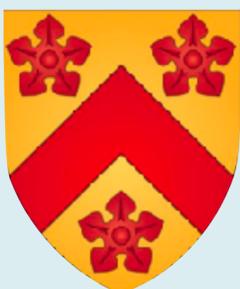


HIGGS PHYSICS IN THE PRECISION ERA

Higgs 2021, Stony Brook, via Zoom, 18 October 2021

Gavin Salam
Rudolf Peierls Centre for Theoretical Physics &
All Souls College, University of Oxford



What are we trying to achieve?

Higgs is the last particle of the SM.

So the SM is complete, right?

The Lagrangian and interactions: two out of three qualitatively new!

$$\mathcal{L}_{\text{SM}} = \dots + |D_{\mu}\phi|^2 + \psi_i y_{ij} \psi_j \phi - V(\phi)$$

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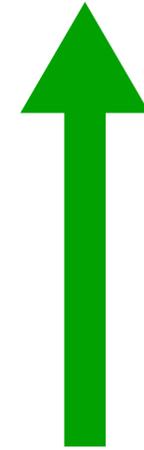
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like those in QED, QCD, EW,
studied for many decades
(but now with a scalar)

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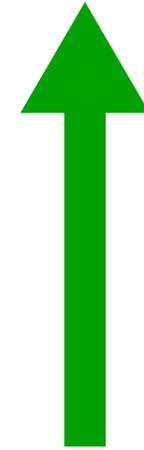
Yukawa interactions. Responsible for fermion masses, and induces “fifth force” between fermions. **Direct study started only in 2018!**

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Higgs potential (\rightarrow self-interaction). Holds the SM together. **Unobserved**

Why do Yukawa couplings matter to everyone?

Because, within SM **conjecture**, they set quark and electron masses

$\psi_i \gamma_j \psi_j \phi$

Up quarks (mass ~ 2.2 MeV) are lighter than
down quarks (mass ~ 4.7 MeV)

proton (up+**up**+down): $2.2 + 2.2 + 4.7 + \dots = 938.3$ MeV
neutron (up+**down**+down): $2.2 + 4.7 + 4.7 + \dots = 939.6$ MeV

So protons are **lighter**
than neutrons,
 \rightarrow protons are stable,
giving us hydrogen

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Bohr radius

$$a_0 = \frac{4\pi\epsilon_0\hbar^2}{m_e e^2} = \frac{\hbar}{m_e c \alpha} \propto \frac{1}{y_e}$$

electron Yukawa
determines size of all
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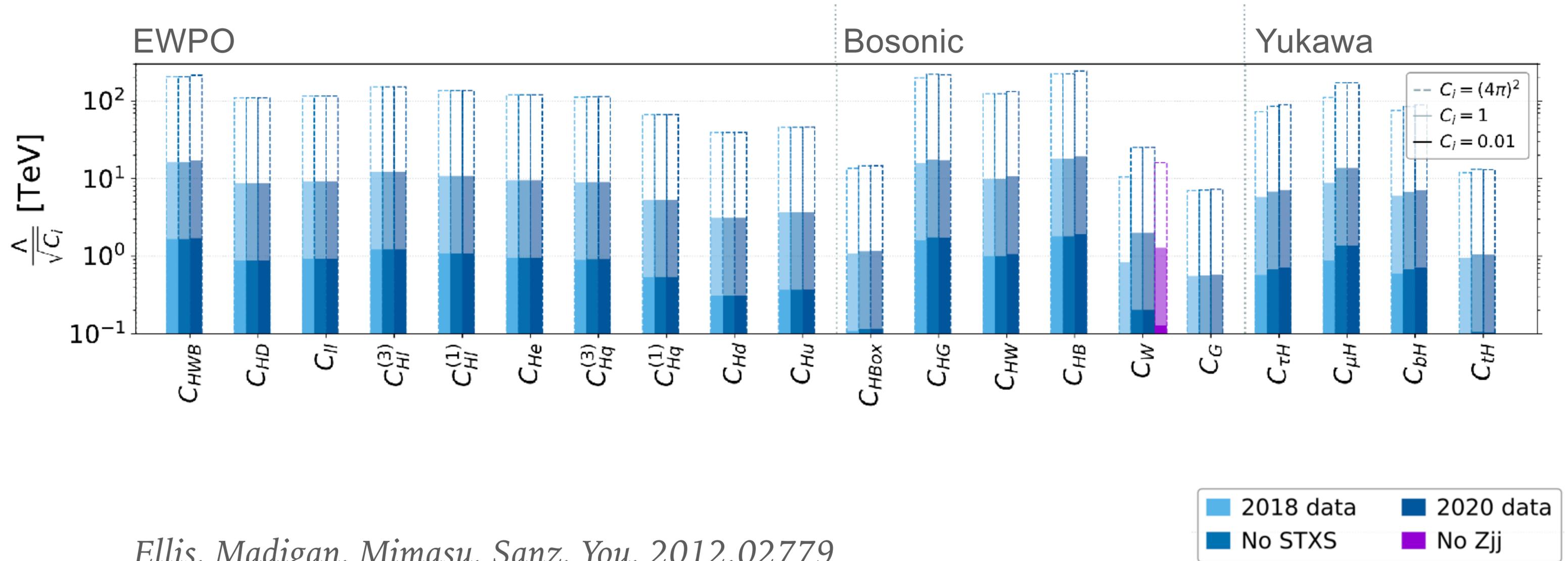
**We are establishing the existence of crucial new interactions
We wouldn't consider QED established if we'd only tested it to 0(10%)**

