Cross section re-design

- 1. Cross sections are needed: for tracking (elastic, inelastic, total, fission, capture) for models (internal)
- 2. For an unification of hadron-nucleus (nucleus-nucleus?) cross section access re-design is needed
- 3. K-meson-nucleus cross sections
- 4. Anti-proton-nucleus cross section
- 5. Nucleus-nucleus cross sections

Various cross sections parameterization at high energies



Hadronic cross section re-design G. Folger, V. Ivanchenko, D. Wright

- Too many classes in cross_sections subdirectory
- Function of classes are mixed
- Not easy identify what class can be used in Physics List, what inside a model, what combined class
- Not possible to exchange internal cross section between models
- Authorship problem

Simplest and backward compatible solution was chosen:

- Directory structure is not changed
- G4VCrossSectionDataSet interface was extended
- New G4VComponentCrossSection interface is added
- It is agreed that cross_section library should not depend on any model library
- Model library can depend on cross_section library

Extension of G4VCrossSectionDataSet

• New methods are added:

G4VComponentCrossSection interface

K-meson-nucleus cross sections



There are parameterization alternative to CHIPS. There are problems in other approaches.

Anti-proton-proton cross sections



There are parameterization alternative to CHIPS. There are problems in other approaches.

Glauber model can be used for hA cross section estimations with following Grichine's parameterization

A simple model for integral hadron-nucleus and nucleus-nucleus cross-sections. V.M. Grichine, Nucl. Instrum. Meth. B267: 2460 (2009).

A simplified Glauber model for hadron-nucleus cross sections. V.M. Grichine, Eur. Phys. J. C62:399 (2009).

