
- **Failing probe**: Probe failing the ID requirements
Lepton Efficiencies: Tag and Probe

- Background is subtracted by performing a likelihood fit (always larger in failing probe case)
- Typically with analytic functional form for the background (erf*exp in this case), analytic (e.g Breit Wigner * Crystal Ball) or smeared MC templates for signal
- Efficiency and corresponding statistical uncertainty can be extracted directly from a simultaneous fit of passing and failing probes
  \[ \epsilon = \frac{N_{\text{pass}}^{\text{sig}}}{N_{\text{pass}}^{\text{sig}} + N_{\text{fail}}^{\text{sig}}} \]
- Systematic uncertainties typically from alternate signal or background models, alternate fitting range, variations in tag selection, etc
- Typically applied to analysis as **scale factors** to MC: \( \epsilon_{\text{data}}/\epsilon_{\text{MC}} \) to exploit e.g. mostly correct MC modelling of efficiency variations within a bin
In this particular case (electrons in a region with a large amount of material), MC models the pt-dependence of the efficiency qualitatively, but some corrections are still needed.
Defining a suitably inclusive probe selection for reconstruction efficiency can sometimes be challenging:

- e.g. for muons use inner tracks as probes to measure muon chamber efficiency and vice versa
- e.g. for electrons use tracks as probes to measure EM calorimeter cluster efficiency and vice versa
- need to carefully consider possibly correlated sources of uncertainty in such cases
Lepton Efficiencies: Tag and Probe: Caveats

- **Sometimes inefficiencies can be correlated with poor energy/momentum measurement** (e.g. electrons incident on gaps or cracks in the calorimeter which are more likely to fail shower profile cuts, but also more likely to have their energy undermeasured)
  - These effects must be accounted for in the signal model and/or associated systematic uncertainties
  - Dedicated or multivariate energy corrections can sometimes mitigate these effects (but be careful about deriving energy corrections on tight objects and applying them to looser ones)
Lepton efficiencies: Tag and Probe: Caveats

- Lepton efficiency may be **dependent on event topology**
  - Must control associated extrapolation/variation of efficiencies when measuring in one process/phase space and applying to another
  - Example shown here concerns orientation of muon with respect to hadronic recoil in drell-yan events, but the effect may be even larger e.g. in $t\bar{t}$ events with more additional jet activity
Sources of inefficiency which are correlated between the tag and the probe cannot be measured by this method and must be accounted for by other means.

**Pathological example:** Cut on $d_{xy}$ (muon, beamspot), but beamspot is mismeasured or otherwise incorrect in the reconstruction → corresponding failing probes will be missing from the tag and probe sample because the tag will also fail the cut!

**Real life example:** Trigger pre-firing: probe lepton is reconstructed by the hardware trigger one bunch crossing too early, correct trigger of the tag is suppressed by trigger rules/deadtime → failing probe won’t appear in the sample because the event is never triggered in the correct bunch crossing.
Precise measurements of the Higgs mass enable more precise consistency tests of the Standard Model using $m_W$ and $\sin^2 \theta_W$.
Production and decay of $Z/\gamma^* \rightarrow \ell^+\ell^-$ or $W \rightarrow \ell\nu$ at the LHC, inclusive in additional hadronic activity, can be characterized by a 5-dimensional differential cross section

$$\frac{d\sigma}{dp_T^Z dy^Z dm^Z d\cos \theta d\phi} = \frac{3}{16\pi} \frac{d\sigma^{U+L}}{dp_T^Z dy^Z dm^Z} \times \left\{ (1 + \cos^2 \theta) \right.$$\left.+ \frac{1}{2} A_0 (1 - 3 \cos^2 \theta) + A_1 \sin 2\theta \cos \phi \right.$$\left.+ \frac{1}{2} A_2 \sin^2 \theta \cos 2\phi + A_3 \sin \theta \cos \phi + A_4 \cos \theta \right.$$\left.+ A_5 \sin^2 \theta \sin 2\phi + A_6 \sin 2\theta \sin \phi + A_7 \sin \theta \sin \phi \right\}.

$\theta$ and $\phi$ are the decay angles of the lepton/neutrino in the rest-frame of the $Z/\gamma^*$ or $W$, defined e.g. in the Collins-Soper frame.
Angular distributions of leptons in Z rest frame are sensitive to weak mixing angle.

