

# Dualities of QCD phase diagram and influence of chiral imbalance on color superconductivity phenomenon



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10th ICNFP 2021



Russian  
Science  
Foundation



Фонд развития  
теоретической физики  
и математики

K.G. Klimenko, IHEP  
T.G. Khunjua, University of Georgia, MSU

**in the broad sense our group stems from**  
Department of Theoretical Physics, Moscow State University  
Prof. V. Ch. Zhukovsky

details can be found in

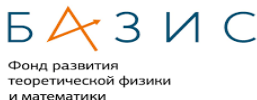
Eur.Phys.J.C 80 (2020) 10, 995 arXiv:2005.05488 [hep-ph]  
JHEP 06 (2020) 148 arXiv:2003.10562 [hep-ph]  
Phys.Rev. D100 (2019) no.3, 034009 arXiv: 1904.07151 [hep-ph]  
JHEP 1906 (2019) 006 arXiv:1901.02855 [hep-ph]  
Eur.Phys.J. C79 (2019) no.2, 151, arXiv:1812.00772 [hep-ph],  
Phys.Rev. D98 (2018) no.5, 054030 arXiv:1804.01014 [hep-ph],  
Phys.Rev. D97 (2018) no.5, 054036 arXiv:1710.09706 [hep-ph]  
Phys.Rev. D95 (2017) no.10, 105010 arXiv:1704.01477 [hep-ph]

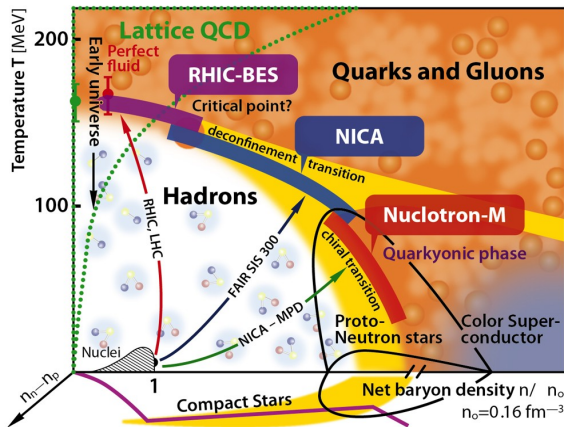
The work is supported by

- ▶ Russian Science Foundation (RSF)  
under grant number 19-72-00077



- ▶ Foundation for the Advancement of Theoretical Physics and Mathematics





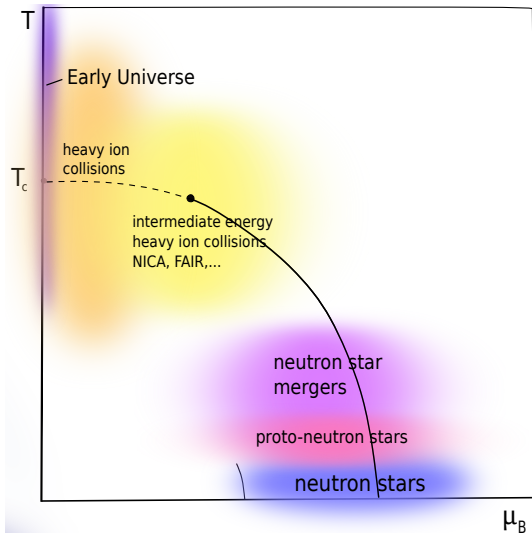
Two main phase transitions

- ▶ confinement-deconfinement
- ▶ chiral symmetry breaking phase—chiral symmetric phase



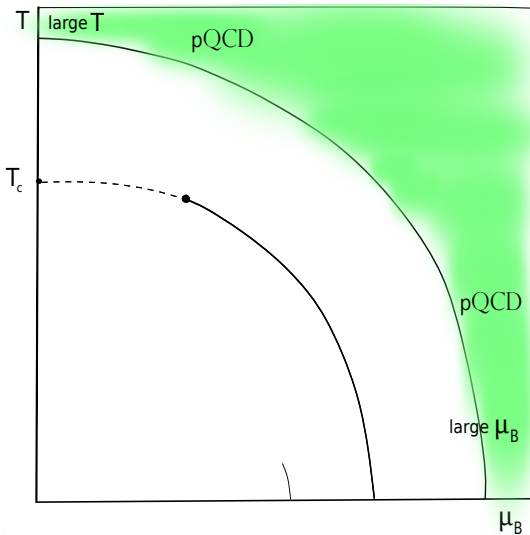
QCD at  $T$  and  $\mu$   
(QCD at extreme conditions)

