

From Quarks to Drugs: A Cross-Disciplinary Journey in Technology Transfer

Pietro Faccioli

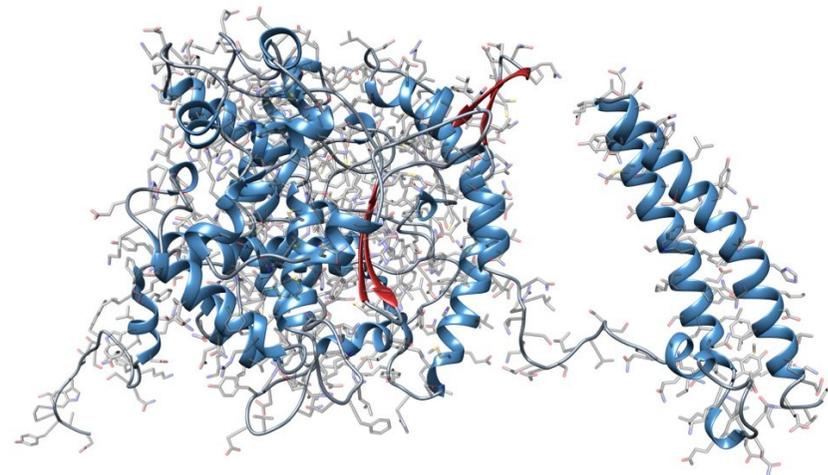


UNIVERSITÀ DEGLI STUDI
DI TRENTO

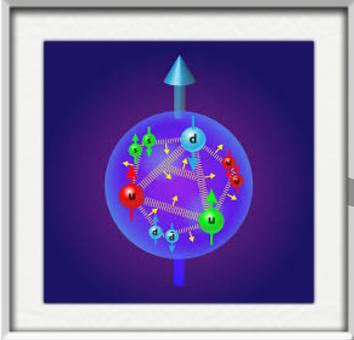
Dipartimento di Fisica



Trento Institute for
Fundamental Physics
and Applications

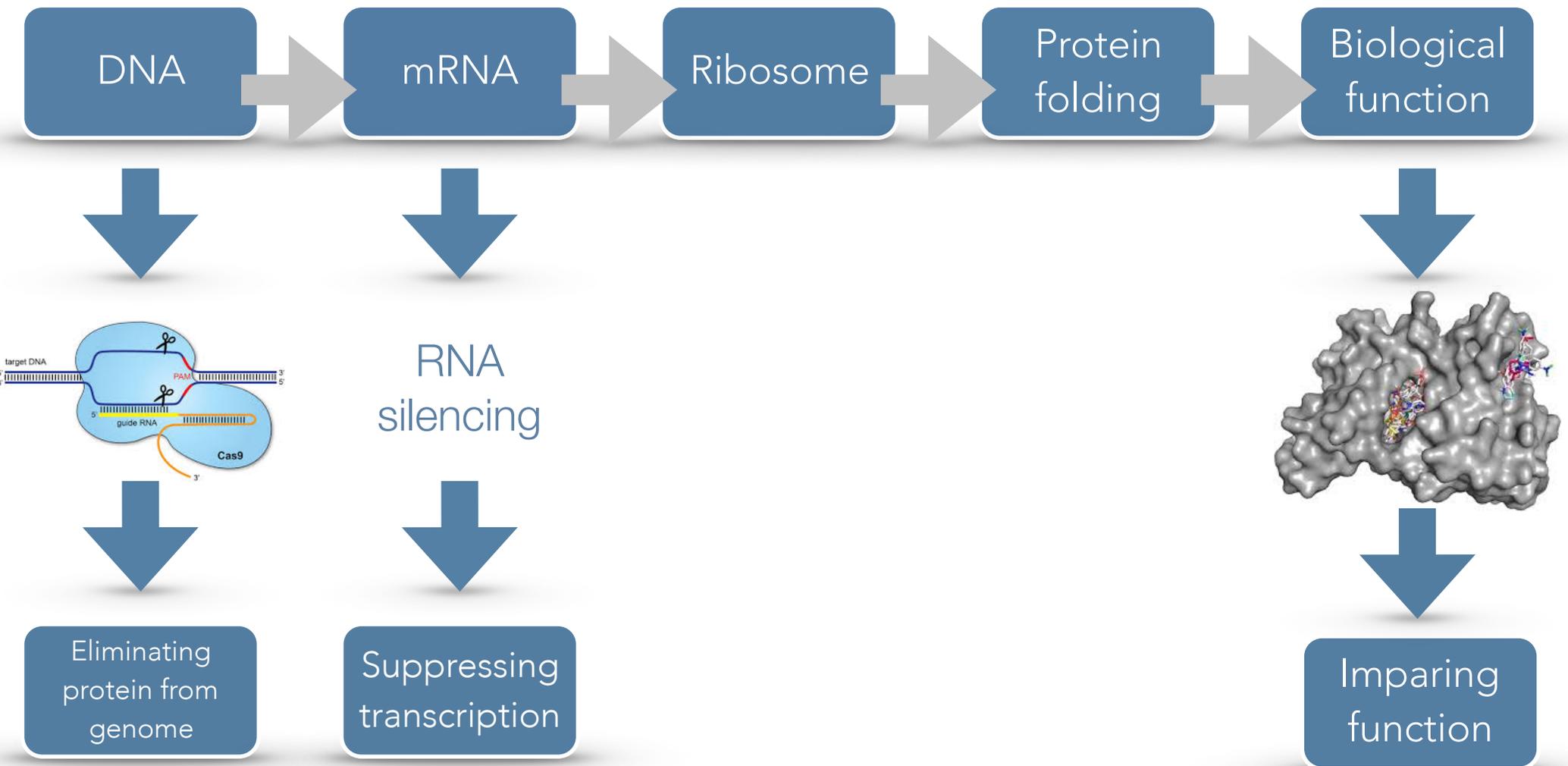


A SCIENTIFIC JOURNEY



FUNDAMENTAL DOGMA OF MOLECULAR BIOLOGY

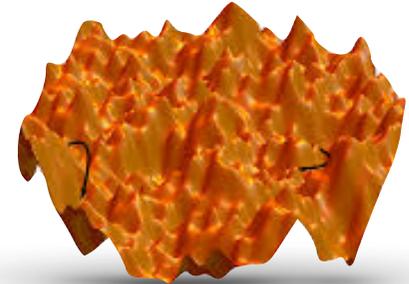
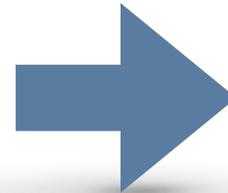
patent file # 102018000007535 (with E. Biasini)



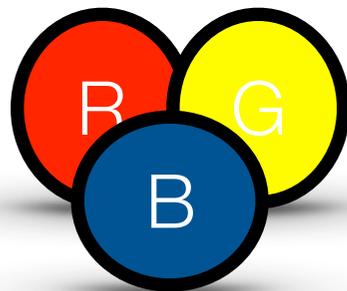
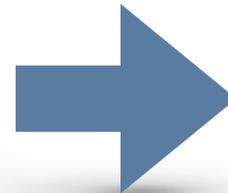
PROTEINS AND HADRONS ARE VERY SPECIAL



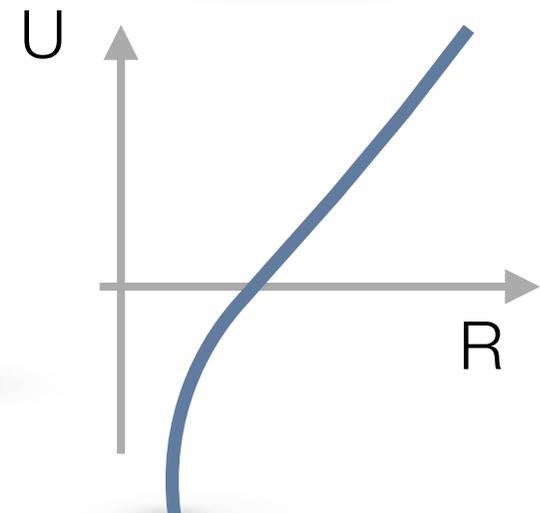
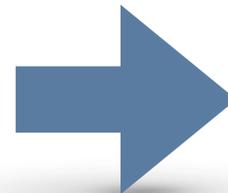
Random polypeptide



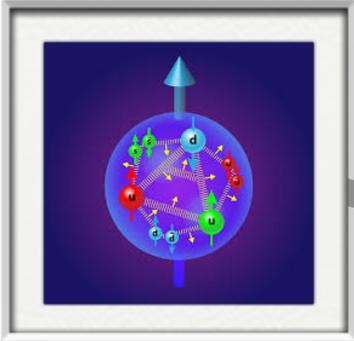
Protein



Baryon



PHASE 1: MATHEMATICAL FORMALISM & HIGH PERFORMANCE COMPUTING

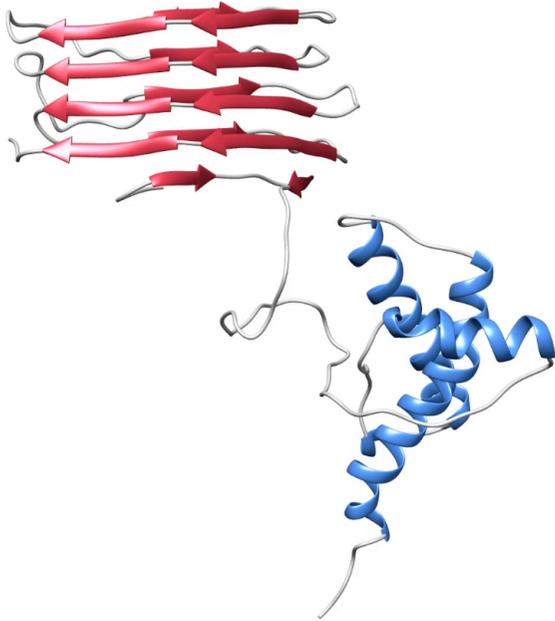


$$\mathcal{L} = \frac{1}{4g^2} G_{\mu\nu}^a G_{\mu\nu}^a + \sum_f \bar{\psi}_f (i \not{\partial} \not{D}_\mu + m_f) \psi_f$$

where $G_{\mu\nu}^a \equiv \partial_\mu A_\nu^a - \partial_\nu A_\mu^a + gf_{abc} A_\mu^b A_\nu^c$
and $D_\mu \equiv \partial_\mu + i t^a A_\mu^a$
That's it!



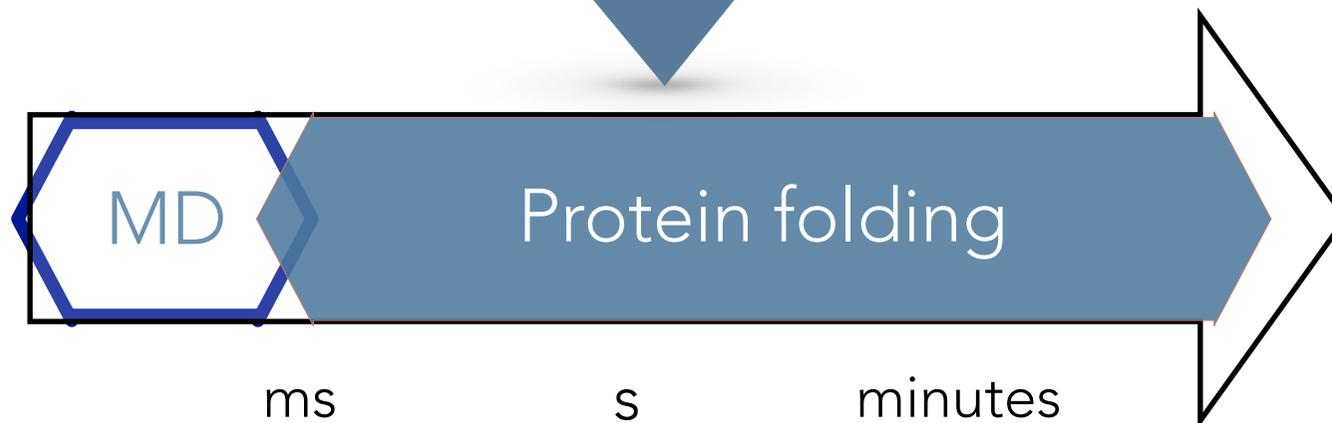
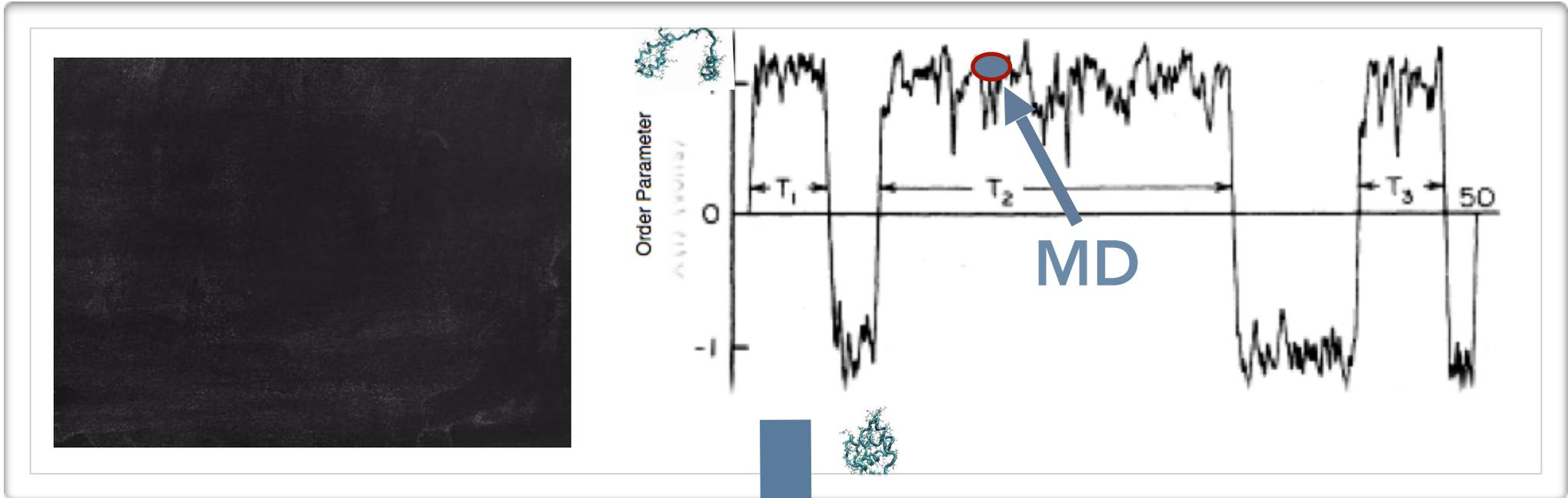
REDUCTIONIST'S APPROACH TO MOLECULAR BIOLOGY



