

# Accelerator Physics - Introduction

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Heavy Ion Therapy School, May 2021



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# Contents

- History and motivation for accelerators
- Beam properties – transverse emittance
- RF acceleration, longitudinal dynamics, phase stability
- Cyclotrons and synchrotrons
- Strong focusing, transverse dynamics beam transport
- *Beam instrumentation* – Wednesday (20 min)
- *Sarajevo Linac project* – Friday (20 min)

## Other presentations this week:

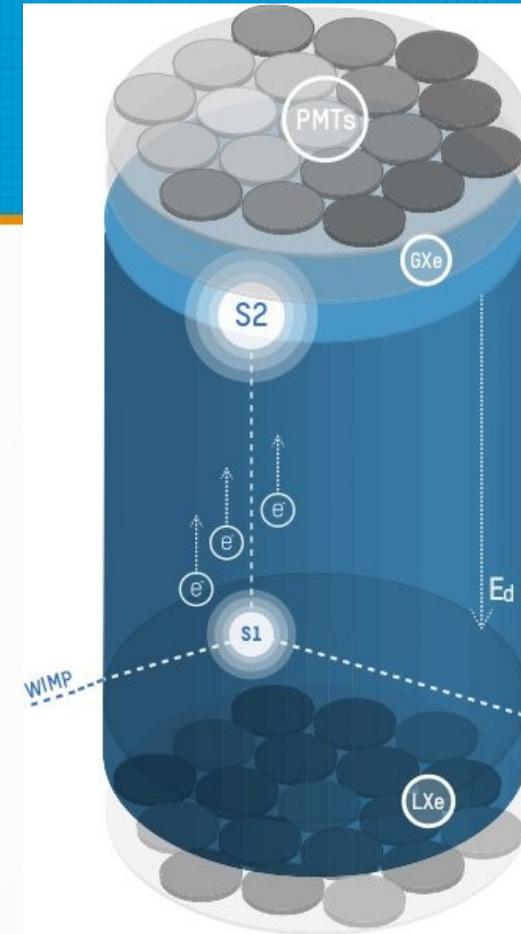
- Introduction to accelerators and medical machines (Maurizio, 1.5 h)
- Linear accelerators (Giovanni, 45 min)
- Injection to synchrotrons (Elena, 20 min)
- Beam extraction (Rebecca, 30 min)
- Ion sources (Nadia, 30 min)
- Gantries and Beam Delivery (Elena, 45 min)
- Low energy accelerators (Milko, 45 min)
- Sarajevo Linac project:
  - Ion Beam Analysis (Fehima, 15 min)
  - Low energy beam transport simulations (Benjamin, 10 min)

# Methods of science

- **Observation of nature:**
  - Astronomy – purely observational science
  - Physics – e.g. Dark Matter search (XENON)
- **Controlled experiments:**
  - Many types of experiment, rich methodology
  - Various tools, including accelerators



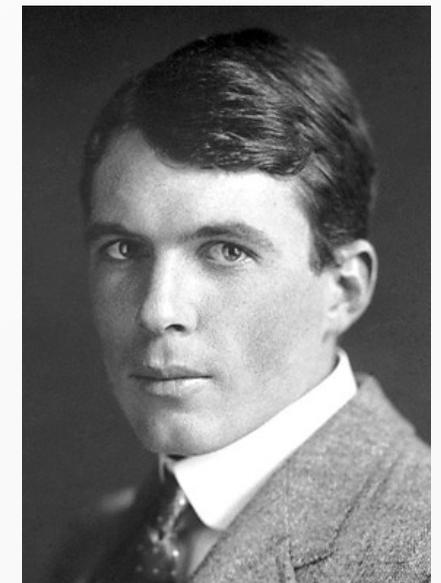
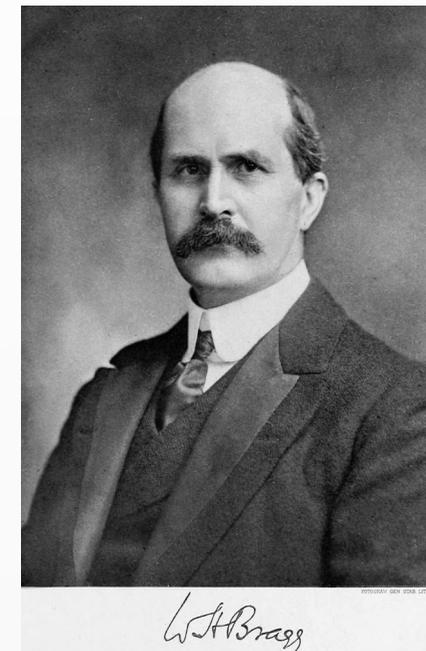
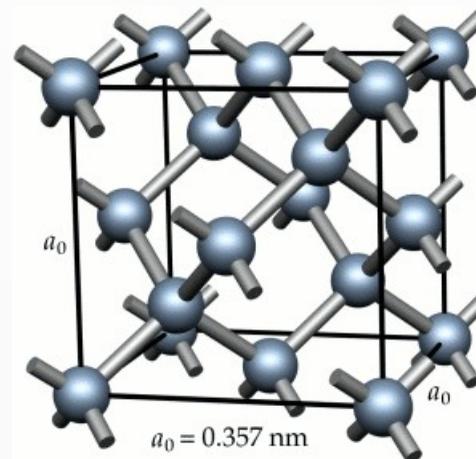
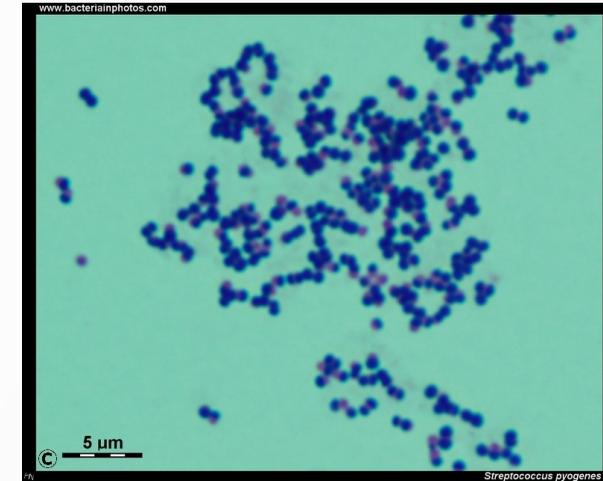
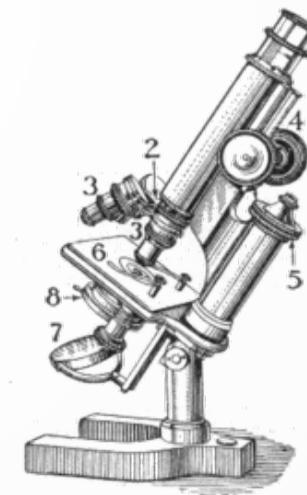
May 18,



e.g. lasers are very important tools of modern physics:  
quantum mechanics, atomic physics, ultra-fast chemistry, etc

# Microscopes

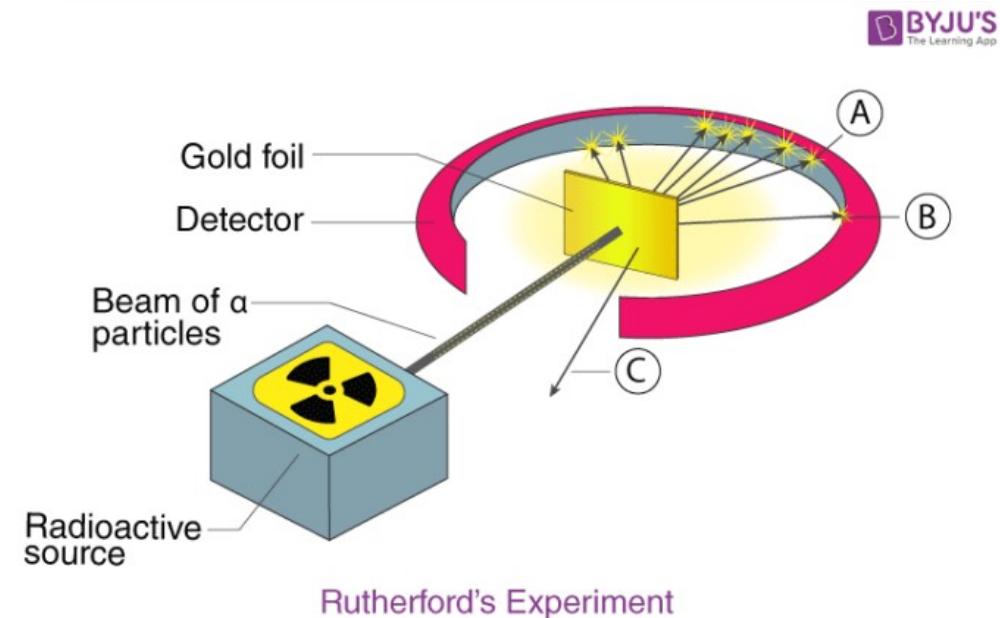
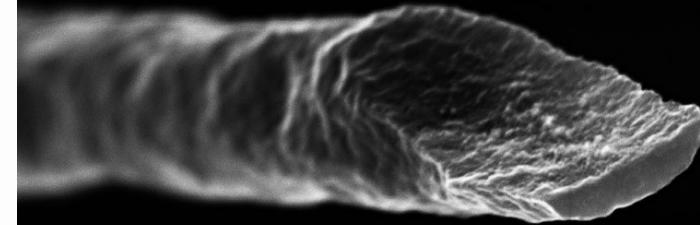
- Optical microscopes invented in 17<sup>th</sup> century
- Resolution 200 nm determined by wavelength length (diffraction limit  $\lambda/2$ , optical  $\lambda=400-800$  nm)
- Bacteria size  $\sim 0.5-2$   $\mu\text{m}$  – seen by light microscopes, but SARS-Cov-2 virus size  $\sim 100$  nm
- Crystalline structures need sub-nm resolution
- One could use shorter wavelengths: X-ray Crystallography (not microscopy – difficult optics)
- Bragg's law 1912, by William Henry - father and William Lorentz - son
- X-ray  $\lambda=0.1$  nm
- BTW Bragg father discovered **Bragg peak** in 1903



# Rutherford experiment

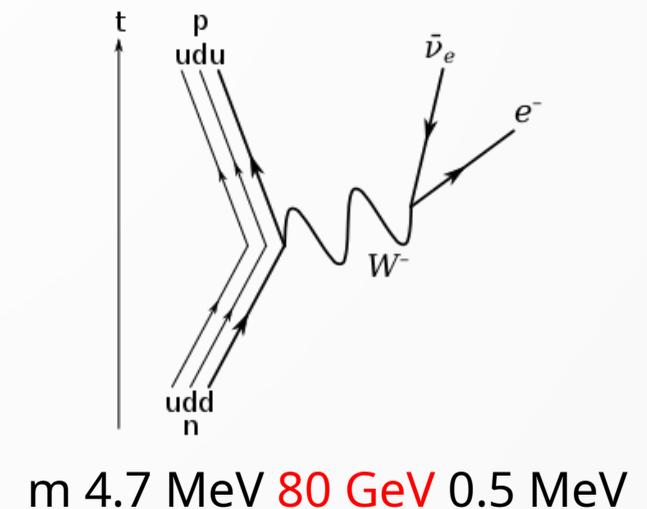
- Electrons are also waves (de Broglie, 1924), use them instead of light (E. Ruska, 1931)
- Electron microscopes can reach: 0.1 nm resolution – atom size; they work on normal objects (not only crystals)
- What if we want to **look inside atoms**? Photons and electrons are interacting with other electrons in the atom.
- Rutherford experiment (1908) – use alpha particles – they are heavy and penetrate through electrons
- Rutherford (+Geiger+Marsden) used Radon-222, which decays emitting alpha with  $E_k=5.5$  MeV
- Experiment lead to discovery of nucleus and further to discovery of protons (1919) and neutrons (1932)
- Note: Beam of particles can be generated from radioactive source, but we have little control on it
- Positrons, muons – discovered in cosmic radiation

Carbon fibre damaged by SPS beam, 5-10  $\mu\text{m}$  diameter.



# Why do we like $E=mc^2$ ?

- Since Rutheford – enormous **progress in nuclear physics thanks to accelerators**
- E.g. discovery of new isotopes, often short-living and non existing in nature, or new particles
- *New matter is produced from energy*
- **Why creation of something what does not exist in nature is important ?**
  - Because those particles really existed during Big Bang shaping our World
  - Because they exist now, in form of virtual particles
- Virtual particles born from vacuum for a short moment
  - Disappear after  $\Delta t \leq h/2E$
  - Unless they interact/decay as e.g. in  $\beta$ -decay ( $^{60}\text{Co}$ )

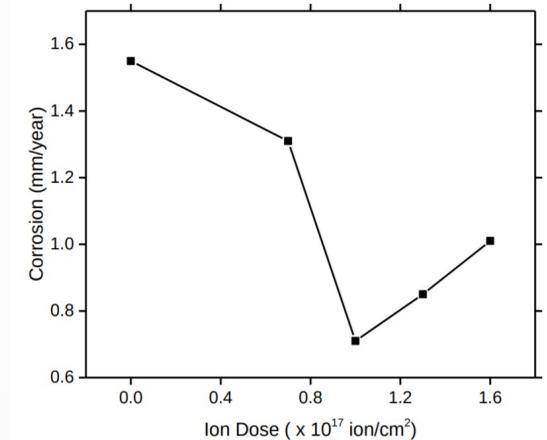


$$E=mc^2$$

We have to do accelerator experiments  
if we want to understand the world around us

# Accelerators in industry and medicine

- Ion implantation
- Ion beam analysis
- Electron beam material processing
- Radioisotope production (also medical)
- Neutron generation
- Radiotherapy, radiosurgery
- Noninvasive diagnostics



Corrosion for implanted 304 SS  
(A. Nikmah et al 2019 IOP Conf. Ser.:  
Mater. Sci. Eng. 515 012018)

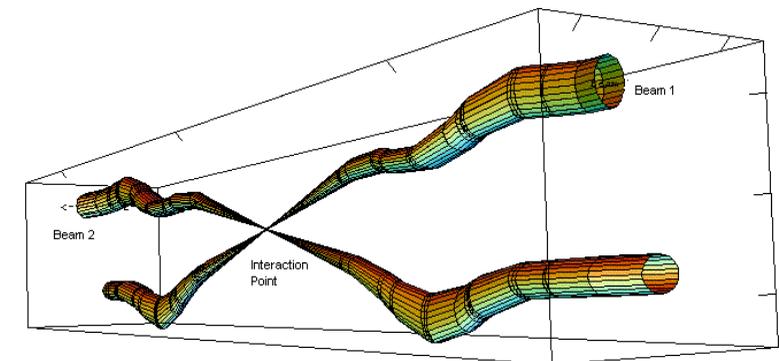
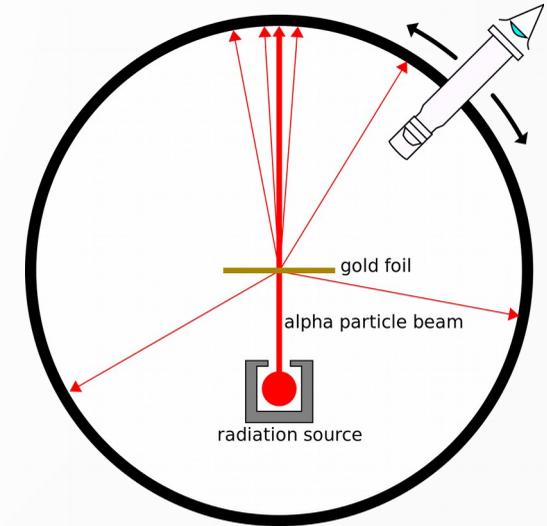


Öztürk, O. (2014). "Structural and Magnetic Characterization of Nitrogen Ion Implanted Stainless Steel and CoCrMo Alloys."