- ▶ Early Universe
- ▶ heavy ion collisions
- ▶ neutron stars
- ▶ proto- neutron stars
- ▶ neutron star mergers



Methods of dealing with QCD

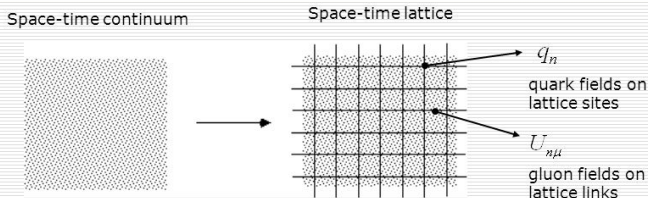
► Perturbative QCD





## QCD on a space-time lattice

*K. G. Wilson 1974*



### □ Feynman path integral

■ Action  $S_{QCD} = \frac{1}{g_s^2} \sum_P n(UUUU) + \sum_f \bar{q}_f (\gamma \cdot U + m_f) q_f$

- Physical quantities as **integral averages**



*Monte Carlo  
Evaluation of  
the path integral*

$$\langle O(U, \bar{q}, q) \rangle = \frac{1}{Z} \int \prod_{n\mu} dU_{n\mu} \prod_n d\bar{q}_n dq_n O(U, (\bar{q}, q)) e^{-S_{QCD}}$$

$$Z = \int D[\text{gluons}] D[\text{quarks}] e^{-S_{\text{QCD}}^E}$$

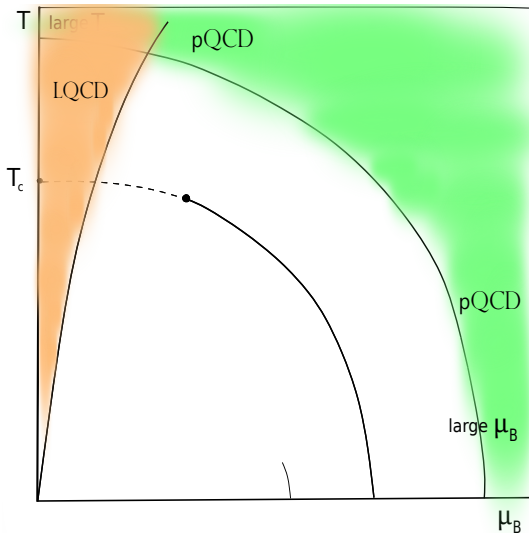
$$Z = \int D[\text{gluons}] \text{Det} D(\mu) e^{-S_{\text{gluons}}^E}$$

It is well known that **at non-zero baryon chemical potential  $\mu_B$  lattice simulation** is quite challenging due to the **sign problem**  
complex determinant

$$(\text{Det}(D(\mu)))^\dagger = \text{Det}(D(-\mu^\dagger))$$

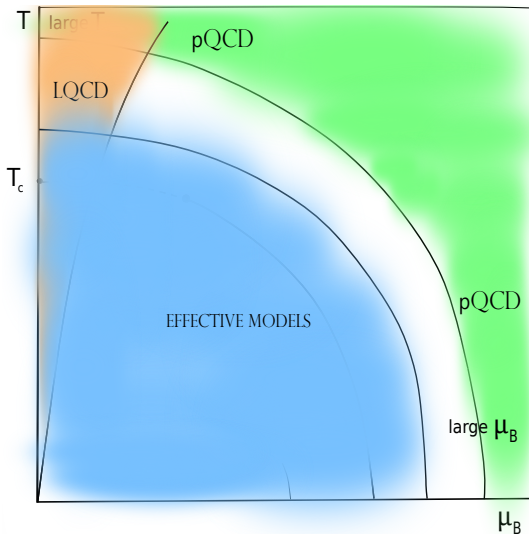
## Methods of dealing with QCD

- ▶ Perturbative QCD
- ▶ First principle calculation  
– lattice QCD



## Methods of dealing with QCD

- ▶ Perturbative QCD
- ▶ First principle calculation  
– lattice QCD
- ▶ Effective models
- ▶ DSE, FRG
- ▶ .....