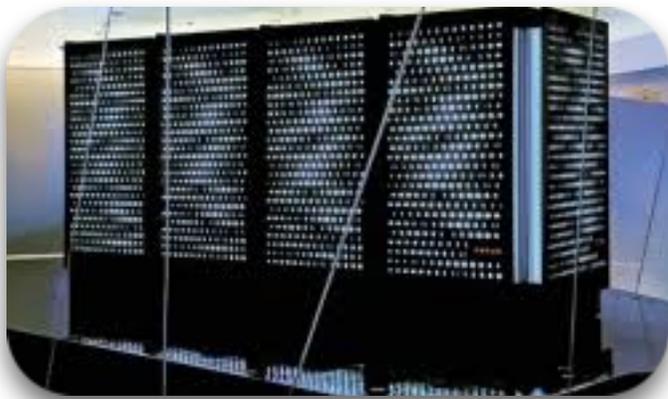
Challenge:

Integrate $\sim 10^6$ coupled
Newton-type equations
looking for **extremely
rare events**

RARE EVENT PROBLEMS



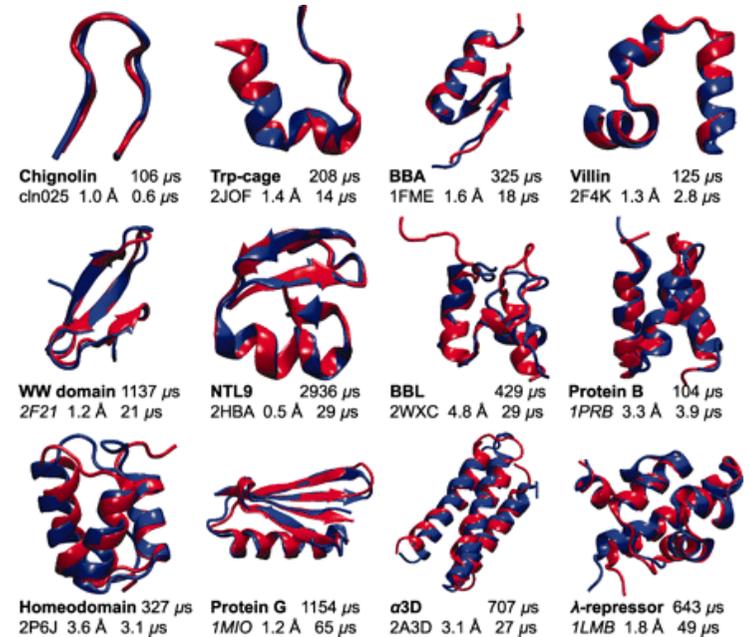
MD YIELDS CORRECT PROTEIN NATIVE STATES



Anton supercomputer
(DES Research)



MD



Atomic-Level Characterization of the Structural Dynamics of Proteins
David E. Shaw, *et al.*
Science **330**, 341 (2010);
DOI: 10.1126/science.1187409

How Fast-Folding Proteins Fold

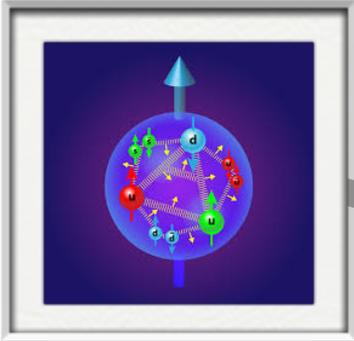
Kresten Lindorff-Larsen,^{1*}† Stefano Piana,^{1*}† Ron O. Dror,¹ David E. Shaw^{1,2†}

ZOOLOGY OF ENHANCED SAMPLING METHODS

Markov State Models, Milestoning, Transition Path Sampling, Transition Interface Sampling, Forward Flux Sampling, Temperature Accelerated Molecular Dynamics, Metadynamics, Umbrella Sampling, Blue Moon Sampling, String Method, Stochastic Difference, ... [and counting]

They are **all too computationally demanding** for many biologically relevant problems.

PHASE 1: MATHEMATICAL FORMALISM & HIGH PERFORMANCE COMPUTING



$$\mathcal{L} = \frac{1}{4g^2} G_{\mu\nu}^a G_{\mu\nu}^a + \sum_f \bar{\psi}_f (i \not{\partial} \not{D}_\mu + m_f) \psi_f$$

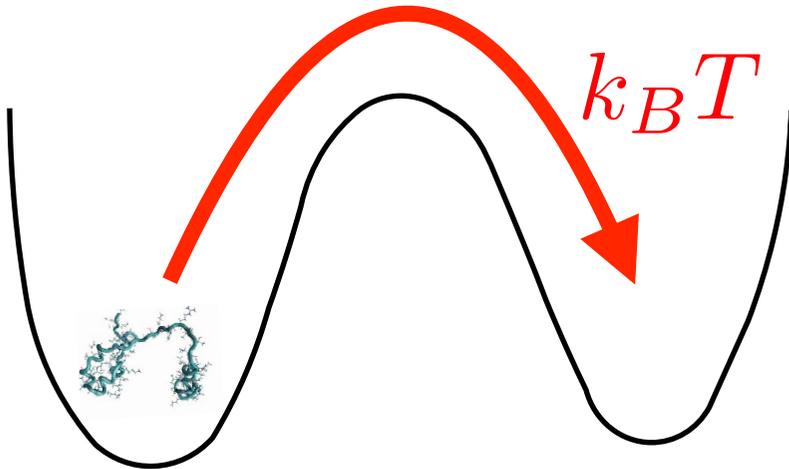
where $G_{\mu\nu}^a \equiv \partial_\mu A_\nu^a - \partial_\nu A_\mu^a + gf_{abc} A_\mu^b A_\nu^c$
and $D_\mu \equiv \partial_\mu + i t^a A_\mu^a$
That's it!



A USEFUL ANALOGY

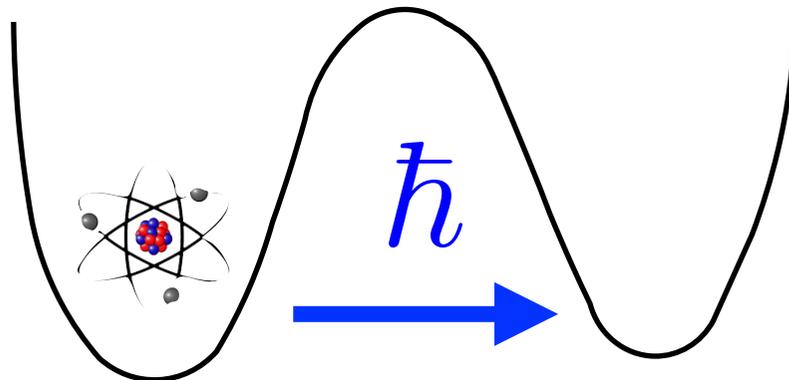
Thermal activation

($\beta = (K_B T)^{-1}$)



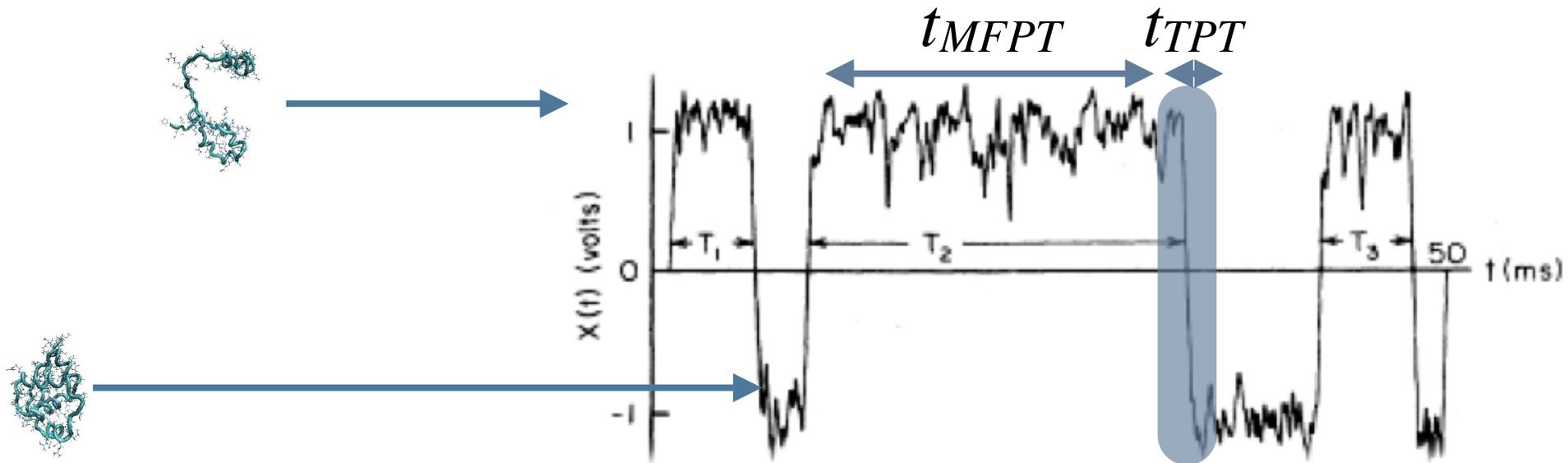
$$P(x_f, t | x_i) = \int_{x_i}^{x_f} \mathcal{D}q e^{-\frac{\beta}{4M\gamma} \int_0^t d\tau (M\ddot{q} + M\gamma\dot{q} + \nabla U(q))^2}$$

Quantum tunneling



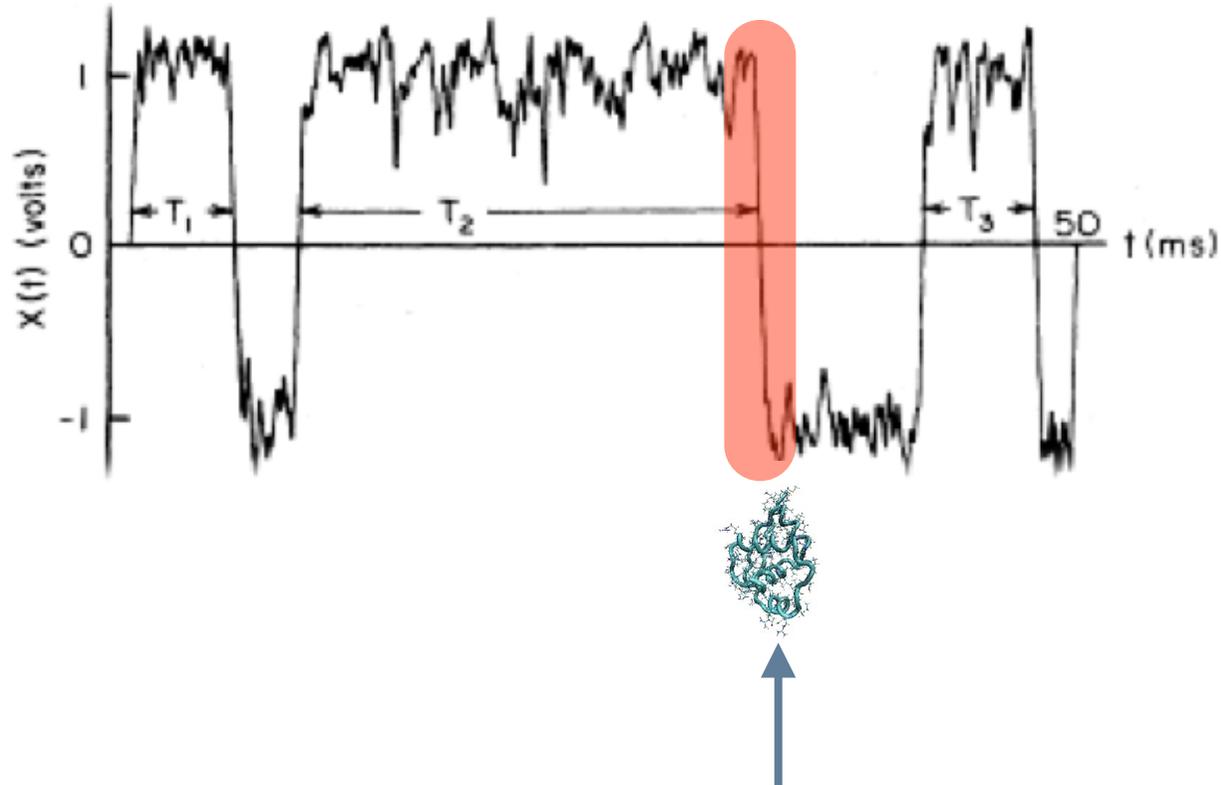
$$K_E(x_f, t | x_i) = \int_{x_i}^{x_f} \mathcal{D}q e^{-\frac{1}{\hbar} \int_0^t d\tau (\frac{M}{2} \dot{q}^2 + U(q))}$$

ADVANTAGES



$$t_{TPT} \sim \tau_0 \log \left[\log \left(\frac{t_{MFPT}}{\tau_0} \right) \right]$$

IS THIS A "FREE LUNCH"?