Leading sensitivity through forward-backward asymmetry or A4 angular coefficient (equivalent up to a constant in the full phase-space).

Sensitivity diluted in p-p collisions due to unknown direction of incoming quark vs anti-quark.
Weak Mixing Angle

- Size of dilution effect is rapidity-dependent and sensitive to PDFs
- $\sin^2 \theta_W$ sensitivity mainly at Z peak, PDF sensitivity mainly above/below
  
  → perform measurements differential in $m_{\ell\ell}$ and $y_{\ell\ell}$

Drell-Yan Angular Distributions

- Unfolded triple-differential \( d^3\sigma/dm_\ell\ell\,d|y_\ell\ell|\,d\cos\theta^* \) cross sections containing information relevant for \( \sin^2\theta_W \) determination and in-situ PDF constraints

J.Bendavid

JHEP 12 (2017) 059
In-situ PDF constraints: Weak Mixing Angle Case

- CMS and ATLAS weak mixing angle measurements exploit in-situ constraints to reduce PDF uncertainties with Bayesian reweighting of Monte Carlo replicas/profiling of nuisance parameters associated with Hessian representation (numerically equivalent in the Gaussian limit)

ATLAS significantly improves sensitivity through inclusion of forward electrons (beyond tracking acceptance) to extend acceptance to higher rapidity.
Both experiments provide results for multiple PDF sets, but choose one nominal set for the final result

Assessing compatibility between different PDF sets non-trivial since largely common input datasets and methodology imply large correlations
## Weak Mixing Angle Measurements

<table>
<thead>
<tr>
<th>$\sin^2 \theta^\ell_{\text{eff}}$</th>
<th>CT10</th>
<th>CT14</th>
<th>MMHT14</th>
<th>NNPDF31</th>
</tr>
</thead>
<tbody>
<tr>
<td>0.23118</td>
<td>0.23141</td>
<td>0.23140</td>
<td>0.23146</td>
<td></td>
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</table>

### Uncertainties in measurements

<table>
<thead>
<tr>
<th></th>
<th>Total</th>
<th>Stat.</th>
<th>Syst.</th>
</tr>
</thead>
<tbody>
<tr>
<td>0.23101 ± 0.00036</td>
<td>39</td>
<td>21</td>
<td>32</td>
</tr>
<tr>
<td>0.23119 ± 0.00018</td>
<td>37</td>
<td>21</td>
<td>31</td>
</tr>
<tr>
<td>0.23140 ± 0.00016</td>
<td>36</td>
<td>21</td>
<td>29</td>
</tr>
<tr>
<td>0.23166 ± 0.00031</td>
<td>38</td>
<td>21</td>
<td>31</td>
</tr>
</tbody>
</table>

$\sin^2 \theta^\ell_{\text{eff}} = 0.23101 \pm 0.00036 \text{ (stat)} \pm 0.00018 \text{ (syst)} \pm 0.00016 \text{ (theo)} \pm 0.00031 \text{ (PDF)}$

### Diagram

ATLAS Preliminary

- LEP-1 and SLD: Z-pole
- LEP-1 and SLD: $A_{FB}^{0,b}$
- SLD: $A_{1}$
- Tevatron
- LHCb: 7+8 TeV
- CMS: 8 TeV
- ATLAS: 7 TeV
- ATLAS: $ee_{CC}+\mu_{CC}$
- ATLAS: $ee_{CF}$
- ATLAS: 8 TeV

- For current results, main contributions to uncertainties are statistical and PDFs
Weak Mixing Angle Prospects

Existing measurements already reduce PDF uncertainties with in-situ constraint

Measurements with full HL-LHC data can reach or surpass LEP+SLD precision, depending also on improved knowledge of PDFs from external sources
**W mass at LHC**

- **W cannot be fully reconstructed due to neutrino** → mass must be inferred from lepton $p_T$ or transverse mass distributions
- Current ATLAS measurement of $m_W$ performed using 1D $p_T^\ell$ and $M_T$ distributions (in bins of $\eta^\ell$)
- Highest possible precision required on lepton momentum and hadronic recoil scale/resolution
- $p_T^\ell$ (and $p_T^\nu$) distributions depend not only on $m_W$ but also critically on $p_T^W$ as well as polarization → strong dependence on QCD calculation and PDFs
- $M_T$ distribution still sensitive to $p_T^W$ and polarization due to finite detector acceptance