Bohr

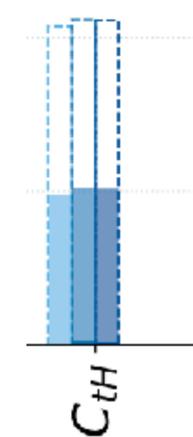
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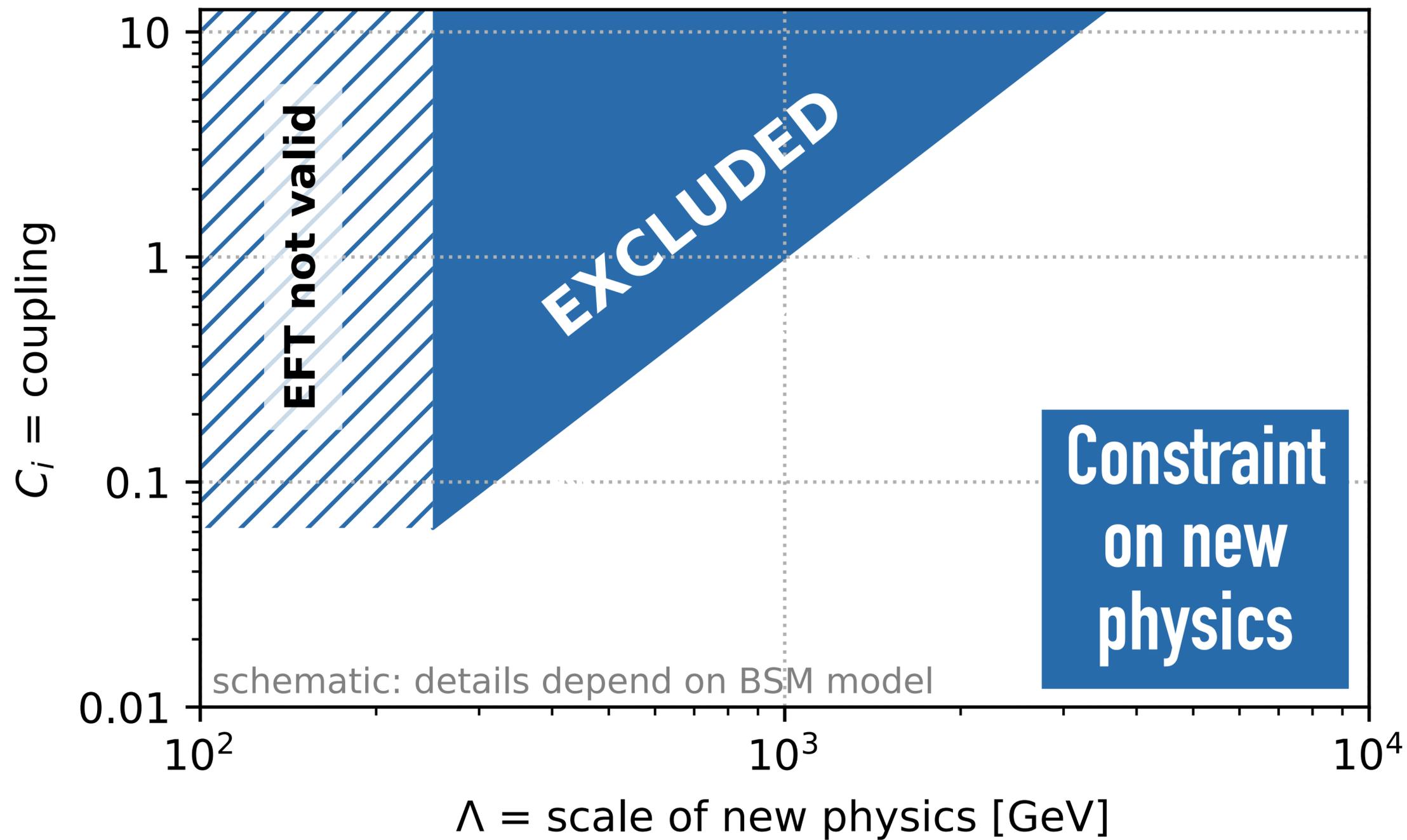
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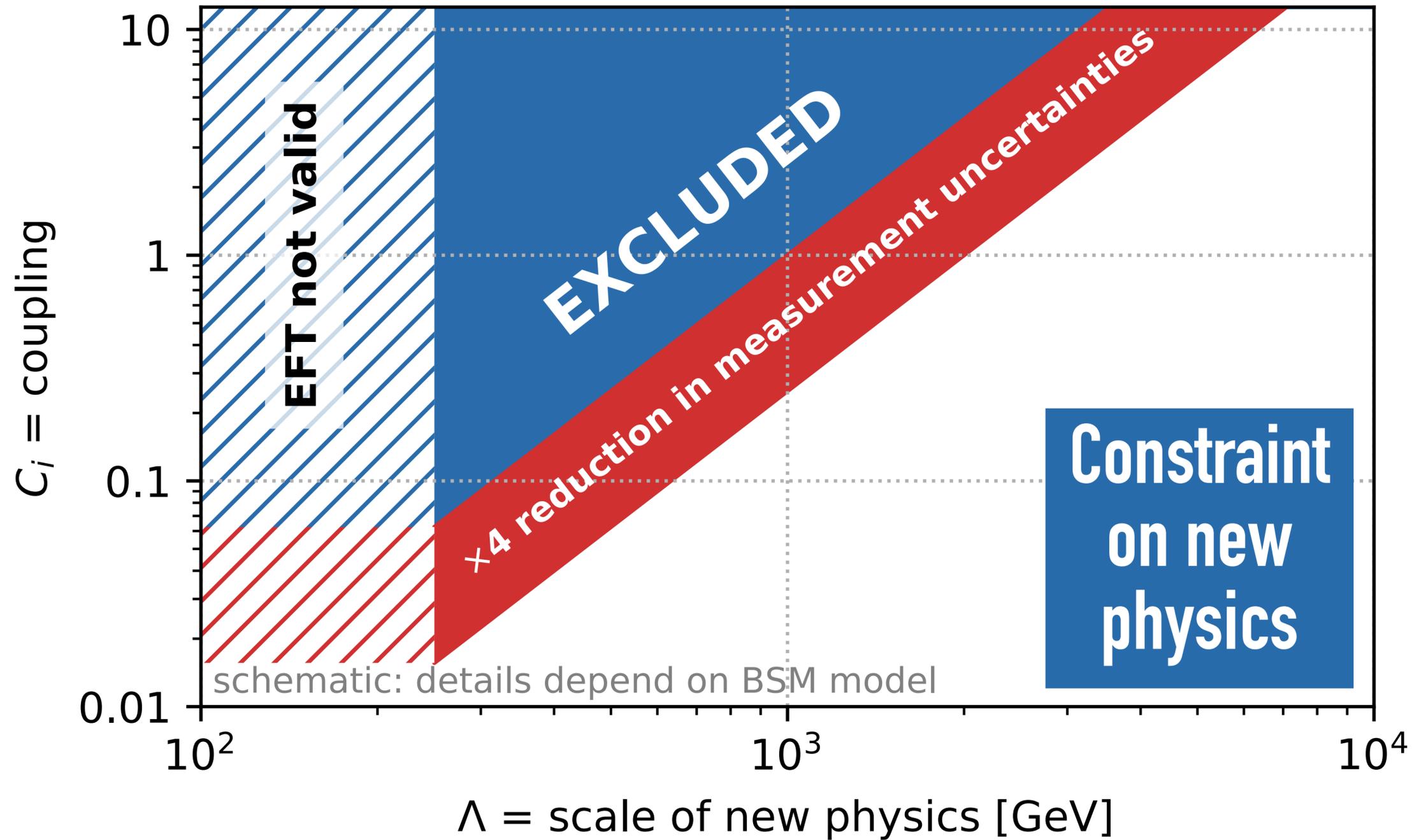
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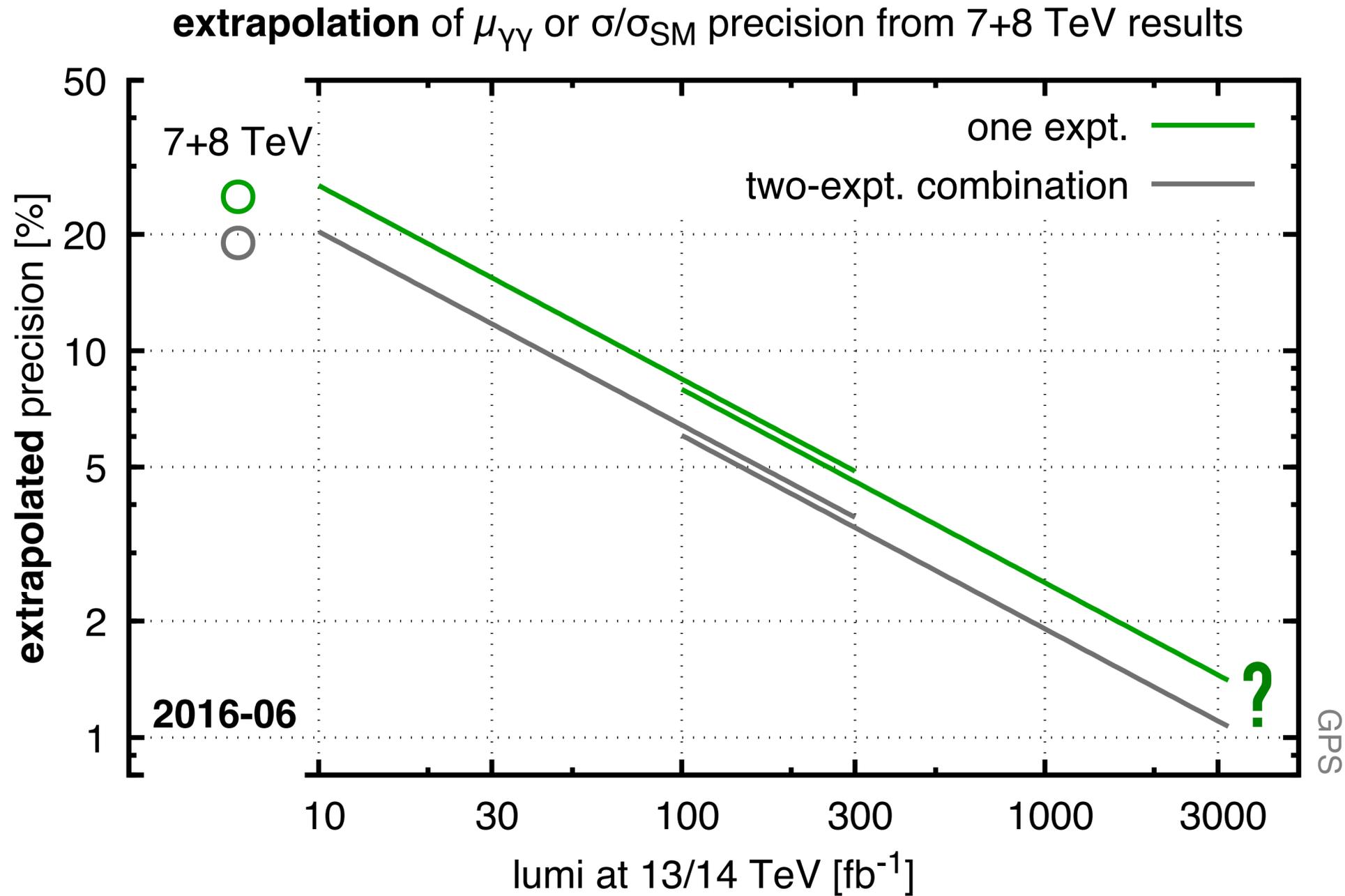
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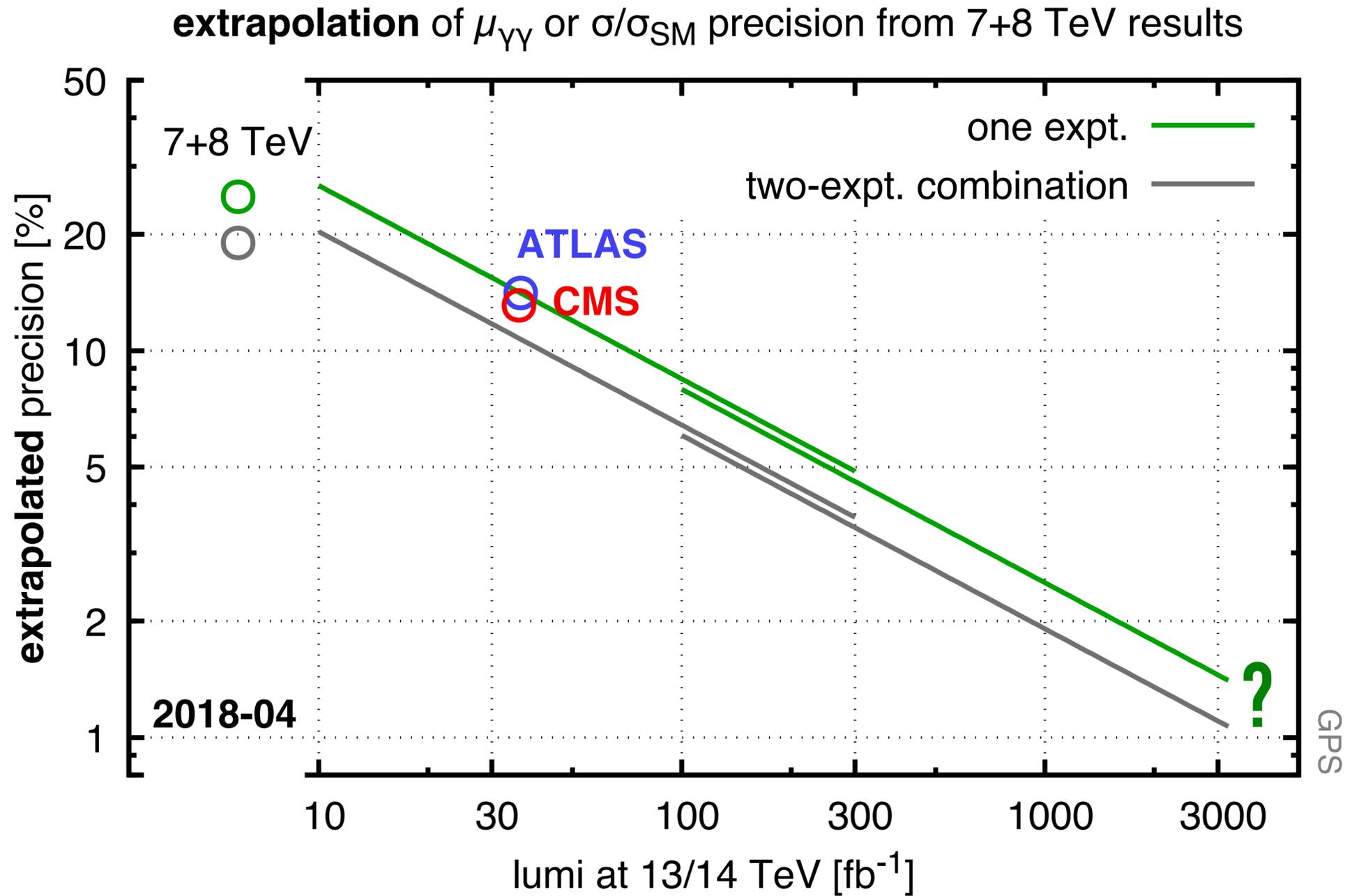
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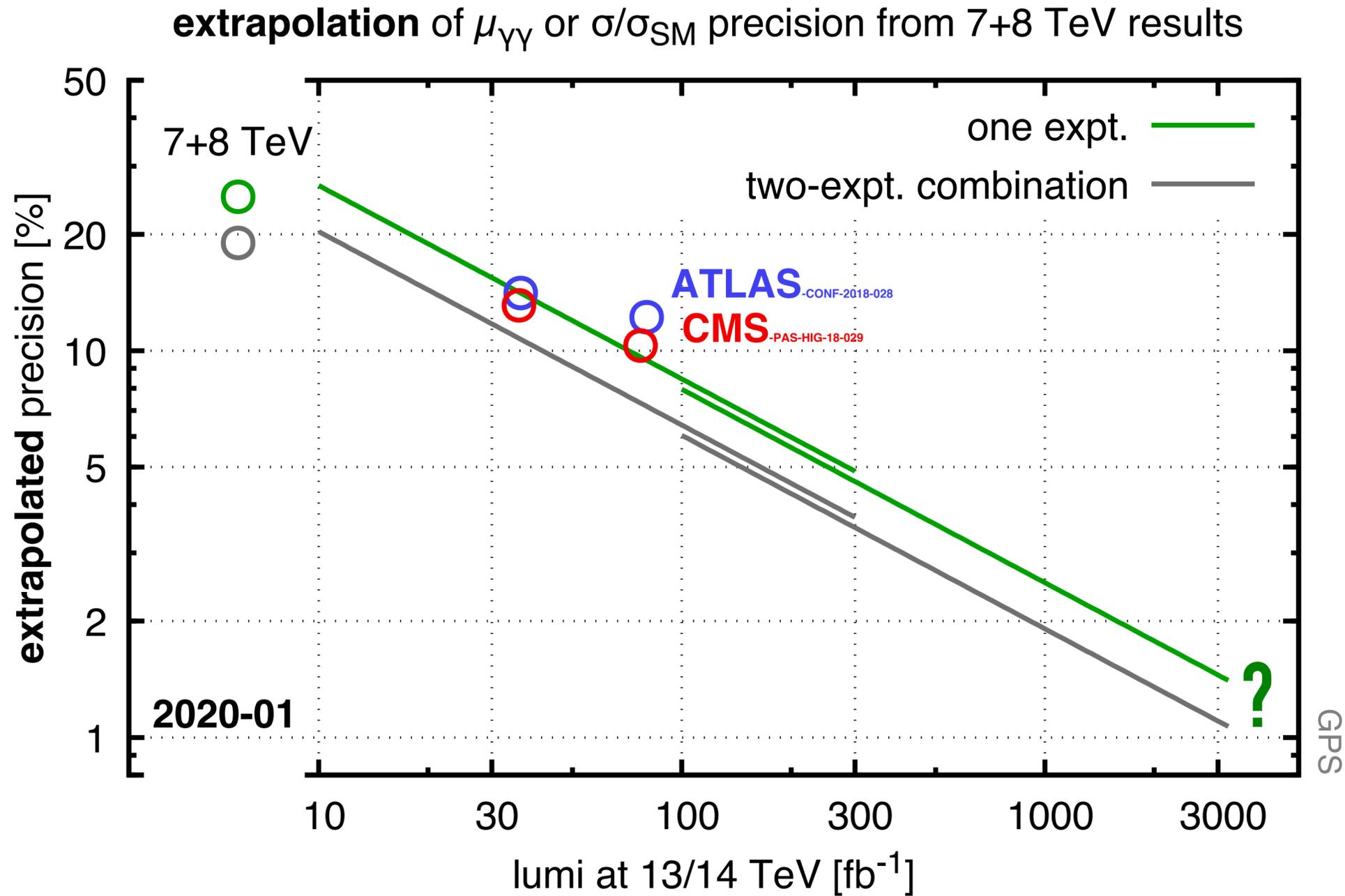
$H \rightarrow \gamma\gamma$, an indirect probe of the top Yukawa, HWW and contact ggH couplings