All atom 3D structure of the native state **are given in input**, not predicted

FULLY EXPLOITING THEORETICAL PHYSICS TOOLS

072336-4 Bartolucci, Orioli, and Faccioli

between the Gibbs distribution and the SCR estimate forward- and backward-committors, as in Eq. (A3). Introducing the distribution

$$P^{(P)}(x, t) \equiv \int dx_i P^{(P)}(x, t | x_i, 0) \rho_0(x_i), \quad (22)$$

the density in Eq. (22) reads

$$m_{SCR}(x) = \frac{1}{t_f - \tau_0} \int_{\tau_0}^{t_f} dt Q^{(R)}(x, t_f - t) P^{(P)}(x, t).$$

Using the detailed balance condition, we find $P^{(P)} = e^{-\beta U(x)} \frac{1}{Z_R} Q^{(P)}(x, t)$. Then, inserting this result into Eq. we find

$$m_{SCR}(x) = \frac{e^{-\beta U(x)}}{Z_R (t_f - \tau_0)} \int_{\tau_0}^{t_f} dt Q^{(R)}(x, t_f - t) Q^{(P)}(x, t).$$

Finally, recalling that $Q^{(R)}(x, t)$ and $Q^{(P)}(x, t)$ are time-independent in the SCR and using Eqs. (17) and we recover a fundamental result of TPT [cf. Eq. (A. Appendix A)],

$$m_{SCR}(x) \propto e^{-\beta U(x)} q_{SCR}^+(x) (1 - q_{SCR}^+(x)).$$

Within the same framework, it is possible to do the reactive current in the SCR in complete analogy Eq. (22),

$$J_{SCR}^i(x) = \frac{-D}{t_f - \tau_0} \int_{\tau_0}^{t_f} dt Q^{(R)}(x, t_f - t) \times (\vec{\nabla} - \overleftarrow{\nabla} + \beta \nabla U(x)) P^{(P)}(x, t).$$

$$\begin{aligned} V_{eff}^R(\mathbf{X}) &\simeq \frac{D_0(1-b)}{\pi b \Omega} \nabla^2 V_{eff}(\mathbf{X}) \\ &+ \frac{1}{2} \left(\frac{D_0(1-b)}{\pi b \Omega} \right)^2 \nabla^4 V_{eff}(\mathbf{X}) \\ &+ \frac{1}{6} \left(\frac{D_0(1-b)}{\pi b \Omega} \right)^3 \nabla^6 V_{eff}(\mathbf{X}) - \frac{D_0^2(1-b^3)}{3\pi(b\Omega)^3} (\partial_i \partial_j V_{eff}(\mathbf{X}))^2. \end{aligned} \quad (24)$$

Note that the first line is the leading order term (i.e. $L = 1$), while the second and third lines display the order $L = 2$ and $L = 3$ corrections, respectively.

We emphasize that the result of the EST construction is a new expression for the *same* path integral (15), in which the UV cutoff has been lowered from Ω to $b\Omega$. Equivalently, the path integral is discretized according to a larger elementary time step, $\Delta t \rightarrow \Delta t/b$:

$$Z^{\Delta t}(t) \equiv \oint_{\Delta t} \mathcal{D}\mathbf{X} e^{-S_{eff}[\mathbf{X}]} \propto \oint_{\Delta t/b} \mathcal{D}\mathbf{X} e^{-S_{eff}[\mathbf{X}] - \int_0^t d\tau V_{eff}^R[\mathbf{X}(\tau)]} \equiv Z_{EST}^{\Delta t/b}(t) \quad (25)$$

In these expressions, the symbol $\oint_{\Delta t}$ denotes the fact that the path integral is discretized according to an elementary time step Δt and we have suppressed the subscript " $<$ ", in the paths. It can be shown that the proportionality factor between $Z^{\Delta t}(t)$ and $Z_{EST}^{\Delta t/b}(t)$

PRL 114, 098103 (2015) PHYSICAL REVIEW LETTERS

$$\mathcal{P}_{\text{bias}}[X] = \int \mathcal{D}Y e^{-S_{\text{OM}}[X, Y] - U(X, Y)/k_B T}, \quad (3)$$

where the functional $S_{\text{bias}}[X, Y]$ is defined as

$$\begin{aligned} S_{\text{bias}} &\equiv \frac{1}{4k_B T} \int_0^t d\tau \left[\sum_{i=1}^N \frac{1}{\gamma_i m_i} (m_i \dot{x}_i + m_i \gamma_i \dot{x}_i + \nabla_i U - \mathbf{F}_i^{\text{bias}})^2 \right. \\ &\quad \left. + \sum_{j=1}^N \frac{1}{\gamma_j m_j} (m_j \dot{y}_j + m_j \gamma_j \dot{y}_j + \nabla_j U)^2 \right]. \end{aligned} \quad (4)$$

The Onsager-Machlup functional $S_{\text{OM}}[X, Y]$ entering Eq. (2) is recovered, setting $\mathbf{F}_i^{\text{bias}} = 0$ in Eq. (4).

Let us now return to the problem of computing the reaction pathways in the unbiased Langevin dynamics [Eq. (1)]. Using the standard reweighting trick we can write the variational condition $(\delta/\delta X)\mathcal{P}[X] = 0$ as

$$\frac{\delta}{\delta X} [\mathcal{P}_{\text{bias}}[X] (e^{-S_{\text{OM}}[X, Y]} - S_{\text{bias}}[X, Y])_{\text{bias}}] = 0. \quad (5)$$

We now introduce our main approximation, by restricting the search for the optimum path $X(\tau)$ within an ensemble of trajectories generated by integrating the biased Langevin equation. By definition, these paths have a large statistical weight in the biased dynamics, i.e., they lie in the functional vicinity of some path $\tilde{X}(\tau)$ which satisfies $(\delta/\delta X)\mathcal{P}[\tilde{X}] = 0$. Thus, the typical biased paths approximately satisfy the stationary condition

This equation states that for which the force is least. In derived in the context of solvent-induced... Let us now emphasize that the history-dependent ratchet-and-rhythm developed in Ref... formalism attempts to terms of slow... z. Conversion... To define... (1) with a... $-\frac{k_B}{2} \nabla$
0
 $z_m(t)$ dec... time t (we... obeys the... Let us

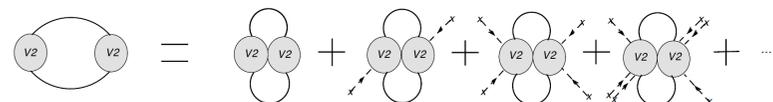


FIG. 3: Diagrammatic representation of the local time-derivative expansion of a non-local diagram —Eq. (49)—. Solid lines are fast-mode propagators, while dashed lines represent a single time derivative acting on the corresponding vertex function.

Notice that each term in the perturbative expansion (35) generates a new vertex, with an increasing power of the $x_{>}(\tau)$ field. The couplings to the fast modes depend implicitly on the time τ , through the slow modes $x_{<}(\tau)$.

By Wick theorem, each term in the series (34) can be related to a Feynman graph with vertices given by (36) and propagators given by —see appendix A —:

$$\langle x_{>}^i(\tau_1) x_{>}^j(\tau_2) \rangle_0 = \sum_{\{|\omega_n|, |\omega_m| \in S_b\}} G_{>}^{ij}(\omega_n, \omega_m) e^{i(\omega_m \tau_1 + \omega_n \tau_2)} = \sum_{\{|\omega_n| \in S_b\}} \delta_{ij} \frac{2}{\beta \gamma t \omega_n^2} e^{i\omega_n(\tau_2 - \tau_1)}. \quad (37)$$

The expansion (34) can be re-organized as the exponent of the sum performed over only connected diagrams:

$$e^{-\beta S_{>}[x_{<}(\tau)]} = e^{\Sigma(\text{all connected diagrams})}. \quad (38)$$

Hence, the path integral (26) for the slow modes can be given the following exact diagrammatic representation

$$Z(t) \equiv \oint \mathcal{D}x_{<} e^{-\beta S_{eff}[x_{<}(t)] + \Sigma(\text{all connected diagrams})}. \quad (39)$$

Below we give a classification of all the connected diagrams that may give a contribution to the expansion above.

064108-3 Orioli, a Beccara, and Faccioli

J. Chem. Phys. 147, 064108 (2017)

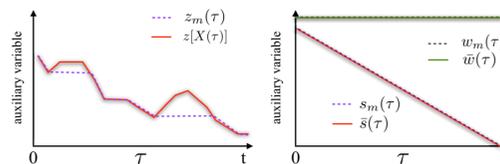


FIG. 1: Illustrative representation of the dynamics of the auxiliary variables introduced in the path integral representation of rMD (left panel) and in the derivation self-consistent path sampling algorithm (right panel).

of such a variable is frozen any time z_m becomes smaller than $z(X)$ and any time the collective coordinate $z(X)$ is increasing. Its time derivative is otherwise set equal to $\dot{z}(X)$. Therefore, by choosing the initial conditions $z_m(0) = z(X(0))$, $z_m(\tau)$ is identically set equal to the minimum value attained by the collective coordinate z until time τ (see left panel of Fig. 1).

The functional $S_{MD}[X, z_m]$ in the exponent of Eq. (8) coincides with an OM action with the addition of the unphysical biasing force \mathbf{F}_i ,

$$S_{MD} = \sum_{i=1}^N \Gamma_i \int_0^t d\tau [m_i \dot{x}_i + m_i \gamma_i \dot{x}_i + \nabla_i U - \mathbf{F}_i]^2. \quad (9)$$

In Eq. (8), $\Phi[z_m, X]$ denotes a Jacobian factor that needs to be introduced in order to ensure that the statistical weight of the paths is not affected by the measure of the $\int \mathcal{D}z_m$ integral, i.e.,

$$\int \mathcal{D}z_m \Phi[z_m, X] \delta \left[z_m(\tau) - \int_0^\tau d\tau' \dot{z}(X(\tau')) \theta(-\dot{z}(X(\tau'))) \right]$$

III. SELF-CONSISTENT PATH SAMPLING

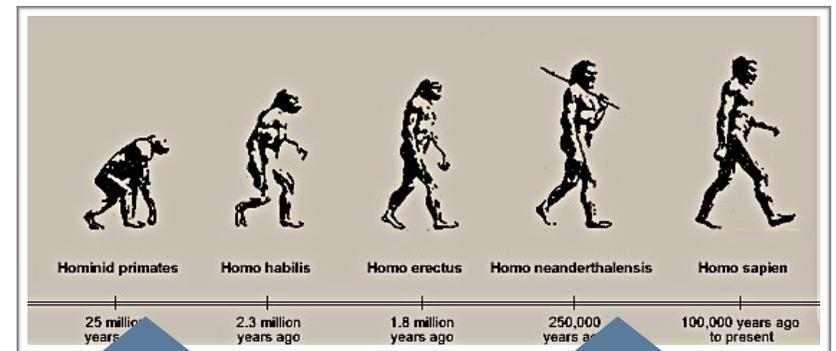
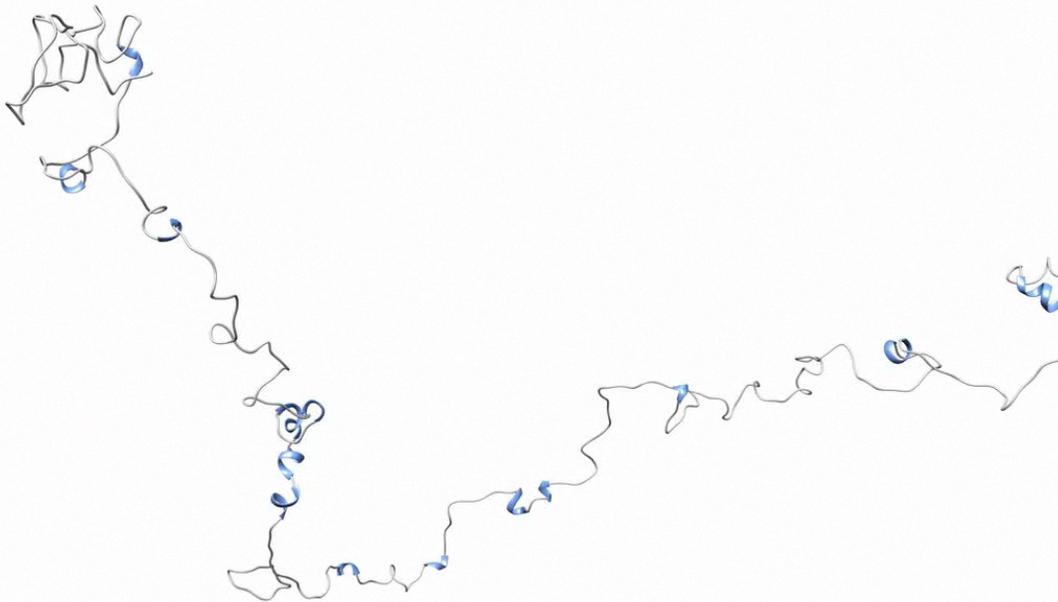
Let us now introduce our new algorithm, which provides major improvement with respect to the rMD and BF schemes discussed in Sec. II A. Indeed, it follows directly from the unbiased Langevin equation and allows us to remove the systematic errors associated to the choice of biasing coordinate.