Sapinski, Accelerator physics

# What is a beam?

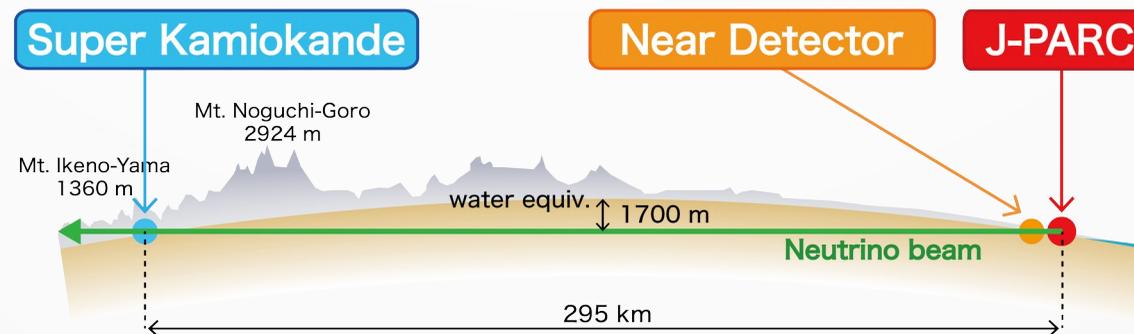
- Accelerators are producing beams, so what is a beam?
- **An ensemble of particles moving in the same direction**
- Characterized by:
  - Particle type (usually monoparticle)
  - Intensity
  - Particle energy and energy spread
  - Transverse size and divergence (emittance)



Relative beam sizes around IP1 (Atlas) in collision

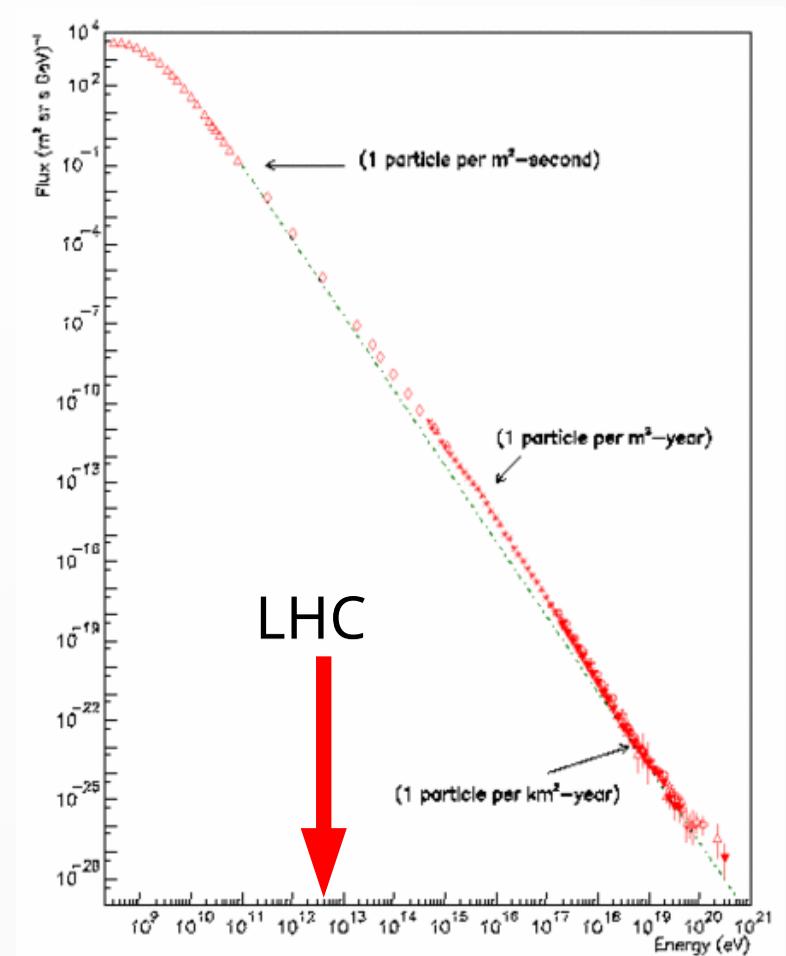
# Particle types

- Electrons (the easiest, e.g. X-ray tube), positrons
- Protons, antiprotons
- Ions, e.g.  ${}^4\text{He}^{2+}$ ,  ${}^{12}\text{C}^{6+}$ , all isotopes and charge states,
- also exotic and radioactive beams eg.  ${}^6\text{He}^{2+}$  ( $\tau_{1/2}=0.8\text{s}$ ) and negative ions (eg.  $\text{H}^-$ )
- Compound particles eg.  $\text{CH}_3^+$
- Neutral particles (eg. neutrons, neutrinos or photons) are produced as secondary beams



# Particle energies

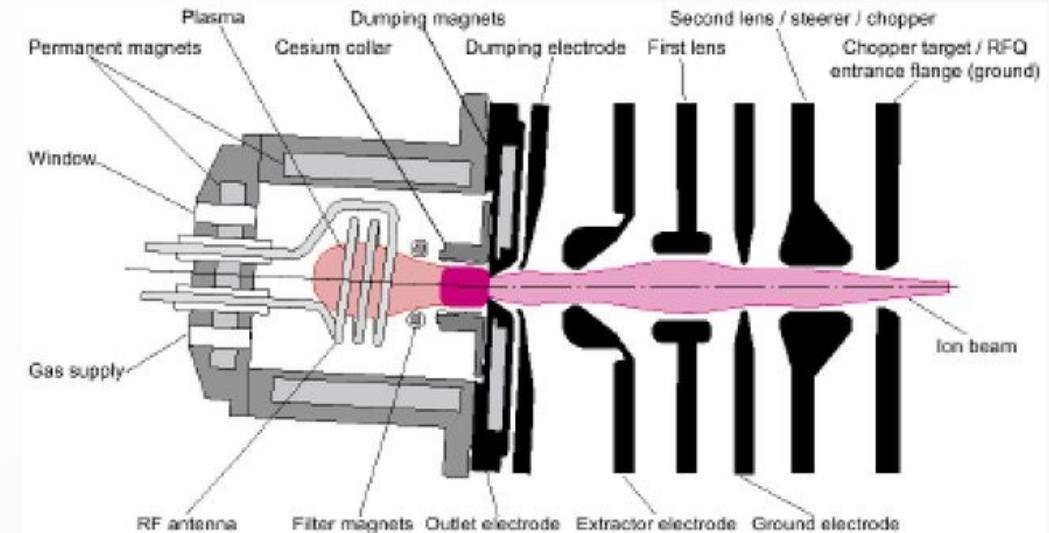
- Energy is conveniently expressed in **electron-volt (eV)** and for ions in eV/u (per nucleon)
- For some studies particles are decelerated down to meV energies and trapped (e.g. antimatter)
- The highest beam energy (per particle) is at LHC: 6.5 TeV proton beams
- Total energy stored in beams: 362 MJ (equivalent of 77,4 kg TNT!)
- Interestingly cosmic rays reach much higher energies: so called cosmic accelerators are probably driven by expanding magnetic field of exploding stars (Fermi acceleration)
- Beams are not monoenergetic; typically we talk about **momentum spread**  $\Delta p/p \sim 10^{-3}$



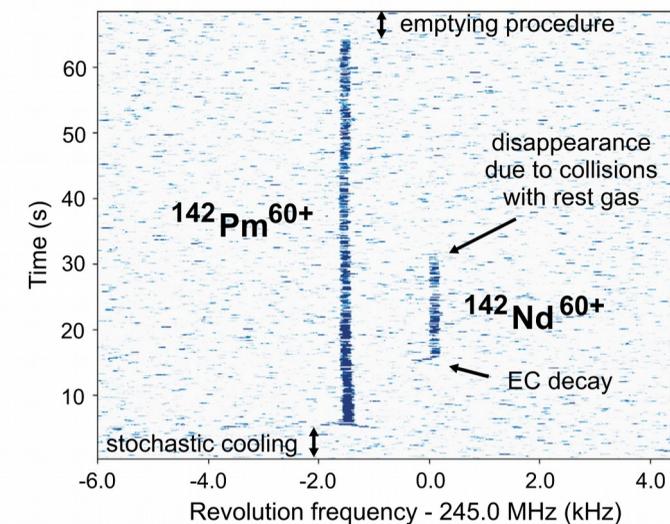
Oh-My-God particle, 50 J

# Beam intensities and time structure

- Beam is typically produced as continuous from the ion source and **bunched** in the accelerating structures
- Therefore ion **source intensity is given mA** of DC current; in the linac it is peak current and pulse duration; in **synchrotron it is easier to talk about number of circulating particles**
- Ion source can reach 65 mA currents, numbers of circulating particles can be in range 1-10<sup>14</sup>
- Bunch length: from DC to 1 ps



M. Steck and Y.A. Litvinov / Progress in Particle and Nuclear Physics 115 (2020) 103811

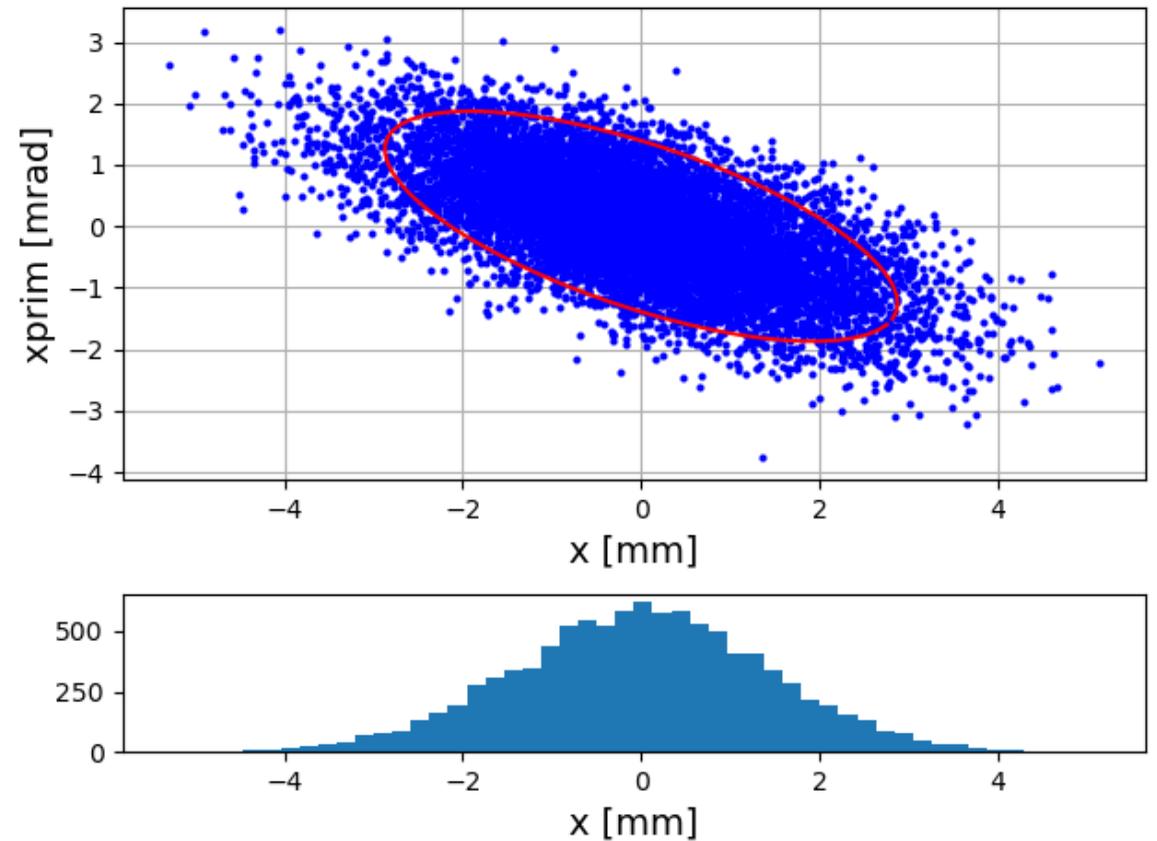


# Transverse size

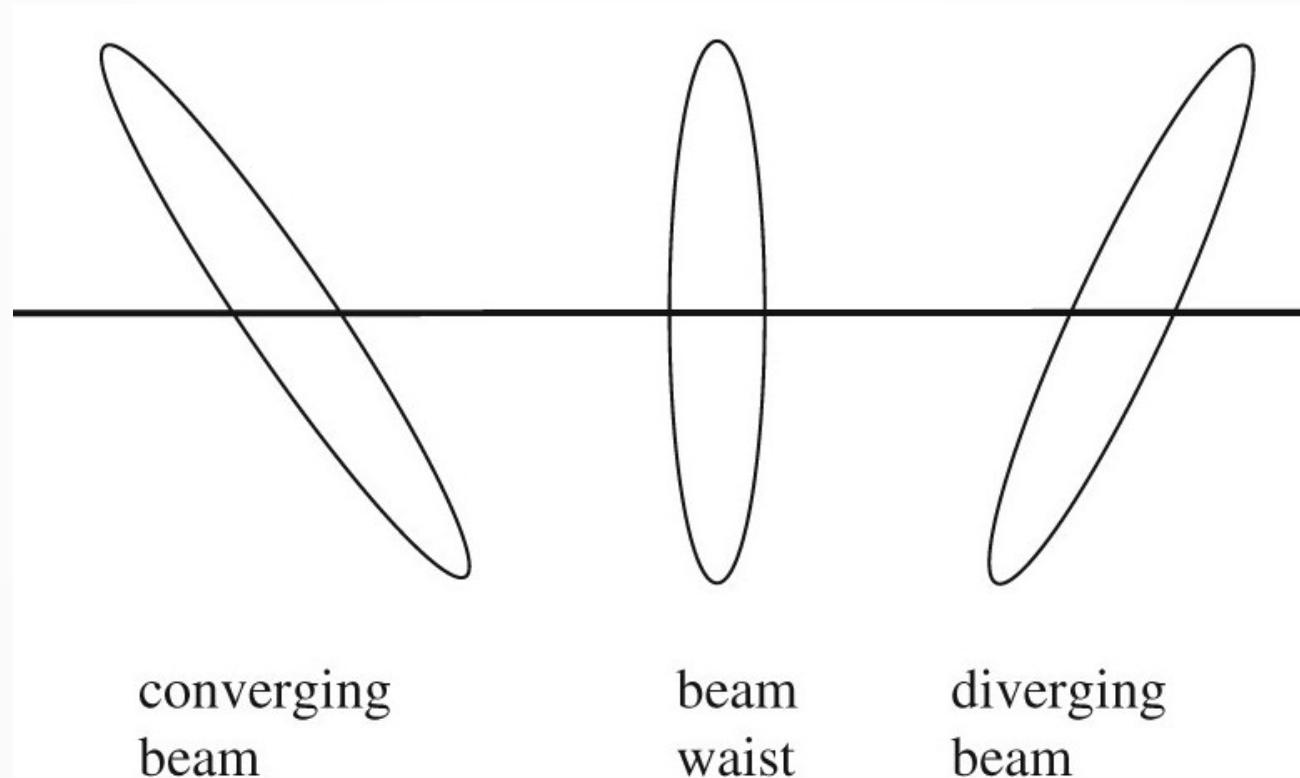
- Transverse sizes of beams can vary:
  - nanometers ( $10^{-9}$  m)- electron beam lithography
  - micrometers ( $10^{-6}$  m) – synchrotron light sources
  - millimeters ( $10^{-3}$  m) – eg. LHC
  - centimeters ( $10^{-2}$  m) – hadron therapy synchrotrons
  - meters – neutron, neutrino beams
- Beam size changes when traveling through accelerator
  - for instance it is usually focused on target
- It is better to use about beam **emittance**

# Phase space and emittance

- **Beam phase space** is defined by its transverse position ( $x$ ) and divergence ( $x'$ )
- Both distributions have usually approximately gaussian shape
- The surface of the ellipse containing 95% of the beam particles is called **emittance** ( $\epsilon_{95\%}$ )
- People also use RMS-emittance ( $\epsilon_{\text{RMS}}$ ), surface of ellipse containing 1 Root Mean Square (RMS) of the particles (40% for 2D gaussian distribution)
- In lack of acceleration and dissipative processes emittance is constant (Liouville's theorem)  $\rightarrow$  ion source must produce good emittance as it cannot be (easily) decreased

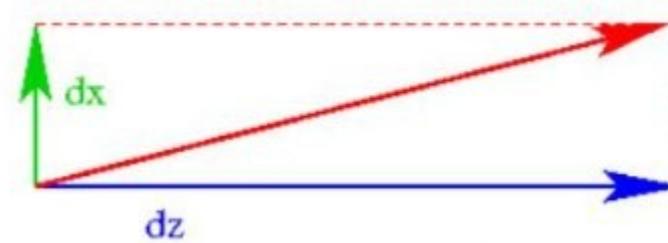
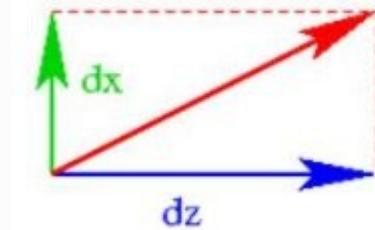


# Understanding the beam phase space

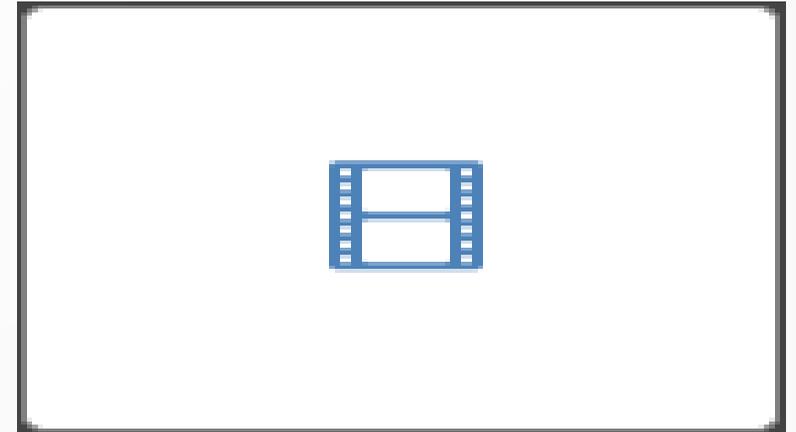


# Emittance and acceleration

- During acceleration (energy ramp) the longitudinal momentum increases while transverse remains the same



- Therefore the divergence of particles decreases, so the emittance shrinks!
- **Normalized emittance is conserved during acceleration:**  
$$\epsilon_n = \beta\gamma\epsilon_{RMS}$$
- Units: [mm\*mrad], [ $\pi$ \*mm\*mrad]
- Typical values for medical ion beams: 0.5-1.0 [mm\*mrad]
- Synchrotron light sources emittance reach  $\sim 1$  nm\*mrad



**I hope at this point you understand the role of accelerators in our civilization and variety of beams produced.**

**You got familiar with a concept of beam emittance.**

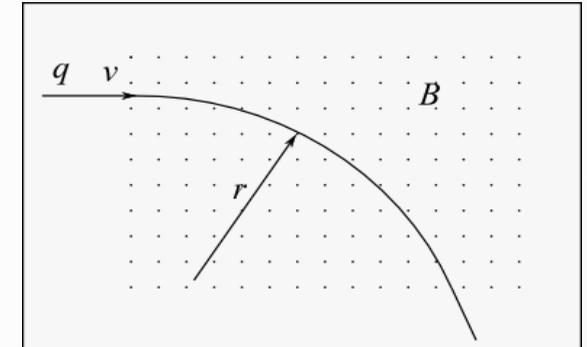
**Let's go to some details.**

# Acceleration techniques

- Which field to use for the acceleration?