NJL model can be considered as **effective model for QCD**.

the model is **nonrenormalizable**

Valid up to  $E < \Lambda \approx 1 \text{ GeV}$

$$\mu, T < 600 \text{ MeV}$$

Parameters  $G, \Lambda, m_0$

**chiral limit**  $m_0 = 0$

in many cases chiral limit is a very good approximation

dof- **quarks**

no gluons only **four-fermion interaction**

attractive feature — dynamical CSB

the main drawback – lack of confinement (PNJL)

Nambu–Jona-Lasinio model

$$\mathcal{L} = \bar{q}\gamma^\nu i\partial_\nu q + \frac{G}{N_c} \left[ (\bar{q}q)^2 + (\bar{q}i\gamma^5 q)^2 \right]$$

$$q \rightarrow e^{i\gamma^5 \alpha} q$$

continuous symmetry

$$\tilde{\mathcal{L}} = \bar{q} \left[ \gamma^\rho i\partial_\rho - \sigma - i\gamma^5 \pi \right] q - \frac{N_c}{4G} \left[ \sigma^2 + \pi^2 \right].$$

**Chiral symmetry breaking**

$1/N_c$  expansion, leading order

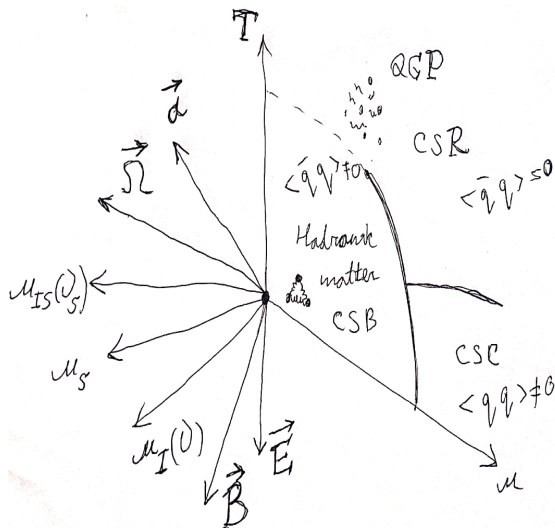
$$\langle \bar{q}q \rangle \neq 0$$

$$\langle \sigma \rangle \neq 0 \quad \longrightarrow \quad \tilde{\mathcal{L}} = \bar{q} \left[ \gamma^\rho i\partial_\rho - \langle \sigma \rangle \right] q$$



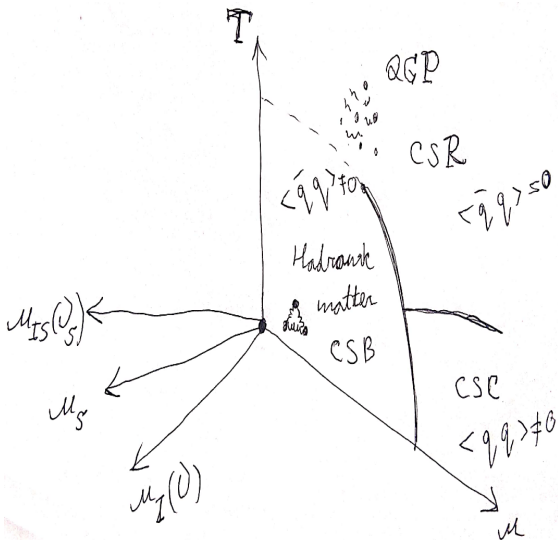
More than just QCD at  $(\mu, T)$

- ▶ more chemical potentials  $\mu_i$
- ▶ magnetic fields
- ▶ rotation of the system  $\vec{\Omega}$
- ▶ acceleration  $\vec{a}$
- ▶ finite size effects (finite volume and boundary conditions)



More than just QCD at  $(\mu, T)$

- ▶ **more chemical potentials**  $\mu_i$
- ▶ magnetic fields
- ▶ rotation of the system
- ▶ acceleration
- ▶ finite size effects (finite volume and boundary conditions)



**Baryon chemical potential  $\mu_B$** 

Allow to consider systems with non-zero baryon densities.

$$\frac{\mu_B}{3} \bar{q} \gamma^0 q = \mu \bar{q} \gamma^0 q, \quad n_B = \frac{1}{3}(n_u + n_d)$$

**Baryon chemical potential  $\mu_B$** 

Allow to consider systems with non-zero baryon densities.

$$\frac{\mu_B}{3} \bar{q} \gamma^0 q = \mu \bar{q} \gamma^0 q, \quad n_B = \frac{1}{3} (n_u + n_d)$$

**Isotopic chemical potential  $\mu_I$** 

Allow to consider systems with isospin imbalance ( $n_n \neq n_p$ ).

$$\frac{\mu_I}{2} \bar{q} \gamma^0 \tau_3 q = \nu (\bar{q} \gamma^0 \tau_3 q)$$

$$n_I = n_u - n_d \quad \longleftrightarrow \quad \mu_I = \mu_u - \mu_d$$

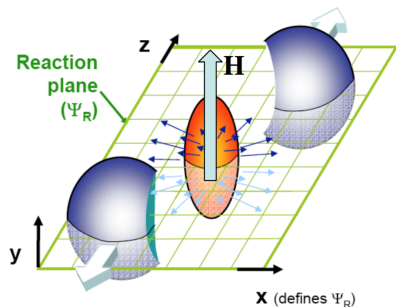
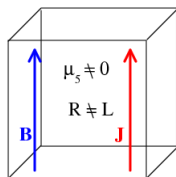
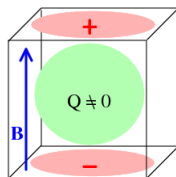
**chiral (axial) chemical potential**

Allow to consider systems with chiral imbalance (difference between densities of left-handed and right-handed quarks).