Our starting point is path integral representation of the unbiased Langevin dynamics (2). We introduce two dumb auxiliary variables $w_m(\tau)$ and $s_m(\tau)$ into this path integral by means of appropriate functional Dirac deltas,

$$\begin{aligned} p(X_N, t | X_0) &= \int_{X_0}^{X_N} \mathcal{D}X \cdot e^{-S[X]} \int_{S(0)} \mathcal{D}s_m \int_{\hat{w}(0)} \mathcal{D}w_m \\ &\cdot \delta \left[w_m(\tau) - \int_0^\tau d\tau' \dot{w}(\tau') \theta(-\dot{w}(\tau')) \theta(w_m(\tau') - \hat{w}(\tau')) \right] \\ &\cdot \delta \left[s_m(\tau) - \int_0^\tau d\tau' \dot{s}(\tau') \theta(-\dot{s}(\tau')) \theta(s_m(\tau') - \bar{s}(\tau')) \right], \end{aligned} \quad (12)$$

where $\bar{s}(\tau)$ and $\hat{w}(\tau)$ are two external time-dependent functions to be defined below. In analogy with the path integral repre-

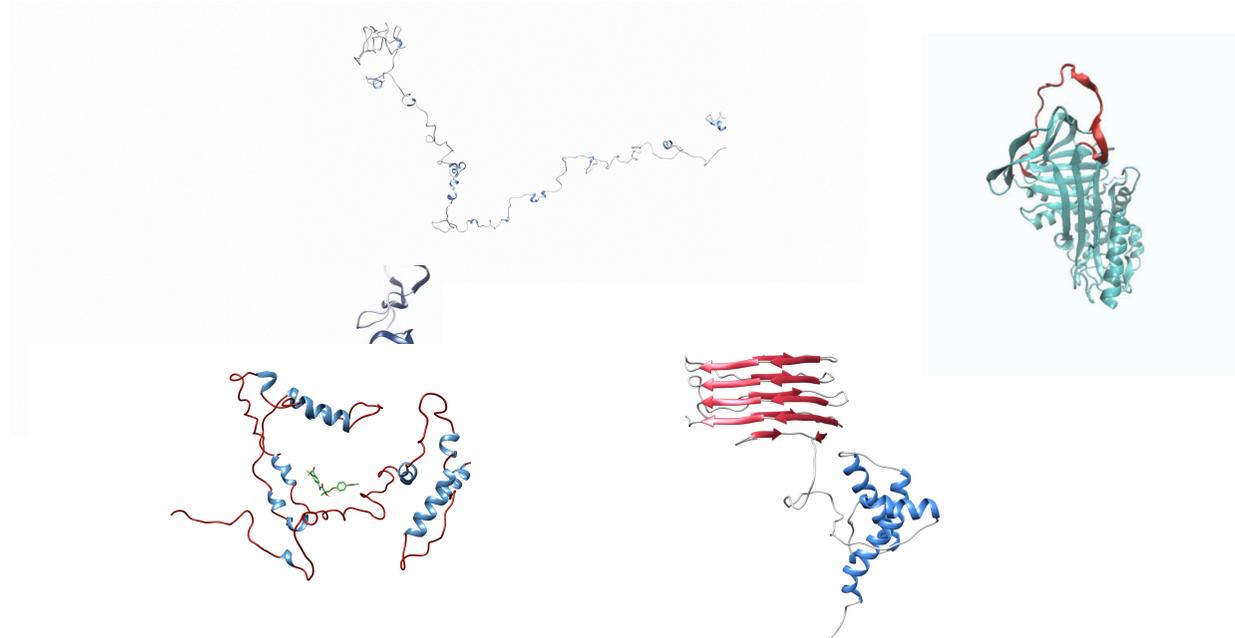
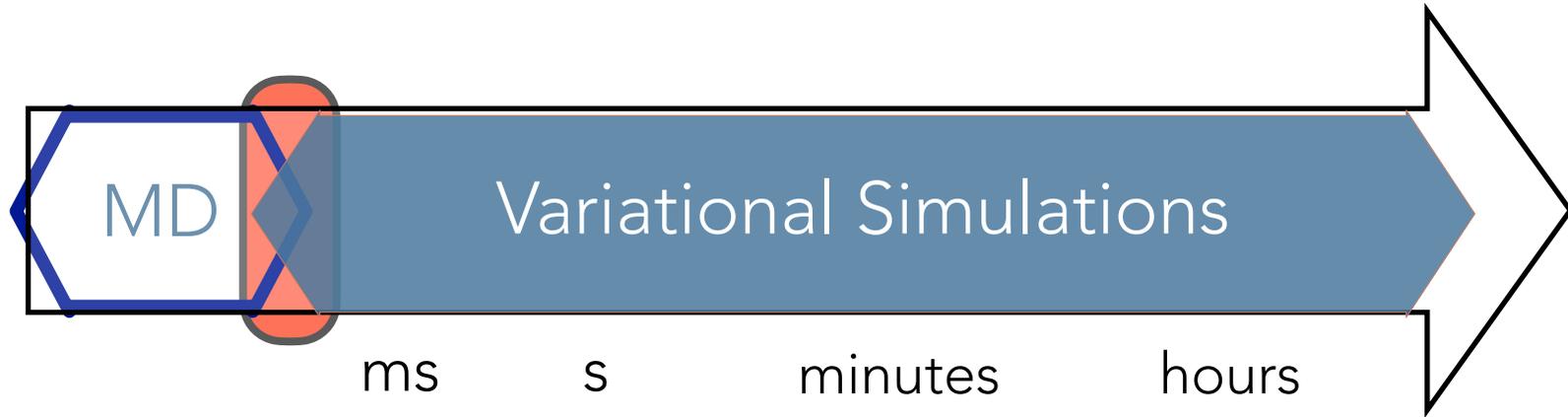
HUGE COMPUTATIONAL GAIN



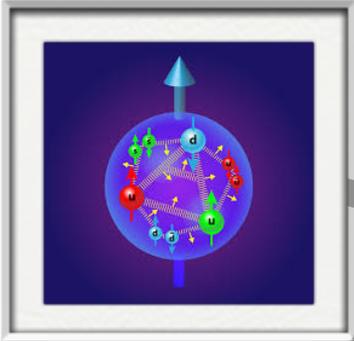
Using top all-
purpose
supercomputers

Using top
special-purpose
supercomputer

VENTURING INTO THE BIO-ZONE

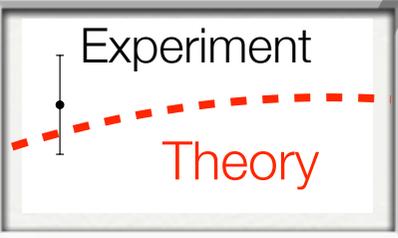


PHASE 2: VALIDATION

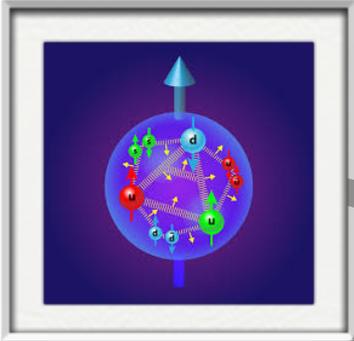


$$\mathcal{L} = \frac{1}{4g^2} G_{\mu\nu}^a G_{\mu\nu}^a + \sum_f \bar{\psi}_f (i \not{\partial} \not{D}_\mu + m_f) \psi_f$$

where $G_{\mu\nu}^a \equiv \partial_\mu A_\nu^a - \partial_\nu A_\mu^a + gf_{abc} A_\mu^b A_\nu^c$
and $D_\mu \equiv \partial_\mu + i t^a A_\mu^a$
That's it!

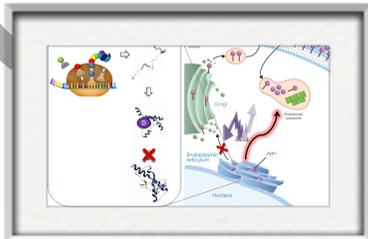
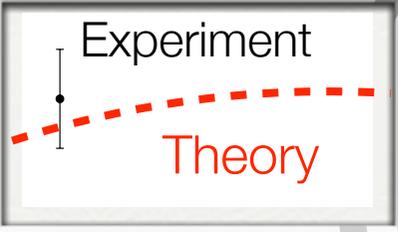


PHASE 3: EXPLOITATION IN MOLECULAR BIOLOGY



$$\mathcal{L} = \frac{1}{4g^2} G_{\mu\nu}^a G_{\mu\nu}^a + \sum_f \bar{\psi}_f (i \not{\partial} \not{D}_\mu + m_f) \psi_f$$

where $G_{\mu\nu}^a \equiv \partial_\mu A_\nu^a - \partial_\nu A_\mu^a + gf_{abc} A_\mu^b A_\nu^c$
and $D_\mu \equiv \partial_\mu + i t^a A_\mu^a$
That's it!



EXPLORING BIOLOGICAL PROCESSES

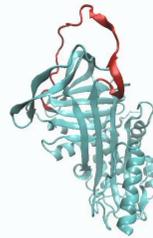
Serpin latency transition at atomic resolution

Giorgia Cazzoli^{2,3}, Fang Wang¹, Silvio a Beccara^{3,4}, Anne Gershenson⁵, Pietro Faccioli^{2,3,1}, and Patrick L. Wintrode^{6,1}

¹Dipartimento di Fisica, Università degli Studi di Trento, 38100 Povo (Trento), Italy; ²Trento Institute for Fundamental Physics and Applications, 38123 Povo (Trento), Italy; ³Department of Pharmaceutical Sciences, School of Pharmacy, University of Maryland, Baltimore, MD 21201; ⁴Interdisciplinary Laboratory for Computational Science, Fondazione Bruno Kessler, 38123 Povo (Trento), Italy; and ⁵Department of Biochemistry and Molecular Biology, University of Massachusetts Amherst, Amherst, MA 01003

Edited by David E. Shaw, D. E. Shaw Research, New York, NY, and approved September 12, 2014 (received for review April 24, 2014)

Protease inhibition by serpins requires a large conformational transition from an active, metastable state to an inactive, stable state for polypeptide chains consisting of nearly 100 amino acids (6), which are considerably smaller than PAI-1. Additionally, the



Biophysical Journal
Article

Biophysical Society

All-Atom Simulations Reveal How Single-Point Mutations Promote Serpin Misfolding

Fang Wang,¹ Simone Orioli,^{2,3} Alan Ianeselli,^{2,3} Giovanni Spagnoli,^{2,3} Silvio a Beccara,^{2,3} Anne Gershenson,^{4,*} Pietro Faccioli,^{2,3,*} and Patrick L. Wintrode^{1,*}

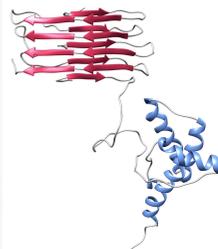


PLOS PATHOGENS

RESEARCH ARTICLE

Full atomistic model of prion structure and conversion

Giovanni Spagnoli^{1,*}, Marta Rigoli^{1,2}, Simone Orioli^{2,3}, Alejandro M. Sevillano⁴, Pietro Faccioli^{2,3}, Holger Wille⁵, Emiliano Biasini^{1,*}, Jesús R. Requena^{6,*}



