Nuclear Astrophysics with Lasers

A. Bonasera
Texas A&M, College Station; INFN-LNS, CATANIA

INTRODUCTION:
Create energetic ion beams under specific physical conditions, for basic nuclear science and applications.

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Results from SGII-Shanghai experiments on CD\(_2\) targets

NEXT

CONCLUSIONS

Measuring S-factors in hot and dense plasma

\[ N = \iiint dV \, dt \, n_1 \, n_2 \, \langle \sigma_{12} \, v \rangle, \]

Measure the number of fusion \( N \), the plasma phase space densities (i.e. \( T \)) \( n_1, n_2 \) (e.g. \( p+B \)) and volumes then recover the cross-section (or S-factor)

\[
(\sigma \, v) = \frac{8}{\mu \pi} \frac{1}{(k_B T)^{3/2}} \int_0^{\infty} S(E) \exp \left( -\frac{E}{k_B T} - \frac{b}{\sqrt{E}} \right) dE
\]

Add a third axis: density!

Alpha tracks from laser \((p+B)\) interaction at ABC-ENEAA (LAPLAFUS coll., W.Sci. in press)
Measuring S-factors in hot and dense plasma

Measure the number of fusion N, the plasma phase space densities (i.e. T) n1,n2 (e.g. p+B) and volumes then recover the cross-section (or S-factor)

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Add a third axis: density!

Alpha tracks from laser (p+B) interaction at ABC-ENEA (LAPLAFUS coll., W.Sci. in press)
1.0 Title: Measuring Cluster Fusion Plasma Temperature and Density from $^3$He(d,p)$^4$He and d(d,p)T Reactions

2.0 PI, Co-PI’s & Affiliation

PIs: Todd Ditmire (UT Austin), Aldo Bonasera (Texas A&M, LNS INFN Catania-Italy)
Co-investigators: W. Bang, G. Dyer, H. Quevedo, A. Bernstein (CHEDS, UT Austin)
M. Barbui, M. Barbarino, G. Giuliani, K. Hagel, Z. Hua, J. Natowitz, K. Schmidt
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A. Caruso (Kore University, Enna-Italy)
J. Sura (Heavy Ion Institute, Warsaw, Poland)
S. Kimura (LNS INFN Catania-Italy)
P. Andreoli, R. DeAngelis, F. Consoli, ENEA

3.0 Experimental Objectives and Concept

We propose to conduct experiments that follow on from the successful deuterium cluster fusion experiments in early 2011 on the TPW. Specifically we propose a detailed investigation of ion temperature in hot exploding cluster plasmas. To do this we will simultaneously measure the experimental yield from two different nuclear reactions. While our first experiments utilized pure deuterium to drive the d(d,p)T and d(d,n)He$^3$ reactions we now propose to mix He$^3$ into the gas jet target to allow us to measure simultaneously yields from the He$^3$(d,p)He$^4$ and the d-d reactions. Because these two reactions have different cross sections, measuring the ratio of the yields between these two reactions will allow a precise determination of the plasma temperature at the time when the reaction occurred (assuming thermalization). The measure of the experimental yield from sequential reactions will also make possible a direct measurement of the plasma density at the time of the reaction. Once the...
High power laser can be used to generate neutrons from the fusion reaction

Nuclear fusion from laser-cluster interaction

- Most of the laser pulse energy is absorbed by the atomic clusters.
- Clusters experience Coulomb explosion after electrons escape.
- DD fusion occurs, and 2.45MeV fusion neutrons are produced.
Expected fusion reactions:

\[ D + D \rightarrow T (1.01 \text{ MeV}) + p (3.02 \text{ MeV}) \] (50%)

\[ D + D \rightarrow ^3\text{He} (0.82 \text{ MeV}) + n (2.45 \text{ MeV}) \] (50%)

\[ D + ^3\text{He} \rightarrow ^4\text{He} (3.6 \text{ MeV}) + p (14.69 \text{ MeV}) \] (100%)

\[ D + T \rightarrow ^4\text{He} (3.5 \text{ MeV}) + n (14.1 \text{ MeV}) \] (100%)
Experimental setup

- Cryo-cooled nozzle
- Laser beam direction
- Proton A
- Proton B
- Proton C
- Neutron 8
- Neutron 13
- Neutron 9
- Neutron 16
- Faraday cup
- 3 plastic scintillators from UT
Measure V for each event
Measured observables