**electric** force:  $\underline{F} = q\underline{E}$  - acts **along the field lines**

**magnetic** force:  $\underline{F} = q(\underline{v} \times \underline{B})$  - acts **perpendicular to field lines**  
and to **particle velocity - no acceleration**



- Force magnitudes:

- electric: 20 MV/m(\*),  $F = 3.2$  pN

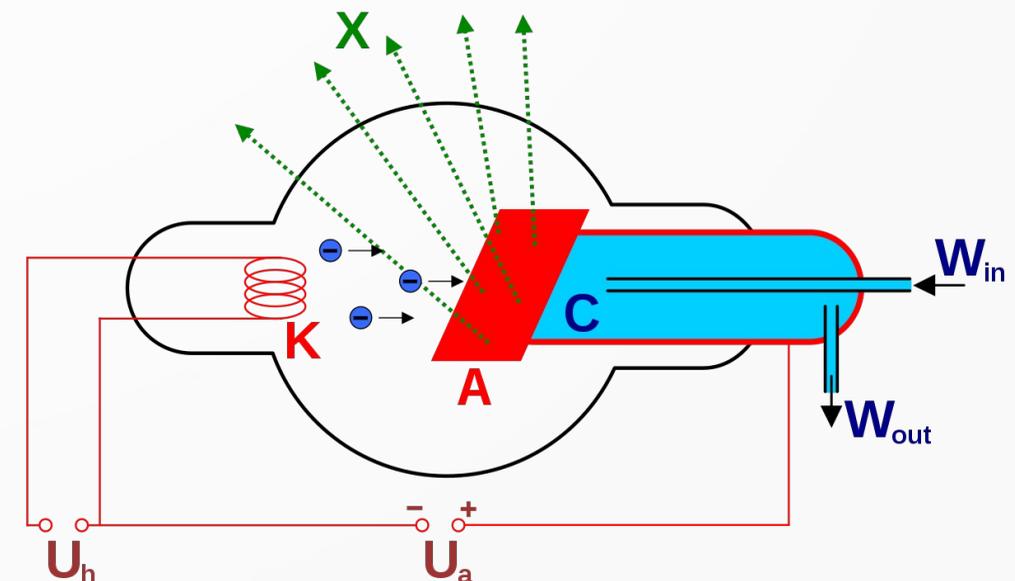
- magnetic: 1.5 T(\*),  $v(p@20 \text{ keV}) = 0.007c$ ,  $F = 0.5$  pN

(\* typical values) but at  $v \rightarrow c$ :  $F = 70$  pN (!)

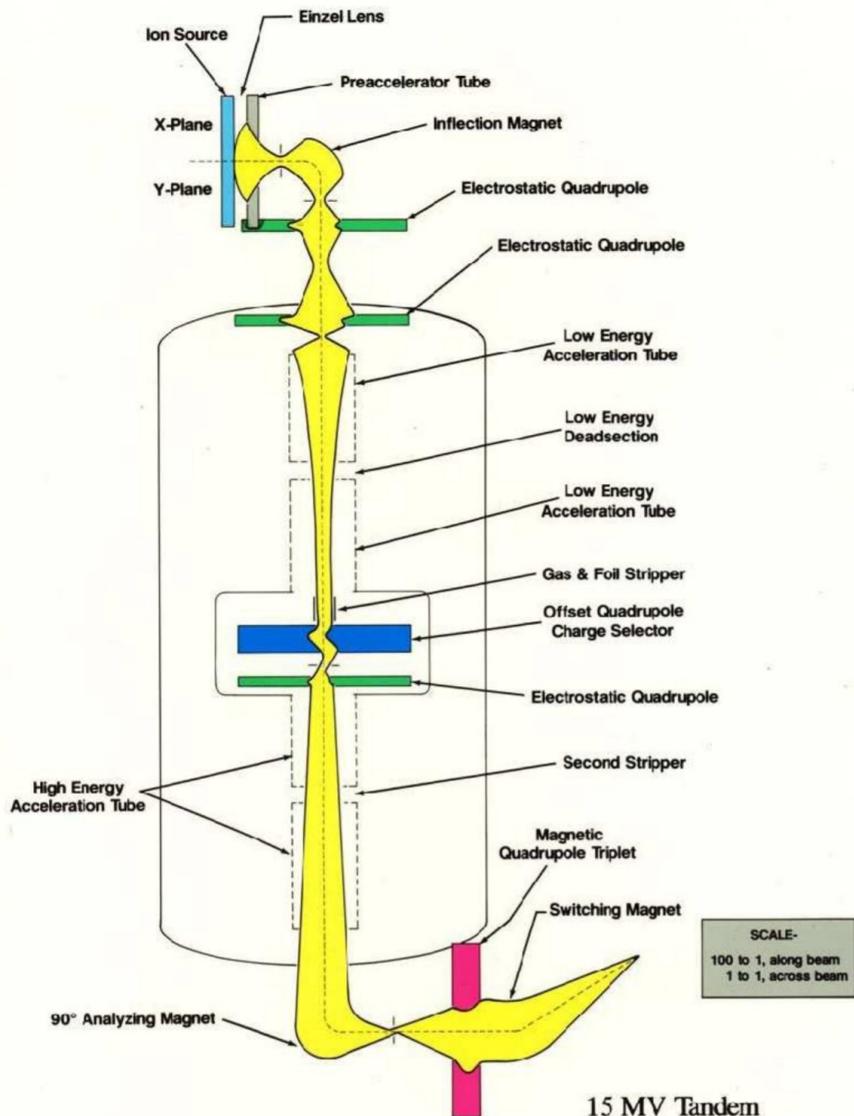
- Electrostatic acceleration:

- Continuous beam, small energy spread

- Easy to tune energy



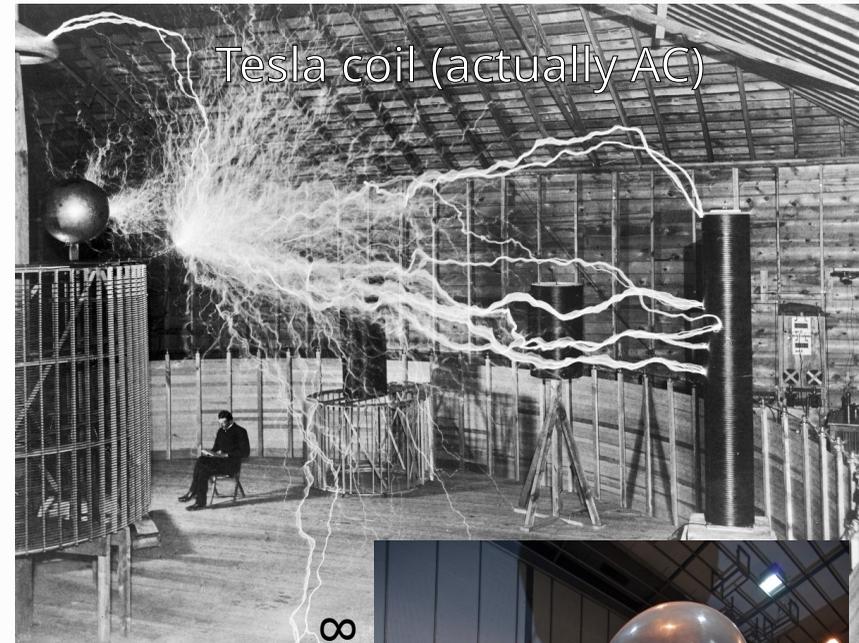
# Electrostatic acceleration



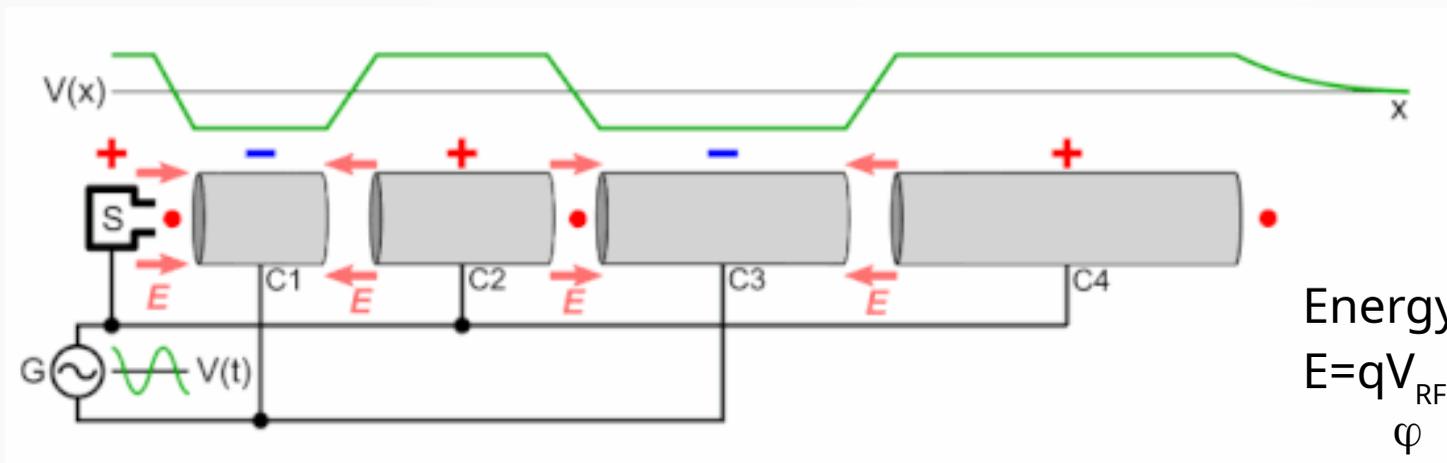
- Tandem accelerator, doubling the energy
- Energies in range 1-40 MeV
- Energy spread  $10^{-4}$
- Electrostatic lenses keep the beam focused (first mention of transverse focusing)
- Still used for instance:
  - Ion Beam Analysis
  - as pre-accelerators for larger facilities
  - ion implantation
- See Giovanni's, Aris'es and Fehima's presentations

# Acceleration techniques: RF

- MV electrostatic generators are huge
- Safe handling these voltages is difficult
- Idea: use oscillating electric field - Gustav Ising (1924)
- First device: Rolf Widreoe (1928)
- Note: beam is bunched and energy tuning is not as easy as for electrostatic machines
- Common name: Drift-Tube Linac (DTL)

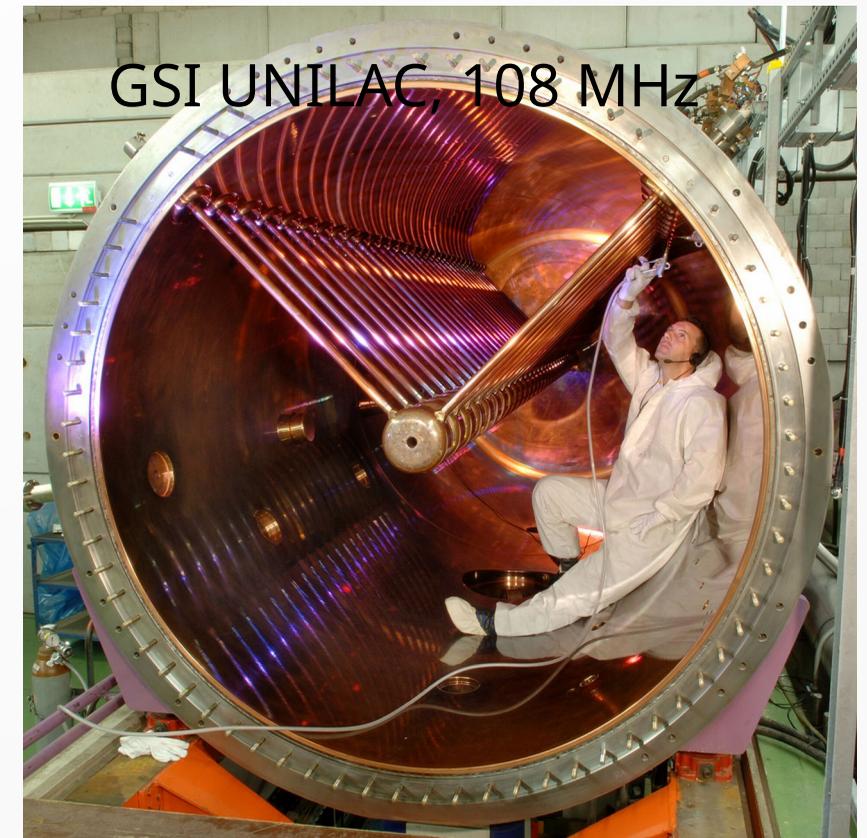
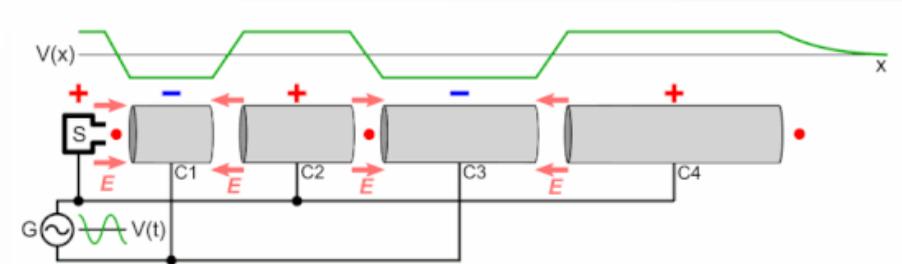


810 kV generator (PSI)