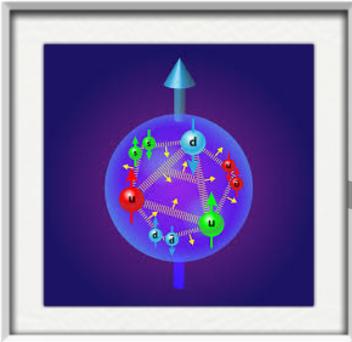
All-Atom Simulation of the HET-s Prion Replication

Luca Terruzzi^{1,2,*}, Giovanni Spagnoli^{2,3**}, Alberto Boldrini^{1,2}, Jesús R. Requena⁴, Emiliano Biasini^{2,3#} and Pietro Faccioli^{5,6#}



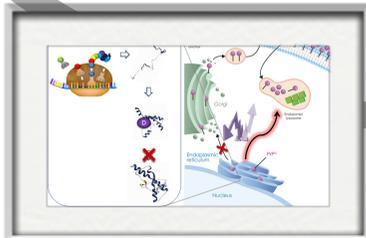
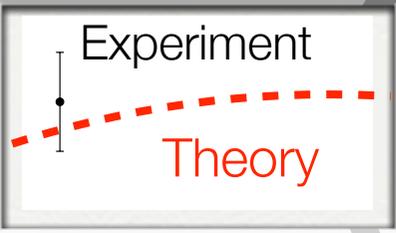
Teaming up with
E. Biasini's lab (DICIBIO)

PHASE 4: PHARMACOLOGOLOGICAL RESEARCH



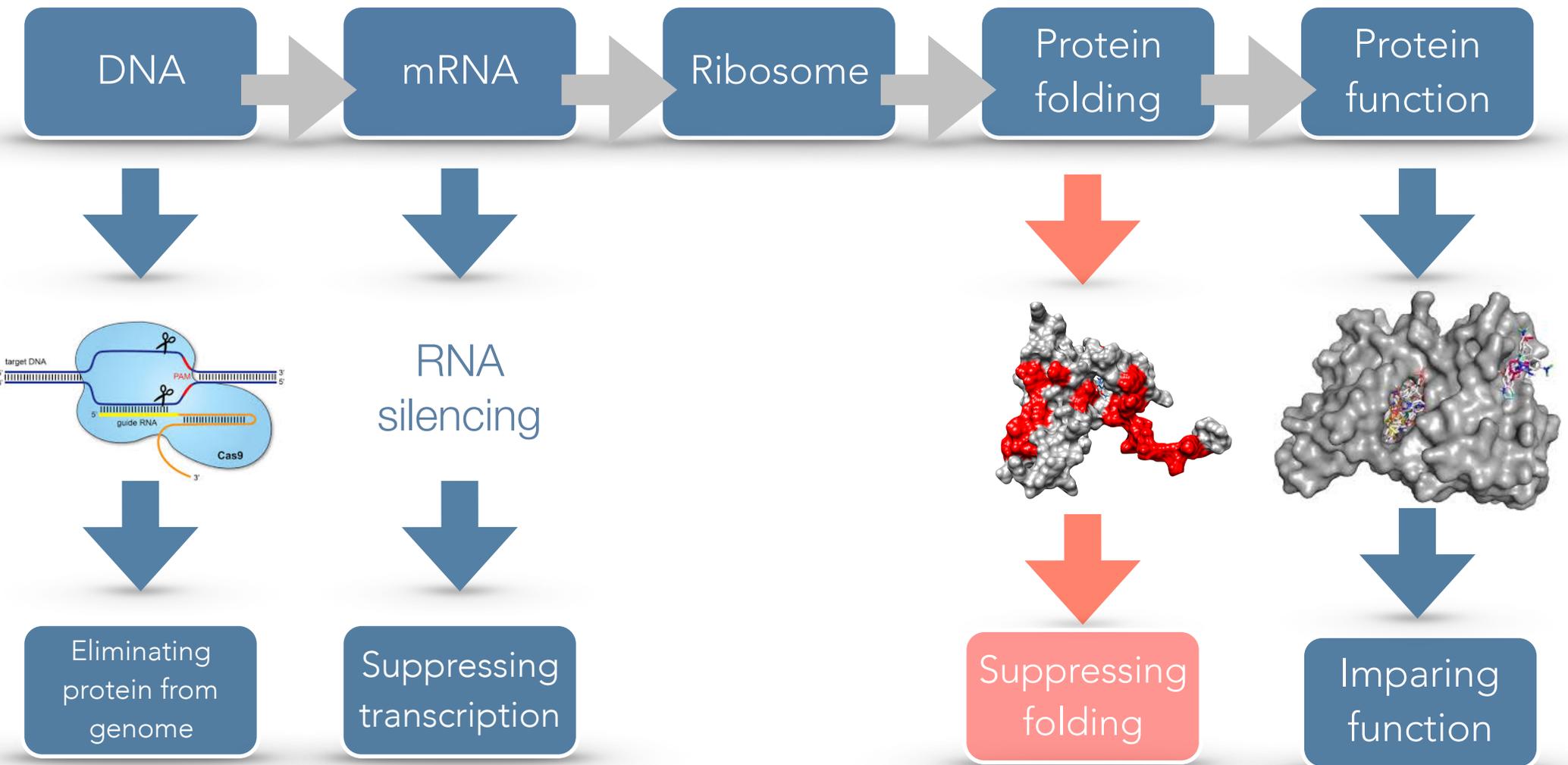
$$\mathcal{L} = \frac{1}{4g^2} G_{\mu\nu}^a G_{\mu\nu}^a + \sum_f \bar{\psi}_f (i\gamma^\mu \partial_\mu + m_f) \psi_f$$

where $G_{\mu\nu}^a \equiv \partial_\mu A_\nu^a - \partial_\nu A_\mu^a + gf_{abc} A_\mu^b A_\nu^c$
and $D_\mu \equiv \partial_\mu + i t^a A_\mu^a$
That's it!

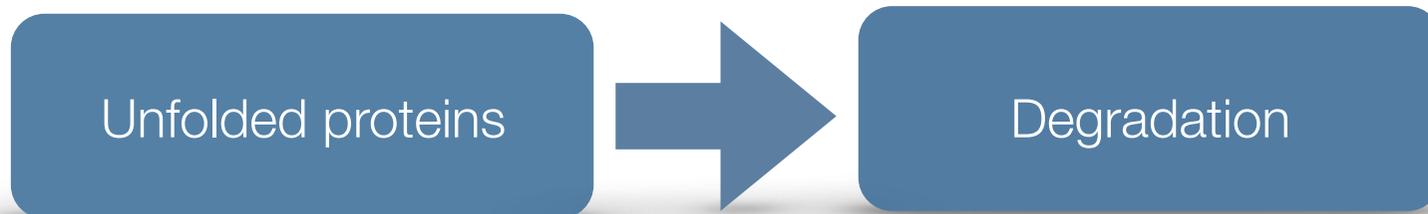
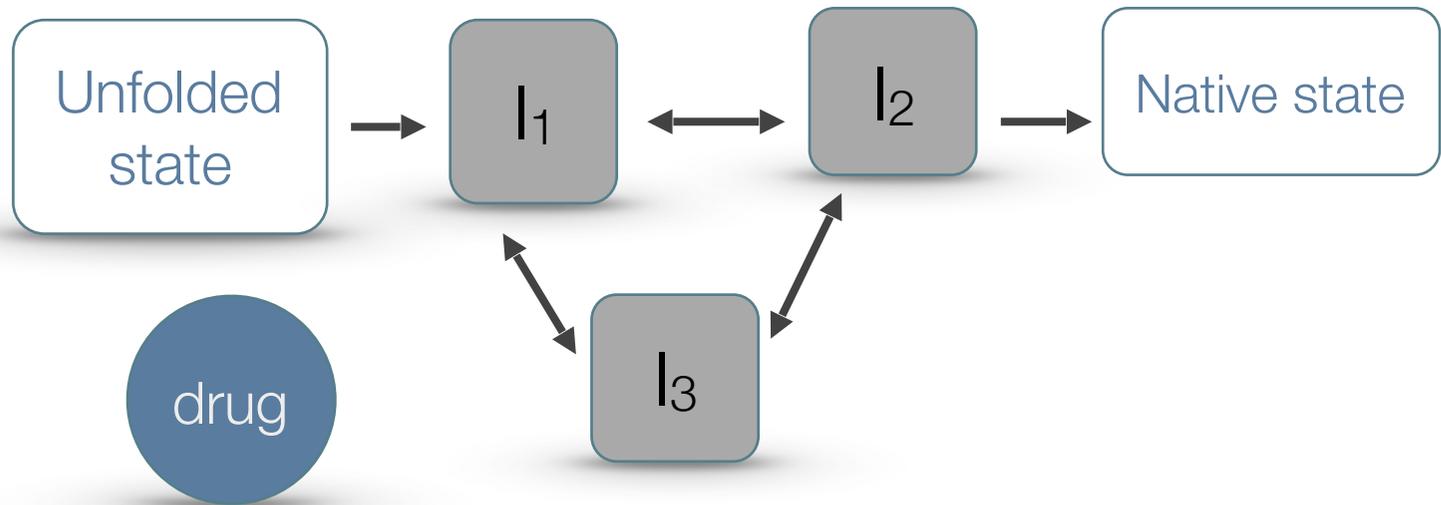


PHARMACOLOGICAL PROTEIN INACTIVATION BY FOLDING INTERMEDIATE TARGETING

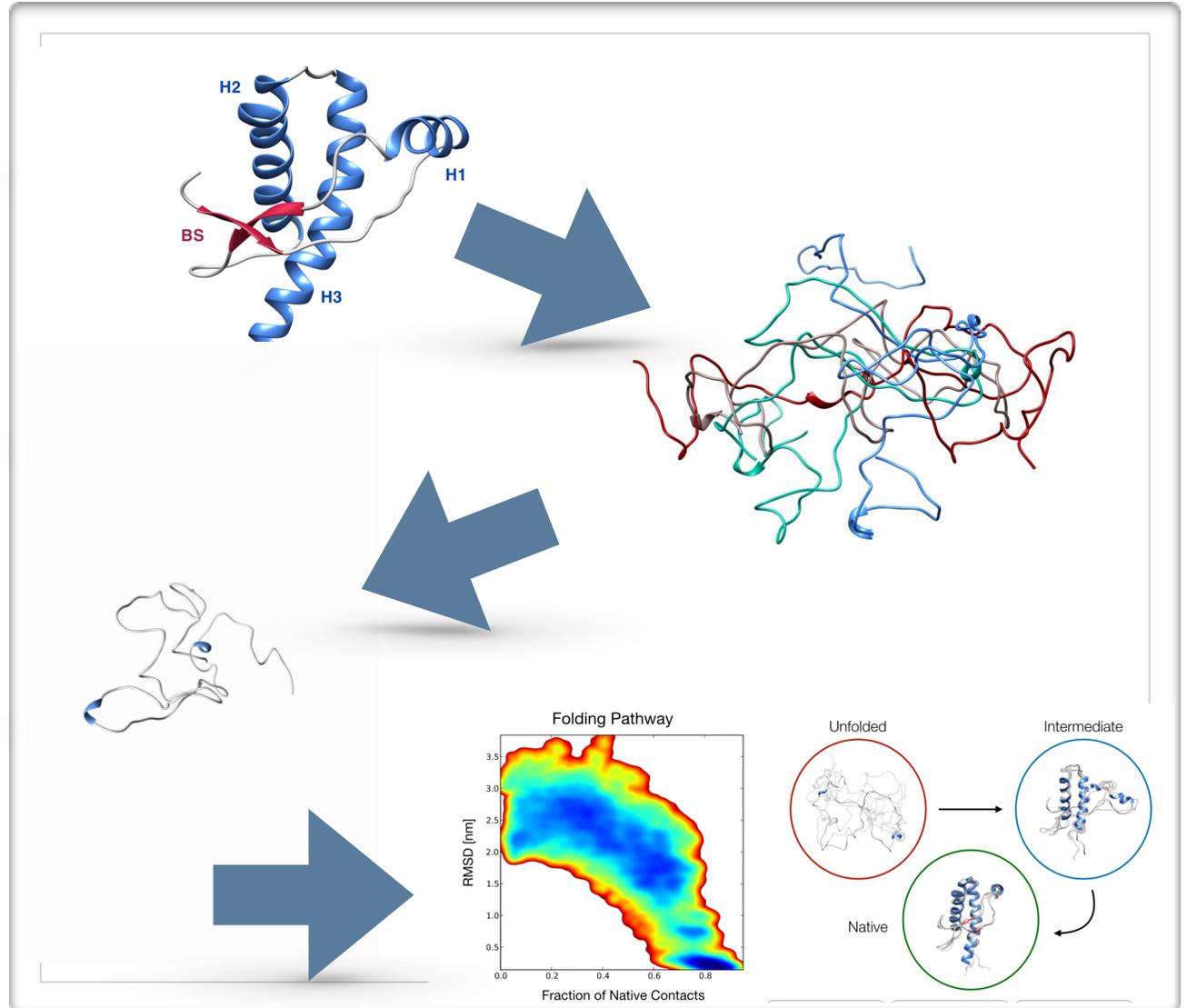
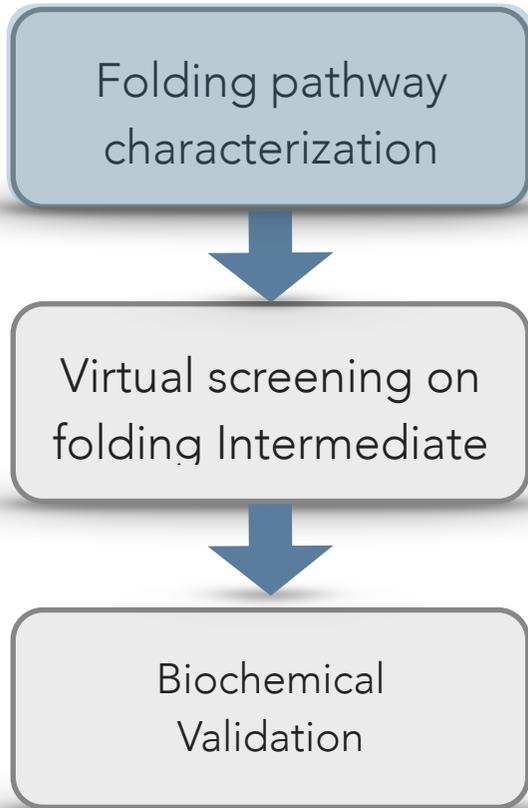
patent file # 102018000007535 (with E. Biasini)