Temperature and the number of the energetic ions

Yield of 2.45 MeV neutrons

Yield of 3.02 MeV protons

Yield of 14.7 MeV protons

Faraday cup

4 liquid scintillators NE213 placed at different angles and 6 plastic scintillators placed at 90 degrees

Thin (254 um) plastic scintillators BC400.
Measured observables

Measure N. of ions

Temperature and the number of the energetic ions

Faraday cup

Yield of 2.45 MeV neutrons

4 liquid scintillators NE213 placed at different angles and 6 plastic scintillators placed at 90 degrees

Yield of 3.02 MeV protons

Thin (254 um) plastic scintillators BC400.

Yield of 14.7 MeV protons

Back ch2 shot 2779
Measure N. of ions

Temperature and the number of the energetic ions

Yield of 2.45 MeV neutrons

Yield of 3.02 MeV protons

Yield of 14.7 MeV protons

Measure N. of fusions for each channel
Gamow peak approximation near strong resonances

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(Received 29 November 2012; revised manuscript received 16 April 2013; published 23 May 2013)

We discuss the most effective energy range for charged-particle-induced reactions in a plasma environment at a given plasma temperature. The correspondence between the plasma temperature and the most effective energy should be modified from the one given by the Gamow peak energy, in the presence of a significant incident-energy dependence in the astrophysical $S$ factor as in the case of resonant reactions. The suggested modification of the effective energy range is important not only in thermonuclear reactions at high temperature in the stellar environment, e.g., in advanced burning stages of massive stars and in explosive stellar environments, as has been already claimed, but also in the application of nuclear reactions driven by ultra-intense laser-pulse
Here maybe high repetition laser low energy fs pulse (BUDAPEST?)

PHYSICAL REVIEW C 87, 058801 (2013)

Gamow peak approximation near strong resonances

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Here maybe high repetition laser low energy fs pulse (BUDAPEST?)

Here maybe the 10PW @ ELI-NP/ALPS/MED
Temperature Measurements of Fusion Plasmas Produced by Petawatt-Laser-Irradiated \(D_2 - ^3\text{He}\) or \(CD_4 - ^3\text{He}\) Clustering Gases


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Two different methods have been employed to determine the plasma temperature in a laser-cluster fusion experiment on the Texas Petawatt laser. In the first, the temperature was derived from time-of-flight data of deuterium ions ejected from exploding \(D_2\) or \(CD_4\) clusters. In the second, the temperature was measured from the ratio of the rates of two different nuclear fusion reactions occurring in the plasma at the same time: \(\text{D}(d, ^3\text{He})n\) and \(^3\text{He}(d, p)^4\text{He}\). The temperatures determined by these two methods agree well, which indicates that (i) the ion energy distribution is not significantly distorted when ions travel in the disassembling plasma; (ii) the kinetic energy of deuterium ions, especially the “hottest part” responsible for nuclear fusion, is well described by a near-Maxwellian distribution.

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PACS numbers: 52.50.Jm, 25.45.-z, 36.40.Wa

Nuclear fusion from explosions of laser-heated clusters has been an active research topic for over a decade [1–11]. Researchers have used explosions of cryogenically cooled deuterium (\(D_2\)) cluster targets or near-room-temperature deuterated methane cluster plasmas produced by the irradiation of a clustering gas jet by 150 fs petawatt peak power laser pulses. We find that the effective ion temperature produced can be in excess of 25 keV.
Measurement of the Plasma Astrophysical S Factor for the $^3\text{He} + ^2\text{H} \rightarrow ^4\text{He}$ Reaction in Exploding Molecular Clusters


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(Received 1 July 2013)

Synopsis: Nuclear Reactions in Lab Plasma

Many low-energy nuclear reactions in astrophysics occur in plasmas, in which the nuclei are free of electrons. By contrast, most nuclear experiments involve neutral targets, whose bound electrons produce a “screening effect.” A new technique uses lasers to remove these unwanted electrons so that low-energy nuclear reactions can be studied directly in laboratory plasma. The authors demonstrate their approach in Physical Review Letters on the deuterium/helium-3 interaction that helped synthesize elements in the early Universe and could potentially be used to power a future nuclear fusion reactor.