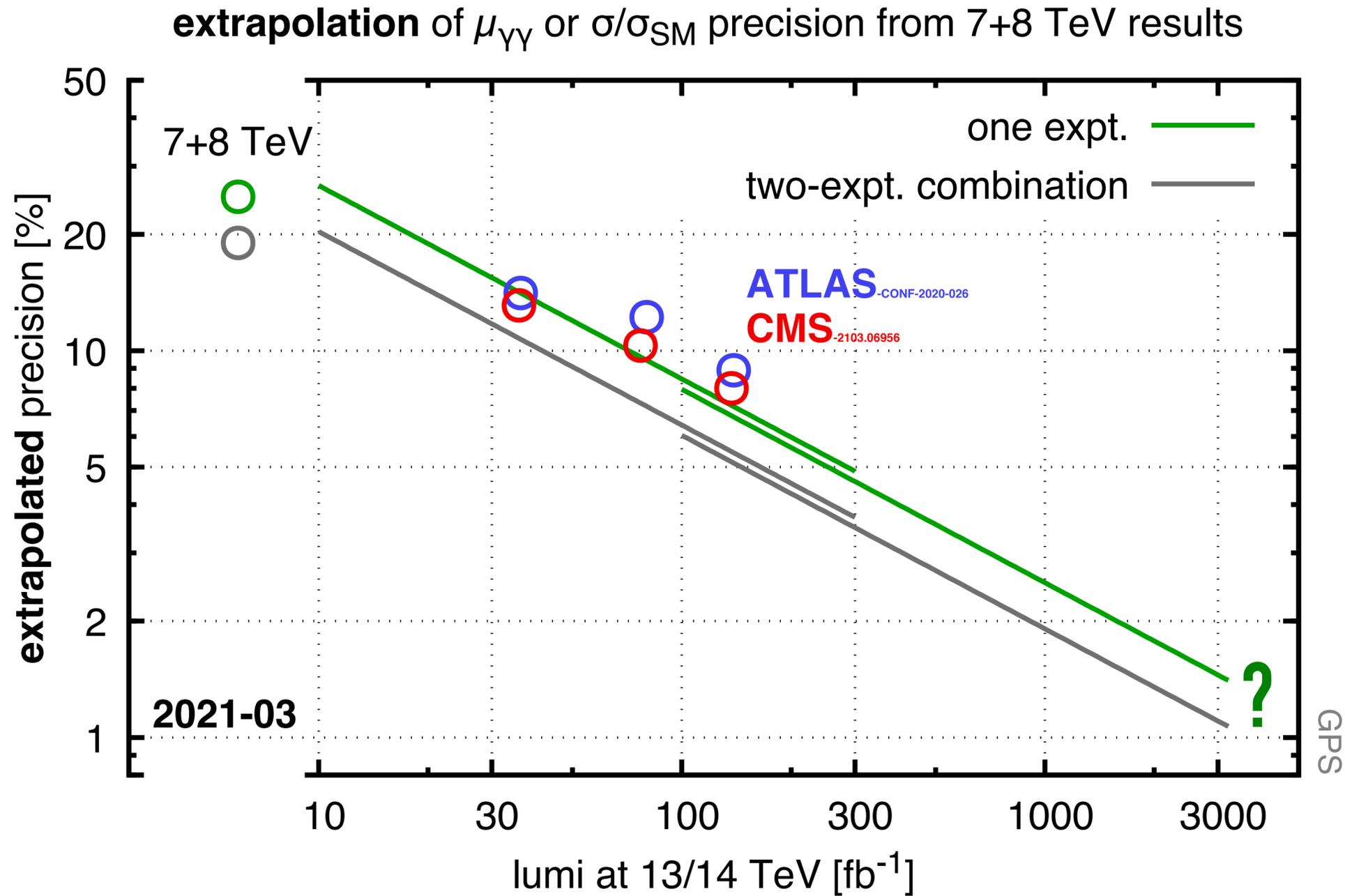
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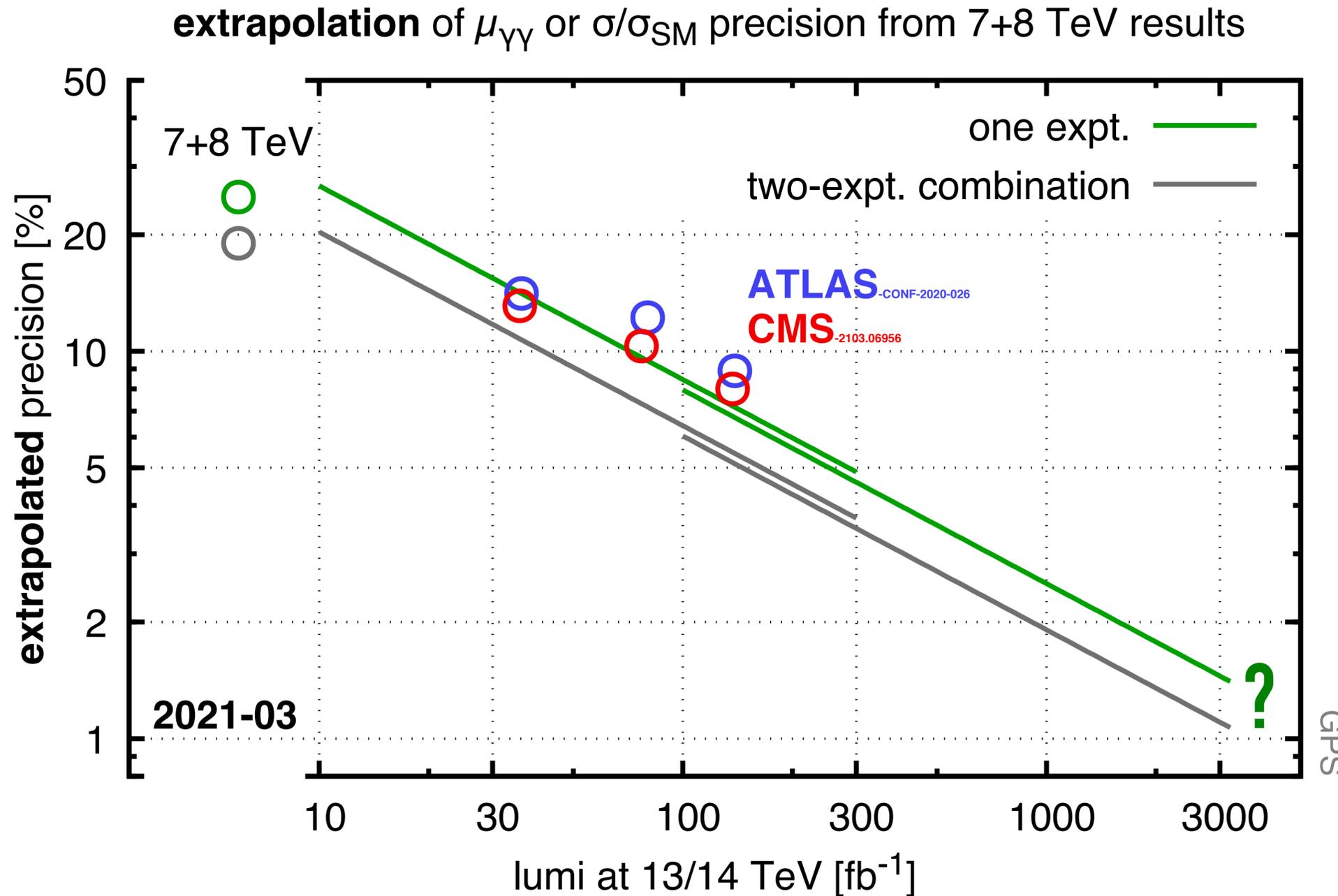
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H → γγ, an indirect probe of the top Yukawa, HWW and contact ggH couplings



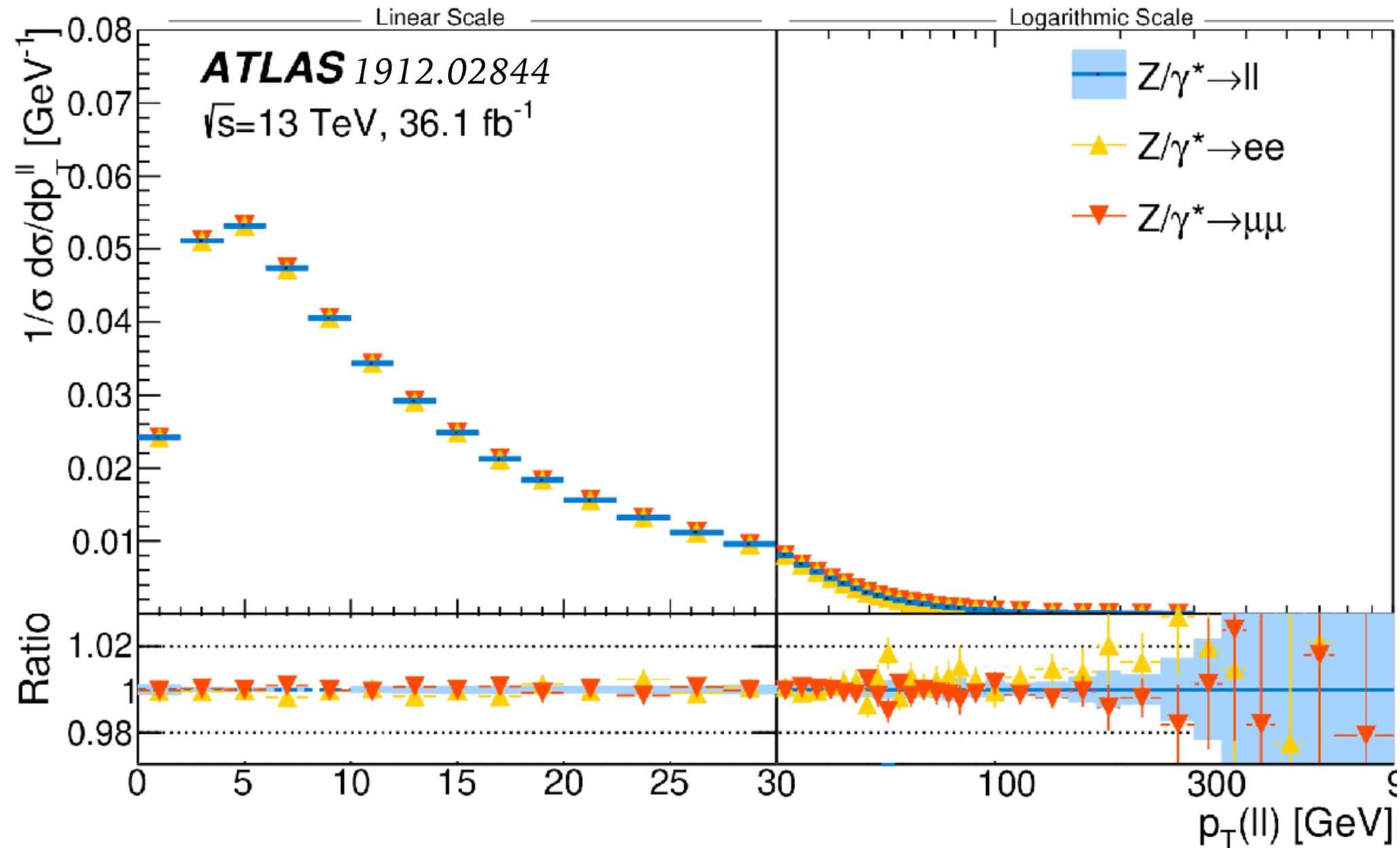
today's ATLAS and CMS total uncertainties (ratio to SM) are at the 8-9% level

- 5-6% stat.
- 3-6% syst.
- ~5% theo.

what is possible experimentally?

[in a quasi-ideal world]

Z p_T distribution — a showcase for LHC precision



Normalised distribution's statistical and systematic errors well below 1% all the way to p_T ~ 200 GeV

Largest normalisation err is luminosity then lepton ID

$$\sigma_{\text{fid}} = 736.2 \pm 0.2 \text{ (stat)} \pm 6.4 \text{ (syst)} \pm 15.5 \text{ (lumi)} \text{ pb}$$

Precision luminosity measurement in proton-proton collisions at $\sqrt{s} = 13$ TeV in 2015 and 2016 at CMS

Table 4: Summary of contributions to the relative systematic uncertainty in σ_{vis} (in %) at $\sqrt{s} = 13$ TeV in 2015 and 2016. The systematic uncertainty is divided into groups affecting the description of the vdM profile and the bunch population product measurement (normalization), and the measurement of the rate in physics running conditions (integration). The fourth column indicates whether the sources of uncertainty are correlated between the two calibrations at $\sqrt{s} = 13$ TeV.

Source	2015 [%]	2016 [%]	Corr
Normalization uncertainty			
<i>Bunch population</i>			
Ghost and satellite charge	0.1	0.1	Yes
Beam current normalization	0.2	0.2	Yes
<i>Beam position monitoring</i>			
Orbit drift	0.2	0.1	No
Residual differences	0.8	0.5	Yes
<i>Beam overlap description</i>			
Beam-beam effects	0.5	0.5	Yes
Length scale calibration	0.2	0.3	Yes
Transverse factorizability	0.5	0.5	Yes
<i>Result consistency</i>			
Other variations in σ_{vis}	0.6	0.3	No
Integration uncertainty			
<i>Out-of-time pileup corrections</i>			
Type 1 corrections	0.3	0.3	Yes
Type 2 corrections	0.1	0.3	Yes
<i>Detector performance</i>			
Cross-detector stability	0.6	0.5	No
Linearity	0.5	0.3	Yes
<i>Data acquisition</i>			
CMS deadtime	0.5	<0.1	No
Total normalization uncertainty	1.3	1.0	—
Total integration uncertainty	1.0	0.7	—
Total uncertainty	1.6	1.2	—

Luminosity: the systematic common to all measurements

- has hovered around 2% for many years (except LHCb)
- CMS has recently shown that they can get it down to 1.2%
- a major achievement, because it matters across the spectrum of precision LHC results

the master formula

$$\sigma = \sum_{i,j} \int dx_1 dx_2 f_{i/p}(x_1) f_{j/p}(x_2) \hat{\sigma}(x_1 x_2 s) \times [1 + \mathcal{O}(\Lambda/M)^p]$$

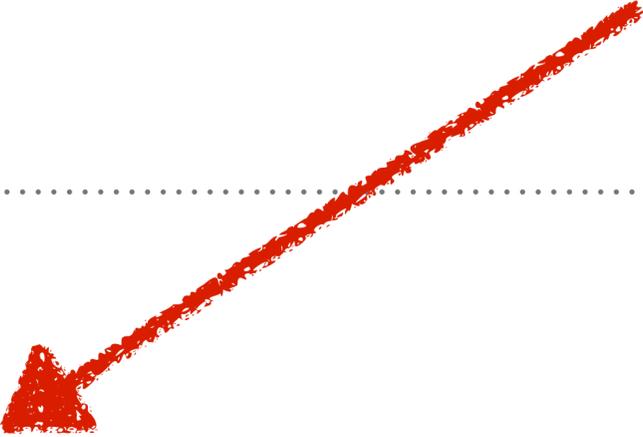
m_H (GeV)	Cross Section (pb)	TH Gaussian %	\pm PDF %	$\pm\alpha_s$ %
125.00	4.858E+01	± 3.9	± 1.9	± 2.6