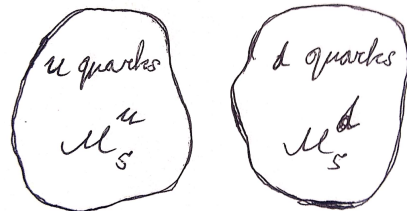
$$n_5 = n_R - n_L \quad \longleftrightarrow \quad \mu_5 = \mu_R - \mu_L$$

The corresponding term in the Lagrangian is

$$\mu_5 \bar{q} \gamma^0 \gamma^5 q$$



$$\vec{J} \sim \mu_5 \vec{B},$$



$$\mu_5^u \neq \mu_5^d \quad \text{and} \quad \mu_{I5} = \mu_5^u - \mu_5^d$$

Term in the Lagrangian —  $\frac{\mu_{I5}}{2} \bar{q} \tau_3 \gamma^0 \gamma^5 q = \nu_5 (\bar{q} \tau_3 \gamma^0 \gamma^5 q)$

$$n_{I5} = n_{u5} - n_{d5}, \quad n_{I5} \longleftrightarrow \nu_5$$

- ▶ Chiral isospin imbalance and chiral imbalance  
 $\mu_{I5}$  and  $\mu_5$  can be generated in parallel magnetic and electric fields  $\vec{E} \parallel \vec{B}$
  
- ▶ Chiral imbalance could appear in dense matter
  - ▶ Chiral separation effect  
*(Thanks for the idea to Igor Shovkovy)*
  
  - ▶ Chiral vortical effect



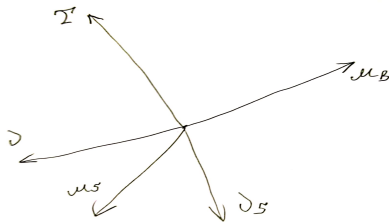
Different chemical potentials and matter content

$$\mu = \frac{\mu_B}{3}, \quad \nu = \frac{\mu_I}{2}, \quad \mu_5, \quad \nu_5 = \frac{\mu_{I5}}{2}$$

- ▶ **QCD phase diagram**  
with **different chemical potentials**  
and matter content including **chiral imbalance**
  
- ▶ **QC<sub>2</sub>D phase diagram** and **diquark condensation** phenomenon  
with different chemical potentials,  
including  $\mu_5$

## Different chemical potentials and matter content

$$\mu = \frac{\mu_B}{3}, \quad \nu = \frac{\mu_I}{2}, \quad \mu_5, \quad \nu_5 = \frac{\mu_{I5}}{2}$$



# Dualities

It is not related to holography or gauge/gravity  
duality

it is the dualities of the phase structures of  
different systems

# Dualities

Chiral symmetry breaking  $\iff$  pion condensation

Isospin imbalance  $\iff$  Chiral imbalance

The TDP

$$\Omega(T, \mu, \mu_i, \dots, \langle \bar{q}q \rangle, \dots)$$

The TDP

$$\Omega(T, \mu, \mu_i, \dots, \langle \bar{q}q \rangle, \dots)$$

$$\Omega(T, \mu, \nu, \nu_5, \dots, M, \pi, \dots)$$

The TDP

$$\Omega(T, \mu, \mu_i, \dots, \langle \bar{q}q \rangle, \dots) \qquad \Omega(T, \mu, \nu, \nu_5, \dots, M, \pi, \dots)$$

