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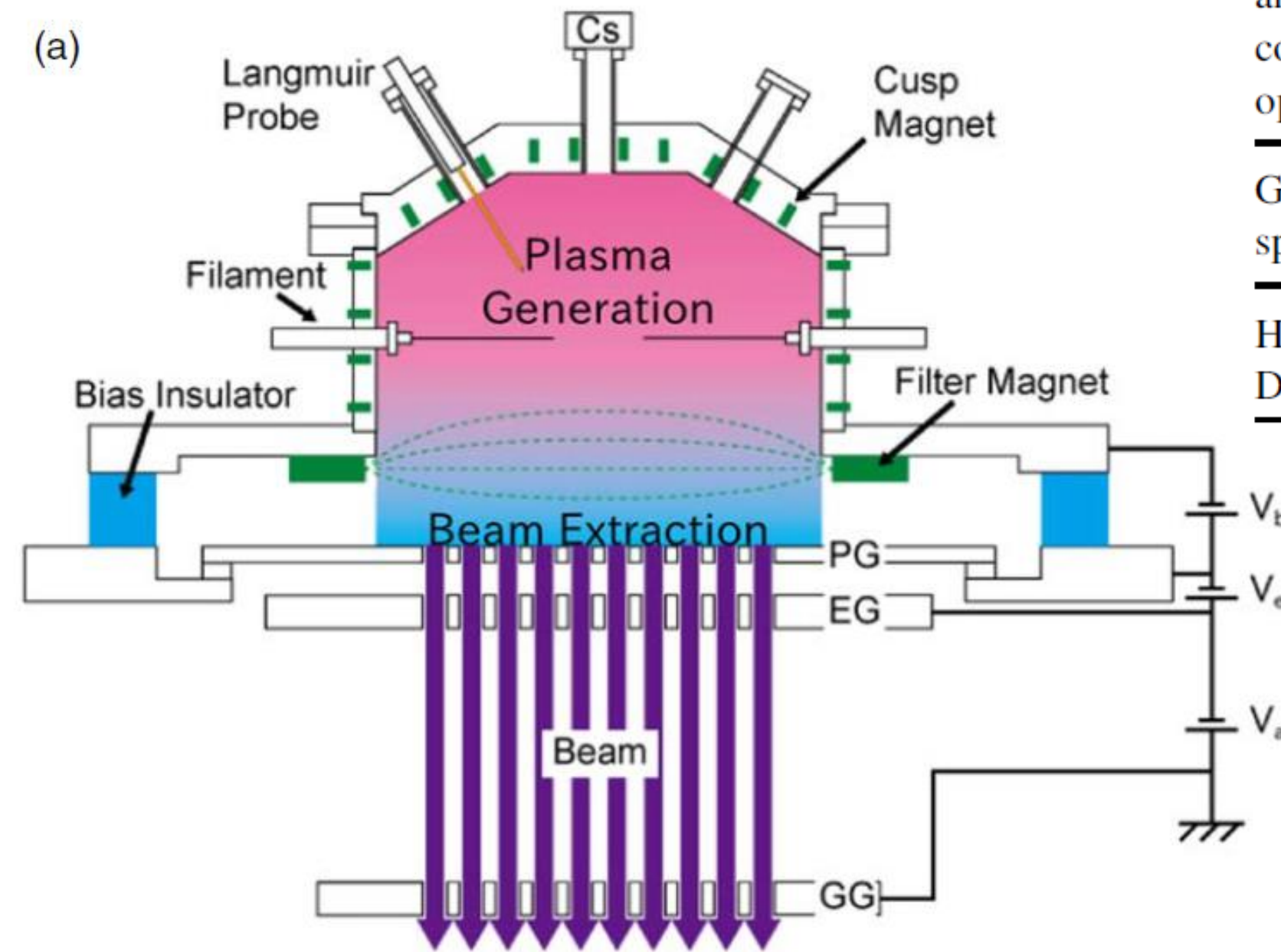
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Abstract

Sequences of hydrogen (H) and deuterium (D) experiments have been done by NIFS research and development negative ion source (RNIS) for the deuterium NBI development. In the experiments, the co-extracted electron current with the negative ions and the electron density in the plasma generation region in the D experiment have been around three times higher than that in the H experiment. To explain the difference of the electron density in the RNIS driver region, a zero-dimensional numerical model is developed in the present study. The model only focuses on the isotope effect for vibrationally excited level of electronically grounded state molecules and its relevant cross-sections. The calculation results show that difference of the ionization channel numbers via molecular vibrationally excited states could be a reason to enhance ionization rate in D plasma.