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Alex Huss's talk tomorrow
(including a conceptual surprise)

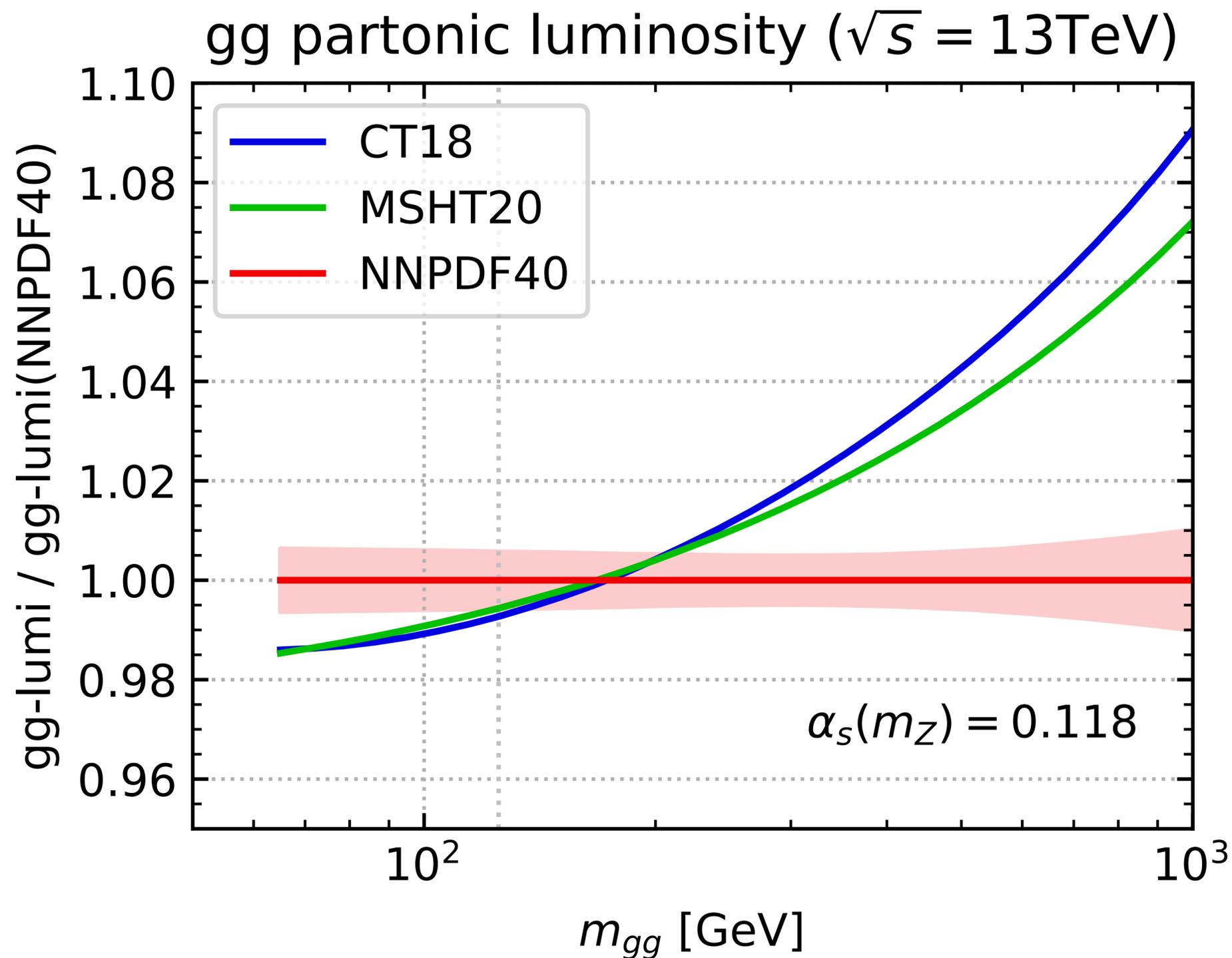
HXSWG YR 4 $gg \rightarrow H$ uncertainties

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Comparing modern PDF sets



gg-lumi, ratio to PDF4LHC15 @ m_H

PDF4LHC15	1.0000	\pm	0.0184	$\times 3$
CT18	0.9914	\pm	0.0180	
MSHT20	0.9930	\pm	0.0108	
NNPDF40	0.9986	\pm	0.0058	

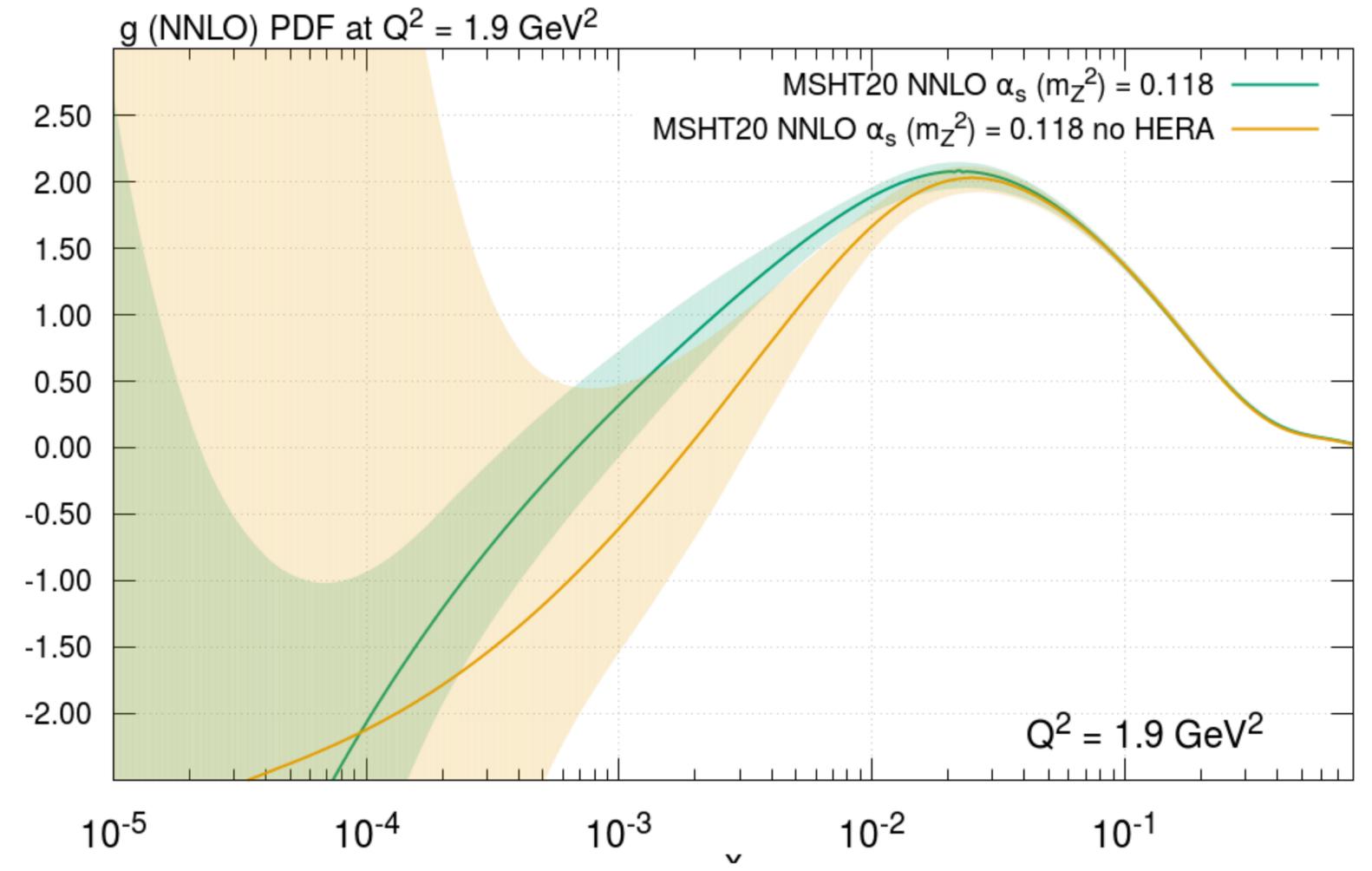
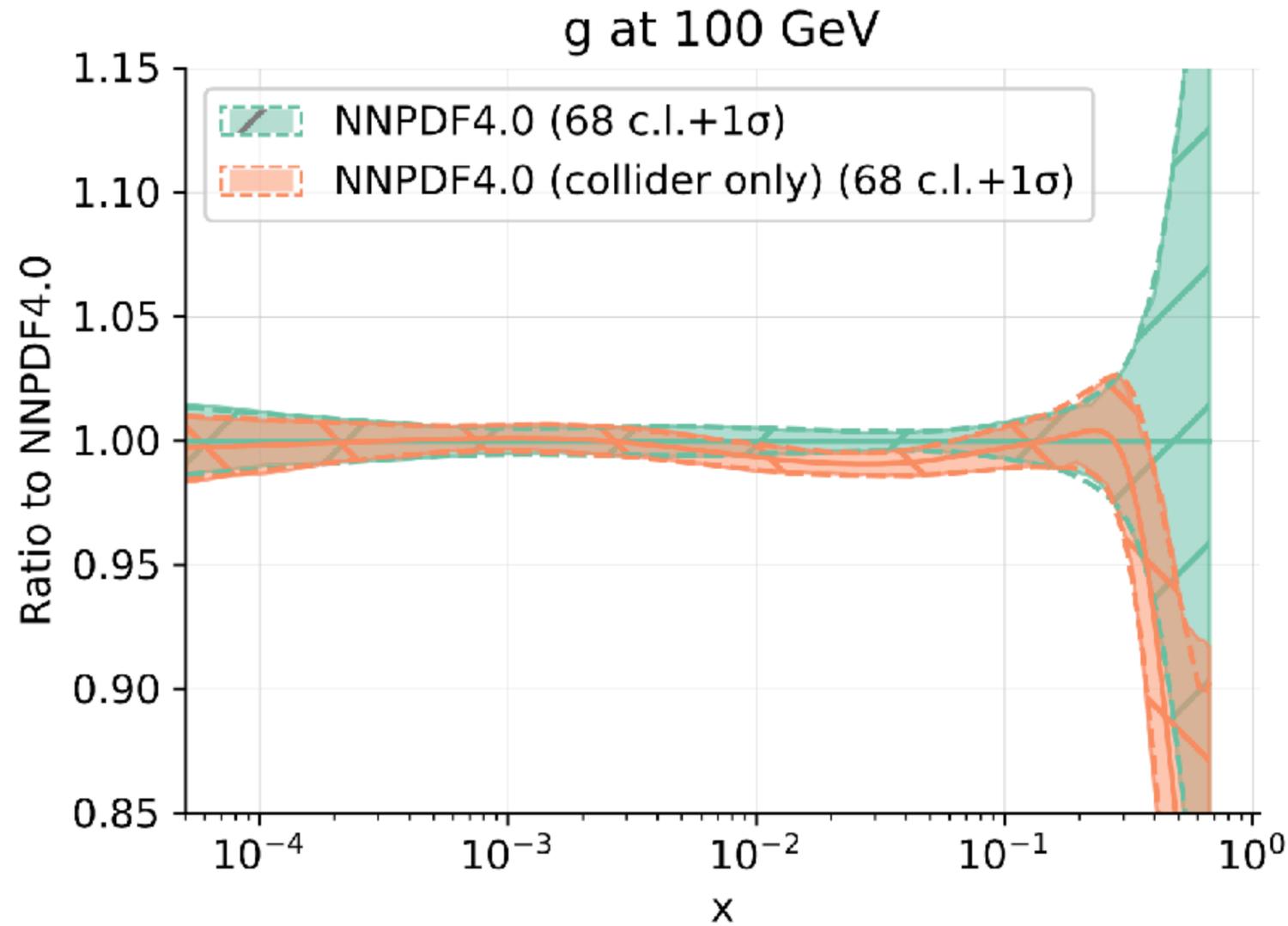
Amazing that MSHT20 & NNPDF40 are reaching %-level precision

Differences include

- methodology (replicas & NN fits, tolerance factors, etc.)
- data inputs
- treatment of charm