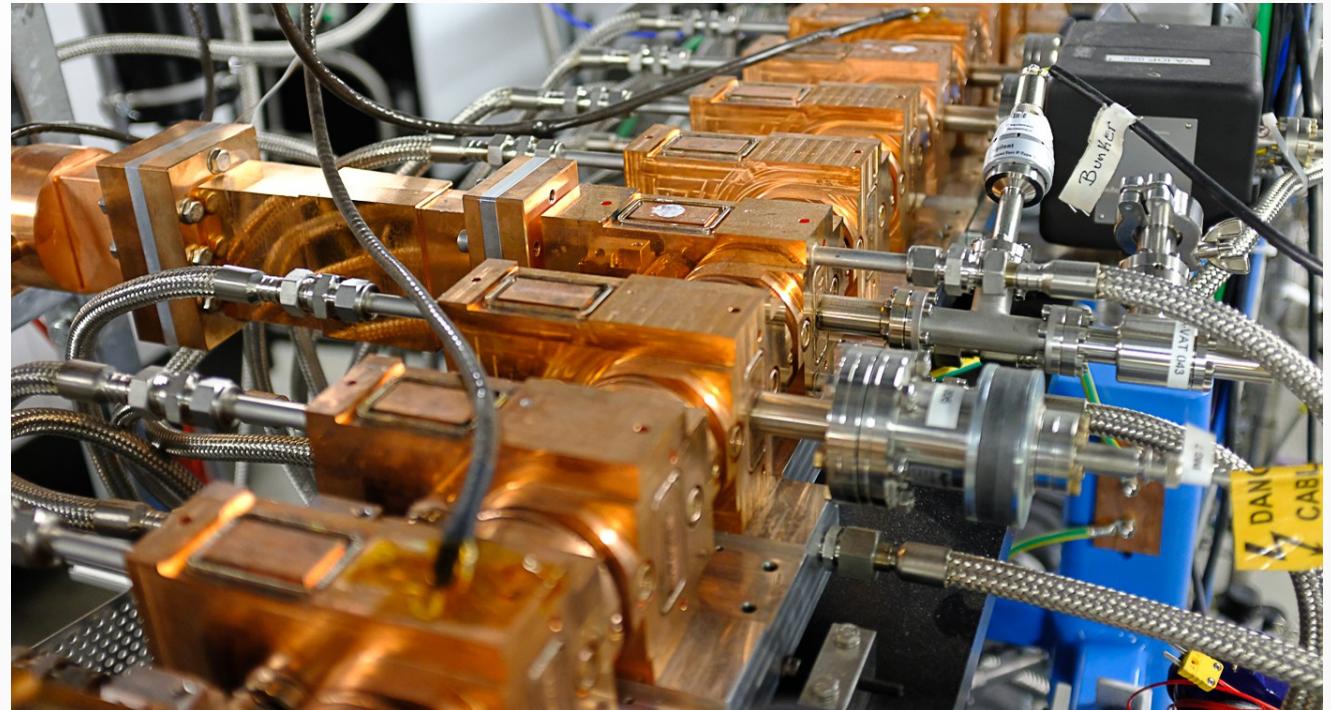
# Acceleration techniques: RF

- With increase of energy, the drift tubes gets longer, higher frequency allows them to be shorter (careful, the first drift tubes can be too short)
- Typically MeV ion beams require frequencies 36-750 MHz, and elements of the system work like antennas emitting most of the energy
- Therefore the accelerator is enclosed in resonant tank and fed by RF source (no need to make electrical connections to the drift tubes)



# Very high frequencies

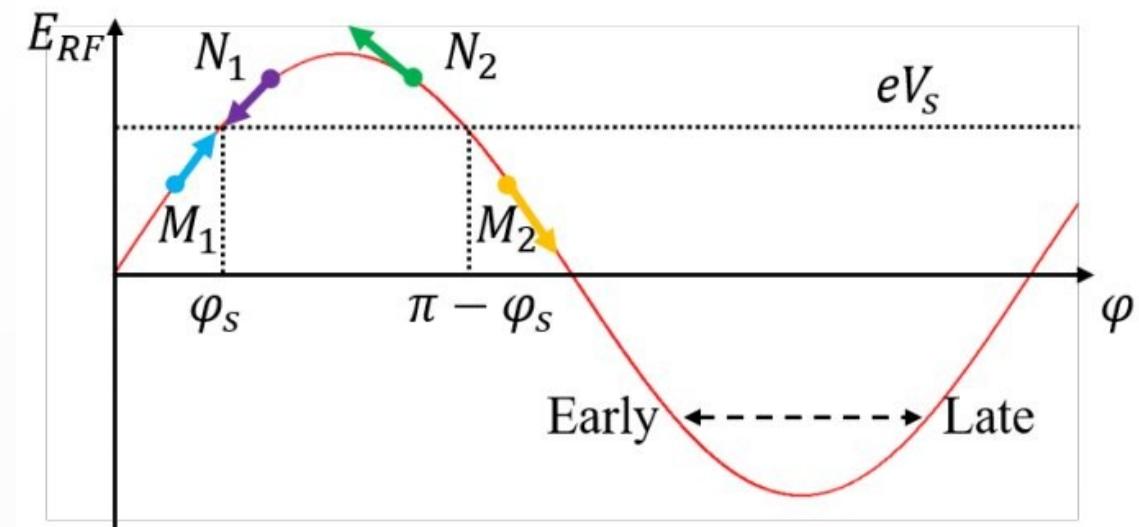
- Higher frequencies allow for smaller linac and higher acceleration gradients
- Maximum frequency depend on particle velocity and accuracy of machining of structures
- For electrons, which are fast relativistic (1 MeV – 95% c) the “golden standard” frequency is 3 GHz (SLAC)
- CLIC developments at CERN (e.g. new industrial CNC machines developed for this project) pushed it to 12 GHz
- For protons new developments:
  - 750 MHz (CERN, low energy)
  - 3 GHz (AVO-ADAM, proton therapy, 70-230 MeV)
  - Accelerating gradient  $>30$  MV/m



AVO-ADAM  
SCDTL

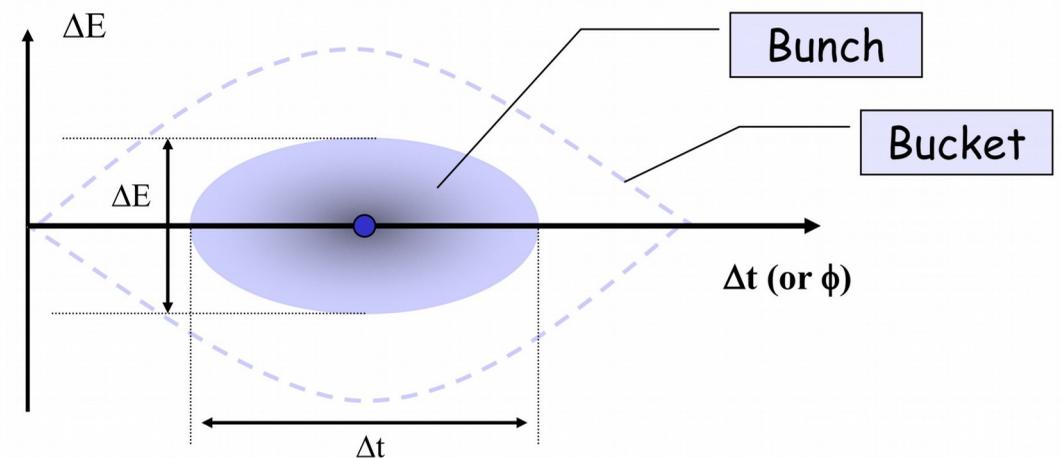
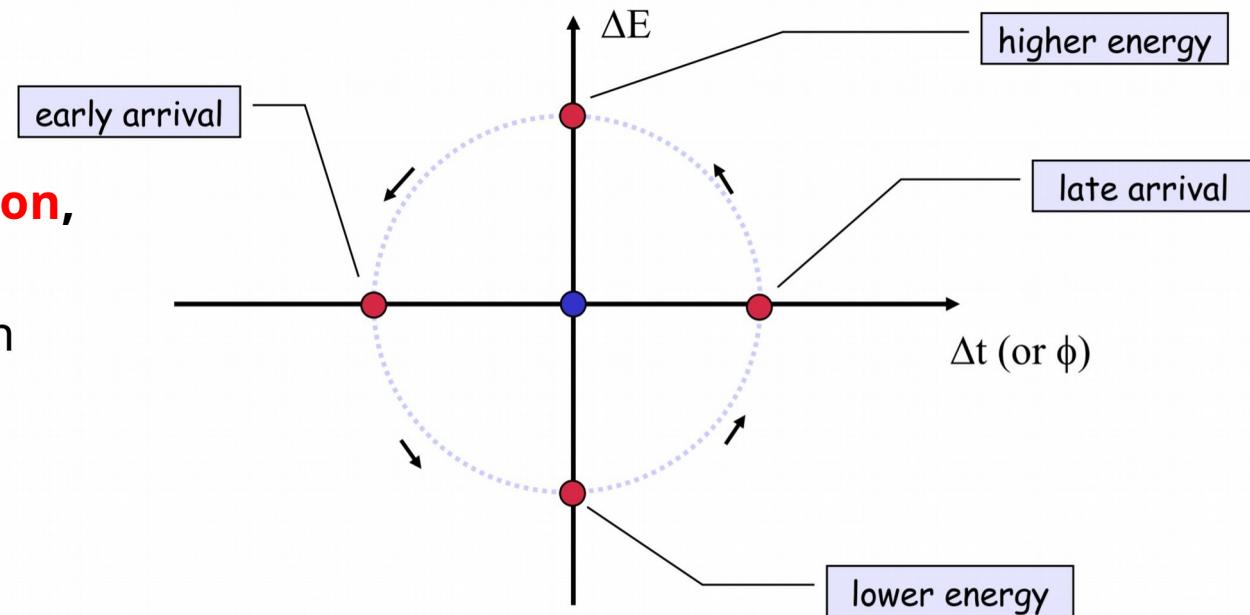
# Tuning the RF cavity

- **Phase stability = longitudinal focusing**
- Particles arriving too early (higher energy) to the accelerating gap experience smaller accelerating field
- Particles arriving too late (lower energy) experience higher accelerating field
- For synchrotrons it is a bit more complex because the different energy means different orbit and depends on beam energy and machine lattice (transition) –  $M_2$  becomes stable
- The stable particle position (RF set-point, optimal phase) is usually found by optimizing the beam transmission through the linac

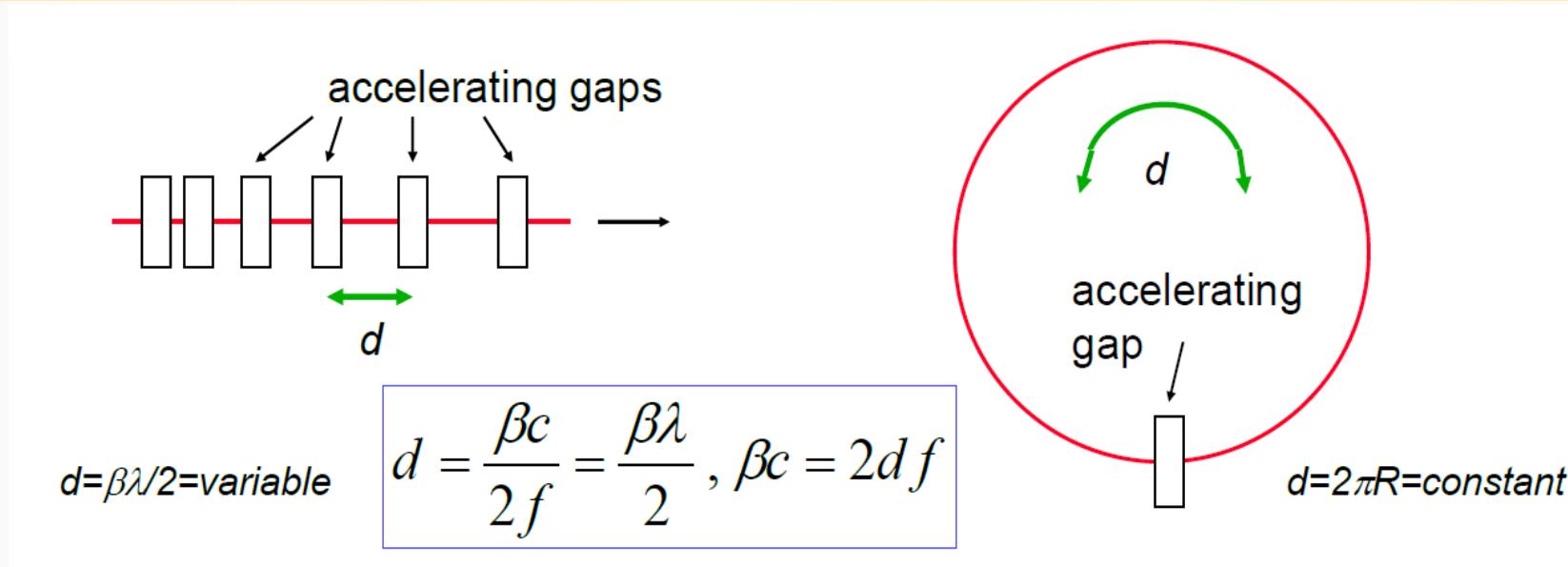


# Longitudinal phase space

- Particles oscillate around the synchronous particle position
- This oscillation is called **synchrotron oscillation**, mainly because it is rather slow oscillation so particle must be circulating in a synchrotron in order to observe it
- Particles stay within separatrix
- Longitudinal phase space: energy-phase
- **Longitudinal emittance** is the phase space area including all particles  $4 \cdot \pi \cdot \sigma_{\Delta E} \cdot \sigma_{\Delta t}$ 
  - Unit [eV.s]



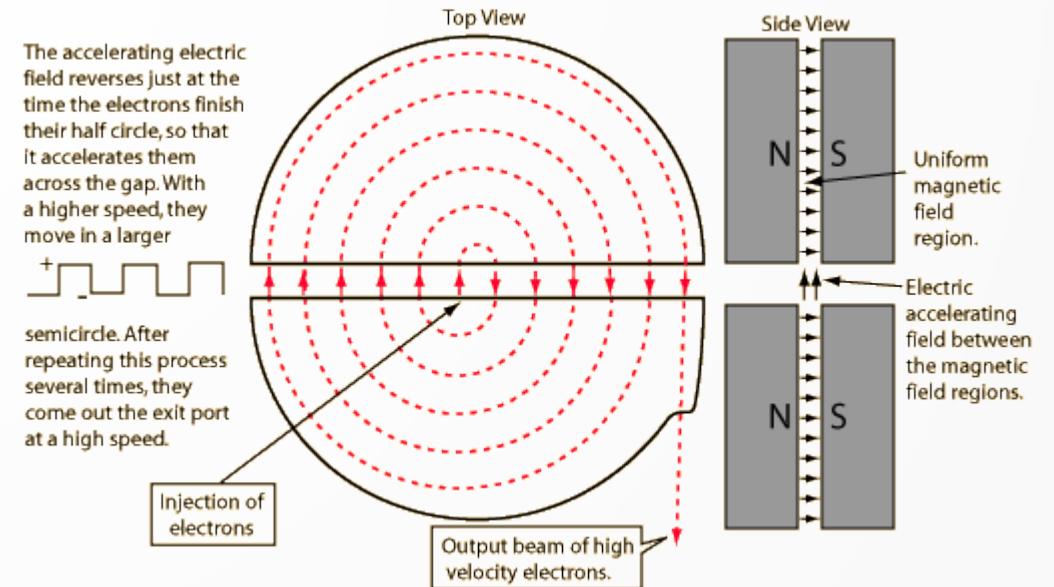
# Linear versus circular machines



- In linear machine each accelerating gap is used once, high accelerating gradient is important; in circular machine beam comes back to the same cavity multiple times, gradient is not so crucial
- Linear machine: distance between gaps increases; circular: frequency of the cavity increases (in non relativistic regime)

# Circular machines: cyclotrons

- Proposed by E.O. Lawrence (1929) and build by Livingstone (1931)
- Vertical magnetic field bends the particle trajectory
- Gap between the dees is used for acceleration
- Radius of the particle increases with its energy
- Lorentz and centrifugal forces balance:
  - $qvB=mv^2/r$
  - $\omega = v/r = qB/m$   
(Larmor frequency)
- Modern cyclotrons: multiple cavities, RF frequency  $\sim 100$  MHz



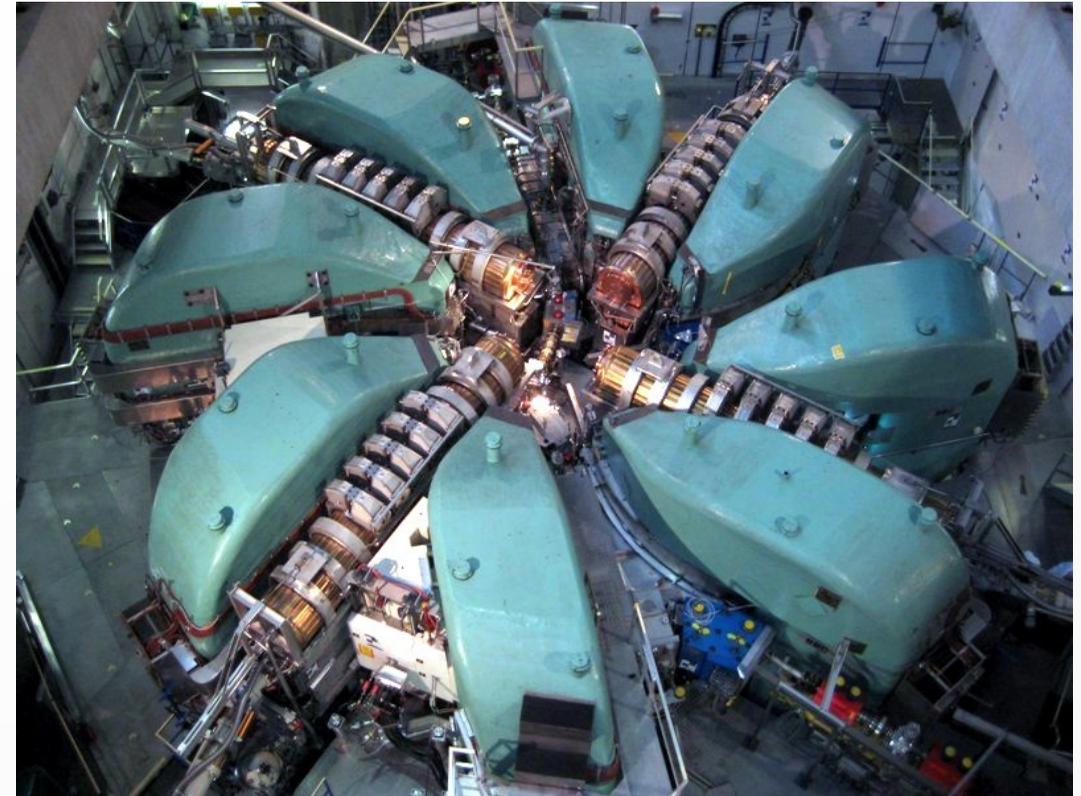
Livingstone cyclotron, diam. 10 cm.  
80 keV protons