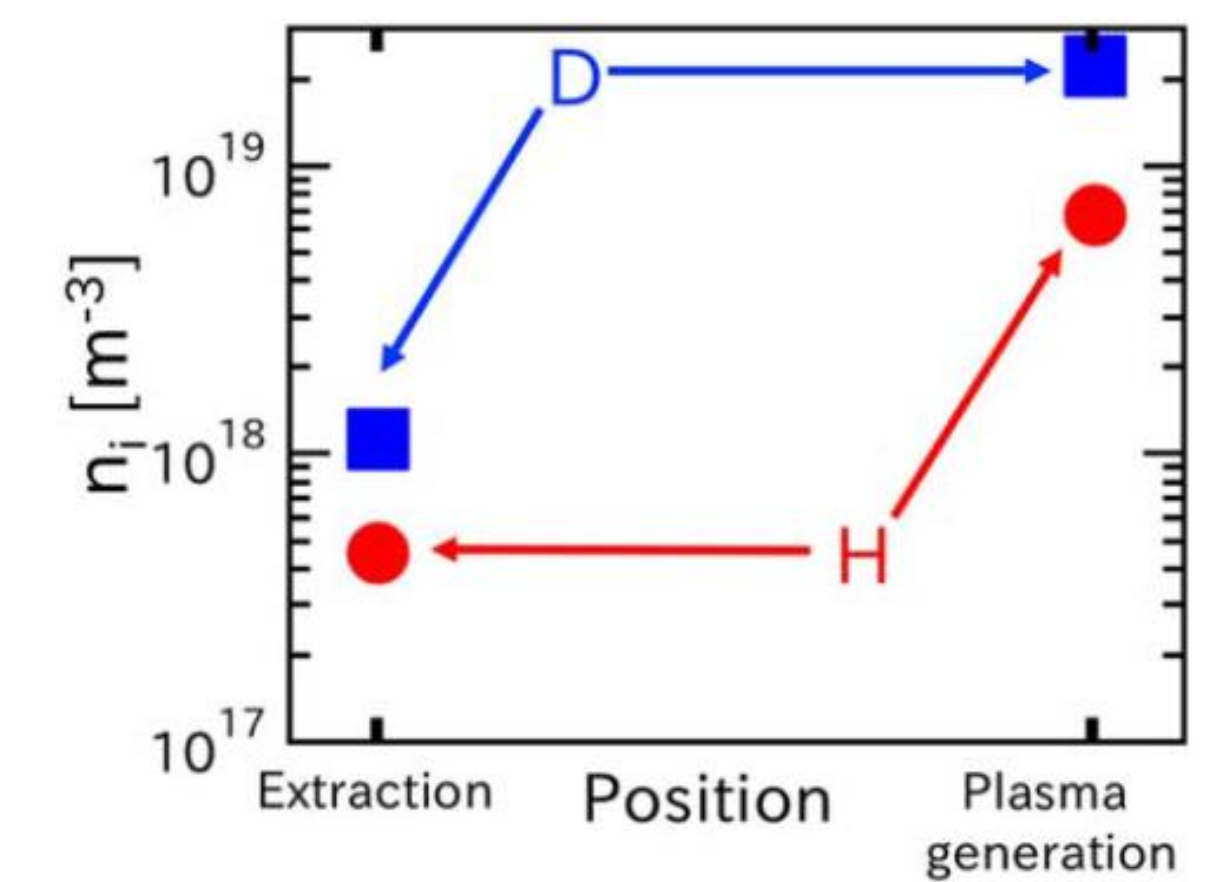
NIFS LHD-NBI and RNIS experiments



Operation parameters of RNIS	
Arc current	600 A ($P_{arc} = 50$ kW)
Gas flow rate	850 sccm (for H2 exp.)

Table 1. Comparison of extracted current characteristics between hydrogen and deuterium operations. The I_e , $I(H^-)$ and $I(D^-)$ are currents of co-extracted electron, H^- beam, and D^- ion beam, respectively. Other operation parameters except discharge power were the same in both operations.

Gas species	Discharge power [kW]	I_e [A]	$I(H^-), I(D^-)$ [A]	$I_e/I(H^-), I_e/I(D^-)$
H ₂	165	9.3	24.6	0.38
D ₂	149	22.6	16.5	1.37



[1] H. Nakano, et al., J. J. Appl. Phys. **59**, SHHC09 (2020).