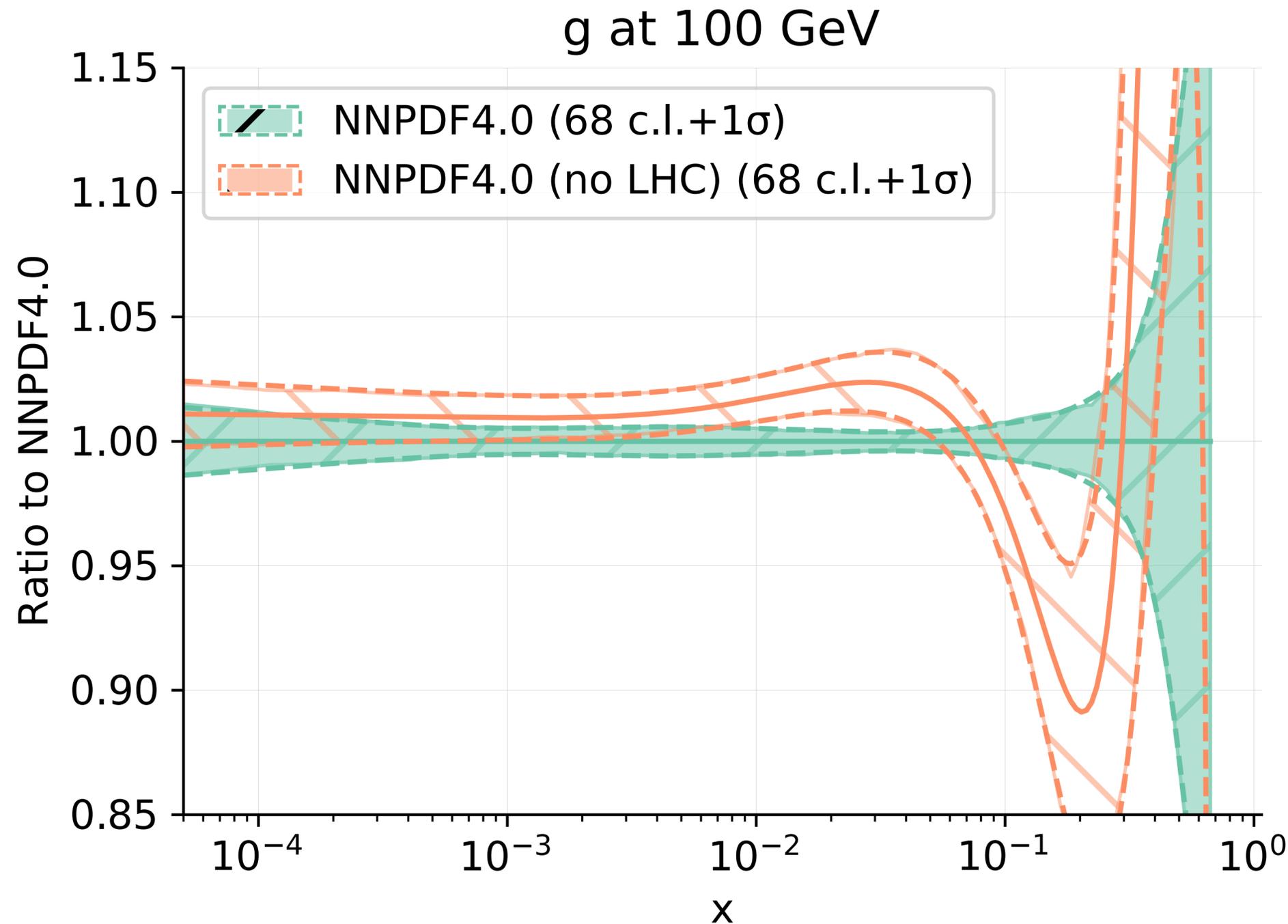
At this level, QED effects probably no longer optional

Removing DIS data (and associated worries about sizeable Λ^2/Q^2 corrections)



Reassuring indications that results are not (substantially) affected by Λ^2/Q^2 corrections from low- Q^2 DIS part of fit

Removing LHC data



- LHC data appears to be dominant in constraining the gluon
- One clear question is how to interpret gg-lumi uncertainties $\lesssim 1\%$ when all input cross sections @ hadron colliders have larger theory uncertainties.

HXSWG YR 4 $gg \rightarrow H$ uncertainties

m_H (GeV)	Cross Section (pb)	TH Gaussian %	\pm PDF %	$\pm\alpha_s$ %
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The strong coupling

HXSWG YR4	0.1180 ± 0.0015
PDG 2019	0.1179 ± 0.0010
ALPHA lattice (step scaling)	0.1185 ± 0.0008

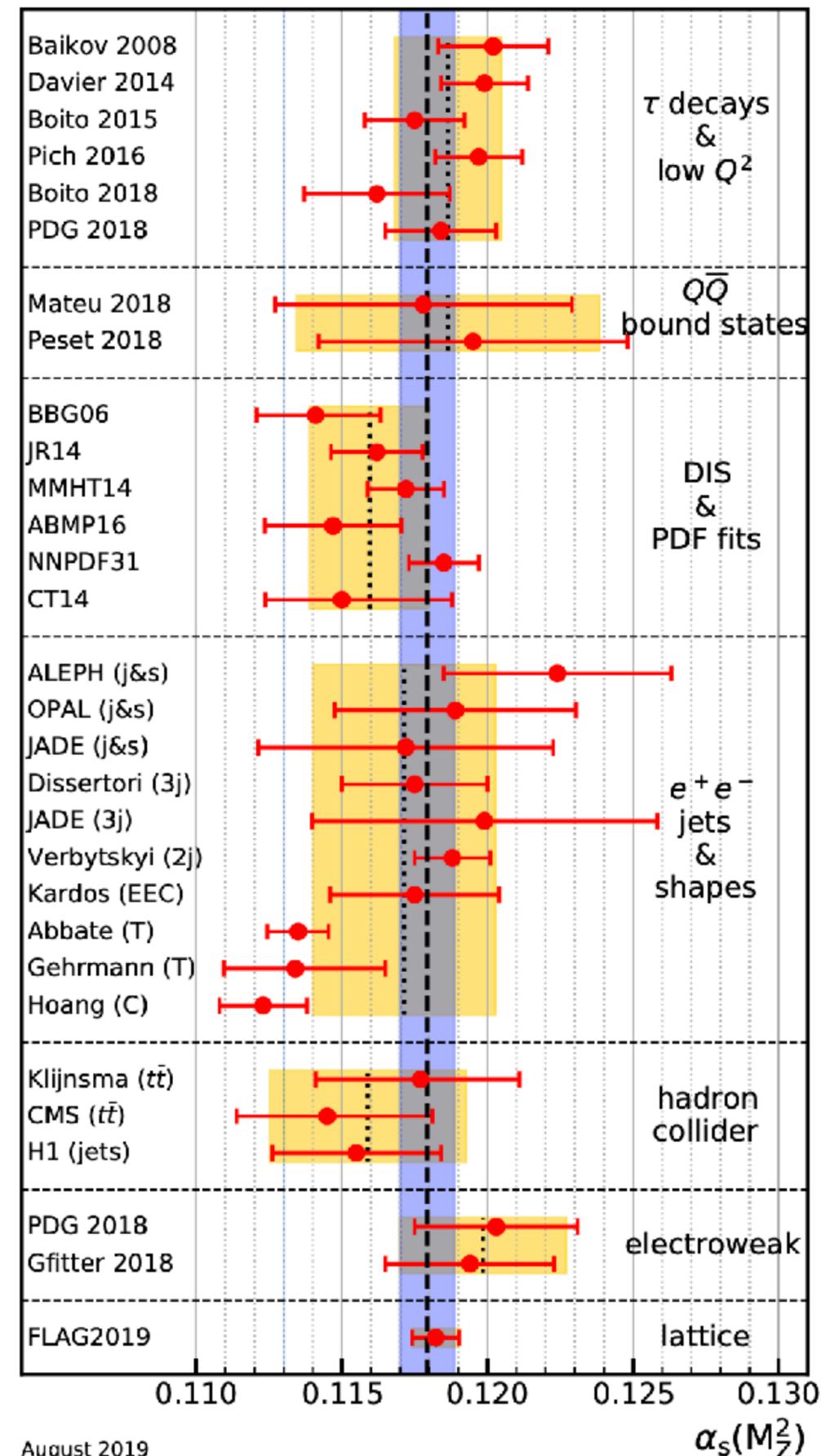
Impact of ± 0.0010 on $\sigma_{gg \rightarrow H}$ is $\pm 2.1\%$ (NNPDF40+ihixs)

Until we get FCC-ee Z hadronic width measurement, I don't see any way forward that isn't (step scaling) lattice-based

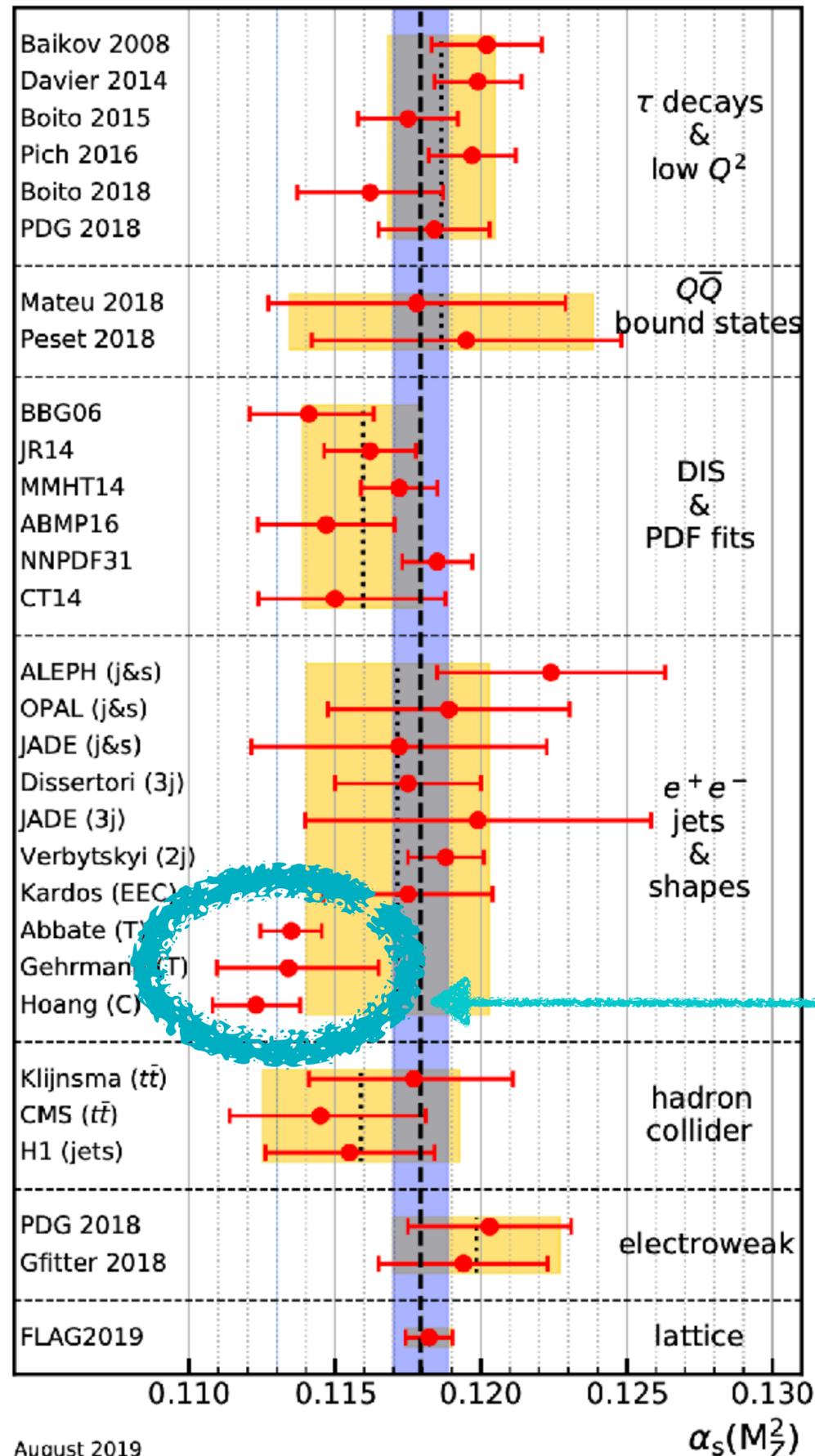
Lattice determinations of the strong coupling