In a typical nuclear reaction experiment, an ion beam is directed at a target containing neutral atoms. The bound electrons provide a screen that reduces the Coulomb repulsion between the positive nuclei. Therefore, laboratory measurements tend to predict higher reaction rates than would be expected between ionized nuclei. To obtain astrophysically relevant parameters, researchers try to correct their data by estimating the screening effect of the bound electrons.
Model-independent determination of the astrophysical S factor in laser-induced fusion plasmas

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In this work, we present a new and general method for measuring the astrophysical S factor of nuclear reactions in laser-induced plasmas and we apply it to \(^3\)He. The experiment was performed with the Texas Petawatt Laser, which delivered 150–270 fs pulses of energy ranging from 90 to 180 J to D2 or CD2 molecular clusters (where \(D\) denotes \(^{2}\)H). After removing the background noise, we used the measured time-of-flight data of energetic deuterium ions to obtain their energy distribution. We derive the S factor using the measured energy distribution of the ions, the measured volume of the fusion plasma, and the measured fusion yields. This method is model independent in the sense that no assumption on the state of the system is required, but it requires an accurate measurement of the ion energy distribution, especially at high energies, and of the relevant fusion yields. In the \(^3\)He(d,n)\(^4\)He and \(^3\)He(d,p)\(^4\)He cases discussed here, it is very important to apply the background subtraction for the energetic ions and to measure the fusion yields with high precision. While the available data on both ion distribution and fusion yields allow us to determine with good precision the S factor in the \(d + d\) case (lower Gamow energies), for the \(d + ^3\)He case the data are not precise enough to obtain the S factor using this method. Our results agree with other experiments within the experimental error, even though smaller values of the S factor were obtained. This might be due to the plasma environment differing from the beam target conditions in a conventional accelerator experiment.

DOI: 10.1103/PhysRevC.93.045808

I. INTRODUCTION

The nuclear reactions between light nuclei in the low energy region (\(\sim \text{keV}\)),

\[
d + d \rightarrow ^3\text{He}(0.82\ \text{MeV}) + n(2.45\ \text{MeV}),
\]

\[
d + d \rightarrow p(3.02\ \text{MeV}) + t(1.01\ \text{MeV}),
\]

\[
d + ^3\text{He} \rightarrow p(14.7\ \text{MeV}) + ^4\text{He}(3.6\ \text{MeV}).
\]

have been studied for many decades [1–10]. The role of low-energy nuclear physics is crucial in both astrophysics, playing a key role in the determination of primordial abundances in Big Bang nucleosynthesis (BBN) models, and applied (plasma) physics, as it lies in the energy region of interest for the operation and design of future fusion power plants. Direct and with bare nuclei and with the ones occurring in astrophysical plasmas [1,6,13,14].

Other physical conditions are possible which might decrease the astrophysical factor, dubbed the dissipative limit (DL) in [11,12]. In a hot plasma, due to the large number of positive and negative charges, fusions occurring in an “electron” cloud might be enhanced. If, however, a large number of positive charges is present in the region where fusion occurs, then the cross section might decrease. In laser-cluster interactions we might be able to create such conditions, thus it would represent a good chance to study the fusion cross sections within stellar plasmas in a laboratory. In particular, we can explore temperatures ranging from few keV up to few tens of keV and a density just above \(10^{18}\) atoms/cm\(^3\). These conditions are similar to those prevailing in the BBN era, and to those occurring in fusion plasmas.
$d + D \rightarrow ^3\text{He} + n$
Range of plasma ions in cold cluster gases near the critical point


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Ion range in systems prepared near a liquid–gas phase transition

ABSTRACT

We measure the range of plasma ions in cold cluster gases by using the Petawatt laser at the University of Texas-Austin. The produced plasma propagated in an ambient gas mixing the cold cluster gas not illuminated by the laser. From the ratio of the measured ion distributions at different angles we can estimate the range of the ions in the cold cluster gas. It is much smaller than estimated using popular models, which take only into account the slowing down of charged particles in uniform matter. We discuss the ion range in systems prepared near a liquid–gas phase transition.