Isotope Effects for vibrationally excited levels of H₂ and D₂

- **Mass Scaling of vibrationally excited molecular levels :**

$$v' = \frac{\omega_{D_2}}{\omega_{H_2}} \times v_i = \sqrt{\frac{\mu_{H_2}}{\mu_{D_2}}} \times v_i$$

- **Franck-Condon Factor:** $F_{v_i, v_f}^{N^1, X^1} = \left| \int_0^\infty \phi_{v_i}^{N^1, X^1}(r) \phi_{v_f}^{X^1}(r) dr \right|^2$
- **Einstein Coefficient (for spontaneous radiation and EV process):**

$$A_{v_i, v_f}^{N^1, X^1} = \frac{16 \pi^3}{3 \epsilon_0 h \lambda^3} \left| \int_0^\infty \phi_{v_i}^{N^1}(r) D_{el}(r) \phi_{v_f}^{X^1}(r) dr \right|^2$$

#	Reactions
1	$H_2(X^1\Sigma_g^+; v_i) + e^- \rightarrow H_2(N^{1,3}\Lambda_\sigma) + e^- \rightarrow 2H + e^-$
2	$H_2(X^1\Sigma_g^+; v_i) + e^- \rightarrow H_2(B^1\Sigma_u^+ \text{ or } C^1\Pi_u) + e^- \rightarrow H_2(X^1\Sigma_g^+; v_f) + e^- + hv$
3	$H_2(X^1\Sigma_g^+; v_i) + e^- \rightarrow H_2(X^1\Sigma_g^+; v_f = v_i \pm 1) + e^-$
4	$H_2(X^1\Sigma_g^+; v_i) + e^- \rightarrow H_2^+(X^2\Sigma_g^+) + 2e^-$
5	$H_2(X^1\Sigma_g^+; v_i) + e^- \rightarrow H_2^+(X^2\Sigma_g^+) + 2e^- \rightarrow H + H^+ + 2e^-$
6	$H_2(X^1\Sigma_g^+; v_i) + e^- \rightarrow H_2^+(X^2\Sigma_g^+) \rightarrow H^- + H$
7	$H_2^+(X^2\Sigma_g^+) + e^- \rightarrow H_2^+(2p\sigma_u; ^2\Sigma_u^+) + e^- \rightarrow H + H^+ + e^-$
8	$H_2^+(X^2\Sigma_g^+) + e^- \rightarrow H_2^{Ryd}(N^{1,3}\Lambda_\sigma) \rightarrow 2H$
9	$H_2^+(X^2\Sigma_g^+) + e^- \rightarrow 2H^+ + 2e^-$
10	$H_3^+ + e^- \rightarrow H_3^+(N^{1,3}\Lambda_\sigma) + e^- \rightarrow H^+ + H_2(X^1\Sigma_g^+; v_f) + e^-$
11	$H_3^+ + e^- \rightarrow H_3^{Ryd}(N^{1,3}\Lambda_\sigma) \rightarrow 3H \text{ or } H_2(X^1\Sigma_g^+; v_f) + H$
12	$H_2(X^1\Sigma_g^+; v_i) + H_2^+ \rightarrow H_3^+ + H$
13	$H + e^- \rightarrow H^+ + 2e^-$
14	$H^+ + e^- \rightarrow H + hv$

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[4] U. Fantz, et al., Atom. Data Nucl. Data Table **92**, 853-973 (2006).

[5] I. I. Fabrikant, Physica Scripta **T96**, 45-51 (2002).

[6] L.C. Johnson, Astrophys. J., **174**, 227 (1971).

- **The isotope effects** taken into account in this model is **only "difference of vibrationally excited levels of H₂ and D₂ molecules"** and the relevant reaction cross-section (with the Franck-Condon factor and the A coefficient for spontaneous radiation).
- **A zero-dimensional (0D) model** solves rate equations for **molecules** (each vibrationally excited levels in electronically grounded state), **atoms**, **positive ions (H⁺, H₂⁺, H₃⁺ or D⁺, D₂⁺, D₃⁺)** and **electrons**.
- The reactions (in the table) for the H plasma and the D plasma are applied to the 0D model individually to obtain the steady-state electron density.
- **The reaction rates** for electron impact processes are calculated from EEDF which is assumed to be Maxwellian distribution with $T_e = 5.0$ eV.
- **Characteristic transport loss time** for each particle species are fixed and given in the rate equations in the form of " $-n/\tau$ ".
- **The H atoms and H⁺ ions** (or D atoms and D⁺ ions) **incoming to the wall are reflected** as atoms or molecules. We assume that the wall is fully covered by H atoms. Most part of the reflection is **"kicking out of wall absorbed atoms"** with probability over 90%. The other part is **"surface recombination"** given in the probability $(1 - R_{N+})$ and $(1 - R_{N0})$, which is below 10%.