![Graphs showing ATLAS Simulation with nominal and $\pm 50$ MeV variations in $m_W$.](image)

\[J.\text{Bendavid} \quad \text{Higgs/SM Experiment} \quad 37\]
$m_W = 80370 \pm 7(\text{stat.})\pm11(\text{exp. syst})\pm14(\text{mod. syst.}) \text{ MeV}$

$m_W = 80370 \pm 7(\text{stat.})\pm11(\text{exp.})\pm8.3(\text{QCD})\pm5.5(\text{EWK})\pm9.2(\text{PDF}) \text{ MeV}$

<table>
<thead>
<tr>
<th>PDF Uncertainty (MeV)</th>
</tr>
</thead>
<tbody>
<tr>
<td>per $</td>
</tr>
<tr>
<td>per-charge</td>
</tr>
<tr>
<td>full combination</td>
</tr>
</tbody>
</table>

- PDFs determine the W rapidity spectrum and lepton decay angles through W polarization.
- Well-defined correlations between phase space regions and processes which are already partly exploited in present measurement to reduce uncertainty.
- Can be further exploited in the future.
$m_W = 80370 \pm 7\text{(stat.)} \pm 11\text{(exp. syst.)} \pm 14\text{(mod. syst.)} \text{ MeV}$

$e + e^(-) \rightarrow \gamma + X$, $\sqrt{s} = 2000 \text{ GeV}$

- $W_pT$ spectrum in relevant region driven by large logarithms in QCD calculation
- Relatively large theoretical uncertainties, and ambiguities in correlations across phase space and processes
- Current measurement using $Z_pT$ spectrum to constrain $W$, assuming strong correlations between $Z$ and $W$ production across $p_T$, but decorrelating contribution of different quark flavours
W mass: QCD Modelling Uncertainties


- Measured hadronic recoil (missing energy) distribution has some sensitivity to $W p_T$ distribution, appears to disfavour more advanced calculations of $W/Z p_T$ ratio

- Future directions for $W p_T$ spectrum:
  - Better direct measurement (special low pileup runs)
  - In-situ constraints
  - Reducing theoretical uncertainties (higher logarithmic accuracy)
  - Better understanding of heavy-flavour contributions
  - More systematic correlations of theory uncertainties across phase space and between $W$ and $Z$
Muon Momentum Scale (and Resolution)

- $Z \rightarrow \mu\mu$ can also be used as a standard candle for the muon momentum scale and resolution, since the mass (and width) are known very precisely from the LEP beam energy scan and calibration (mass is known to $2.3 \times 10^{-5}$ relative precision).

- $J/\psi$ and $\Upsilon$ can also be used ($\Upsilon$ mass is known to similar precision, and $J/\psi$ mass to $2 \times 10^{-6}$).

- **To first order calibration is trivial**: Match the $Z$ peak (+ width) between data and MC (in bins of $\eta$ for example).

- **More complicated**: Account for possible charge/$p_T$ dependence of any momentum scale or resolution bias.

![Graph showing data and prediction for $Z \rightarrow \mu\mu$]
For curvature $k \equiv 1/p_T$, the momentum scale bias can be written as
\[ \delta k/k \sim A + qM/k - \epsilon k \]
(e.g. CMS PAS SMP-14-007)

The three terms correspond to magnetic-field bias, misalignment (e.g. from weak modes in the global alignment procedure), and the average effect of material mis-modelling on the energy loss assumed in the track reconstruction.

Resolution can be written as:
\[ \sigma_k^2/k^2 \sim a + c/k^2 \]
Where the two terms correspond to average contributions from multiple scattering and hit resolution.

For CMS W-like measurement, all 5 terms are explicitly determined/corrected for using the $J/\psi$.
Muon Momentum Scale (and Resolution) $p_T$ dependence

(a) Alignment-like bias

- In the ATLAS measurement, the alignment and b-field like biases are explicitly corrected for (using the $Z$) together with the hit resolution contribution to the resolution.

(b) Material-like bias

- Material-like bias is checked (again with $Z$) and upper bound is propagated as a systematic uncertainty (also cross-checked with explicit $\pm 10\%$ variation of material model).
In the LHCb $m_W$ measurement, alignment-like bias is first corrected in a fine-grained binning using the $Z$.