Nuclear fusion from explosions of laser-heated clusters has been an active research topic for over a decade [1–14]. The explosions of cryogenically cooled deuterium (D 2 ) cluster targets or near-room-temperature deuterium-methane (D 3 ) cluster targets drive fusion reactions. A high intensity femtosecond laser pulse irradiates the cluster gas. This produces energetic explosions of the clusters and ions of keV ion plasma temperature results. DID fusion occurring within this high temperature plasma combined with beam-target fusion, between the ejected ions of the cluster and surrounding cold cluster gas, leads to a burst of fusion neutrons and protons. Following these experiments, we have modified some aspects in order to be able to measure the range of energetic ions in the cold cluster gases. Recall that the range of ions is crucial to estimate the fusion rates in the plasma. We have opportunely focused the laser in such a way that the high-energy pulse drills a “hole” in the target. We found that less than 1% of the laser energy went through the cluster gas for each shot.

A schematic view of this scenario is plotted in Fig. 1a, while an actual experimental result is given in Fig. 1b. Two Faraday cups (FC) were opportune located: the first one (UTCFC) as close as possible to the incoming laser direction (−67.5° minimum) thus measuring essentially the hot plasma only; the second one (CTFC) was located at an angle around 45°, see Fig. 1, and comparable to the physical constraints of the laboratory (walls). The ratio of the FC signals gives an indication of the range of the ions in the surrounding cold cluster gas. Our experimental results show that the range of the ions in the cluster gas is almost independent of their energies and it is much shorter than the range calculated using popular SRM code for instance [15]. The physics included in SRM or similar models, is the slowing down of keV ions due to the interaction with electrons in the uniform gas. In our case, the gas has not an uniform density distribution but it is made of drops of different sizes, well explained by a log-normal distribution [16,17], formed during the free expansion into vacuum after the opening of the pulsed value. Drops can present already inside the value before the expansion, if the gas is prepared for instance near the critical point of the liquid–gas (LG) phase transition. It is of great interest to study what happens in those cases after the gas expands. Near the second order LG phase transition, the mass distribution of the clusters follows the Fisher’s law and in particular it is a power law at the critical point [18]. The free expansion might change such distribution. Theoretical calculations of a classical interacting gas, which freely expands after has been prepared near the critical point [19], do not display much variation from the predicted Fisher’s cluster distribution. Thus, the cluster size distribution obtained from the
Neutron enhancement from laser interaction with a critical fluid

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ABSTRACT

We discuss experimentally and theoretically neutron production from the laser driven explosion of gas clusters prepared near the liquid-gas critical point. We let deuterated methane that was prepared very close to its critical temperature and pressure expand through a conical nozzle to create clusters, and then irradiated those clusters with a high intensity pulse from the Texas Petawatt Laser. After ionization, the clusters explode producing energetic ions, some of which fuse with resultant neutron emission. We show that the critical fluctuations present in the nozzle before the expansion influence the dynamics of neutron production. Neutron production near the critical point follows a power law, which is a signature of a second order phase transition and it is consistent with the Fisher model. This result might be relevant for energy production from fusion reactions.

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Cryogenic d+t Target or solid $^6\text{Li}(t)$

Measure plasma ions kinetic energy distribution: determine non-equilibrium and thermal components

$d+t \rightarrow a+n$

$d+d \rightarrow ^3\text{He}+n(t+p)$

$t+\alpha \rightarrow ^7\text{Li}+\gamma$

From these informations we can derive the astrophysical S-factor for the main channels and some of the higher order reactions as function of temperature, density and collective energy of the plasma.

In the 18 months we plan to perform 8 shots @ NIF with various targets-laser beams combinations and preliminary data analysis completed.
operating since 2011
8 beams output 40 kJ/3 ns/1 ω, 24 kJ/3 ns/3 ω
PW laser (1.5kJ, 2ps, 2011)

for SINAP\SIOM\TAMU\INFN\IMUN Collaboration
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- Zhangguoqiang 2017-04-13 @SINAP
Cross sections\S-factor\reaction rate in plasmas

The effect of ternary fusion reactions

Highly compressed and not so hot plasma

\[ N_{f1} = N_{i} \rho <\sigma \nu > / 2 = N_{i} \rho <\sigma 1 > r 1 / 2 \]  
\[ N_{f2} = N_{f1} \rho <\sigma 2 > r 2 \]

\[ d + d \rightarrow ^{3}\text{He} (0.82 \text{ MeV}) + n (2.45 \text{ MeV}) \]
\[ d + d \rightarrow t (1.01 \text{MeV}) + p (3.02 \text{ MeV}) \]

\[ d + t \rightarrow ^{4}\text{He} (3.5\text{MeV}) + n (14.1 \text{ MeV}) \]
\[ d + ^{3}\text{He} \rightarrow ^{4}\text{He} (3.6\text{MeV}) + p (14.7 \text{ MeV}) \]