# Limitations of cyclotrons

- For relativistic particles mass increases:

$$\omega = v/r = qB/m(\mathbf{E})$$

- Need to increase magnetic field (isochronous) or frequency (synchrocyclotrons)
- At high energies large vacuum chamber becomes difficult (large disc with vacuum)
- Most of proton therapy machines are based on cyclotrons (Varian, IBA)
- The extraction energy is constant (e.g. 230 MeV), must be degraded if needed (e.g. for shallow tumors)



PSI cyclotron, isochronous, sectored, AVF (azimuthally varied field)

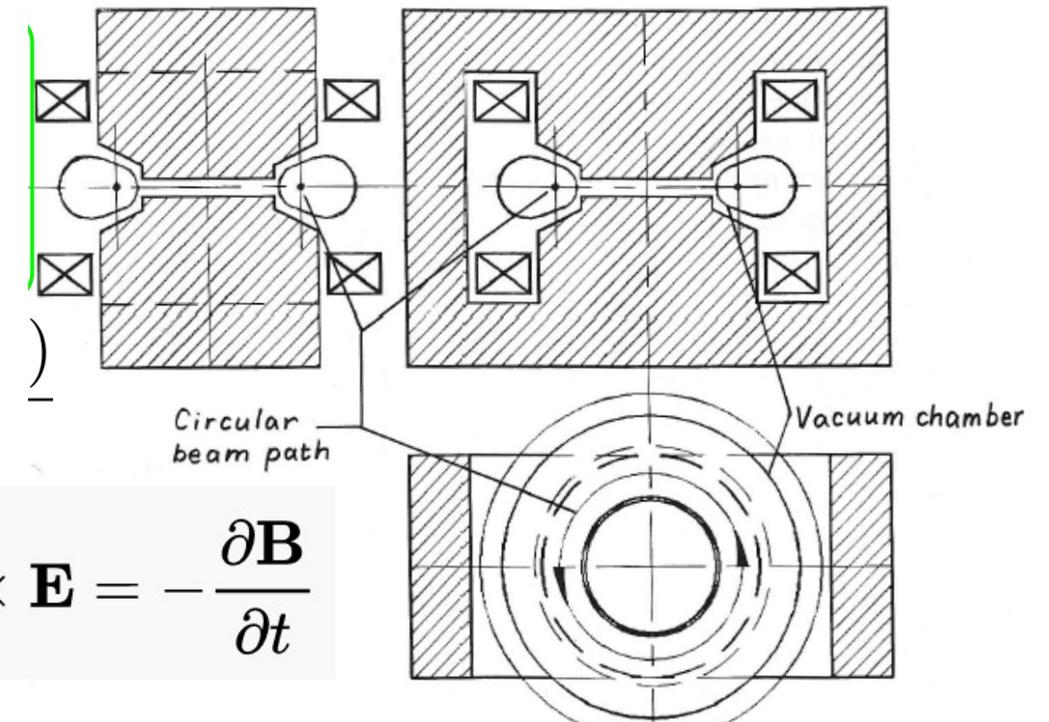
protons at 590 MeV,  $P_{\text{beam}} = 1.3 \text{ MW}$

Diameter 15 m

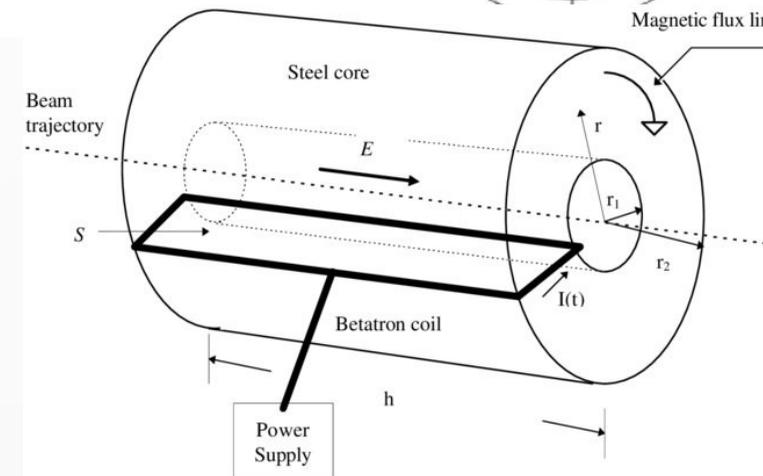
4 RF cavities, 8 magnets

# Betatron

- Electrons are relativistic at 500 keV  
– classical cyclotron not useful
- **Magnetic field increases with energy, orbit is constant**
- Energy is transmitted through transformer effect:  
increasing magnetic field generates vortex  
electric field which accelerates electrons
- Acceleration takes place over  $\frac{1}{4}$  of the RF cycle
- Betatrons were used to produce electron  
beams up to 300 MeV
- Similar idea of final smooth acceleration is used for  
slow extraction in CNAO and MedAustron  
(betatron core)

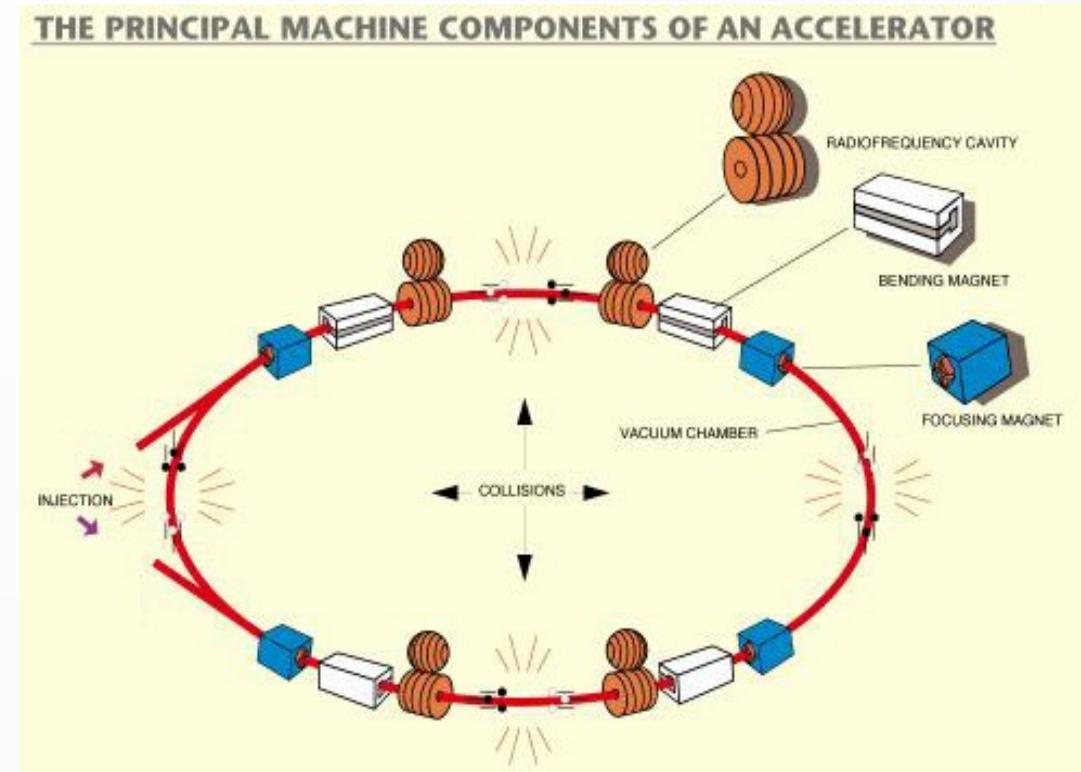


$$\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t}$$



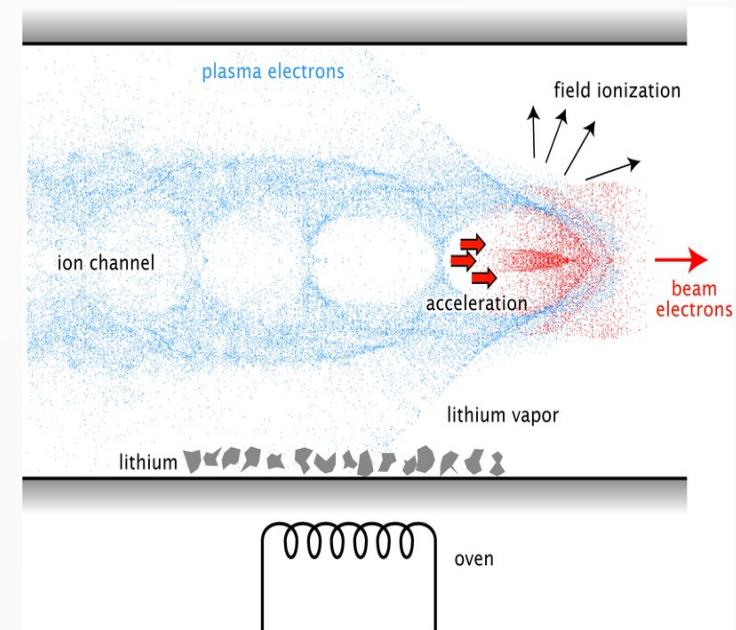
# Synchrotron

- Betatron acceleration method is limited by magnet size and iron saturation
- For larger energies need to split the magnets
- And use much more efficient RF-acceleration
- Synchrotron idea: Vladimir Veksler (1944)
- Origin of the name: **synchronous change of RF frequency with magnets' current**
- First electron synchrotron: Edwin McMillan (1945, independently from Veksler)
- First proton synchrotron: Marcus Oliphant (1952)



# Future of acceleration techniques

- The best cavities reach 50 MV/m (less in regular operation e.g. 30 MV/m European XFEL, DESY, Germany)
- Vacuum breakdown limits possible fields
- Idea: **use plasma** – it is already broken down
- Separate electrons from ions using strong laser pulse, generate locally fields of **100 GV/m** (factor 5000!)
- Currently plasma acceleration, dielectric acceleration, laser ion sources are very active fields of research



**Accelerators started with electrostatic machines.**

**The crucial development was RF resonant acceleration.**

**Phase stability keeps beam bunched.  
Cyclotrons and Synchrotrons.**

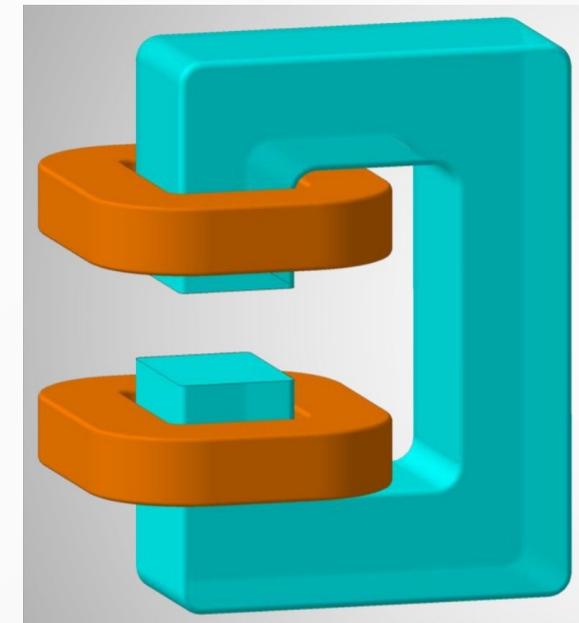
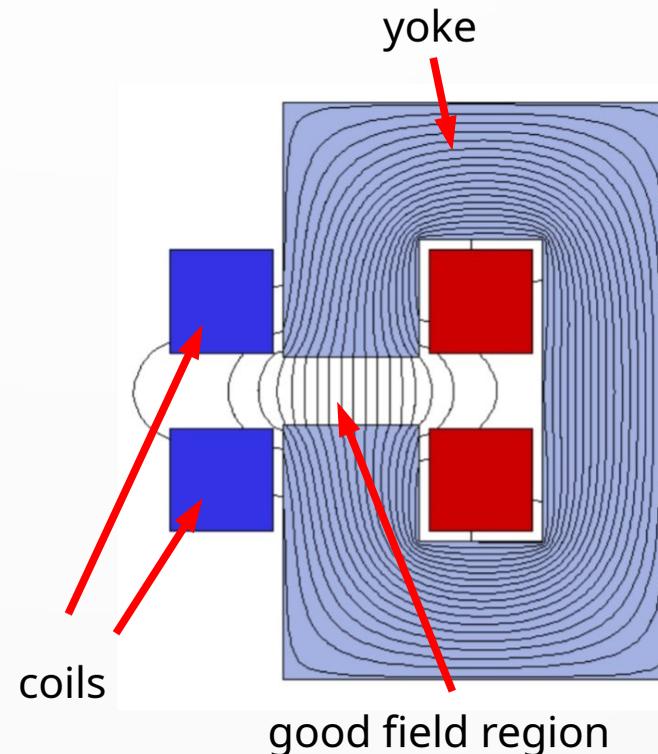
**The future may be in plasma accelerators.**

# Beam trajectory

- Beam can be steered using electric or magnetic fields
- Magnetic field is more effective for high velocities, because:

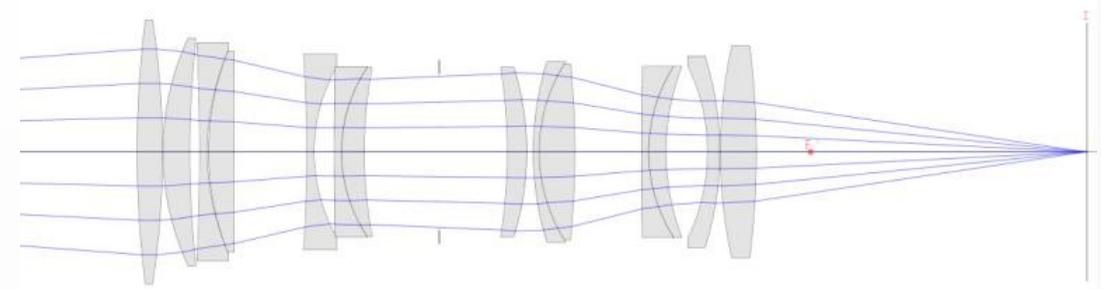
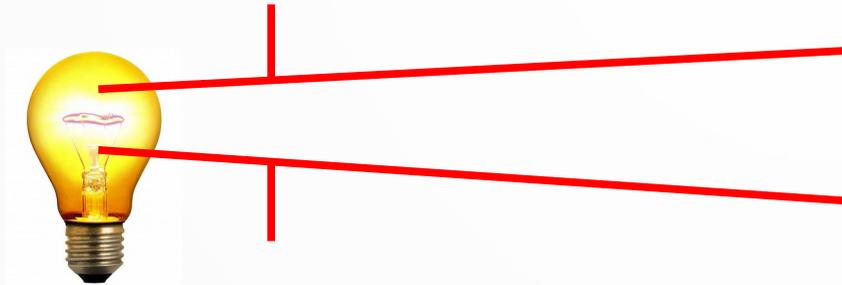
$$\underline{E} = q(\underline{v} \times \underline{B})$$

- Dipole magnets steer the beam
- Particles of the same **magnetic rigidity** have the same trajectory:
  - $p/q = B\rho$ ,  $\rho$ -bending radius
  - $^{12}\text{C}^{6+}$  and  $^4\text{He}^{2+}$  can circulate in the same machine having the same kinetic energy per unit mass eg. 430 MeV/u

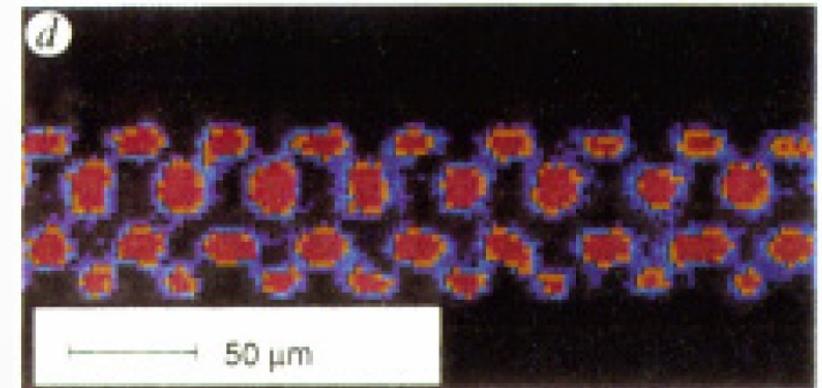


# Beam stability

- Beam is naturally divergent; think about a “beam of light” (from a lamp) and a diaphragm:

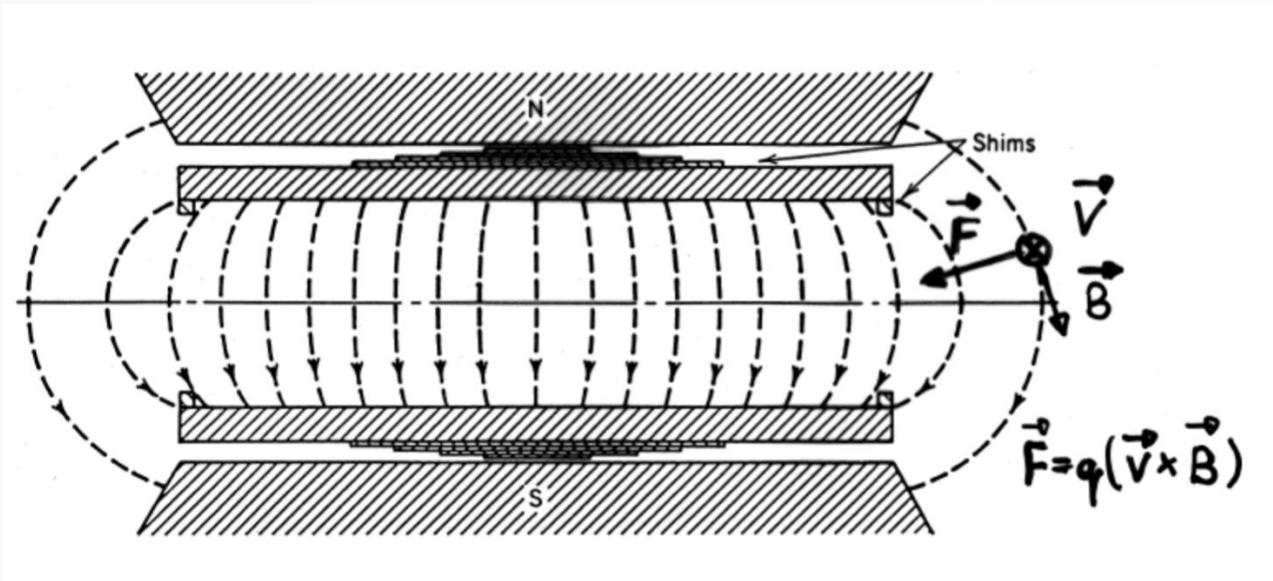


- Set of lenses focus the beam of light
- BTW laser light stays “collimated” without lenses: this is because of spatial coherence of photons in the laser beam
- Can we do similar with ion beams? Not really! Ions are fermions not bosons, they are charged (Coulomb repulsion forces);
- The crystalline ion beams reach limits of ion beam emittance, beam density

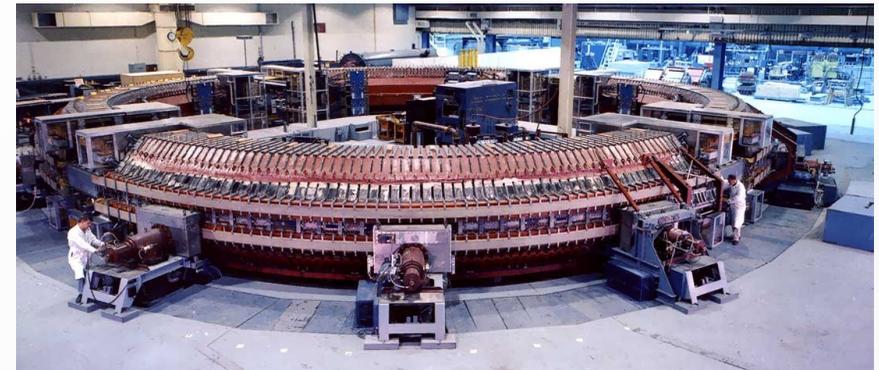


# Weak focusing

- The principle – radial field gradient leads to forces which focus the beam
- This mechanism is called weak focusing
- Every dipole magnet gives vertical focusing at its edges
- In synchrotrons, which store the beam over long time, **weak focusing is not enough!**



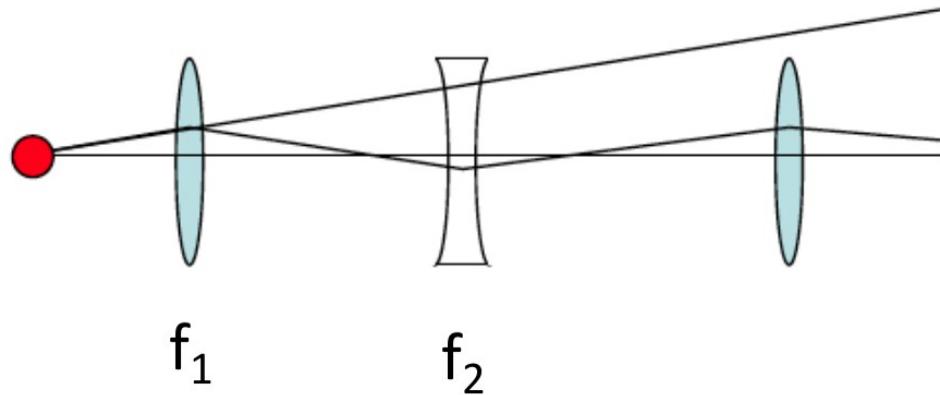
Weak focusing in cyclotron



Cosmotron – 3 GeV proton synchrotron, BNL 1953

# Strong focusing

- Idea: N. Christofilos, 1949 (patented but not published), rediscovered independently by E. Courant, M. Livingston, H. Snyder in 1952
- Strong focusing principle: the net effect on a particle beam of charged particles passing through alternating field gradients is to make the beam converge



$$\frac{1}{F} = \frac{1}{f_1} + \frac{1}{f_2} - \frac{d}{f_1 f_2}$$

Consider  $f_1=f$ ,  $f_2 = -f \rightarrow F = f^2/d > 0$

# Quadrupoles

- Quadrupole magnet provides focusing in one plane and defocusing in other

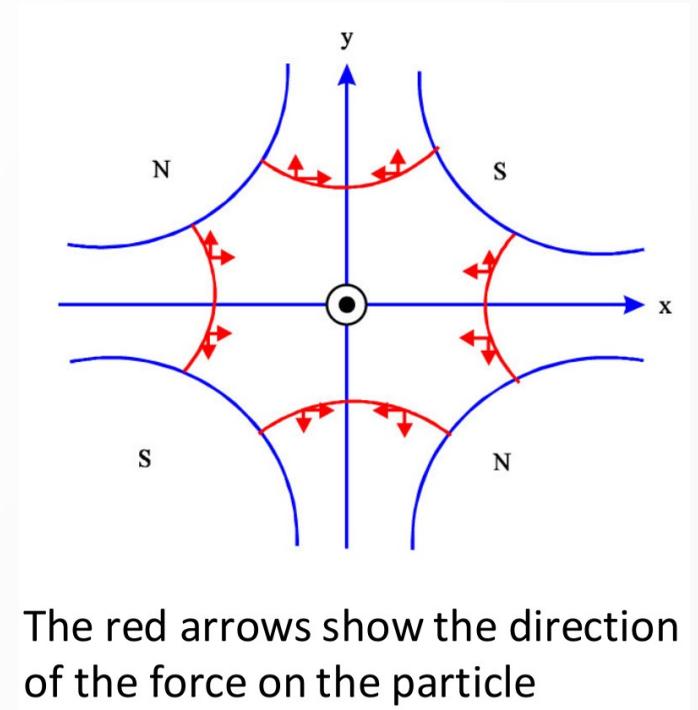
- $F = qvB(x) = qv(g \cdot x)$

- Magnetic field gradient:

$$g = \frac{2\mu_0 n I}{r^2} \left[ \frac{T}{m} \right]$$

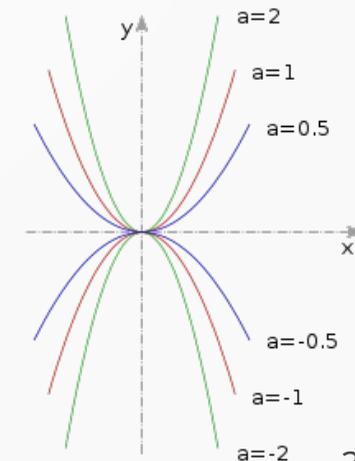
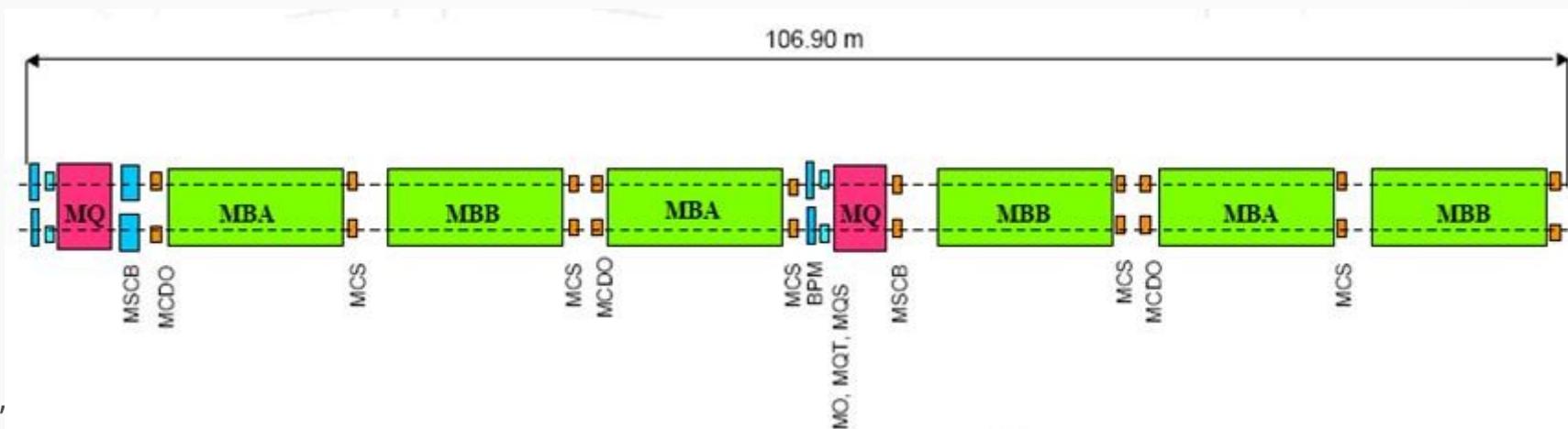
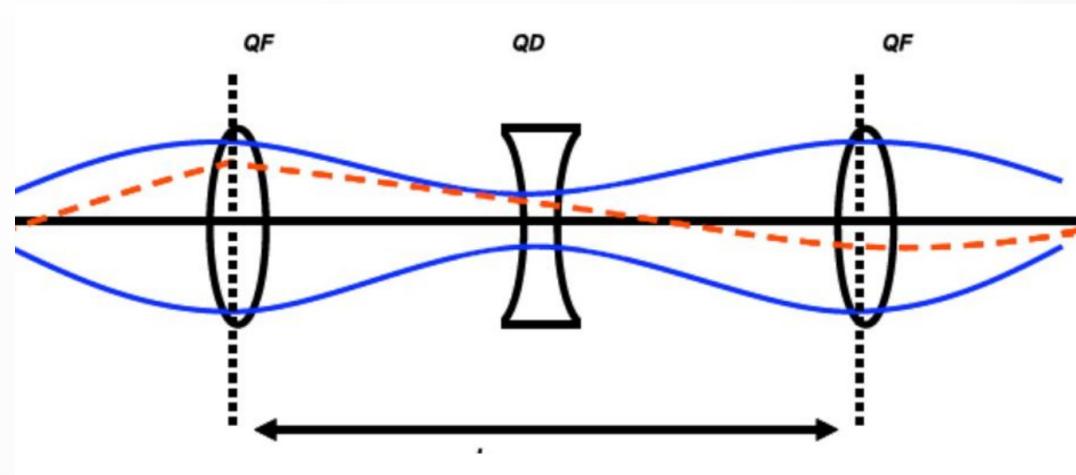
- Gradient normalized to rigidity:

$$k = \frac{g}{p/q} \left[ m^{-2} \right]$$



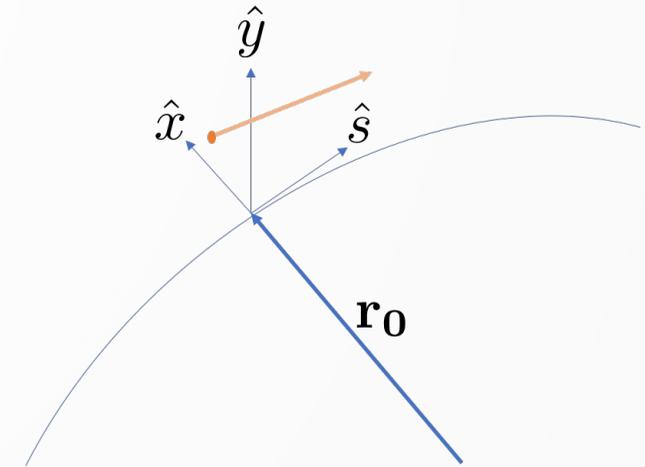
# Synchrotron cells

- As we've seen we need a system of lenses (i.e. of quadrupoles)
- Focusing-Defocusing (FODO) – the easiest elementary cell layout of a synchrotron
- Dipoles, placed between quadrupoles  
- add weak focusing
- Real example: LHC FODO  
184 FODO cells in LHC arcs
- Other cells often used (e.g. multi-bend achromat to minimize beam size in light sources)



# Equation of transverse motion (I)

- First: reference system
- Equation of motion: position of particles in function of time  $x(t)$ ,  $x'(t)$
- Particles move through lattice with constant velocity, so we replace: time ( $t$ )  $\rightarrow$  position along the machine ( $s$ )
  - $dx/dt = dx/ds * ds/dt = (dx/ds) * v$
  - $d^2x/dt^2 = d/dt(dx/ds) * ds/dt + dx/ds * d^2s/dt^2 = (d^2x/ds^2) * v^2$
- Equation of motion  $F=ma=qvB$ :
  - $d^2x/ds^2 = qv(g * x) / mv^2$ ;  $k = -g / (p/q)$ ;  $p/q$  - rigidity
  - $d^2x/ds^2 = -kx$  (focusing, harmonic oscillator!)
- Solution is periodic:  $x(s) = x(0) * \cos(k^{1/2} s) + x'(0) * \sin(k^{1/2} s)$   
(focusing)  $x'(s) = x(0) * k^{1/2} * \sin(k^{1/2} s) + x'(0) * k^{1/2} * \cos(k^{1/2} s)$