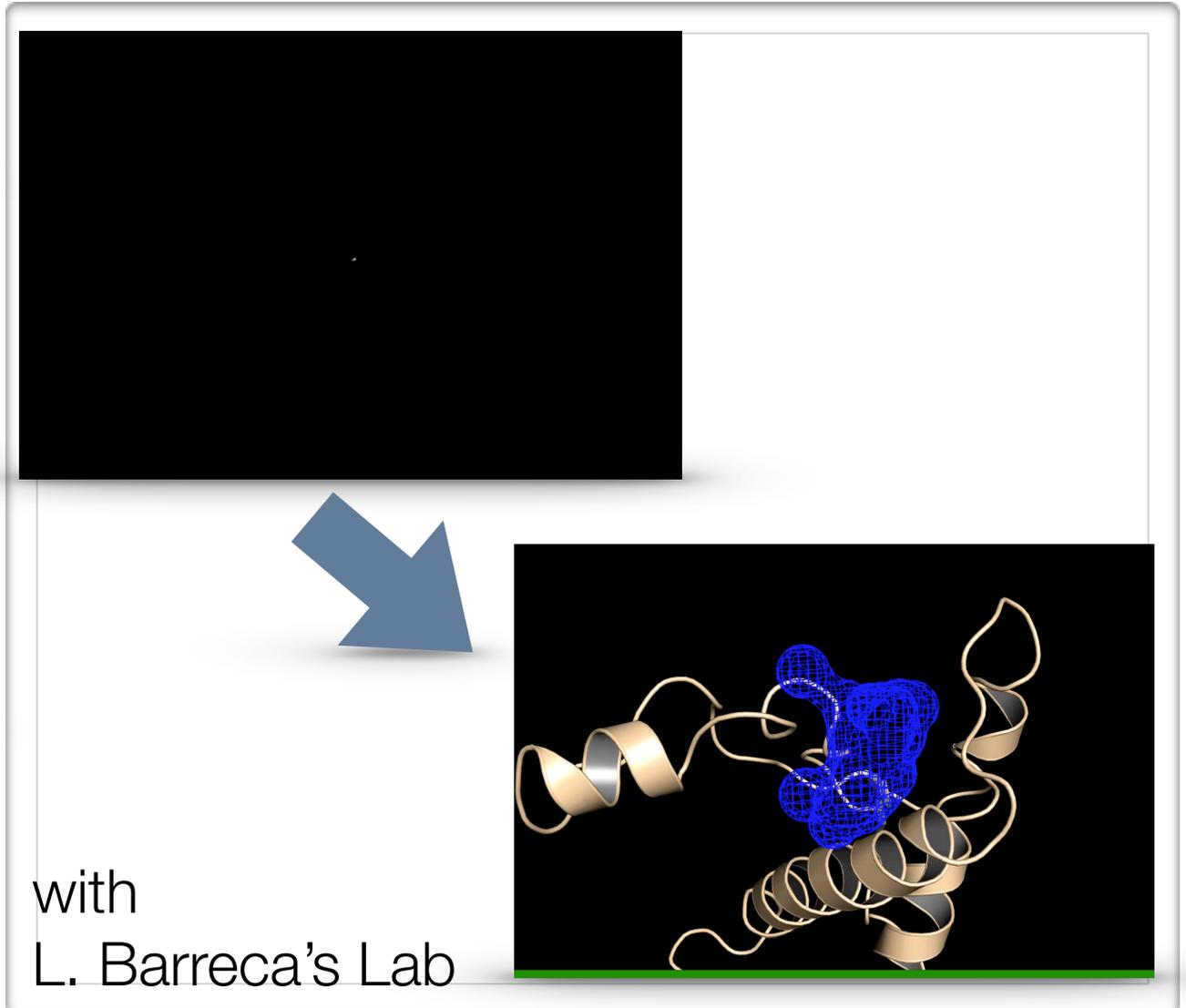
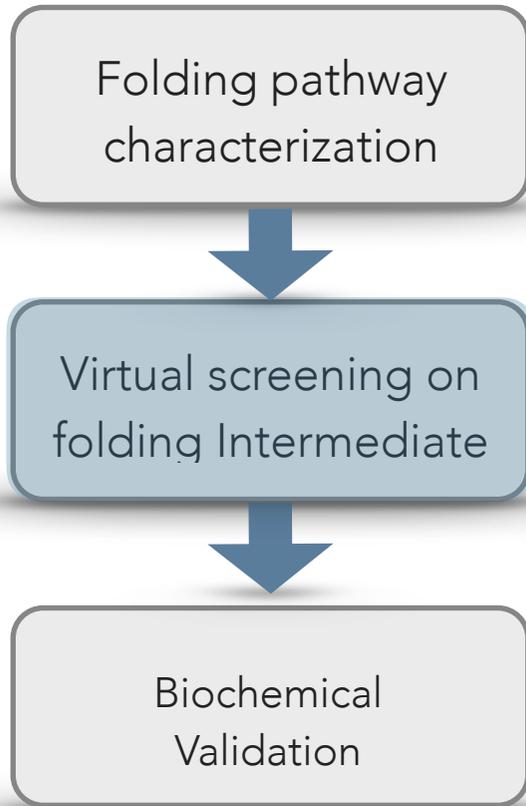
PHARMACOLOGICAL PROTEIN INACTIVATION BY FOLDING INTERMEDIATE TARGETING



PPI-FIT PIPELINE



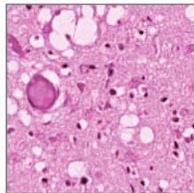
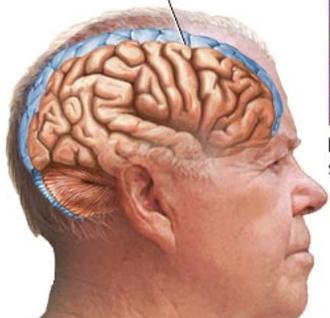
PPI-FIT PIPELINE



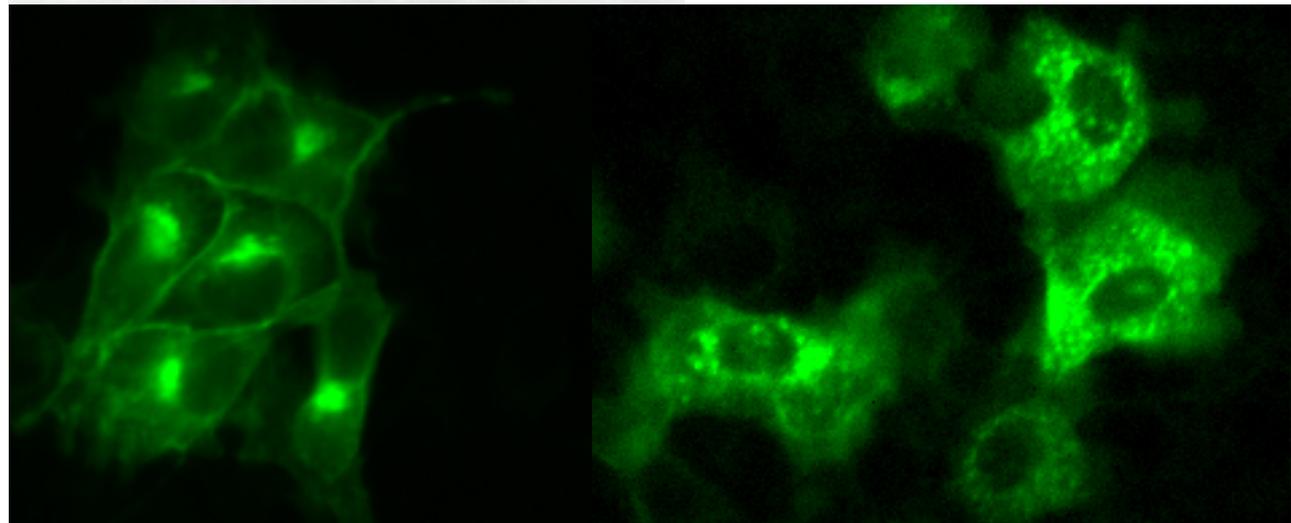
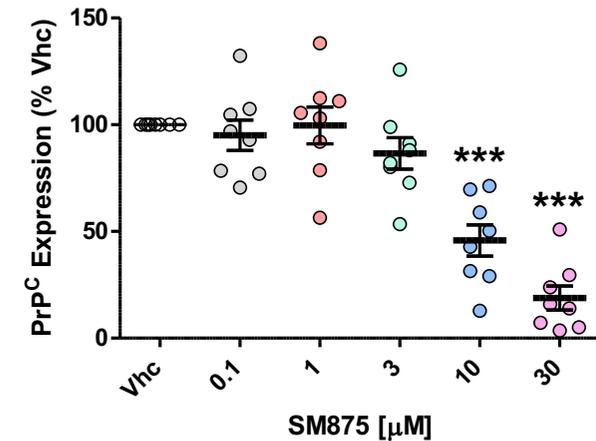
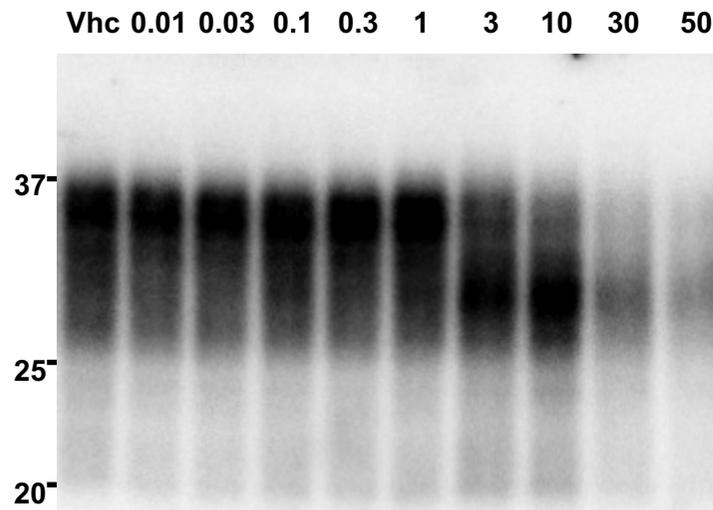
A FIRST VALIDATION

Inactivation of Cellular Prion protein

Brain shrinkage and deterioration occurs rapidly



Brain section showing spongiform pathology characteristic of Creutzfeldt-Jakob



PHARMACOLOGICAL PROTEIN INACTIVATION BY TARGETING FOLDING INTERMEDIATES

© Giovanni Spagnolli, Tania Massignan, Andrea Astolfi, Silvia Biggi, Paolo Brunelli, Michela Libergoli, Alan Ianeselli, Simone Orioli, Alberto Boldrini, Luca Terruzzi, Giulia Maietta, Marta Rigoli, Nuria Lopez Lorenzo, Leticia C. Fernandez, Laura Tosatto, Luise Linsenmeier, Beatrice Vignoli, Gianluca Petris, Dino Gasparotto, Maria Pennuto, Graziano Guella, Marco Canossa, Hermann Clemens Altmeppen, Graziano Lolli, Stefano Biressi, Manuel Martin Pastor, Jesús R. Requena, Ines Mancini, Maria Letizia Barreca, Pietro Faccioli, © Emiliano Biasini

doi: <https://doi.org/10.1101/2020.03.31.018069>



A few facts:

- * Founded in 2017 & funded with 2.4 MEUR by VERTIS SCR in 2019.
- * Spinoff of U.Trento, U.Perugia and INFN
- * Exclusive license of the PPI-FIT technology
- * Developing internal pipelines as well R&D service with a major Pharma
- * International academic partnerships in USA and Europe