A zero-dimensional model for H₂ and D₂ plasma

RATE EQUATIONS

Electron:

$$\frac{dn_e}{dt} = S_{IOZ-H} + S_{IOZ-H_2} - S_{RCMB} + S_{DI} + S_{filament} - n_e/\tau_e$$

Positive ions:

$$\frac{dn_{H^+}}{dt} = S_{DI} - S_{RCMB} + S_{IOZ-H} + S_{DE-H_2+} - n_{H^+}/\tau_+$$

$$\frac{dn_{H_2^+}}{dt} = S_{IOZ-H_2} - S_{DR} - S_{DE} - S_{DI-H_2+} + S_{DE-H_3+} - S_{H_3+PROD} - n_{H_2^+}/\tau_+$$

$$\frac{dn_{H_3^+}}{dt} = -S_{DR} - S_{DE-H_3+} + S_{H_3+PROD} - n_{H_3^+}/\tau_+$$

Atom:

$$\frac{dn_H}{dt} = \sum_v S_{DISS-H_2}(v) + S_{RCMB-H+} - S_{IOZ-H} + S_{DISS} + S_{DI} + S_{DE} + S_{DR} - \frac{n_H}{\tau_H} + R_{N+} \times \frac{n_{H^+}}{\tau_+} + R_{N0} \times \frac{n_{H0}}{\tau_{H0}}$$

Mol.:

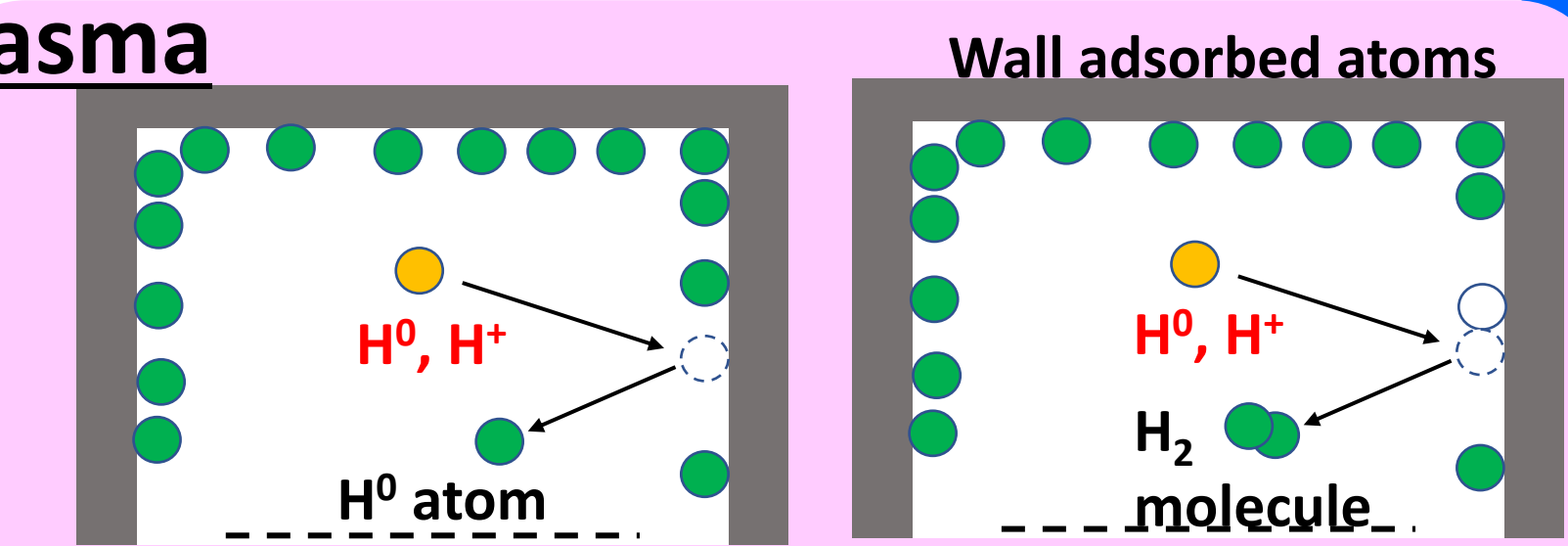
$$\frac{dn_{H_2}(v)}{dt} = -S_{DISS}(v) - S_{IOZ-H_2}(v) - \sum_{v'} S_{EVBC}(v, v') + \sum_{v'} S_{EVBC}(v', v) + S_{eV}(v \pm 1) - S_{eV}(v) + S_{gas}(v=0) - n_{H_2}(v)/\tau_{H_2} + 0.5 \times (1 - R_{N+}) \times \frac{n_{H^+}}{\tau_+} + 0.5 \times (1 - R_{N0}) \times \frac{n_{H0}}{\tau_{H0}}$$