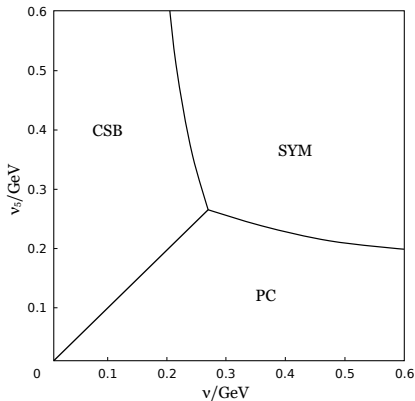
The TDP (phase diagram) is invariant under  
Interchange of - condensates - matter content

$$\Omega(M, \pi, \nu, \nu_5)$$

$$M \longleftrightarrow \pi, \qquad \nu \longleftrightarrow \nu_5$$

$$\Omega(M, \pi, \nu, \nu_5) = \Omega(\pi, M, \nu_5, \nu)$$





$$\mathcal{D} : M \longleftrightarrow \pi, \quad \nu \longleftrightarrow \nu_5$$

Duality between chiral  
symmetry breaking and pion  
condensation

$$\text{PC} \longleftrightarrow \text{CSB} \quad \nu \longleftrightarrow \nu_5$$

Figure: NJL model results

$$\mathcal{L}_{\text{QCD}} = \sum_{f=u,d} \bar{q}_f (i\not{D} - m_f) q_f - \frac{1}{4} \mathcal{G}_{\mu\nu,a} \mathcal{G}_a^{\mu\nu}.$$

$$\mathcal{L}_{\text{NJL}} = \sum_{f=u,d} \bar{q}_f \left[ i\gamma^\nu \partial_\nu - m_f \right] q_f + \frac{G}{N_c} \left[ (\bar{q}q)^2 + (\bar{q}i\gamma^5 \vec{\tau}q)^2 \right]$$

$m_f$  is current quark masses

**In the chiral limit  $m_f = 0$  the Duality  $\mathcal{D}$  is exact**

$$m_f : \frac{m_u + m_d}{2} \approx 3.5 \text{ MeV}$$

In our case typical values of  $\mu, \nu, \dots, T, \dots \sim 10 - 100$ s MeV, for example, 200-400 MeV

One can work in the chiral limit  $m_f \rightarrow 0$

$$m_f = 0 \quad \rightarrow \quad m_\pi = 0$$

physical  $m_f$  a few MeV  $\rightarrow$  physical  $m_\pi \sim 140$  MeV

Duality between CSB and PC is **approximate** in **physical point**

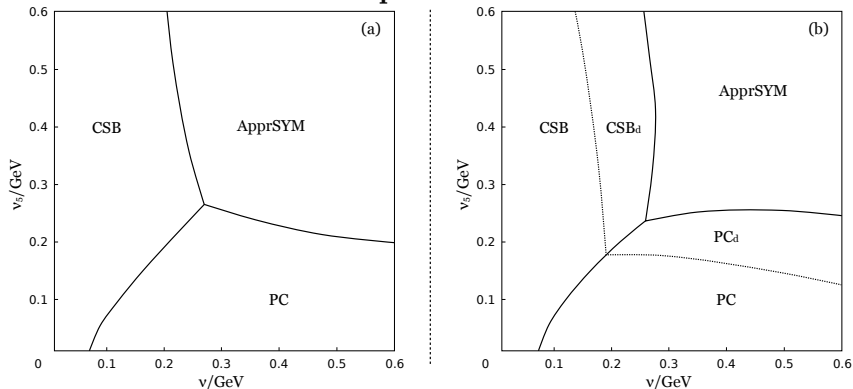


Figure:  $(\nu, \nu_5)$  phase diagram

$\mu_B \neq 0$  impossible on lattice but if  $\mu_B = 0$

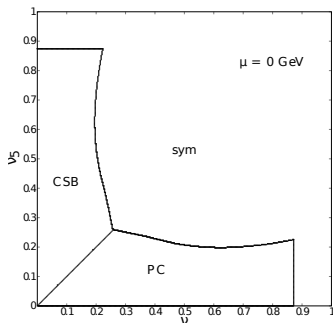
Duality is shown to take place in particular cases in lattice QCD

► **QCD at  $\mu_5$**  —  $(\mu_5, T)$

V. Braguta, A. Kotov et al, ITEP lattice group

► **QCD at  $\mu_I$**  —  $(\mu_I, T)$

G. Endrodi, B. Brandt et al, Emmy Noether junior research group, Goethe-University Frankfurt, Institute for Theoretical Physics ()



# Uses of Dualities

A few rather interesting uses of dualities

*discussed in Particles 2020, 3(1), 62-79*

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Two colour QCD case

$QC_2D$

There are a lot similarities:

- ▶ similar phase transitions:

*confinement/deconfinement, chiral symmetry  
breaking/restoration at large  $T$  and  $\mu$*

- ▶ A lot of physical quantities coincide up to few dozens percent

*Critical temperature  $T_c/\sqrt{\sigma}$ , topological susceptibility  
 $\chi^{\frac{1}{4}}/\sqrt{\sigma}$  shear viscosity  $\eta/s$*