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\[ ^{3}\text{He} (0.82 \text{ MeV}) n (2.45 \text{ MeV}) t (1.01\text{MeV}) p (3.02 \text{ MeV}) \]

\[ ^{4}\text{He} (3.5\text{MeV}) n (14.1 \text{ MeV}) ^{4}\text{He} (3.6\text{MeV}) p (14.7 \text{ MeV}) \]
Figure 2: (color online) a) Deuterium energy distribution from the top (full circles) and bottom (open circles) FC, and TP (full triangles). b) TP spectra relative to the laser irradiation of CD₂ target. c) TOF results for 14.1 MeV and 2.45 MeV neutrons from plastic scintillator detector. The two arrows indicate the 14.1 MeV and 2.45 MeV neutron energies.
Figure 3: (color online) Fusion yield as function of laser energy. Different experimental results Ditmire-2004\textsuperscript{35}, UT-2011\textsuperscript{20}, UT-2016\textsuperscript{19}, Fu-2015 SGII\textsuperscript{40}, Dittrich-1994\textsuperscript{44}, NIF-2014\textsuperscript{43}, Osaka -2001\textsuperscript{41}, Osaka-2004\textsuperscript{42}, OMEGA-shot5241\textsuperscript{36} and SGIIIpro2017\textsuperscript{37} are indicated in the inset.
Figure 5: (color online) The average cross section as function of temperature with Maxwell-Boltzmann distribution, expressed by eq. (6). The red points are the experimental cross section data from eq.(5).
Figure 4: (color online) $\Lambda \rho \sigma / \ln 2$ obtained from eq.(4) vs T from eq.(1). Omega and NIF data are derived from the experiments$^{25}$, using the Down Scatter Ratio$^{21,23}$. Our results using the DSR method (N4/N3) are given by the open triangle symbols in good agreement with the N3/N2 ratios.
Nuclear probes of an out-of-equilibrium plasma at the highest compression

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ABSTRACT

We report the highest compression reached in laboratory plasmas using eight laser beams, \( F_{\text{max}} \approx 12 \) kJ, \( \tau_{\text{FWHM}} = 2 \) ns in third harmonic on a \( \text{CD}_2 \) target at the ShenGuang-II Upgrade (SGII-Up) facility in Shanghai, China. We estimate the deuterium density \( n_D = 2.0 \pm 0.9 \) kg/cm\(^3\), and the average kinetic energy of the plasma ions less than 1 keV. The highest reached areal density \( \Lambda n = 4.8 \pm 1.5 \) g/cm\(^2\) was obtained from the measured ratio of the sequential ternary fusion reactions (\( \text{dd} \rightarrow t+p \) and \( t+d \rightarrow \alpha+n \)) and the two body reaction fusions (\( \text{dd} \rightarrow ^3\text{He} + n \)). At such high densities, sequential ternary and also quaternary nuclear reactions become important as well (i.e. \( \text{n} + ^{14}\text{C} \rightarrow n' + ^{12}\text{C} \) etc.) resulting in a shift of the neutron (and proton) kinetic energies from their birth values. The Down Scatet Ratio (DSR-quaternary nuclear reactions) method, i.e. the ratio of the 10-12 MeV neutrons divided by the total number of 14.1 MeV neutrons produced, confirms the high densities reported above. The estimated lifetime of the highly compressed plasma is 52 ± 9 ps, much smaller than the lasers pulse duration.

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COMPRESS
Nuclear Astrophysics with Lasers

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A systematic program of investigations of nuclear reactions in laser-produced plasmas is addressed. Such reactions provide an important diagnostic tool for probing the dynamics and thermodynamics in the plasma and understanding laser ion acceleration and neutron production mechanisms. The goal will be to reach the level of knowledge that allows the measurement of fundamental nuclear cross sections at low and high particle densities. The quantitative measurement of fusion probabilities in hot and dense plasmas will contribute significantly to our comprehension of stellar composition and evolution and will provide important information for development of fusion energy production and applications such as medical isotope production and compact neutron source development. All of these are some of the main goals of the European Extreme Light Infrastructure (ELI), the Shanghai Superintense Ultrafast Laser Facility (SULF), the Station of Extreme Light (SEL) in China and similar projects in other countries.
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