2101.04762

Luigi Del Debbio^a, Alberto Ramos^{b,1}



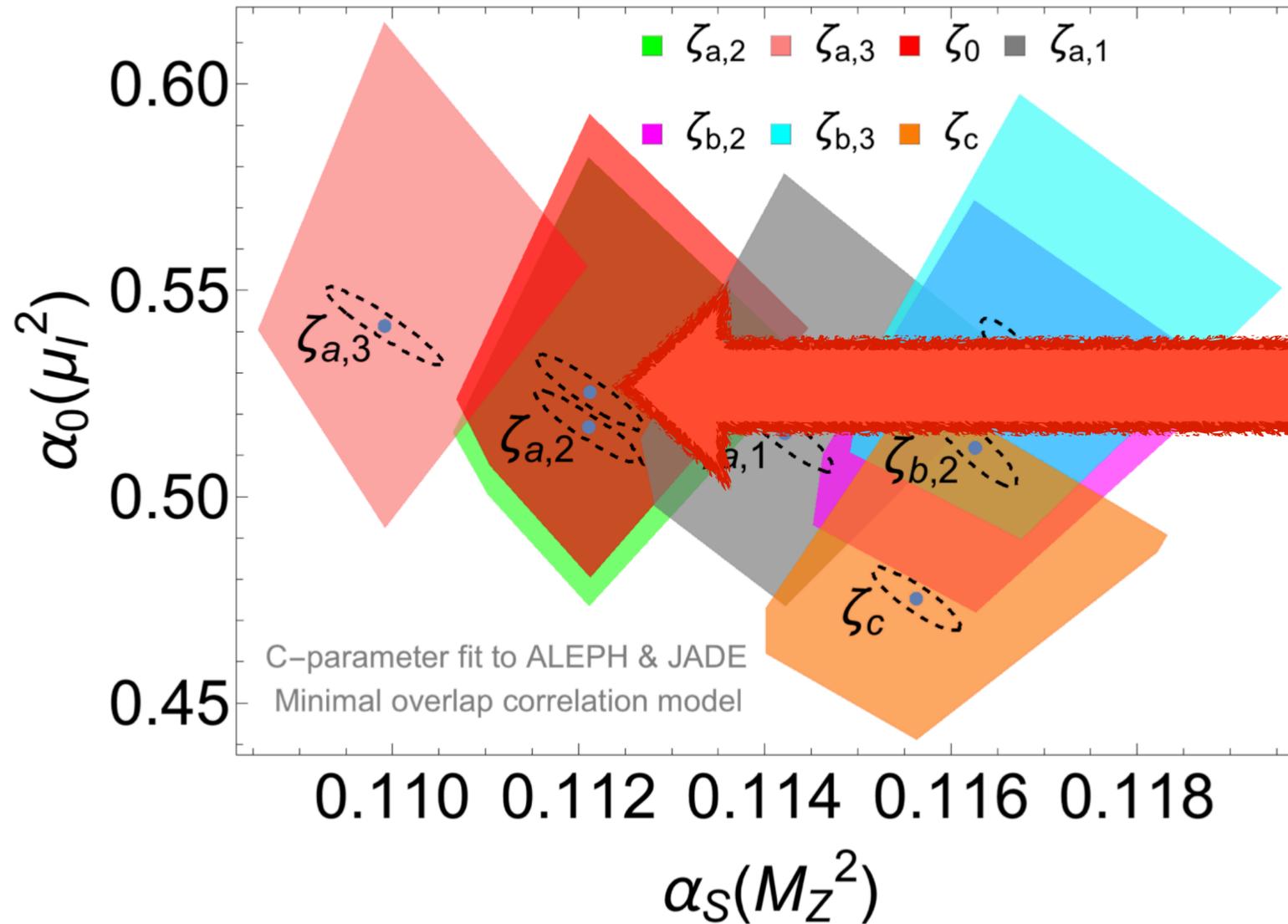
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e^+e^- C-parameter [SCET]	0.1123 ± 0.0015
e^+e^- Thrust [SCET]	0.1135 ± 0.0011

Aside from EW fit and ALPHA lattice, most determinations depend, in some way or other, on measurements that are uncomfortably close / sensitive to non-perturbative physics

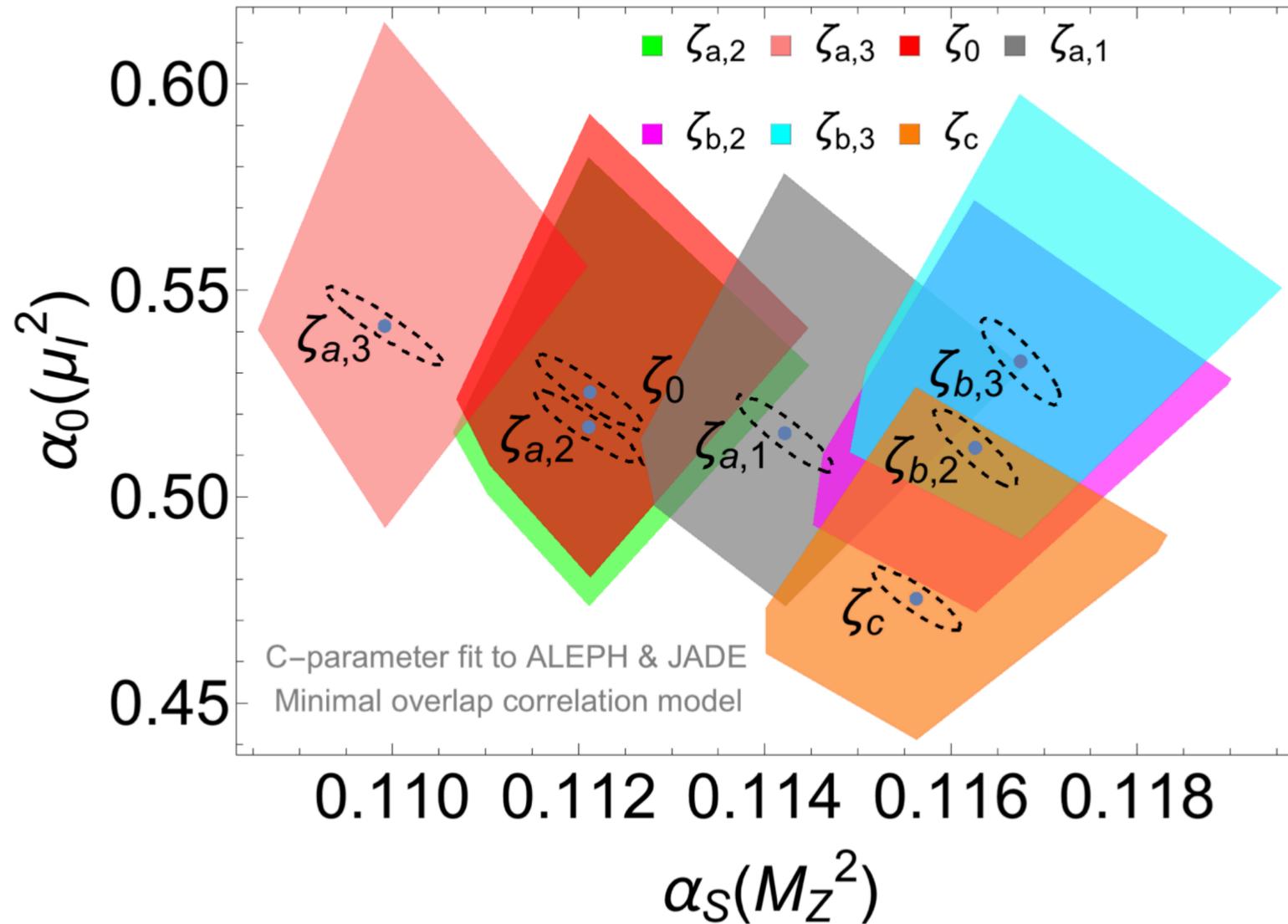
C-param fits with different assumptions for Λ/Q correction (between 2 & 3-jet limits)



- measurement essentially looks at rate of 3rd jet emission in $e^+e^- \rightarrow q\bar{q}$
- $0.1123 \pm 0.0015 \Leftrightarrow$ assumption about the structure of Λ/Q corrections, based on the 2-jet limit

Luisoni, Monni & GPS, [2012.00622](#)

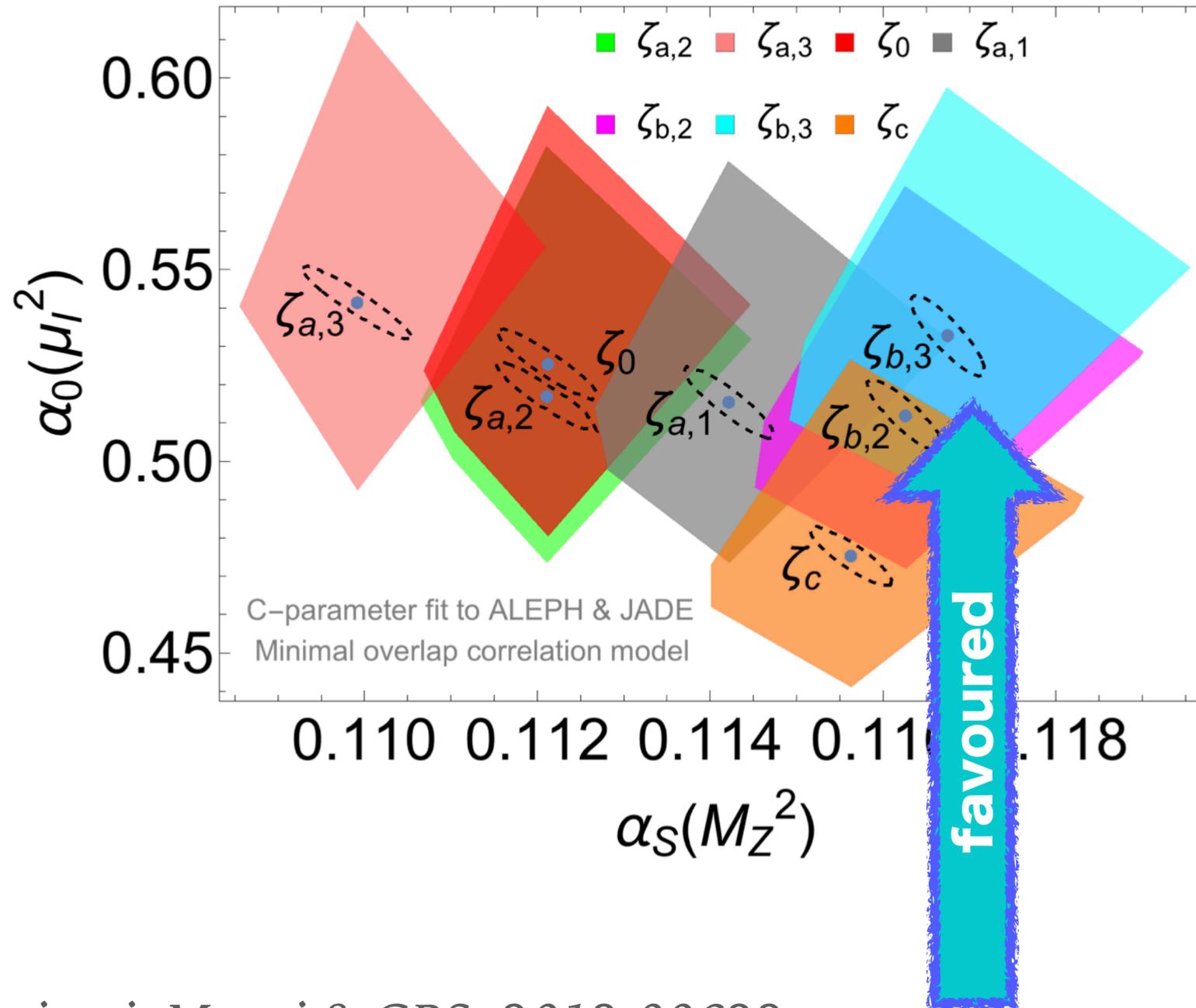
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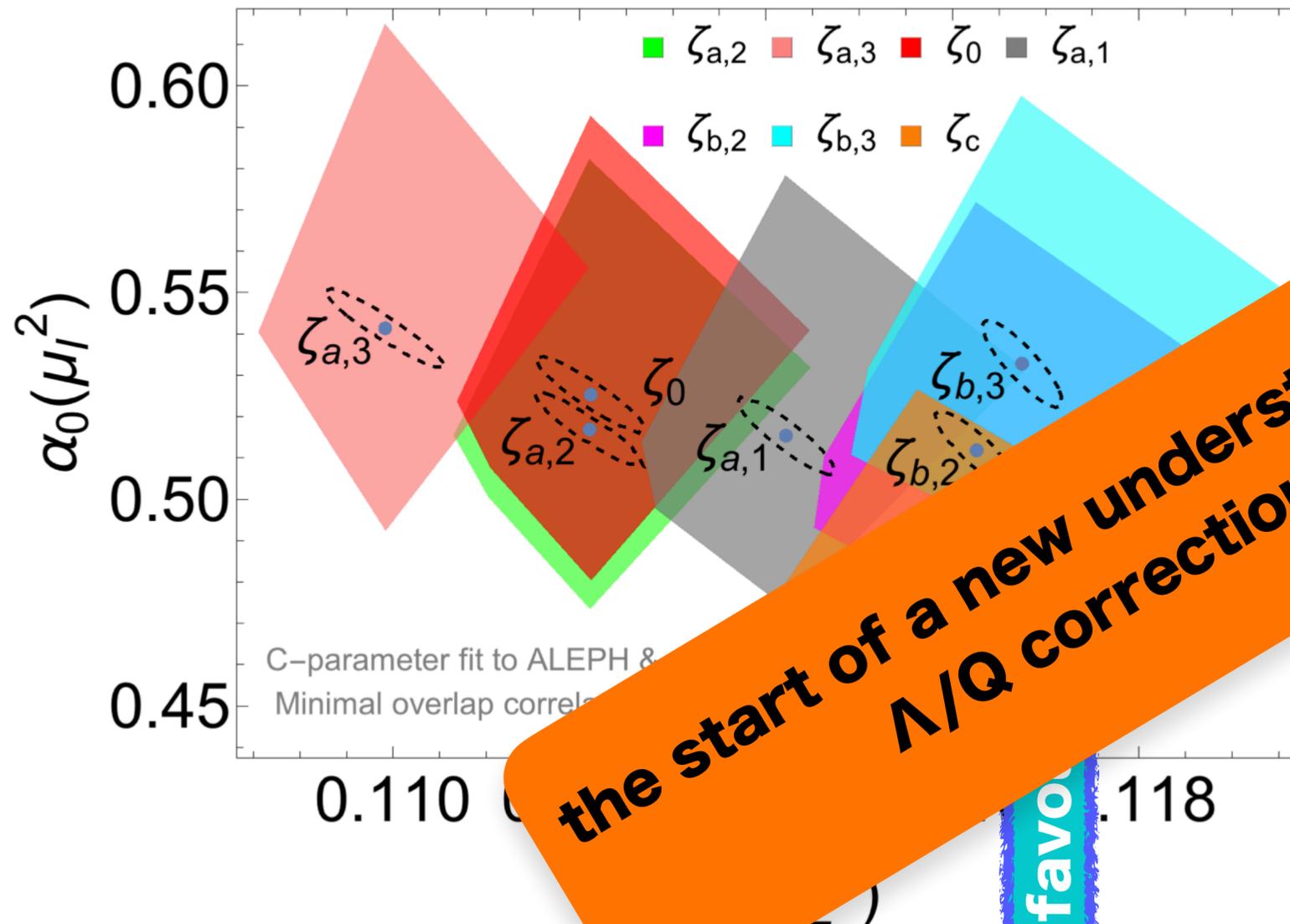