There are **no sign problem** in SU(2) case

$$(Det(D(\mu)))^\dagger = Det(D(\mu))$$

and lattice simulations at non-zero baryon  
density are possible

It is a great playground for studying dense matter



# Phase diagram of $\text{QC}_2\text{D}$

## Condensates and phases

$$M = \langle \sigma(x) \rangle \sim \langle \bar{q}q \rangle,$$

CSB phase:  $M \neq 0$ ,

$$\pi_1 = \langle \pi_1(x) \rangle = \langle \bar{q}\gamma^5\tau_1q \rangle,$$

PC phase:  $\pi_1 \neq 0$ ,

$$\Delta = \langle \Delta(x) \rangle = \langle qq \rangle = \langle q^T C \gamma^5 \sigma_2 \tau_2 q \rangle,$$

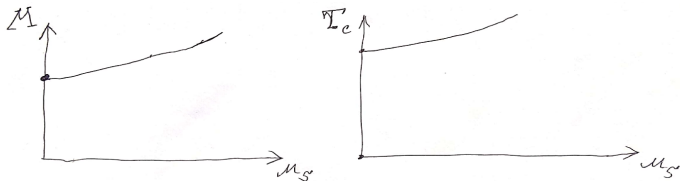
BSF phase:  $\Delta \neq 0$ .

A number of papers predicted **anticatalysis** ( $T_c$  decrease with  $\mu_5$ ) of dynamical chiral symmetry breaking

A number of papers predicted **catalysis** ( $T_c$  increase with  $\mu_5$ ) of dynamical chiral symmetry breaking

lattice results show the **catalysis**

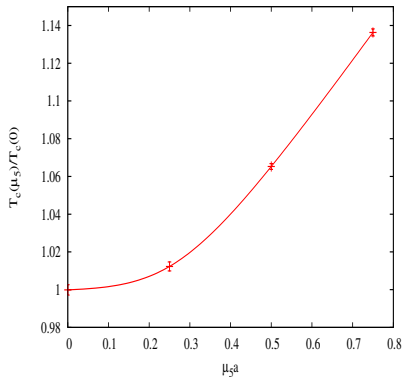
(ITEP lattice group, V. Braguta, A. Kotov, et al)

QCD at non-zero  $\mu_5$ 

catalysis of CSB by chiral imbalance:

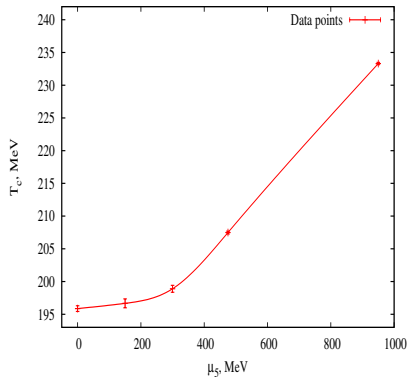
- ▶ increase of  $\langle \bar{q}q \rangle$  as  $\mu_5$  increases
- ▶ increase of critical temperature  $T_c$  of chiral phase transition (crossover) as  $\mu_5$  increases