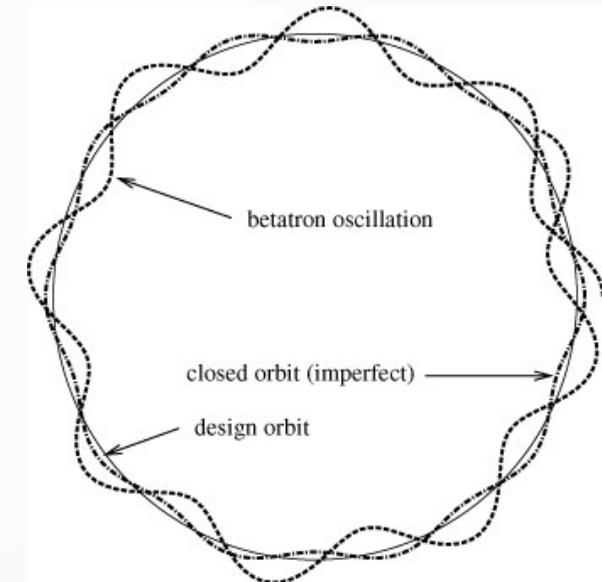


reference orbit  
or trajectory

# Equation of transverse motion (II)

- For defocusing quadrupole:  $x(s) = x(0) \cdot \cosh(k^{1/2}s) + x'(0) \cdot \sinh(k^{1/2}s)$   
 $x'(s) = x(0) \cdot k^{1/2} \cdot \sinh(k^{1/2}s) + x'(0) \cdot k^{1/2} \cdot \cosh(k^{1/2}s)$
- General equation of motion (Hill's equation):  
 $x''(s) + K(s)x(s) = 0$   
where  $K_x = 1/\rho + k$  (includes weak focusing)
- $K(L+s) = K(s)$  – where L is lattice period (eg. length of FODO cell)
- General solution describes quasi-harmonic movement called **betatron oscillations**:

$$x(s) = \sqrt{2J_x \beta_x(s)} \cos(\psi(s) + \phi)$$

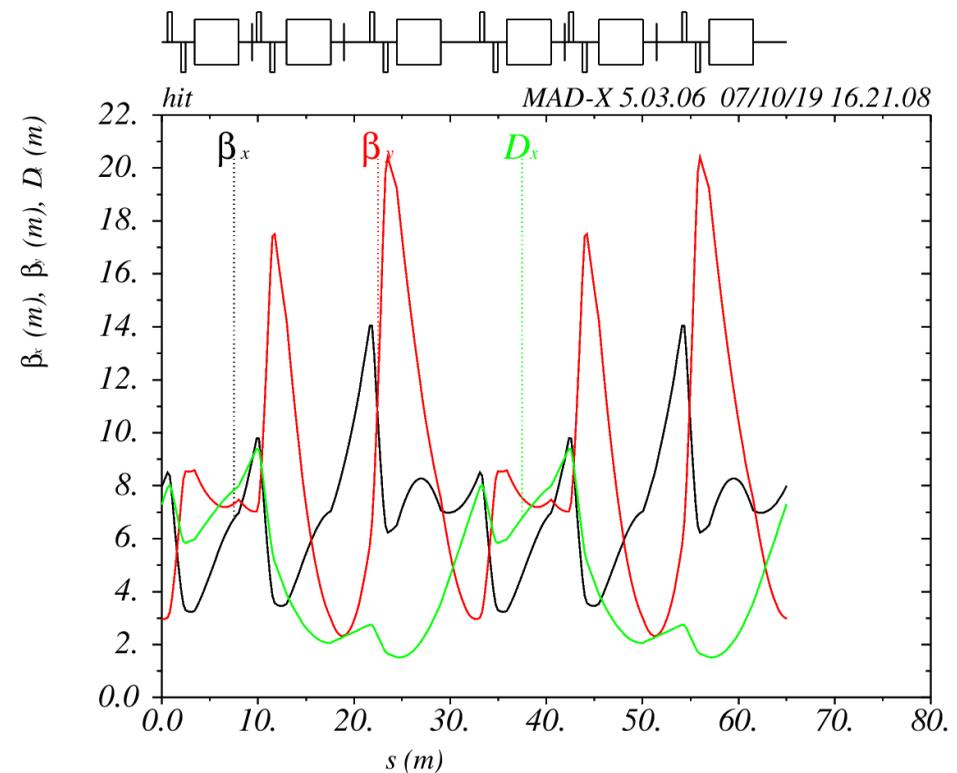
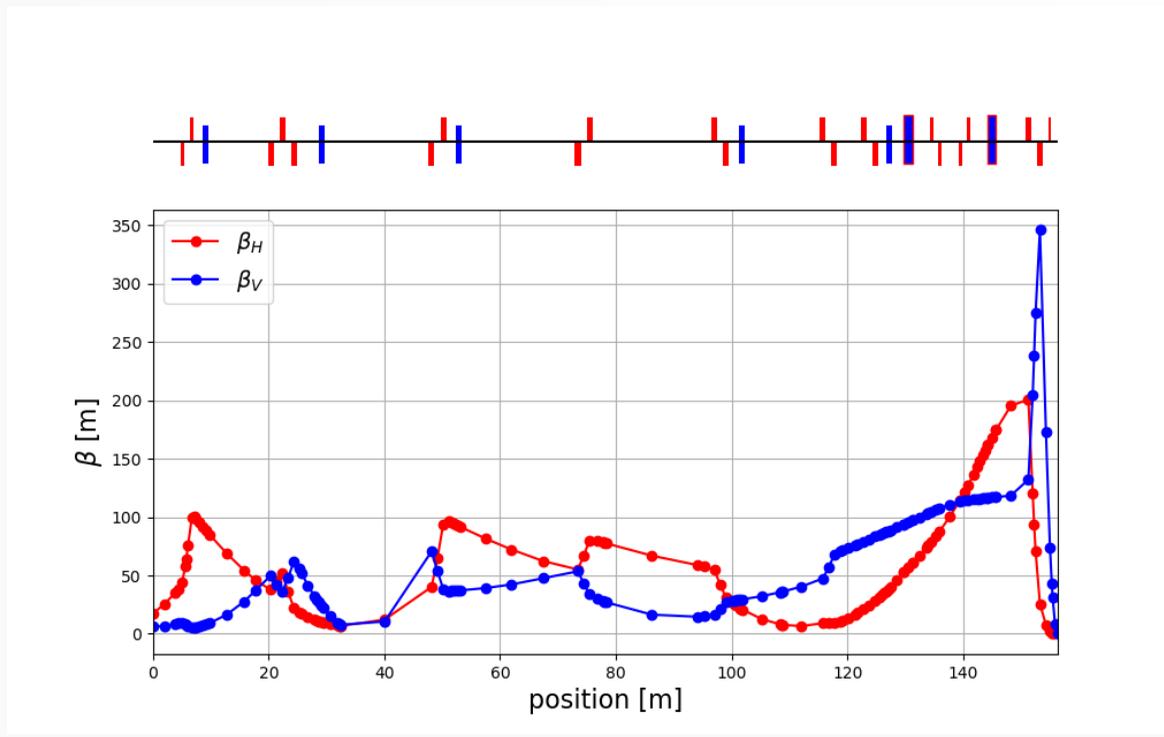


$J_x$  and  $\phi$  – depend on initial conditions

# Beta function

- **betatron oscillations** – transverse oscillations of particle in the beam around the design orbit (reminder: synchrotron oscillations are longitudinal oscillations around the stable RF phase)
- **Beta function  $\beta(s)$**  describes amplitude of betatron oscillations along the accelerator or transfer line, often called *beam optics*
- Beam size (often called beam envelope):

$$\sigma = \sqrt{\beta \epsilon}$$



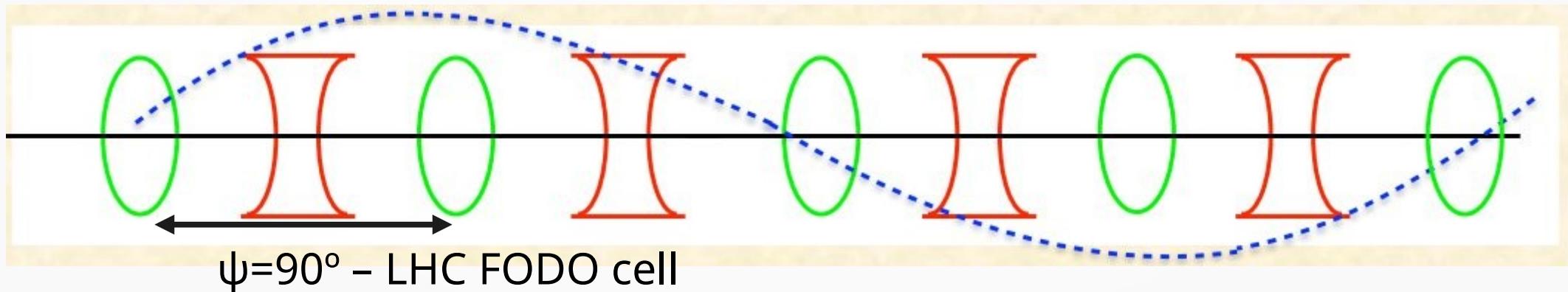
# Phase advance and tune

- The difference of betatron motion phase between two points is called **phase advance**:

$$\psi(s) = \int_0^s \frac{ds}{\beta(s)}$$

- Number of betatron oscillations per turn is called **tune**:

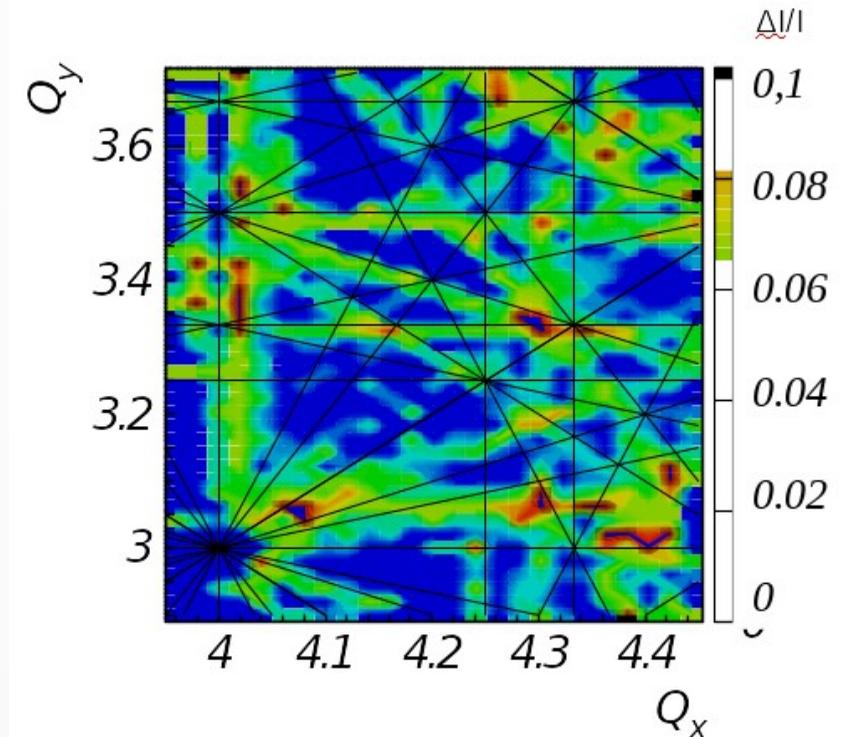
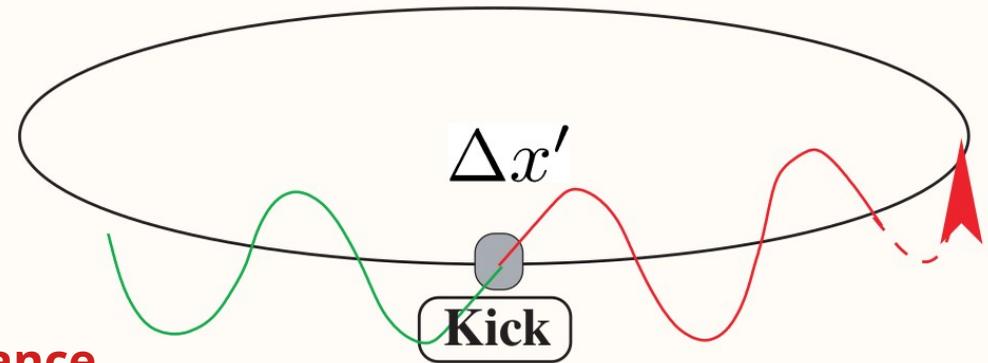
$$Q = \frac{\psi(L_{turn})}{2\pi} = \frac{1}{2\pi} \oint \frac{ds}{\beta(s)}$$



# Tune and resonances

- **Tune** depends on optics (setting of quadrupoles) and can be regulated
- Each small field errors or magnet misalignments create perturbation of the beam trajectory
- If tune is integer (N) or N/2, N/3... the effect of those perturbations add up every turn, **machine is in resonance** and operation is unstable
- This is bad for storage rings but is also a basics of **resonant slow extraction** used in medical machines to extract beam to the patient
- e.g. CNAO/MedAustron **working point:  $Q_{x,y}=(1.672,1.72)$**
- $Q_x=1.666$  is third order resonance which can be excited by sextupole magnets (see later)

GSI SIS18 tune diagram:



# Dispersion

- Transverse and longitudinal motions are not independent; they are coupled via dispersion
- **Dispersion is a deviation of the particle trajectory due to momentum difference:**

$$D_x(s) = dx(s) / (dp/p)$$

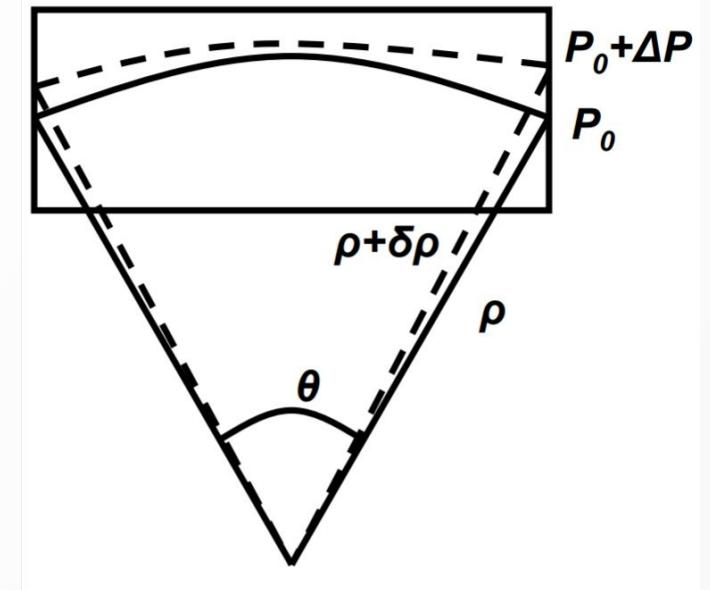
- Similar for angle:

$$D'_x = dx'(s) / (dp/p)$$

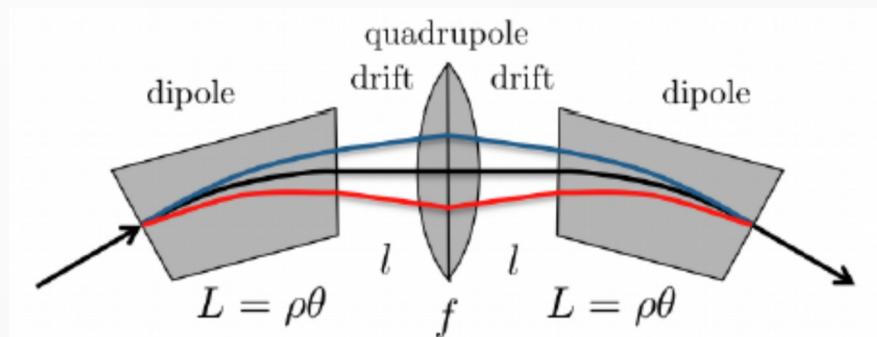
- Dispersion leads to increase of beam size:

$$\sigma = \sqrt{\beta\varepsilon + D^2 \left(\frac{\Delta p}{p}\right)^2}$$

- Dispersion-free regions are often needed: minimize beam size and movement on the patient, maximize luminosity, measure emittance

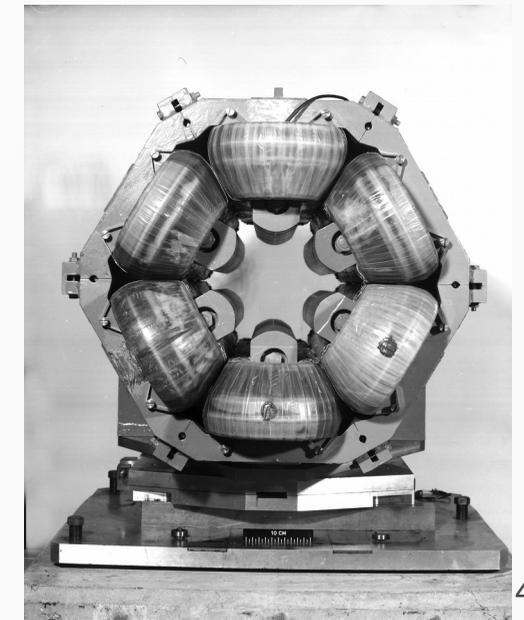
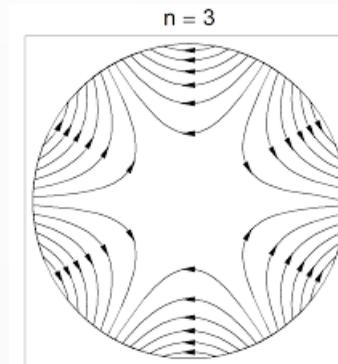
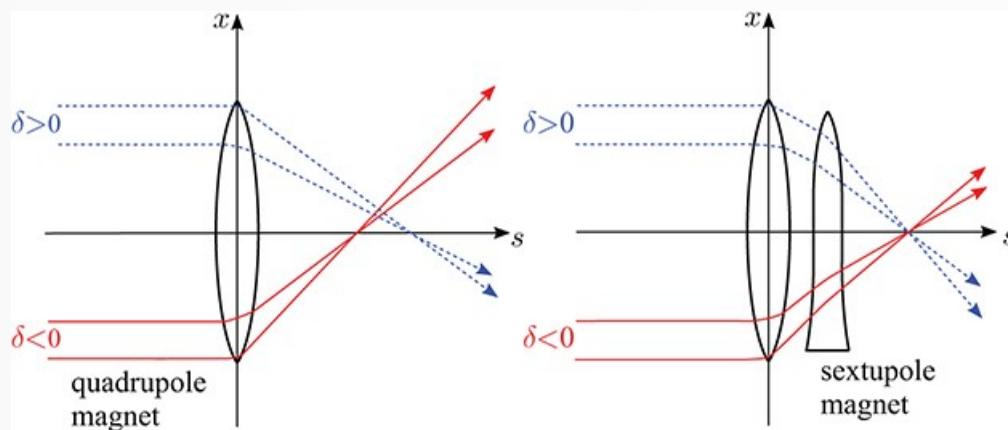