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Caola, Ferrario Ravasio, Limatola, Melnikov & Nason, [2108.08897](#)

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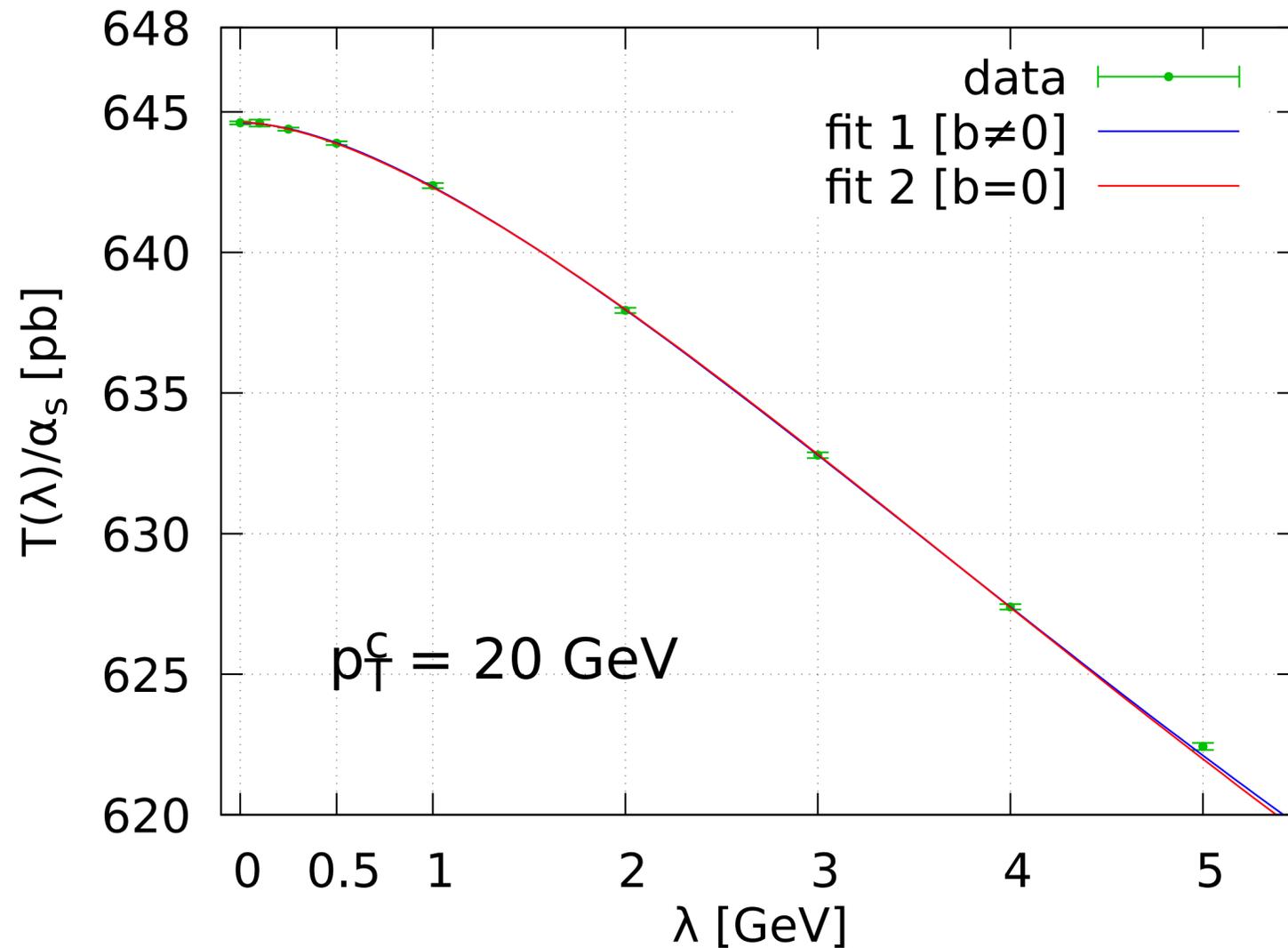
the non-perturbative part

$$\sigma = \sum_{i,j} \int dx_1 dx_2 f_{i/p}(x_1) f_{j/p}(x_2) \hat{\sigma}(x_1 x_2 s) \times \left[1 + \mathcal{O}(\Lambda/M)^p \right]$$

What is value of p in $(\Lambda/Q)^p$?

- ▶ LEP event-shape (C-parameter, thrust) fit troubles came about because $p = 1$
 $\Lambda \sim 0.5 \text{ GeV} \rightarrow (\Lambda/20\text{GeV}) \sim 2.5 \%$
- ▶ Jet physics at LHC is dirty because $p = 1$ (hadronisation & MPI)
- ▶ Hadron-collider inclusive and rapidity-differential Drell-Yan cross sections are believed to have $p = 2$ (Higgs hopefully also), so leptonic / photonic decays should be clean, aside from isolation.
 $\Lambda \sim 0.5 \text{ GeV} \rightarrow (\Lambda/125\text{GeV})^2 \sim 0.002 \%$
[Beneke & Braun, hep-ph/9506452; Dasgupta, hep-ph/9911391]
- ▶ But at LHC, we're also interested in Z, W and Higgs production with non-zero p_T
Nobody knew if we have $(\Lambda/p_T)^p$ with $p = 1$ (a disaster) or $p = 2$ (all is fine)

What is value of p in $(\Lambda/Q)^p$?



- Explicit calculations with an effective gluon mass (λ) can provide an answer
- Flatness in plot for $\lambda \rightarrow 0$ indicates **absence of $p = 1$** (linear) contribution
- arguably the most important result of the year, because it lays foundations for precision physics at non-zero p_T

Ferraro Ravasio, Limatola & Nason, 2011.14114

+ analytic demonstration in Caola, Ferrario Ravasio, Limatola, Melnikov & Nason, [2108.08897](#)

beyond the fixed-order formula

parton shower Monte Carlos

Take example of ATLAS boosted VH — stat (28%) ~ syst (24%)

ATLAS VH: 2008.02508,

Source of uncertainty	Avg. impact	
Total	0.372	
Statistical	0.283	
Systematic	0.240	
Experimental uncertainties		
Small- R jets	0.038	
Large- R jets	0.133	
E_T^{miss}	0.007	
Leptons	0.010	
b -tagging	b -jets	0.016
	c -jets	0.011
	light-flavour jets	0.008
	extrapolation	0.004
Pile-up	0.001	
Luminosity	0.013	
Theoretical and modelling uncertainties		
Signal	0.038	
Backgrounds	0.100	
$\leftrightarrow Z + \text{jets}$	0.048	
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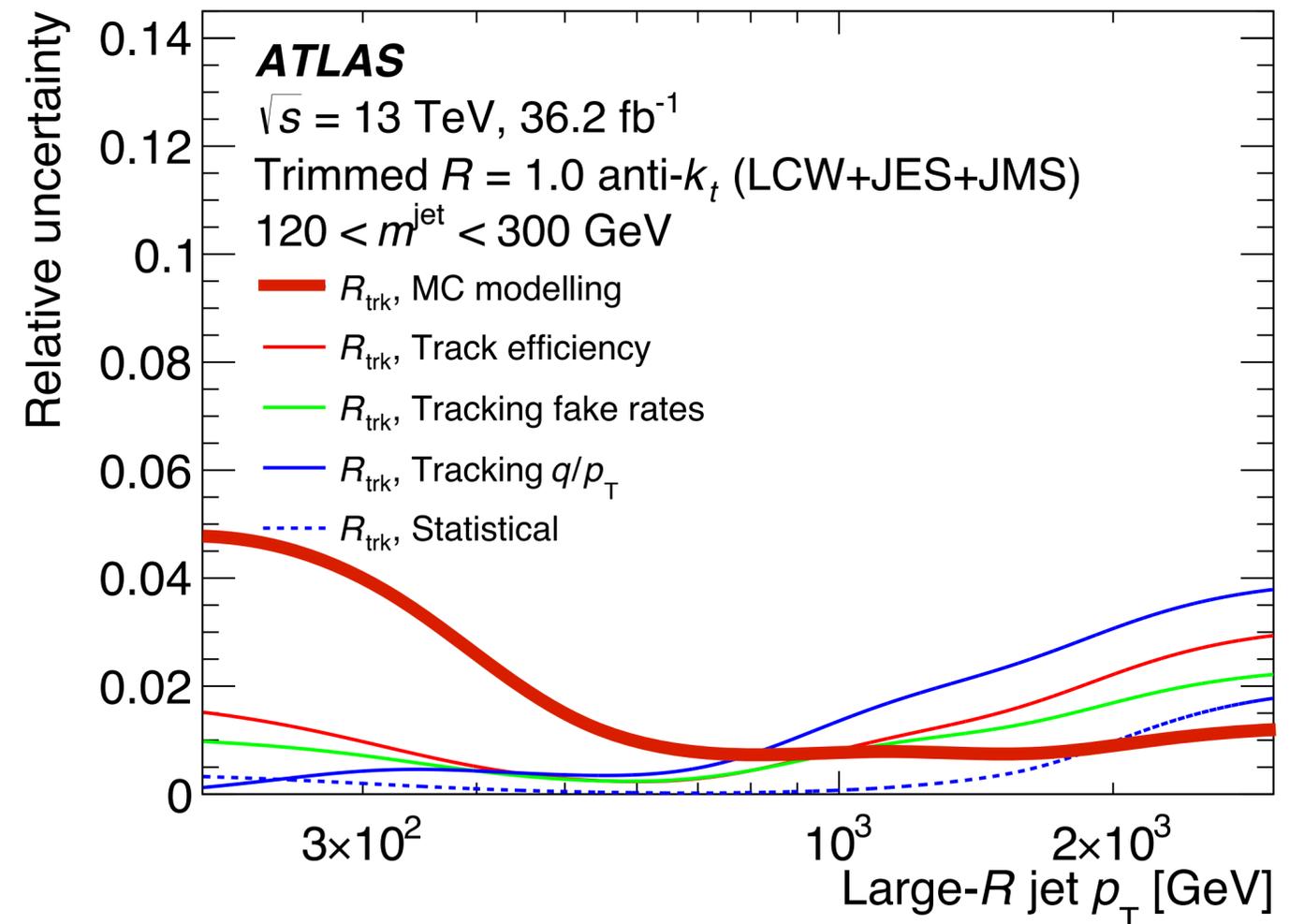
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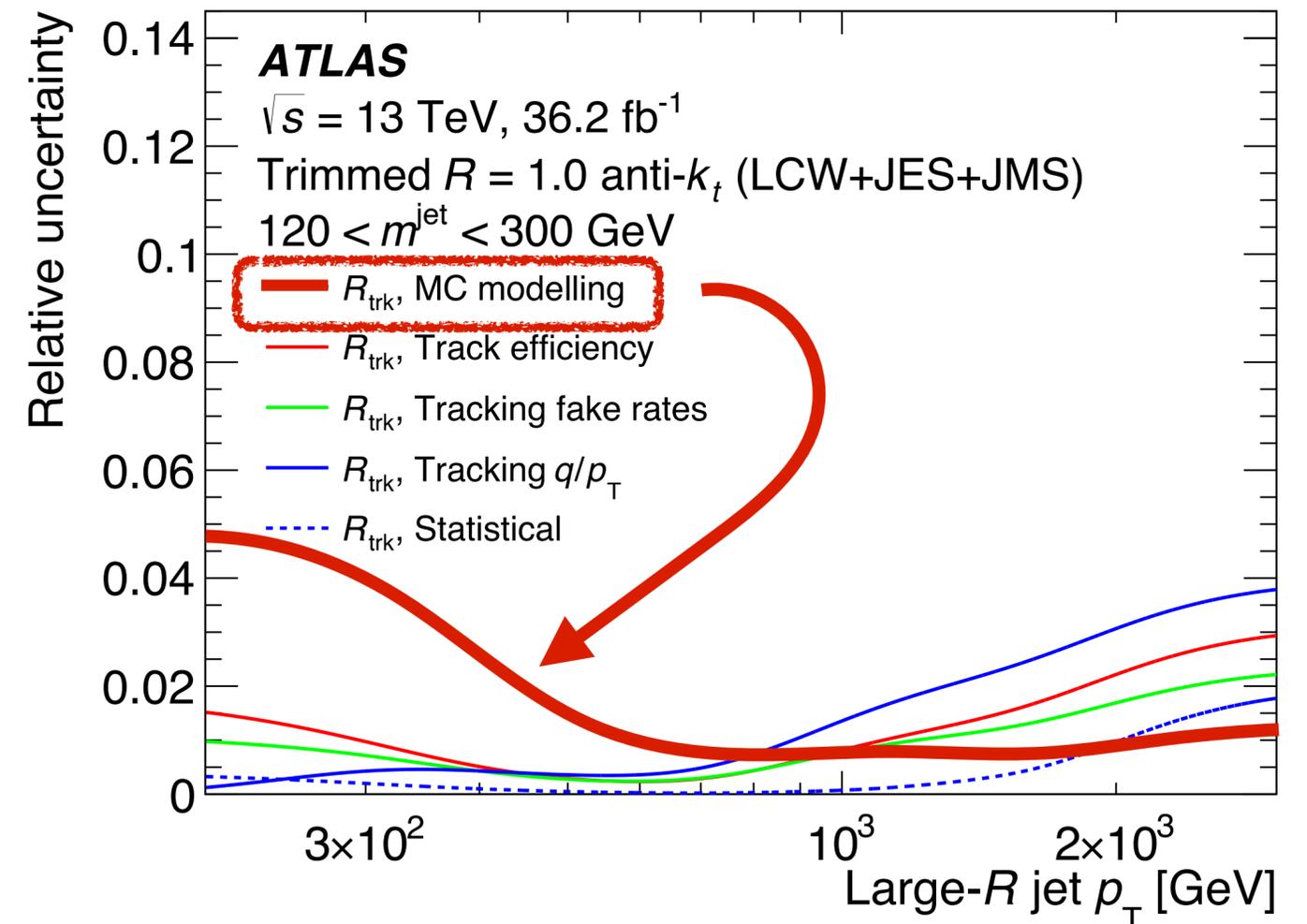