## SU(2)

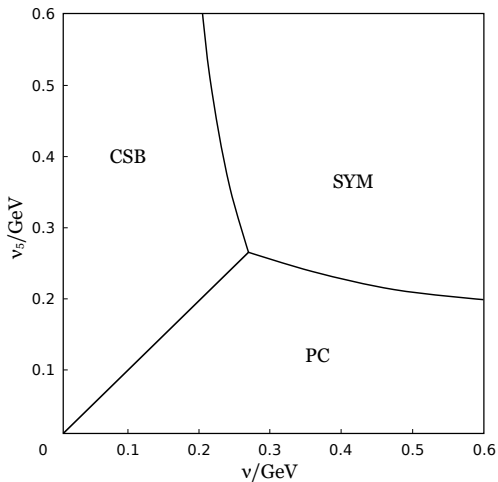


V. Braguta, A. Kotov et al, *JHEP* 1506, 094  
(2015), *PoS LATTICE 2014*, 235 (2015)

## SU(3)

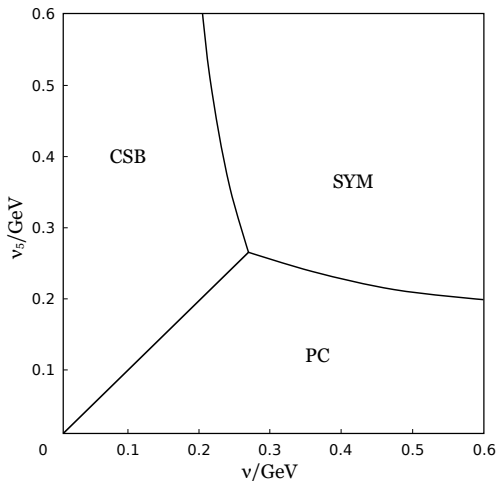


V. Braguta, A. Kotov et al, *Phys. Rev. D* 93,  
034509 (2016), *arXiv:1512.05873 [hep-lat]*



*Isospin imbalance  $\nu$   
generates pion condensation*

*Chiral imbalance  $\nu_5$   
generates chiral symmetry  
breaking*

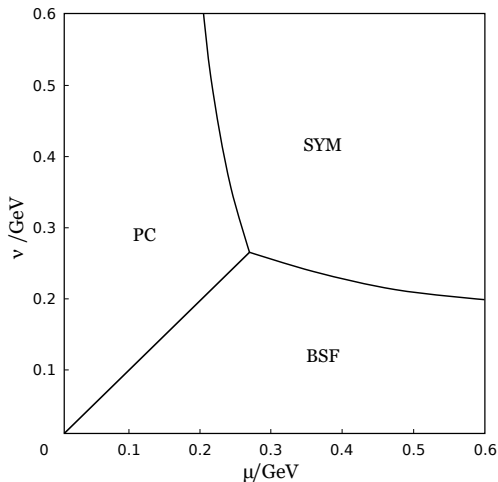


Duality between chiral  
symmetry breaking and  
pion condensation

$$\text{PC} \longleftrightarrow \text{CSB} \quad \nu \longleftrightarrow \nu_5$$

$$\mathcal{D}_3 : \quad M \longleftrightarrow \pi_1, \quad \nu \longleftrightarrow \nu_5$$

$$\Omega(M, \pi, \dots) = \Omega(\pi, M, \dots)$$



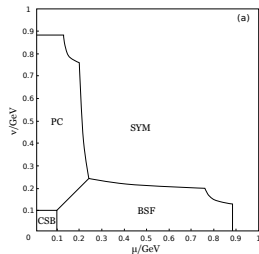
*Baryon density  $\mu$  generates  
color superconductivity*

$$\text{PC} \longleftrightarrow \text{BSF} \quad \nu \longleftrightarrow \nu_5$$

$$\mathcal{D}_1 : \quad \pi_1 \longleftrightarrow |\Delta|, \quad \mu \longleftrightarrow \nu$$

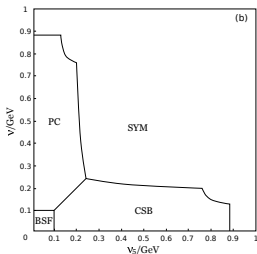
$$\Omega(\pi, \Delta, \dots) = \Omega(\Delta, \pi, \dots)$$





$$(a) \quad \mathcal{D}_1 : \quad \mu \longleftrightarrow \nu, \quad \pi_1 \longleftrightarrow |\Delta|, \quad PC \longleftrightarrow BSF$$

*J. Andersen, T. Brauner, D. T. Son, M. Stephanov, J. Kogut, ...*

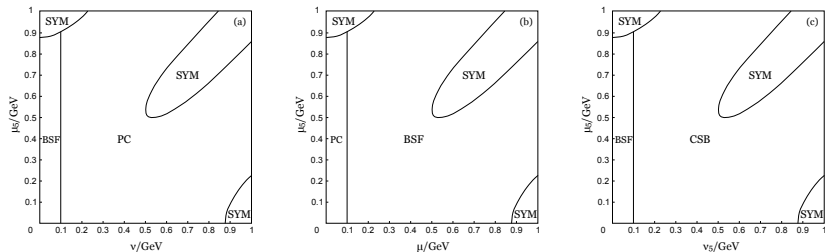


$$(b) \quad \mathcal{D}_3 : \quad \nu \longleftrightarrow \nu_5, \quad M \longleftrightarrow \pi_1, \quad PC \longleftrightarrow CSB$$

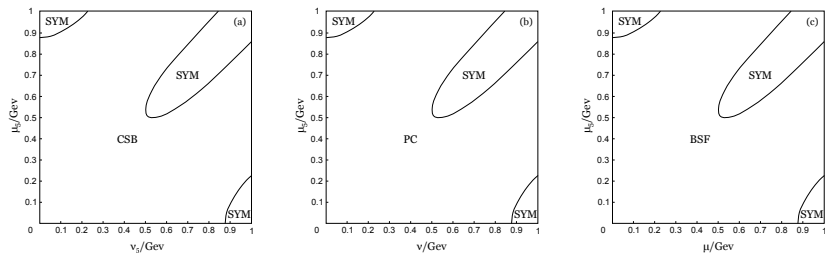
$$(c) \quad \mathcal{D}_2 : \quad \mu \longleftrightarrow \nu_5, \quad M \longleftrightarrow |\Delta|, \quad CSB \longleftrightarrow BSF$$