CEO:

Dr. Lidia Pieri

Chief Scientists:

Dr. Giovanni Spagnolli, Dr. Tania Massignan

Business Developer:

Dr. Federica Mantovani

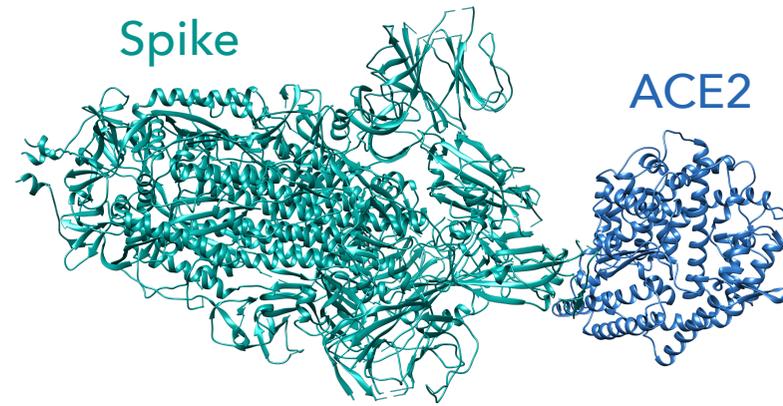
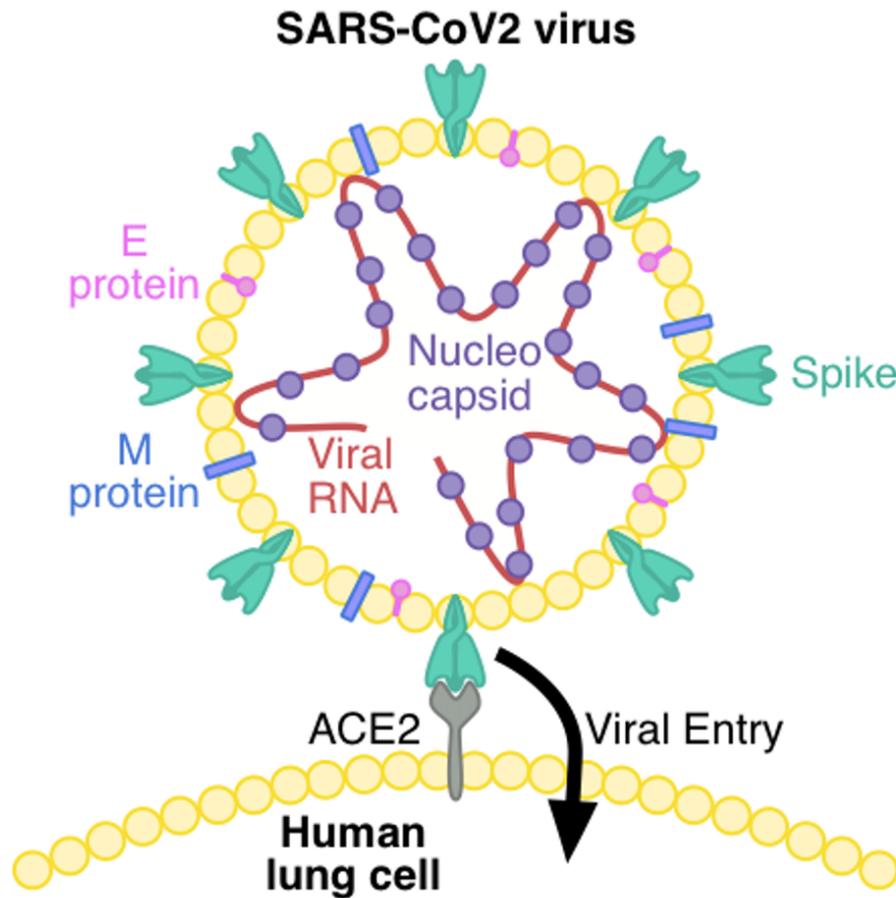
Scientific Staff:

Dott. Alberto Boldrin, Dott. Luca Terruzzi, Dott. Enrica Colasurdo,
Dr. Serena Bonifati , Dr. Andrea Astolfi

Scientific Consultant:

Prof. Vincenzo Summa,

SARS-CoV-2 Replication



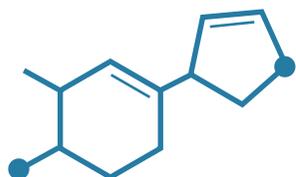
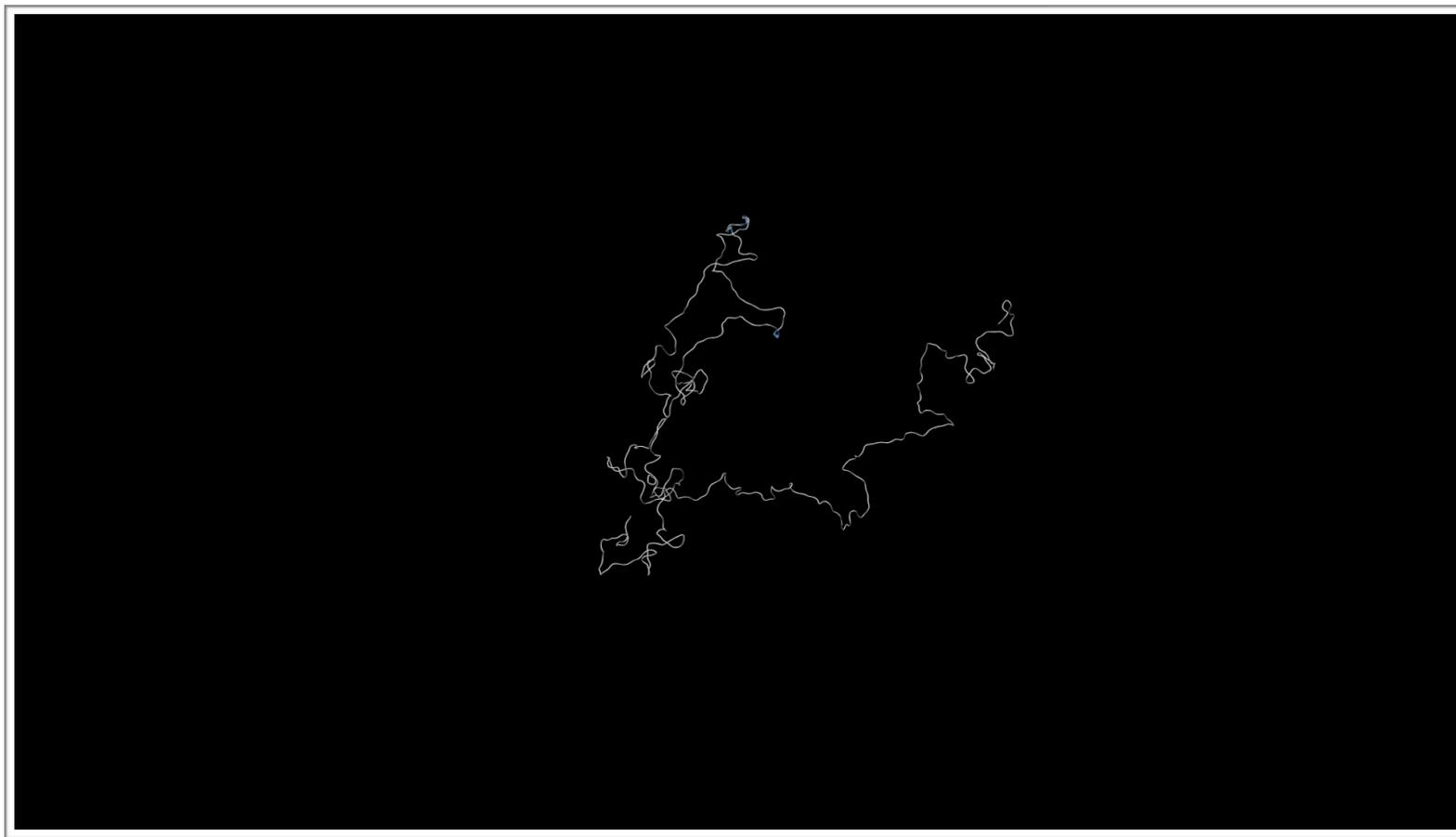
Goals:

Repurposing of approved drugs!
Looking for suppressors of ACE2
expression levels

Figure taken from:

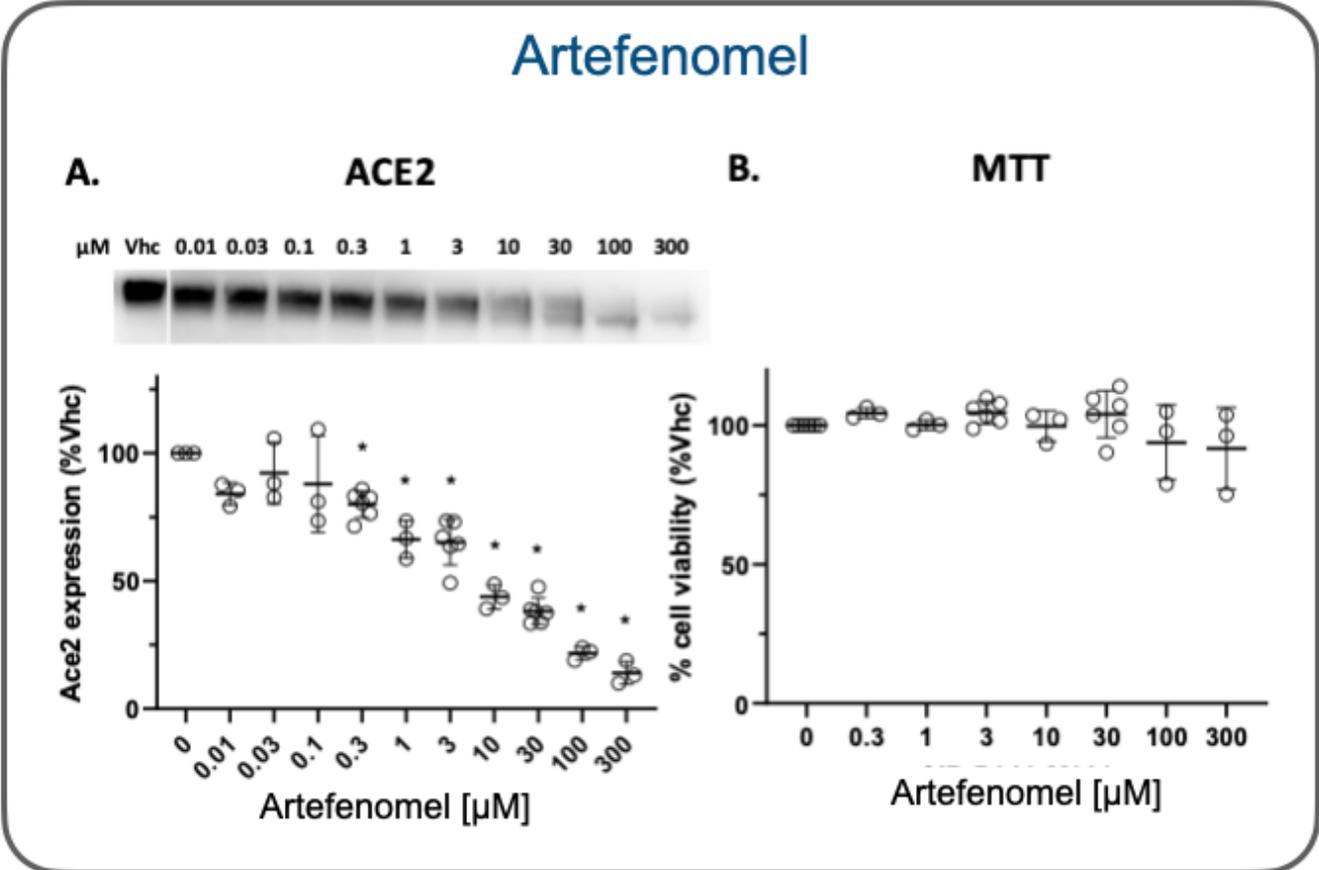
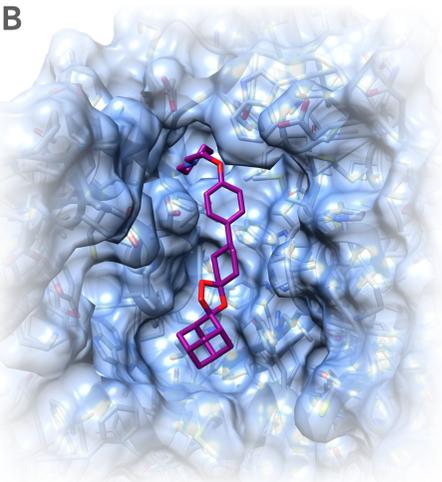
<https://theconversation.com/where-are-we-at-with-developing-a-vaccine-for-coronavirus-134784>

PPI-FIT ON ACE2



Out of 9000 candidates, we found 35 molecules binding in-silico the intermediate. Validation experiments on cellular bio-assays are ongoing.