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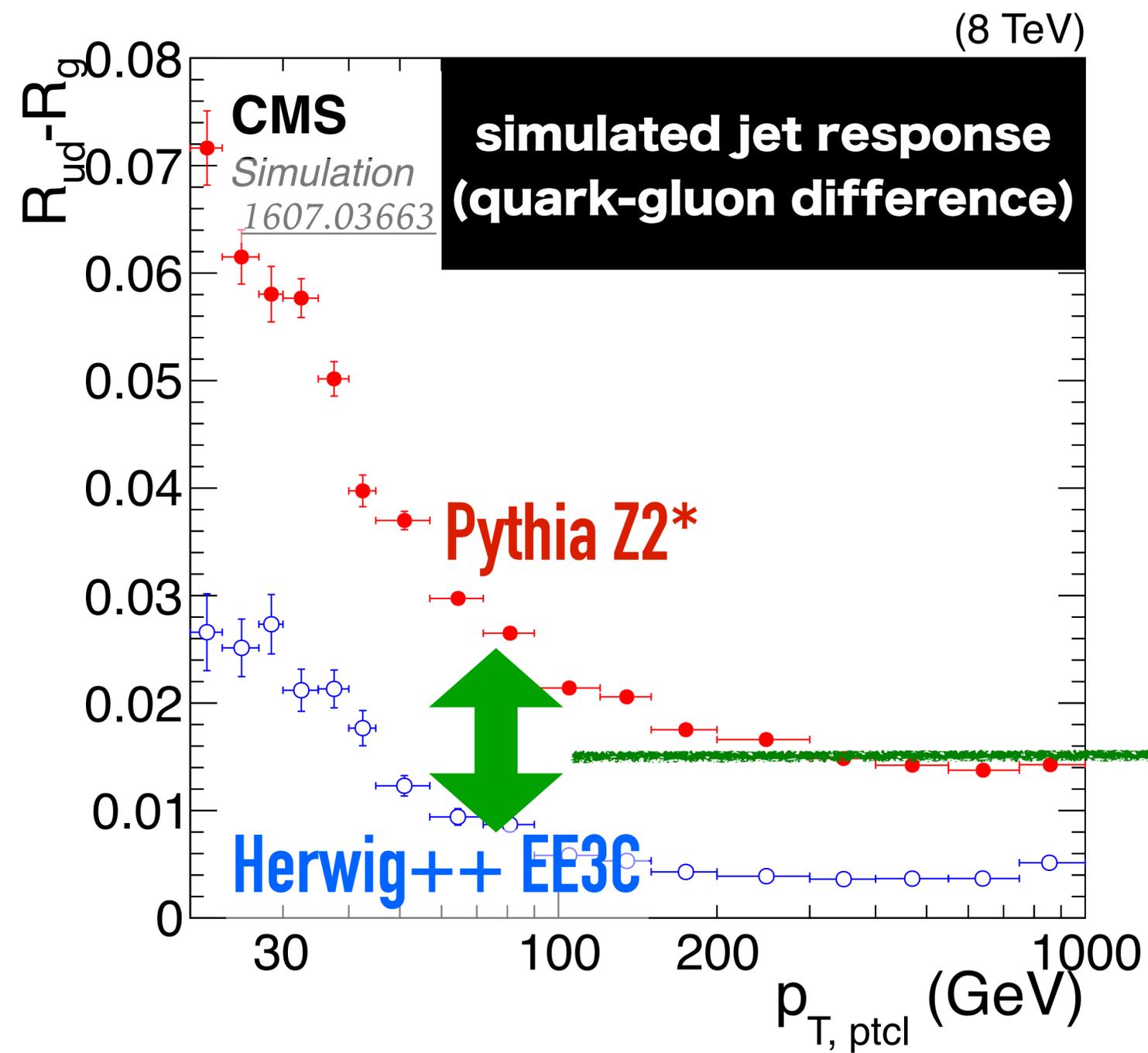
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↔ Multijet	0.009	
MC statistical	0.092	

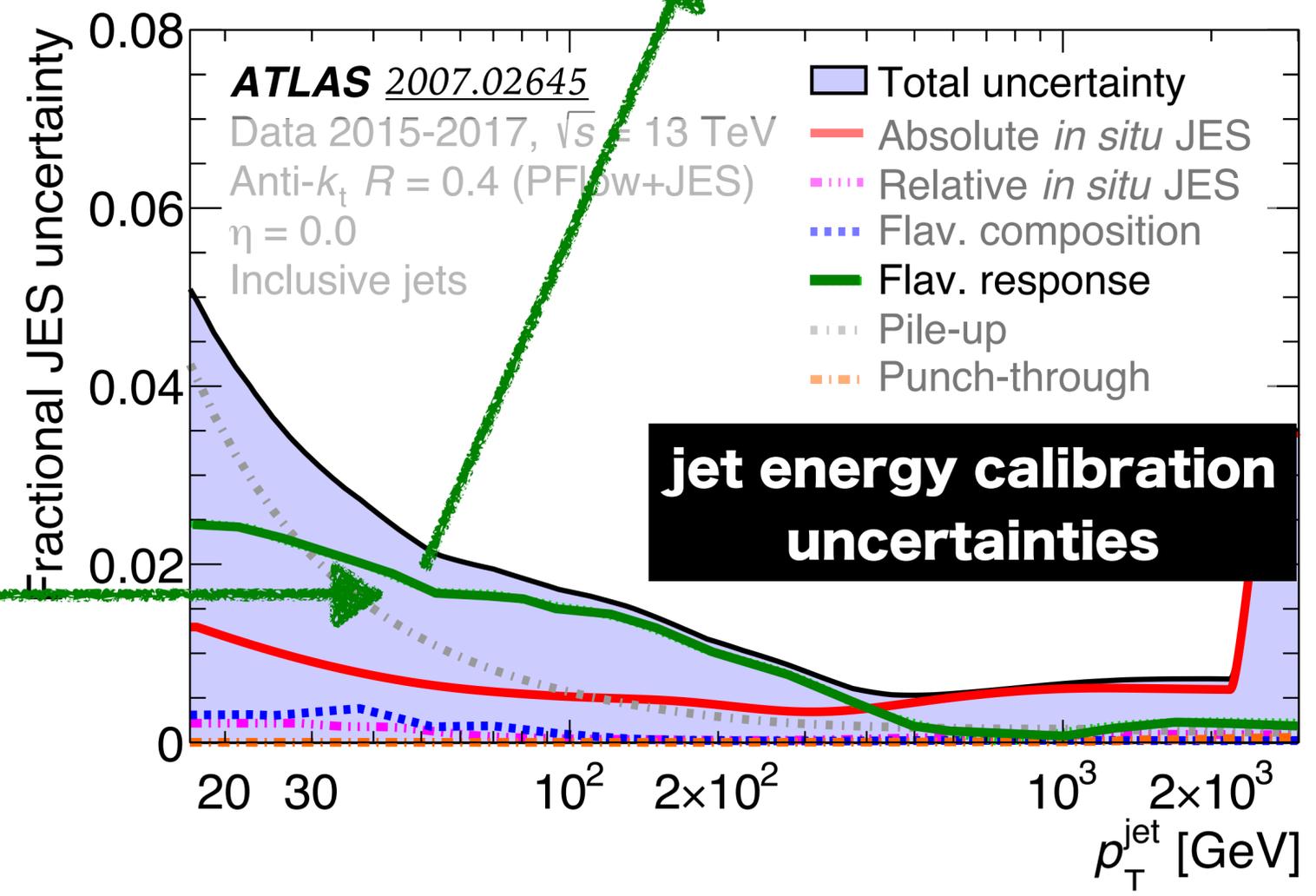
For large- R jets, the uncertainties in the energy and mass scales are [...] as described in [81]



But imperfections matter: e.g. for jet energy calibration (affects ~1500 papers)



Largest uncertainty source is poor understanding of [parton shower simulations of] quark v. gluon-induced jet responses



Resummation @N3LL, but parton showers only LL? **Now evolving to NLL**

Deductor

$k_t \theta$ (“ Λ ”) ordered

Recoil

\perp : local

+: local

-: global

Tests

analytical /numerical
for thrust

Nagy & Soper

2011.04777 (+past decade)

FHP

k_t ordered

Recoil

\perp : global

+: local

-: global

Tests

analytical
for thrust &
multiplicity

Forshaw, Holguin & Plätzer

2003.06400

PanLocal

$k_t \sqrt{\theta}$ ordered

Recoil

\perp : local

+: local

-: local

Tests

numerical
for many observables

Dasgupta, Dreyer, Hamilton, Monni, GPS & Soyez 2002.11114

Hamilton, Medves, GPS, Scyboz, Soyez, 2011.10054

Karlberg, GPS, Scyboz, Verheyen, 2103.16526

PanGlobal

k_t or $k_t \sqrt{\theta}$ ordered

Recoil

\perp : global

+: local

-: local

Tests

numerical
for many observables

future colliders

Will e^+e^- colliders make precision easy?

Table 1.1. Relative statistical uncertainty on $\sigma_{\text{HZ}} \times \text{BR}(H \rightarrow \text{XX})$ and $\sigma_{\nu\bar{\nu}H} \times \text{BR}(H \rightarrow \text{XX})$, as expected from the FCC-ee data, obtained from a fast simulation of the CLD detector and consolidated with extrapolations from full simulations of similar linear-collider detectors (SiD and CLIC).

\sqrt{s} (GeV)	240		365	
Luminosity (ab^{-1})	5		1.5	
$\delta(\sigma\text{BR})/\sigma\text{BR}$ (%)	HZ	$\nu\bar{\nu}$ H	HZ	$\nu\bar{\nu}$ H
H \rightarrow any	± 0.5		± 0.9	
H \rightarrow bb	± 0.3	± 3.1	± 0.5	± 0.9
H \rightarrow cc	± 2.2		± 6.5	± 10
H \rightarrow gg	± 1.9		± 3.5	± 4.5
H \rightarrow W^+W^-	± 1.2		± 2.6	± 3.0
H \rightarrow ZZ	± 4.4		± 12	± 10
H \rightarrow $\tau\tau$	± 0.9		± 1.8	± 8
H \rightarrow $\gamma\gamma$	± 9.0		± 18	± 22
H \rightarrow $\mu^+\mu^-$	± 19		± 40	
H \rightarrow invisible	< 0.3		< 0.6	

Notes. All numbers indicate 68% CL intervals, except for the 95% CL sensitivity in the last line. The accuracies expected with 5 ab^{-1} at 240 GeV are given in the middle column, and those expected with 1.5 ab^{-1} at $\sqrt{s} = 365 \text{ GeV}$ are displayed in the last column.

- Up to $\sim \times 10$ reduction in uncertainties
- Interpreting 0.3% for $H \rightarrow b\bar{b}$ will require substantial improvements in parametric inputs
- Much of the statistics involves hadronic modes — how well will we be able to exploit them?
- Agreement between e^+e^- and LHC will be powerful validation of hadron colliders as precision machines

conclusions

Conclusions

- Across much of Higgs physics, theory / MC uncertainties are among the dominant systematic uncertainties — addressing them is key to benefitting from $\times 20$ statistics of the next 15 years.
- Perturbative calculations are making amazing strides (cf. Alex Huss's talk tomorrow) → technically immensely challenging, and making remarkable progress
- Other aspects (parameters, PDFs, parton showers, non-perturbative contributions) force us to address conceptually complicated questions, e.g.
 - non-perturbative corrections, with remarkable progress (& good news) this past year!