Subsequently remaining alignment-like, bfield-like, and resolution corrections (both hit resolution and multiple scattering) are determined from combined fit of $J/\psi$, $\Upsilon(1S)$ and $Z$.

Material impact on scale is assessed through explicit variation of material model.
LHCb $m_W$ measurement

- LHCb measurement is complementary because of forward rapidity coverage ($2.2 \leq \eta \leq 4.4$) → PDF uncertainties expected to be anti-correlated with ATLAS and CMS
- Current measurement is statistically limited, but only $\sim 1/3$ of the run 2 dataset is used

### Measurement uncertainty summary

<table>
<thead>
<tr>
<th>Source</th>
<th>Size [MeV]</th>
<th>Notes</th>
</tr>
</thead>
<tbody>
<tr>
<td><strong>Parton distribution functions</strong></td>
<td><strong>9.0</strong></td>
<td>Average of NNPDF31, CT18, MSHT20</td>
</tr>
<tr>
<td><strong>Theory (excl. PDFs) total</strong></td>
<td><strong>17.4</strong></td>
<td>Envelope from five different models</td>
</tr>
<tr>
<td>Transverse momentum model</td>
<td><strong>12.0</strong></td>
<td>&quot;Uncorrelated&quot; 31 point scale variation</td>
</tr>
<tr>
<td>Angular coefficients</td>
<td><strong>9.0</strong></td>
<td>Envelope of Pythia, Photos and Herwig</td>
</tr>
<tr>
<td>QED FSR model</td>
<td><strong>7.2</strong></td>
<td>Test with POWHEGew</td>
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<tr>
<td>Additional electroweak corrections</td>
<td><strong>5.0</strong></td>
<td></td>
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<tr>
<td><strong>Experimental total</strong></td>
<td><strong>10.6</strong></td>
<td></td>
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<tr>
<td>Momentum scale and resolution modelling</td>
<td><strong>7.5</strong></td>
<td>Includes simple statistical contributions,</td>
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<tr>
<td>Muon ID, trigger and tracking efficiency</td>
<td><strong>6.0</strong></td>
<td>dependence on external inputs</td>
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<tr>
<td>Isolation efficiency</td>
<td><strong>3.9</strong></td>
<td>and details of the methods.</td>
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<td>QCD background</td>
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<tr>
<td><strong>Statistical</strong></td>
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<tr>
<td><strong>Total</strong></td>
<td><strong>31.7</strong></td>
<td></td>
</tr>
</tbody>
</table>
$m_W$ summary

![Graph showing the $m_W$ summary with various experimental data points and uncertainties. The graph is labelled LHCb preliminary.]
Precision W/Z Cross Section Measurements

Detector level plots of selected $W$ and $Z$ events

\[ \text{Events / 2 GeV} \]

$ATLAS$

- $\sqrt{s} = 7 \text{ TeV, } 4.6 \text{ fb}^{-1}$
- $W^+ \to e^+\nu$
- $W^+ \to \mu^+\nu$
- $Z \to \mu^+\mu$

- Multijet backgrounds to $W$ determined in this case using combination of $M_T$ distribution and inverted identification and/or isolation criteria (more details on this type of background estimate later in the week)
Going from detector level distributions to unfolded cross sections:

- Backgrounds are subtracted
- Acceptance/efficiency is corrected
- Migration of events between bins due to reconstruction biases and/or resolution effects are corrected for
- (+ propagation of systematic uncertainties)
Unfolding

Going from detector level distributions to unfolded cross sections:

- Backgrounds are subtracted
- Acceptance/efficiency is corrected
- **Migration** of events between bins due to reconstruction biases and/or resolution effects are corrected for
  - (+ propagation of systematic uncertainties)

Migrations can be corrected for via a response matrix (by simple inversion, or an alternative method incorporating some degree of regularization)

Alternatively, backgrounds, acceptance, efficiency and migrations can be corrected for implicitly by means of a maximum likelihood fit, aka likelihood based unfolding
Example shown here for statistical component of uncertainty on muon reconstruction efficiency for ATLAS W/Z measurement

Underlying uncertainty is uncorrelated in bins of single muon $p_T$ and $\eta$ in which efficiencies were measured with tag and probe, leading to non-trivial correlations in particular for $Z/\gamma^* \rightarrow \mu\mu$ measurements