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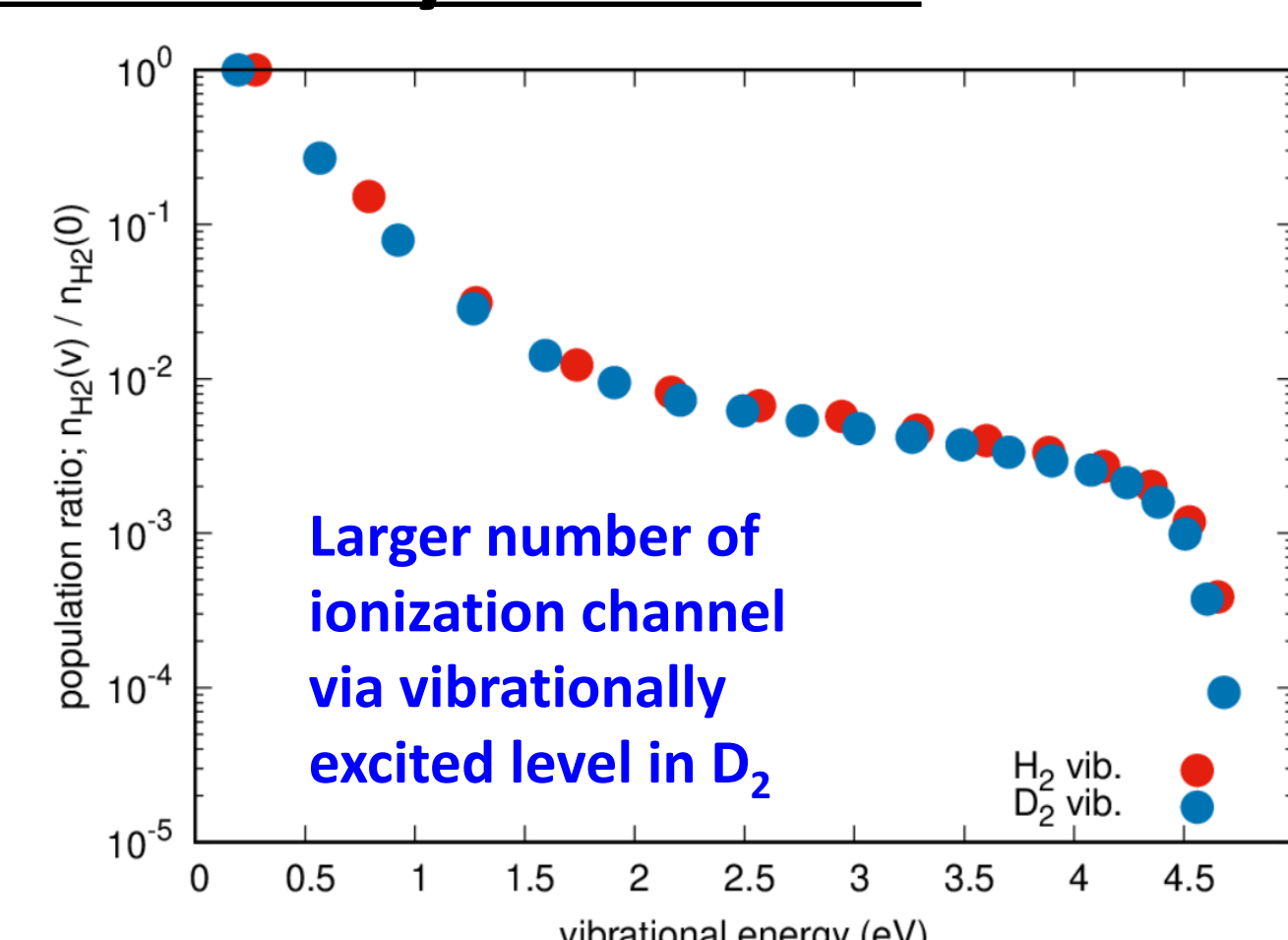