Each chemical potential is connected  
in one-to-one correspondence with some  
phenomenon (condensation)

- ▶ Baryon density  $\mu$   $\iff$  diquark condensation
  - ▶ Isospin imbalance  $\nu$   $\iff$  pion condensation
  - ▶ Chiral imbalance  $\nu_5$   $\iff$  chiral symmetry breaking
-



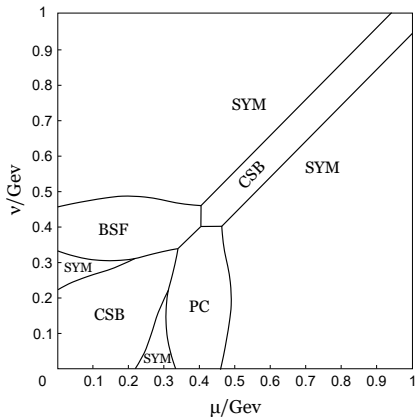
Chiral imbalance  $\mu_5$  catalyzes all the phenomena



Chameleon nature of chiral imbalance  $\mu_5$

$\mu_5$  mimics other chemical potentials  $\mu, \nu, \nu_5$

Chiral imbalance  $\mu_5$  leads to several rather peculiar phases in the system, e. g. the **diquark condensation** in the region of the phase diagram at  $\mu = 0$



$(\mu, \nu, \nu_5)$  phase diagram is highly symmetric due to dualities

and intermingled by dualities at  $\mu_5 \neq 0$

$\mu_5$  deforms the  $(\mu, \nu, \nu_5)$  phase diagram

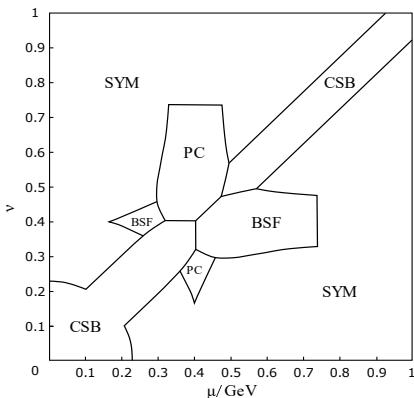
The influence of  $\mu_5$  is constrained by dual properties

- ▶ Chameleon nature of chiral imbalance  $\mu_5$  is also a consequence of dualities
  - ▶ The feature of  $\mu_5$  of being universal catalizer is a consequence of dualities as well
-

There are a lot of features of  $QC_2D$  phase diagram that remains the same in the case of QCD

Including the behaviour of diquark condensation in quark matter with different imbalances

- ▶  $PC_d$  phase has been predicted without possibility of diquark condensation
- ▶ Diquark condensation can take over the  $PC_d$  phase
- ▶ In two colour case diquark condensation is in a sense even stronger than in three colour case and starts from  $\mu > 0$



$PC_d$  phase is unaffected by BSF phase in two color case.  
 Maybe one can infer that it is the case also for 3 color QCD



- ▶  $PC_d$  phase has been predicted without possibility of diquark condensation
  - ▶ Diquark condensation can take over the  $PC_d$  phase
- 

$PC_d$  phase is unaffected by CSC phase in three color case.

Dualities  $\mathcal{D}_1$ ,  $\mathcal{D}_2$  and  $\mathcal{D}_3$  were found in

- In the framework of NJL model
  - In the mean field approximation
-

Dualities are connected with Pauli-Gursey group

Dualities were found in

- In the framework of NJL model  
beyond mean field
- In  $QC_2D$  non-perturbatively (at the level of  
Lagrangian)

Duality  $\mathcal{D}$  is a remnant of chiral symmetry

Duality was found in

- ▶ In the framework of NJL model beyond mean field or at all orders of  $N_c$  approximation
  - ▶ In QCD non-perturbatively (at the level of Lagrangian)
-

- ▶  $(\mu_B, \mu_I, \nu_5, \mu_5)$  phase diagram was studied in two color color case
- ▶ It was shown that there exist dualities in QCD and  $QC_2D$   
*Richer structure of Dualities in the two colour case*
- ▶ There have been shown ideas how dualities can be used  
*Duality is not just entertaining mathematical property but an instrument with very high predictivity power*
- ▶ Dualities have been shown non-perturbatively in the two colour case
- ▶ Duality has been shown non-perturbatively in QCD