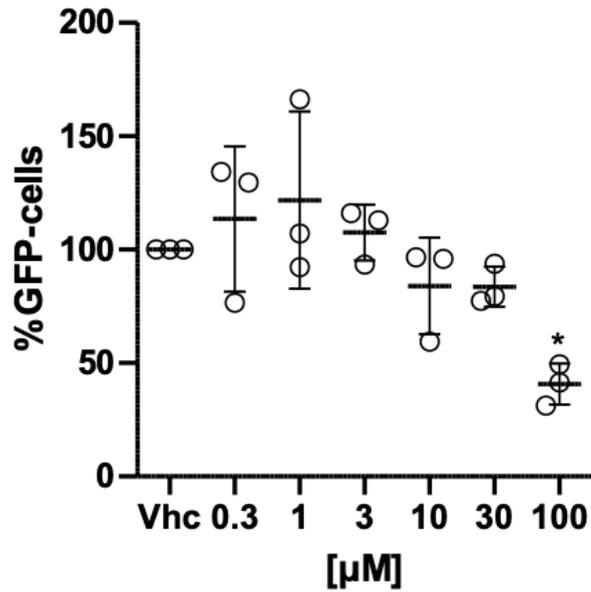
DOSE-DEPENDENT RESPONSE



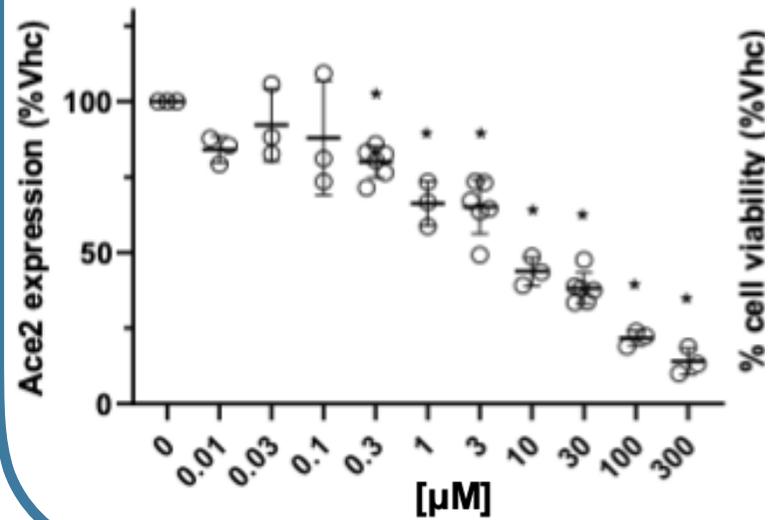
ACTION ON A PSEUDOVIRAL VECTOR



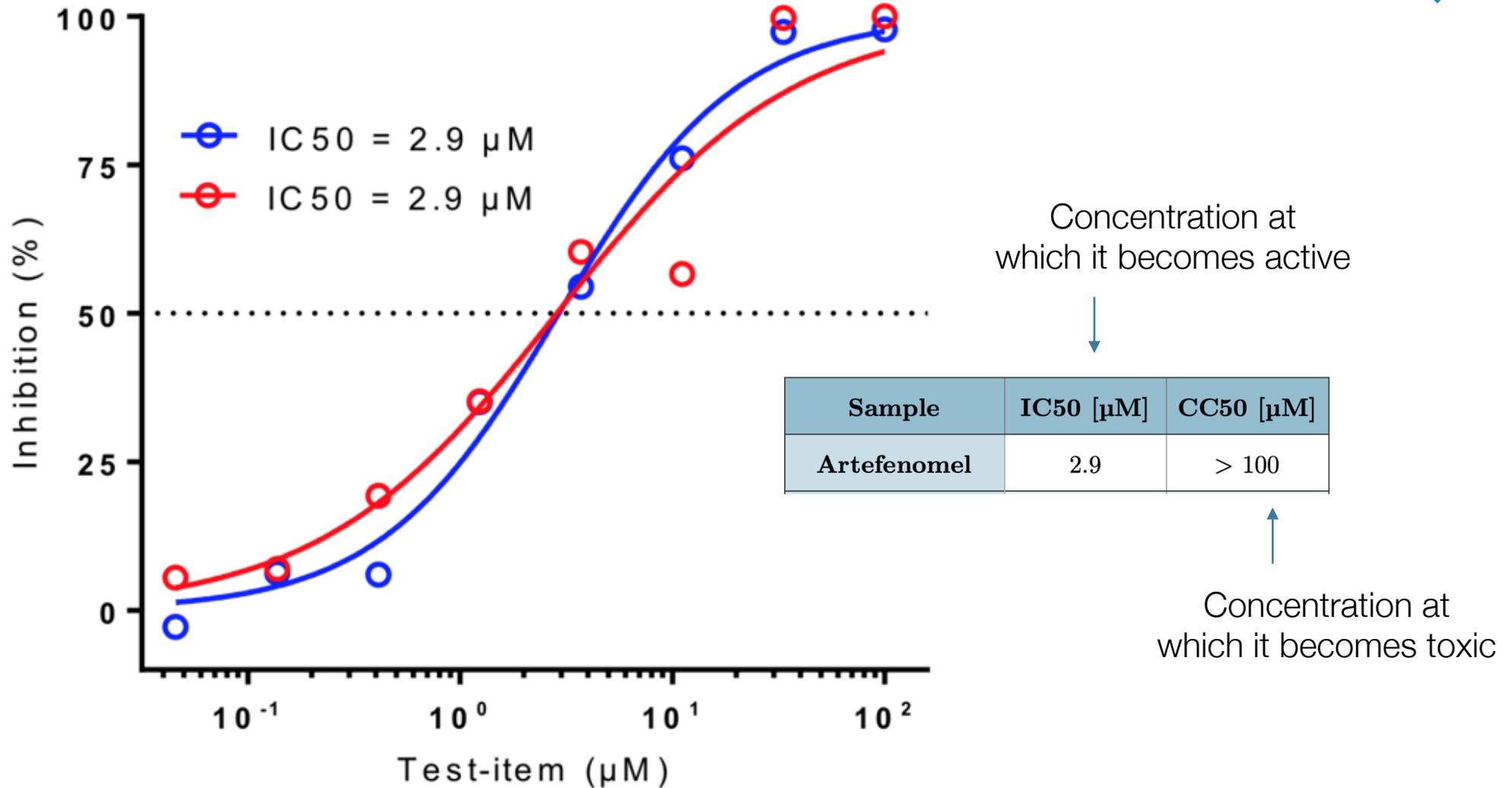
Pseudo-virus Transduction Efficiency



ACE-2 cellular expression level



ANTI-VIRAL ACTIVITY AGAINST LIVE SARS-COV2



This value is *in principle* compatible with the maximum tolerated dose in humans. More to follow...



NEW DIRECTIONS

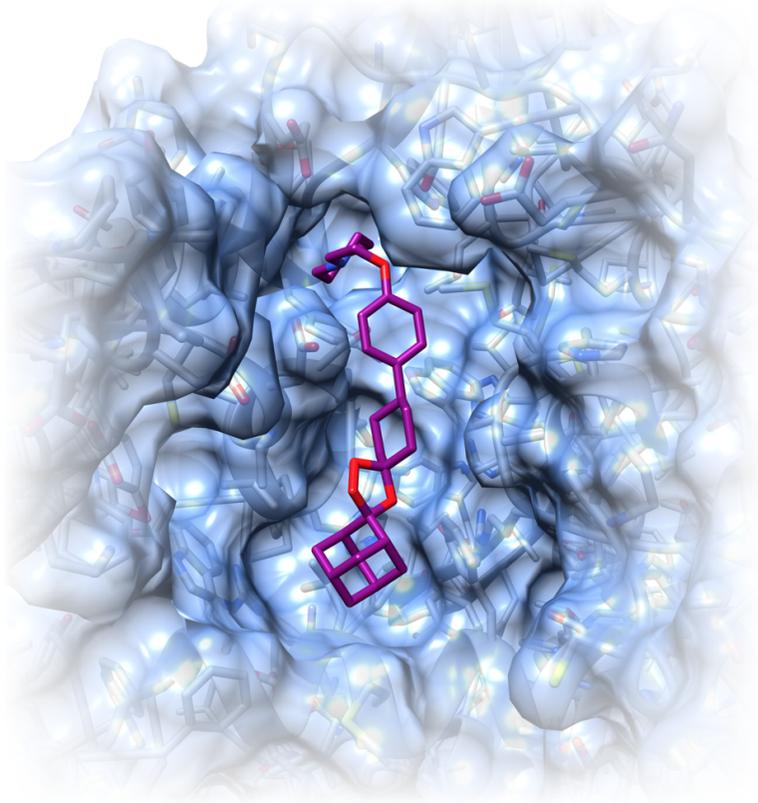
SPACE IS THE NEXT FRONTIER!



Partners:

U. Trento, Space Pharma, CJD Foundation (Israel), U. Tel Aviv, U. Santiago de Compostela, INFN

A MAIN LIMITING FACTOR



Impossible to crystallize
folding intermediates
on Earth



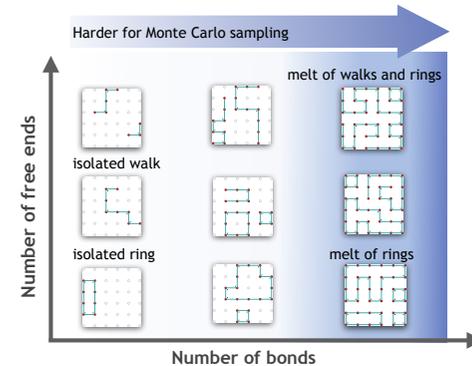
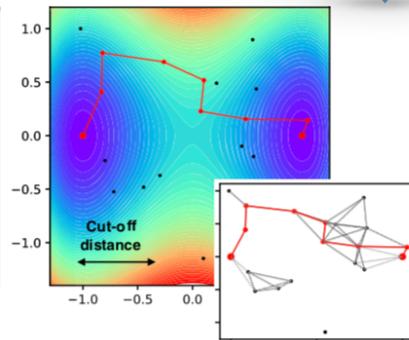
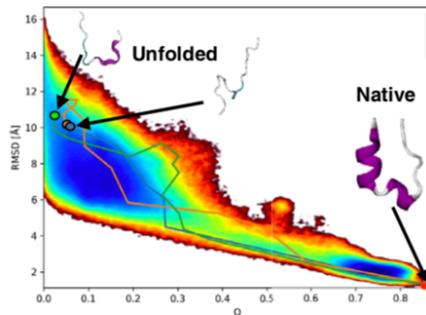
Microgravity
conditions may
provide the solution!

Towards the next generation of molecular simulations

Quantum Computing

Theoretical Physics

Artificial Intelligence



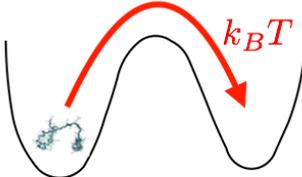
Partners:

U. Trento, Q@Trento, INFN, SISSA, BEC-CNR

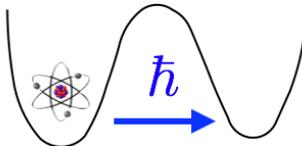
Take home messages

1. Fundamental science breeds innovation

Thermal activation ($\beta = (k_B T)^{-1}$)


$$P(x_f, t|x_i) = \int_{x_i}^{x_f} \mathcal{D}q e^{-\frac{\beta}{4M\gamma} \int_0^t d\tau (M\ddot{q} + M\gamma\dot{q} + \nabla U(q))^2}$$

Quantum tunneling


$$K_E(x_f, t|x_i) = \int_{x_i}^{x_f} \mathcal{D}q e^{-\frac{1}{\hbar} \int_0^t d\tau (\frac{M}{2}\dot{q}^2 + U(q))}$$

2. Cross-disciplinarity is key to tackle complexity.

But in return it requires to redefine expectations

3. The importance of swimming in muddy waters

Acknowledgments



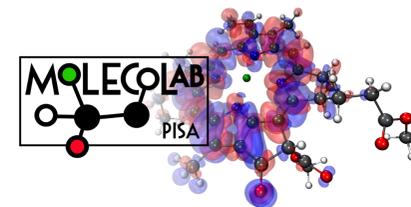
Trento Institute for
Fundamental Physics
and Applications



Italy:

Trento: *E. Biasini*, *A. Ianeselli* (2014-2017), *G. Spagnoli* *S. A Beccara* (2009-2017), *S. Orioli* (2014-2018), *E. Schneider* (2012-2015), *M. Carli* (2017), *M. Turelli* (2018), *F. Mascherpa* (2014), *G. Garberoglio*, *F. Pederiva*, *M. Sega*, *R. Covino* (2012-2015),

Pisa: *B. Mennucci*, *L. Cupellini*, *S. Jurinovich*



Perugia: *L. Barreca*

SISSA: *C. Micheletti*, *A. Laio*

Europe:

U. Zurich: *B. Schuler*

U. Compostela: *J. Raquena*

CEA-Saclay: *H. Orland*

USA: U. Maryland: *P. Wintrode*

U. Mass.: *A. Gershenson*