There are many ways to remove dispersion, eg. double-bend achromat (DBA):



# Chromaticity and sextupoles

- As dispersion is a change of trajectory with the momentum deviation, **chromaticity is change of machine tune with the momentum deviation:  $Q' = dQ / (dp/p)$  [dimensionless]**
- Reminder: typical momentum spread in a synchrotron ( $\Delta p/p \sim 10^{-3}$ )
- **Chromaticity is controlled by sextupole magnets installed in dispersive region**
- Typically small negative chromaticity is needed to make machine stable
- Higher order effects demand octupoles, decapoles to correct



# Twiss parameters and beam ellipse

- Beam ellipse can be described in terms of emittance and **Twiss parameters** (called also Courant-Snyder parameters):

$$\alpha(s) \equiv -\frac{1}{2} \frac{d\beta(s)}{ds}$$

$$\gamma(s) \equiv \frac{1 + \alpha^2(s)}{\beta(s)}$$

- Alpha ( $\alpha$ ) is *slope* of beta;
- “parallel beam”:  $\alpha=0$
- Gamma is dependent parameter and it is *beta for angle*

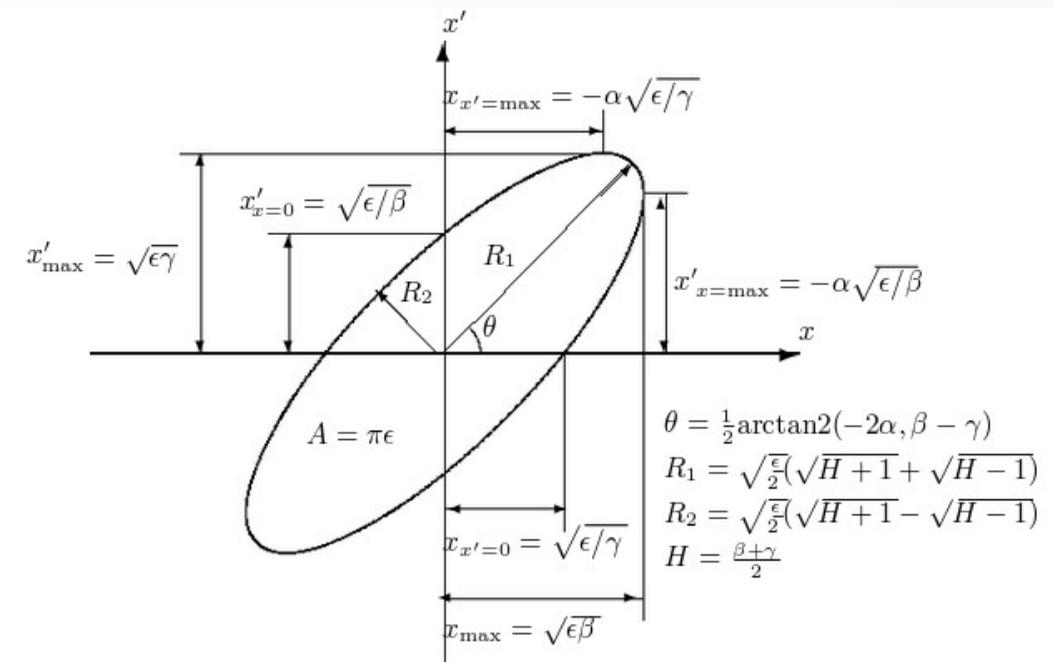


Figure 9: Emittance ellipse geometry with the most important dimensions

# Beam transport

- **Matrix formalism** is used to transfer the beam from one element to another:

$$\begin{pmatrix} x \\ x' \end{pmatrix}_{s_1} = M \begin{pmatrix} x \\ x' \end{pmatrix}_{s_0}$$

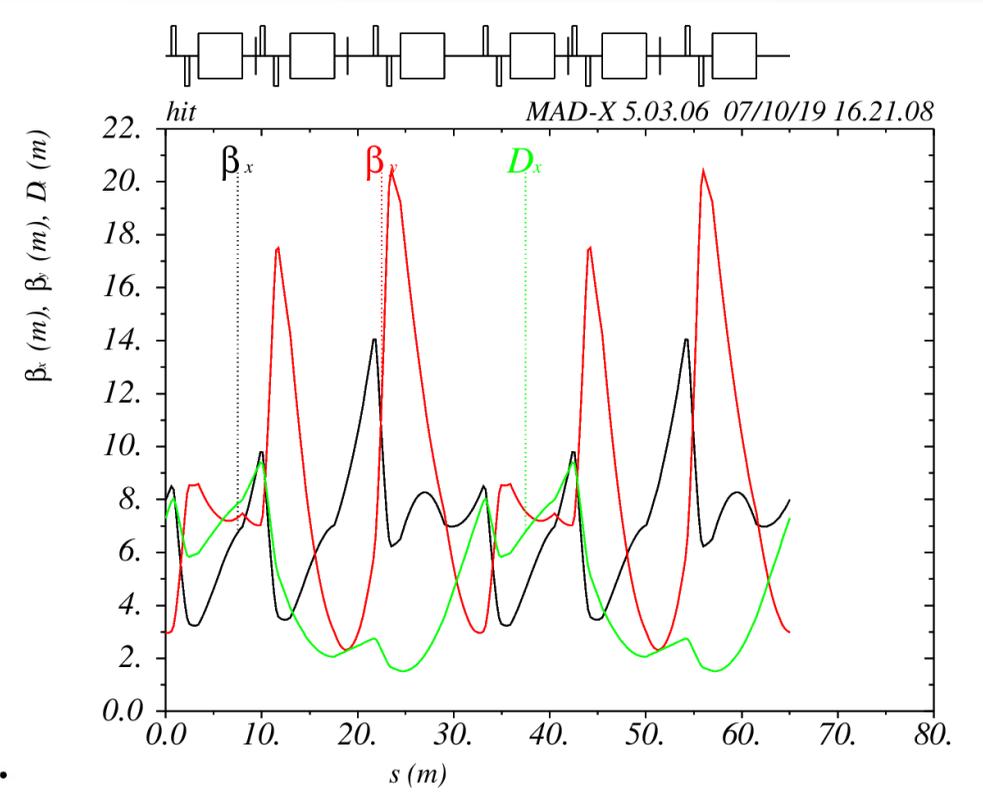
- e.g. transfer matrix for focusing quad:

$$M_{foc} = \begin{pmatrix} \cos(\sqrt{K}s) & \frac{1}{\sqrt{K}} \sin(\sqrt{K}s) \\ -\sqrt{K} \sin(\sqrt{K}s) & \cos(\sqrt{K}s) \end{pmatrix}$$

- Transport through multiple elements:

$$M_{total} = M_{QF} \cdot M_D \cdot M_{Bend} \cdot M_D \cdot M_{QD} \cdot \dots$$

- These are first steps in designing a synchrotron or a beam line



Using this formalism, or tracking of the particles in magnet fields, programs like MAD-X, allow to compute Twiss parameters and dispersion

# Beam transport- 6D

- Equation of ellipse can be also written in form of matrix  $\Sigma$ :

$$[x]^T \Sigma [x] = 1$$

- Beam matrix  $\Sigma(s)$  describes the beam ellipse at a given position; determinant of the ellipse is emittance

- Beam matrix is transformed using matrix formalism:

$$\Sigma(s) = M \Sigma(0) M^T$$

- Beam has 2 independent parameters per dimension, so total 6-D is needed to write full beam matrix
- Transverse-longitudinal coupling via dispersion (D) and D', here included in  $\eta$

$$\Sigma = \begin{pmatrix} \sigma_{11} & \sigma_{12} \\ \sigma_{21} & \sigma_{22} \end{pmatrix} = \epsilon \begin{pmatrix} \beta & -\alpha \\ -\alpha & \gamma \end{pmatrix}$$

$$\Sigma = \begin{pmatrix} \epsilon_x \beta_x & -\epsilon_x \alpha_x & 0 & 0 & 0 & \eta_x \sigma_\delta^2 \\ -\epsilon_x \alpha_x & \epsilon_x \gamma_x & 0 & 0 & 0 & \eta_{p_x} \sigma_\delta^2 \\ 0 & 0 & \epsilon_y \beta_y & -\epsilon_y \alpha_y & 0 & \eta_y \sigma_\delta^2 \\ 0 & 0 & -\epsilon_y \alpha_y & \epsilon_y \gamma_y & 0 & \eta_{p_y} \sigma_\delta^2 \\ 0 & 0 & 0 & 0 & \sigma_z^2 & 0 \\ \eta_x \sigma_\delta^2 & \eta_{p_x} \sigma_\delta^2 & \eta_y \sigma_\delta^2 & \eta_{p_y} \sigma_\delta^2 & 0 & \sigma_\delta^2 \end{pmatrix}$$

**Strong focusing principle allows to construct stable storage rings and transport the beam efficiently.**

**Things to remember:**

**synchrotron cells, Twiss parameters, dispersion, tune, chromaticity, resonances and matrix formalism**

# Conclusions

- Accelerators are one of the most important tools in science, medicine and industry
- They produce beam of particles with a given energy and emittance
- **RF acceleration** allows to reach very high energies; phase stability assures longitudinal focusing
- **Strong focusing** made possible large machines able to produce, transport and store high intensity beams for hours (or days)
- The most important concepts: elementary cell, Twiss parameters ( $\alpha, \beta$ ), beam phase space, beam ellipse, dispersion, tune, resonances, chromaticity and matrix formalism



## Acknowledgments:

- Preparing these slides I used presentations of several CERN Accelerator Schools and summer student lectures

# Thank you for your attention!

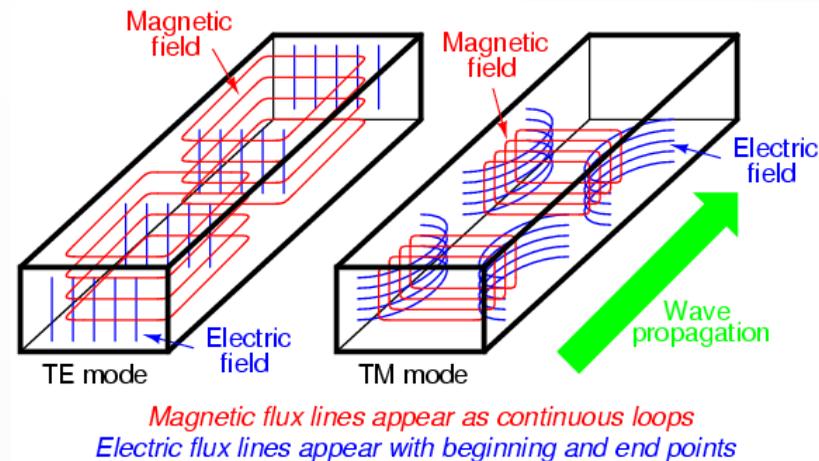
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Please contact me if you have questions concerning this lecture: [mariusz.sapinski@cern.ch](mailto:mariusz.sapinski@cern.ch)

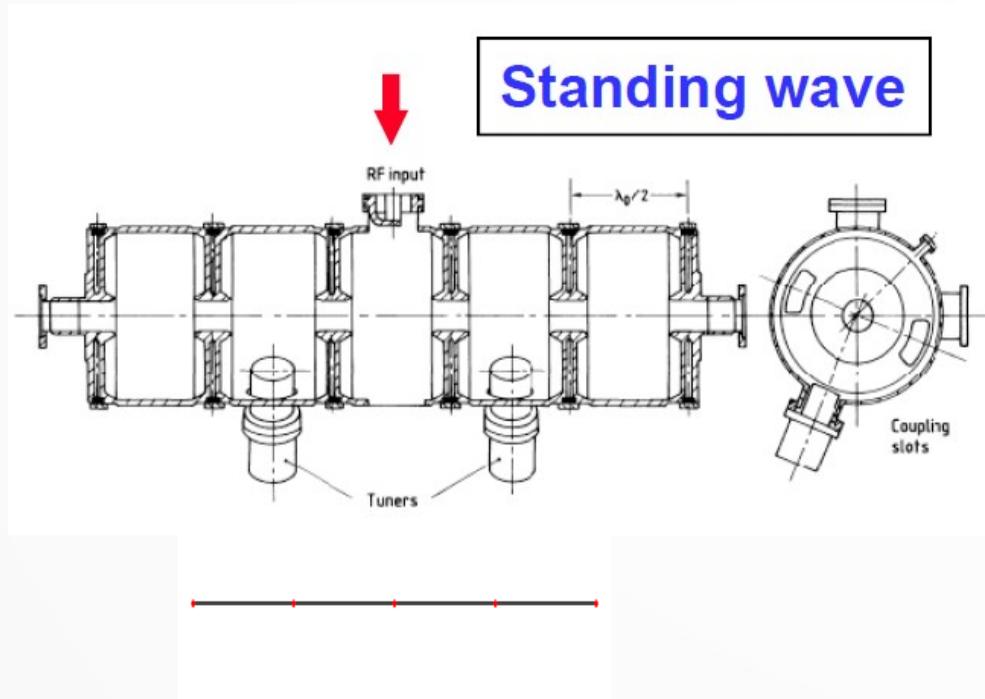


# RF sources

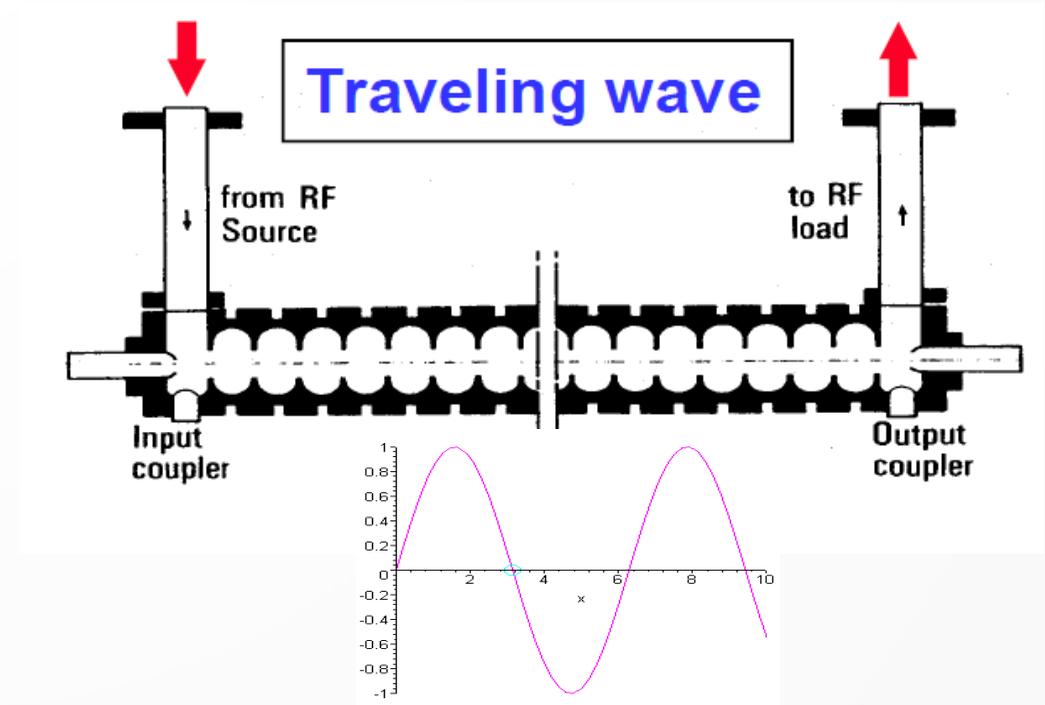
- High-frequency transmission lines are waveguides, not cables, because cables are antennas
- (But... the simplest waveguide is a concentric cable)
- **RF accelerators need powerful RF sources**
- The RF sources are closely related to II Wold War radars
- One of the first devices, still in use, was klystron, developed by Varian brothers (yes, they set up Varian company known for cyclotrons)
- Klystrons by themselves are small electron accelerators
- Trend: solid state RF generators



# Standing and travelling wave



Acceleration  $\sim 5$  MV/m



Acceleration  $\sim 30$  MV/m

# Radio-Frequency Quadrupole

- DTL can accept ion beams from energy of hundreds of keV/u (limits on frequency and size of the tank)
- Ion sources provide ion energies of  $\sim 5\text{-}50$  keV/u
- Acceleration in between is difficult, space charge forces act to disrupt the beam
- Electrostatic acceleration is a valid option, but
- RFQ proves to be a very efficient and compact acceleration element
- It provides focusing and smooth bunching
- Increases the transmission from source to DTL from 50% to 90%
  - crucial for high-power machines