$$\begin{split} \sigma_{tot}^{hA} &= 2\pi R^2 \ln \left[1 + \frac{A\sigma_{tot}^{hN}}{2\pi R^2} \right], \\ \sigma_{in}^{hA} &= \pi R^2 \ln \left[1 + \frac{A\sigma_{tot}^{hN}}{2\pi R^2} \right], \\ \sigma_{in}^{hA} &= \pi R^2 \ln \left[1 + \frac{A\sigma_{tot}^{hN}}{\pi R^2} \right], \\ \sigma_{in}^{hA} &= \pi R^2 \ln \left[1 + \frac{A\sigma_{tot}^{hN}}{\pi R^2} \right], \\ \sigma_{in}^{hA} &= \pi R^2 \ln \left[1 + \frac{A\sigma_{tot}^{hN}}{\pi R^2} \right], \\ \sigma_{in}^{ApAt} &= 2\pi (R_p^2 + R_t^2) \ln \left[1 + \frac{A_p A_t \sigma_{tot}^{NN}}{\pi (R_p^2 + R_t^2)} \right], \\ \sigma_{in}^{ApAt} &= \pi (R_p^2 + R_t^2) \ln \left[1 + \frac{A_p A_t \sigma_{tot}^{NN}}{\pi (R_p^2 + R_t^2)} \right], \\ \sigma_{prod}^{ApAt} &= \pi (R_p^2 + R_t^2) \ln \left[1 + \frac{A_p A_t \sigma_{in}^{NN}}{\pi (R_p^2 + R_t^2)} \right], \\ \sigma_{prod}^{ApAt} &= \pi (R_p^2 + R_t^2) \ln \left[1 + \frac{A_p A_t \sigma_{in}^{NN}}{\pi (R_p^2 + R_t^2)} \right], \\ \sigma_{prod}^{ApAt} &= \pi (R_p^2 + R_t^2) \ln \left[1 + \frac{A_p A_t \sigma_{in}^{NN}}{\pi (R_p^2 + R_t^2)} \right], \\ \sigma_{prod}^{ApAt} &= \pi (R_p^2 + R_t^2) \ln \left[1 + \frac{A_p A_t \sigma_{in}^{NN}}{\pi (R_p^2 + R_t^2)} \right], \\ \sigma_{prod}^{ApAt} &= \pi (R_p^2 + R_t^2) \ln \left[1 + \frac{A_p A_t \sigma_{in}^{NN}}{\pi (R_p^2 + R_t^2)} \right], \\ \sigma_{prod}^{ApAt} &= \pi (R_p^2 + R_t^2) \ln \left[1 + \frac{A_p A_t \sigma_{in}^{NN}}{\pi (R_p^2 + R_t^2)} \right], \\ \sigma_{prod}^{ApAt} &= \pi (R_p^2 + R_t^2) \ln \left[1 + \frac{A_p A_t \sigma_{in}^{NN}}{\pi (R_p^2 + R_t^2)} \right], \\ \sigma_{prod}^{ApAt} &= \pi (R_p^2 + R_t^2) \ln \left[1 + \frac{A_p A_t \sigma_{in}^{NN}}{\pi (R_p^2 + R_t^2)} \right], \\ \sigma_{prod}^{ApAt} &= \pi (R_p^2 + R_t^2) \ln \left[1 + \frac{A_p A_t \sigma_{in}^{NN}}{\pi (R_p^2 + R_t^2)} \right], \\ \sigma_{prod}^{ApAt} &= \pi (R_p^2 + R_t^2) \ln \left[1 + \frac{A_p A_t \sigma_{in}^{NN}}{\pi (R_p^2 + R_t^2)} \right], \\ \sigma_{prod}^{ApAt} &= \pi (R_p^2 + R_t^2) \ln \left[1 + \frac{A_p A_t \sigma_{in}^{NN}}{\pi (R_p^2 + R_t^2)} \right], \\ \sigma_{prod}^{ApAt} &= \pi (R_p^2 + R_t^2) \ln \left[1 + \frac{A_p A_t \sigma_{in}^{NN}}{\pi (R_p^2 + R_t^2)} \right], \\ \sigma_{prod}^{ApAt} &= \pi (R_p^2 + R_t^2) \ln \left[1 + \frac{A_p A_t \sigma_{in}^{NN}}{\pi (R_p^2 + R_t^2)} \right], \\ \sigma_{prod}^{ApAt} &= \pi (R_p^2 + R_t^2) \ln \left[1 + \frac{A_p A_t \sigma_{in}^{NN}}{\pi (R_p^2 + R_t^2)} \right], \\ \sigma_{prod}^{ApAt} &= \pi (R_p^2 + R_t^2) \ln \left[1 + \frac{A_p A_t \sigma_{in}^{NN}}{\pi (R_p^2 + R_t^2)} \right], \\ \sigma_{prod}^{ApAt} &= \pi (R_p^2 + R_t^2) \ln \left[1 + \frac{A_p A_t \sigma_{in}^{NN}}{\pi (R_p^2 + R_t^2)} \right], \\ \sigma_{prod}^{ApAt} &= \pi (R_p^2 + R_t^2) \ln \left[1 + \frac$$

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Hadron-Nucleus and Nucleus-Nucleus Cross Sections



Fig. 2. The integral inelastic cross-section for protons on carbon target versus the proton kinetic energy. Experimental data (open points and squares) from [11, 12]; the line corresponds to the simplified Glauber model.



Fig. 3. The integral inelastic cross-sections for protons on copper target versus the proton kinetic energy. Experimental data (open points and squares) from [11, 12]; the line corresponds to the simplified Glauber model.



Fig. 4. The integral inelastic cross-sections for neutrons on lead target versus the neutron kinetic energy. Experimental data (open points and squares) from [11,12]; the line corresponds to the simplified Glauber model.



Fig. 5. The integral total cross-sections for neutrons on helium target versus the neutron kinetic energy. Experimental data (open points and squares) from [11,12]; the line corresponds to the simplified Glauber model.



- 1. Access to the hadron-nucleus cross section is re-designed
- 2. K-meson-nucleus, anti-proton-nucleus and nucleus-nucleus cross sections would be well to estimate in the Glauber approximation.
- 3. The cross sections can be parameterized using Grichine's approach.