Consistent propagation of correlations of uncertainties is crucial to the (re)-interpretability of the result, its use in PDF fits, etc
$W$ lepton charge asymmetry and PDF constraints

Lepton charge asymmetry is especially sensitive to the ratio of $u$ to $d$ quarks in the proton.

\[ A_\ell = \frac{d\sigma_{W^+}/d|\eta_\ell| - d\sigma_{W^-}/d|\eta_\ell|}{d\sigma_{W^+}/d|\eta_\ell| + d\sigma_{W^-}/d|\eta_\ell|} \]
PDF Constraints from ATLAS Precision W/Z cross sections

- Significant constraints on especially sea quark distributions
Lepton charge asymmetry vs $\eta$ is a convolution of PDF effect with V-A structure of W decay

W charge asymmetry as a function of W rapidity more directly probes the PDFs (but less directly accessible experimentally)

Tevatron experiments historically provided both measurements

n.b. at Tevatron, asymmetries are sensitive to sign of $\eta$ or $y$ due to $p\bar{p}$ collisions $\rightarrow$ final results are “CP” folded $A(-\eta/y) \rightarrow -A(\eta/y)$
Unfolding to $W$ rapidity using missing transverse momentum and $M_W$ constraint

Resolving resulting twofold ambiguity requires assumption about relative fractions of incoming quark vs antiquark in proton beam (plus smaller effect from gluon-initiated production) $\rightarrow$ 10% effect in total, with non-negligible uncertainty from PDF’s $\rightarrow$ some circularity in using data in this form for PDF determination.
On the other hand, lepton charge asymmetry vs $\eta^\ell$ does not contain all available information, since information on $p^\ell_T$, $p_\nu^T$ and $\Delta \phi_{\ell,\nu}$ are lost.
**W Helicity/Rapidity at LHC**

(a) left-handed $W^+$  
(b) right-handed $W^+$  
(c) $W^+$ Rapidity

At tree level:
- All $W$ production at LHC is $q\bar{q}$ induced
- Direction of the $W$ relative to the incoming quark determines the helicity
- Only two helicity amplitudes/polarization states
- $W$ has zero transverse momentum
- **Full information on valence quark PDF’s in the relevant $x$ range contained in $d\sigma/dy$ broken down into the two helicity states**

JHEP12(2017)130 E. Manca, O. Cerri, N. Foppiani, G. Rolandi
W Helicity/Rapidity at LHC

(a) left-handed $W^+$

(b) right-handed $W^+$

(c) $W^+$ Rapidity

- Direction of incoming quark depends even more on PDF’s in $pp$ vs $p\bar{p}$ collisions

- Gluon-induced contribution from higher order effects larger and more uncertain (also due to higher $E_{cm}$ compared to Tevatron)

JHEP12(2017)130 E. Manca, O. Cerri, N. Foppiani, G. Rolandi
**W Helicity/Rapidity at LHC**

- 2D distribution of charged lepton $p_T$ and $\eta$ can discriminate between helicity states as well as rapidity of the $W$. 

![Images showing 2D distributions of charged leptons $p_T$ and $\eta$ for different helicity states of $W$.](image)
2D distribution of charged lepton $p_T$ and $\eta$ can discriminate between helicity states as well as rapidity of the W
Left and right polarization components can be extracted simultaneously as a function of $W$ rapidity, using only charged lepton kinematics (likelihood-based unfolding).

- Avoids dependence on less precisely measured missing transverse momentum (at the cost of some statistical dilution).
- Avoids circular dependence on PDFs since quark vs anti-quark fraction for each rapidity is measured.
Polarized cross sections (+ covariance matrices) contain the full set of information.
Unpolarized xsecs or charge asymmetry can be produced by integrating over polarization (without assuming underlying polarization)
W Helicity/Rapidity at LHC: PDF Constraints

(a) $\bar{u}$  
(b) $\bar{d}$  
(c) $r_s = (s + \bar{s})/(\bar{u} + \bar{d})$

- Strong PDF constraints possible here as well, and a step towards further reduced PDF uncertainty in future $m_W$ measurements
Conclusions

- Questions welcome (now or in discussion sessions)
- Part 2 tomorrow afternoon
- Part 3 on Friday
- If there is particular interest in certain topics or request for clarifications we can cover them in a bit more detail tomorrow/Friday as well