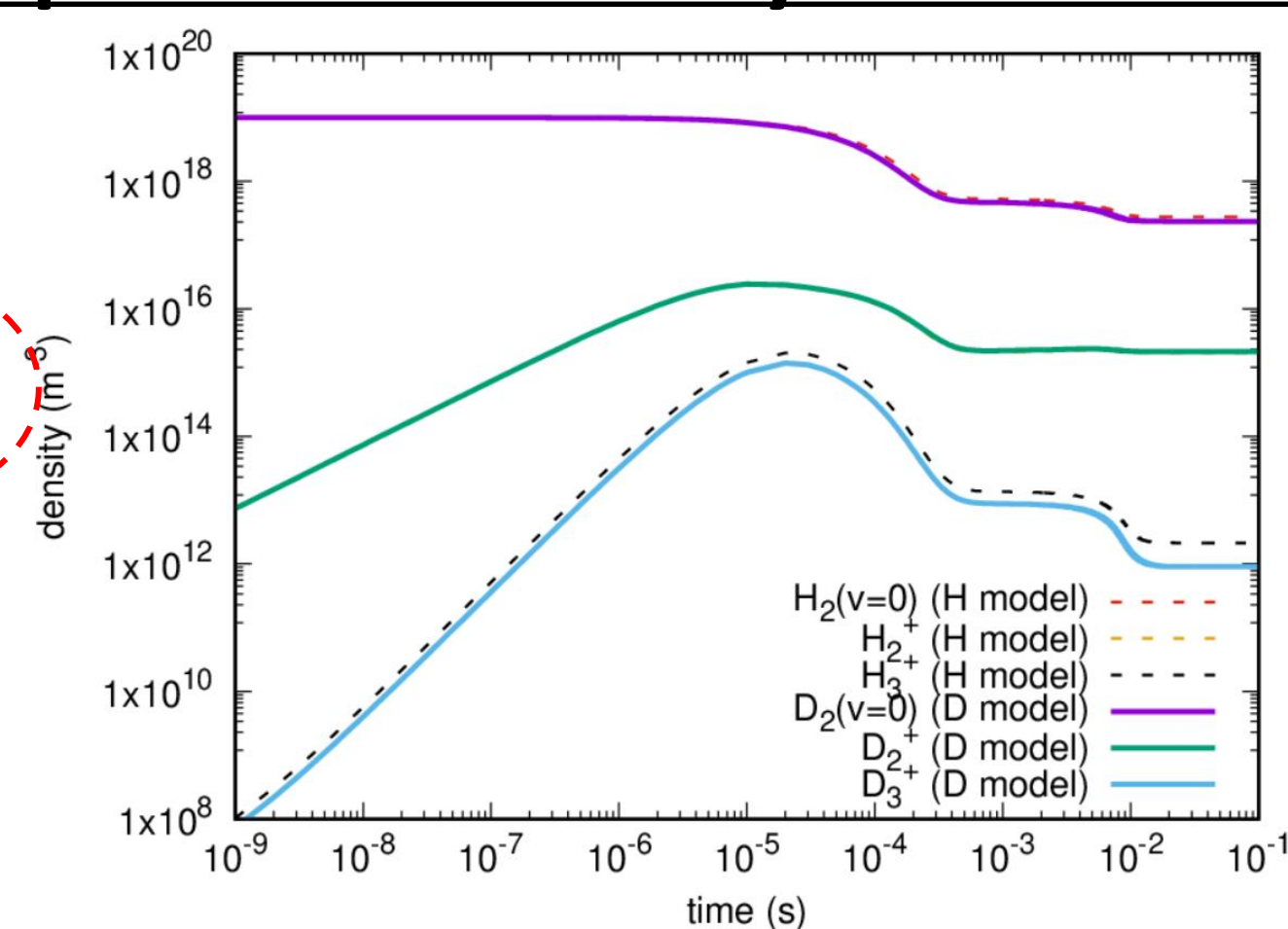
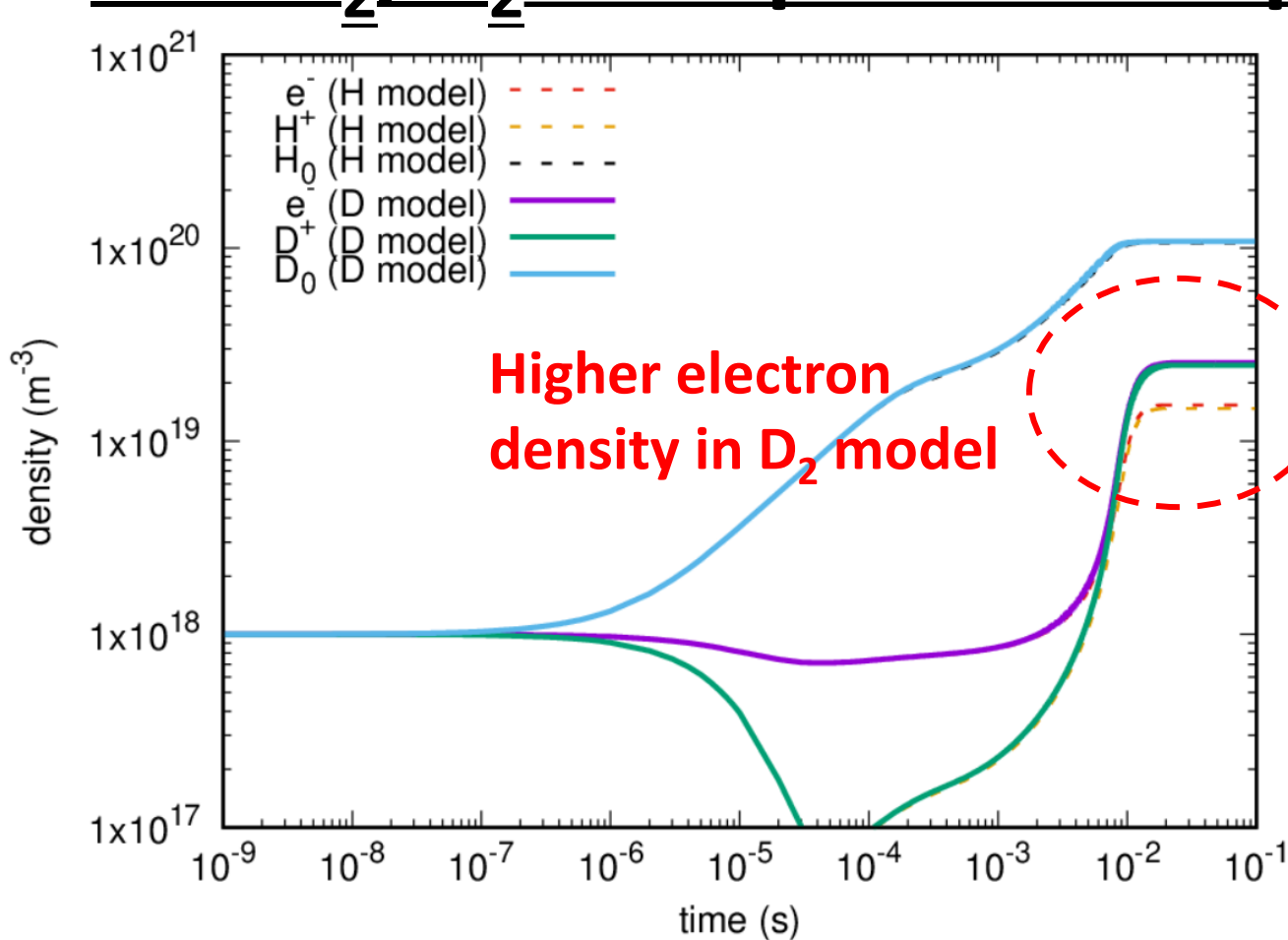
① Kicking-out of wall-adsorbed atom ② Surface Recombination

① Kicking-out of wall-adsorbed atom

② Surface Recombination

Temperature (fixed)	
Electron	5.0 eV
Positive Ions	1.0 eV
Neutrals	0.1 eV (~1000 K)
Characteristic time for transport loss	
Electron & positive ions	$\tau_e = \tau_+ = 1 \times 10^{-5}$ s
H ₂ /D ₂ molecules	$\tau_{H_2} = 5 \times 10^{-3}$ s
H/D atoms	$\tau_H = 1 \times 10^{-3}$ s
Wall Reflection Probability as Atoms	
For Incoming H ⁺ /D ⁺	$R_{N+} = 90.95\%$
For Incoming H/D	$R_{N0} = 95.55\%$

The H₂/D₂ Isotope Effects appear in the steady-state electron density difference



Steady-state density (m ⁻³)	H ₂ model	D ₂ model
e ⁻	1.54E+19	2.54E+19
H ⁺	1.48E+19	2.48E+19
H ⁰	1.06E+20	1.08E+20
H ₂ (v=0)	2.76E+17	2.36E+17
H ₂ ⁺	2.06E+15	2.13E+15
H ₃ ⁺	2.14E+12	9.09E+11

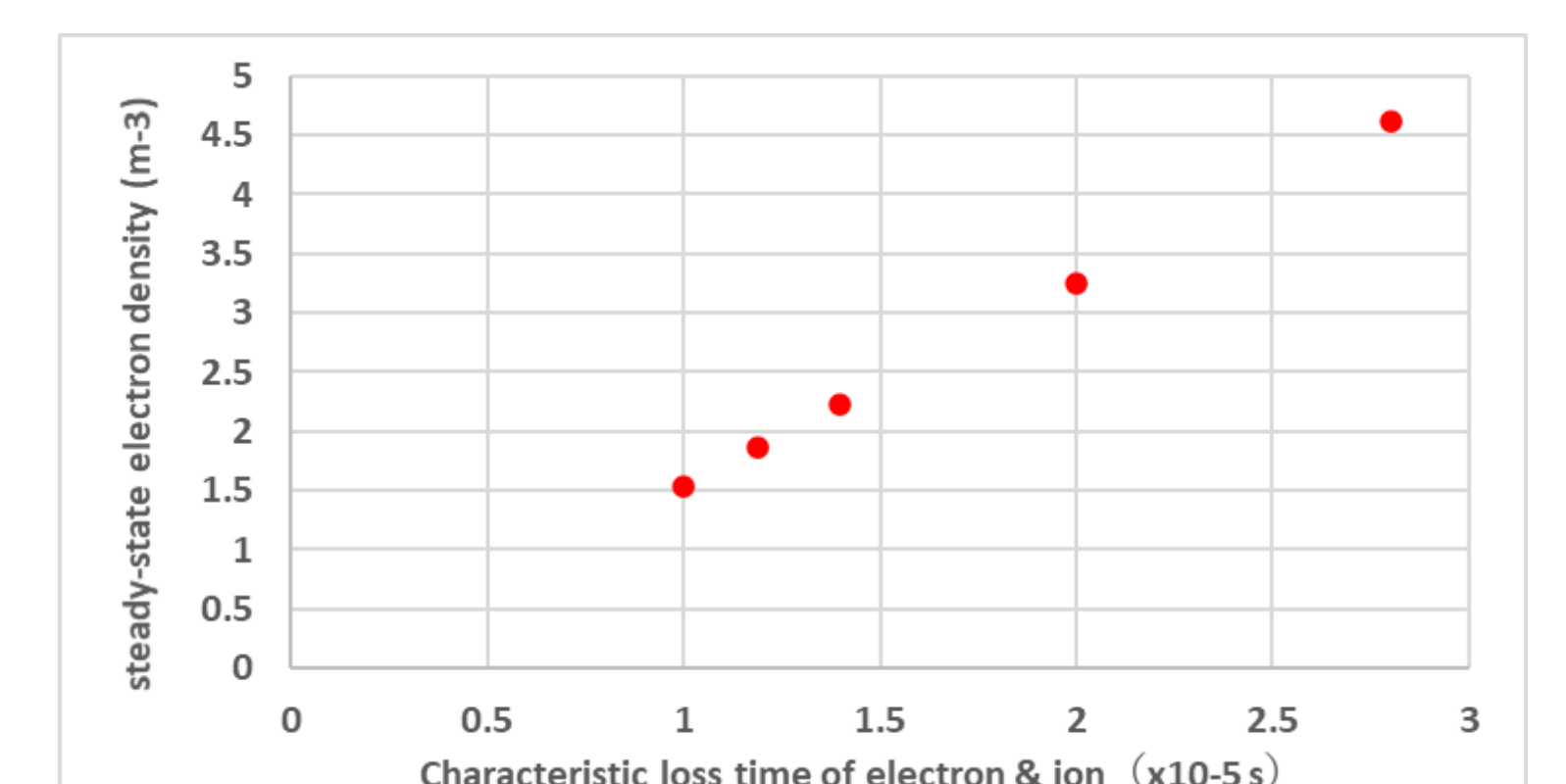
Main Reaction Rate Coefficients (m ³ s ⁻¹) (Rate)/ne/nH ₂ (v=0)	H ₂ model	D ₂ model
Non-Dissociative EXC X ¹ - B ¹	1.56E-15	2.02E-15
Dissociative EXC X ¹ - B ¹	3.53E-17	4.99E-17
Dissociative EXC X ¹ - C ¹	4.00E-17	5.67E-17
Dissociative EXC X ¹ - b ³	1.51E-14	1.77E-14
Non-Dissociative Ionization from X ¹	9.88E-16	1.17E-15
Dissociative Ionization from X ¹	3.32E-18	4.19E-18

- **From the calculation, the higher electron (and D⁺) density are obtained in the D₂ model.**
- The rate coefficient of ionization process (including all vibrationally excited levels) is higher in the D₂ model, however, dependence of the vibrationally excited levels to the vibrational energy matches in the H₂ and D₂ models.
- **The difference of the ionization channel number via vibrationally excited levels is the reason of the different ionization rate** coefficient in the H₂ and D₂ models. This isotope effect appear in the electron density difference.
- However, the ratio of the electron density in H₂ and D₂ models are around 1.65 while the ratio is around 3 in the experiments.
- We conclude that this difference due to the vibrationally excited levels is **one of the isotope effects**.

Other isotope effects

(will be reported in the near future)

1. Characteristic transport loss time for electrons and positive ions



2. Electron Temperature

