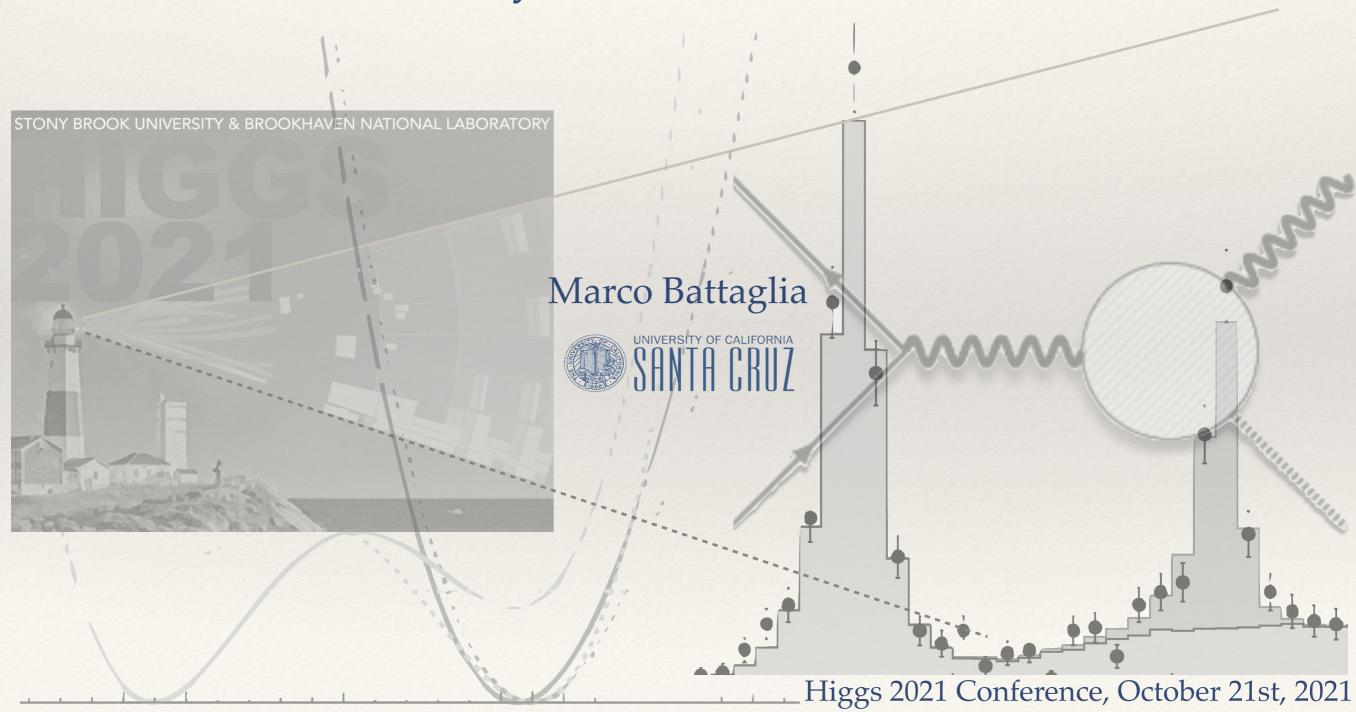




Effective Field Theory Results from Higgs (and beyond) by ATLAS and CMS



Effective Field Theories, the Higgs Sector ...

Leading effects of many BSM scenarios can be parametrized through operators of effective field theories (EFTs), with effects suppressed by new physics scale Λ corresponding to new particle masses;

Higgs measurements bring sensitivity from interplay of tree and loop-mediated couplings involving 3rd generation quarks and leptons, gluons and gauge bosons;

... and beyond

EFTs have become a standard tool allowing study of large experimental data sets, correlating Higgs, top and EW sectors probing same operators, without assuming underlying theory to be valid to arbitrarily high energies;

Interpretations of EFT results brings some caveats since effects of dim>6 operators are not precisely known and EFT validity is restricted within that of decoupling theorem where heavy fields decouple at low momenta.

Very active area for both extraction of new experimental results and their interpretation by theorists (JHEP06 (2018) 146, arXiv:2105.00006, ...) and experimental collaborations;

This talk reviews recent Higgs physics results relevant to studies in EFT framework and related EFT interpretation performed by ATLAS and CMS collaborations with a view on focusing on open problems and highlighting interplay with non-Higgs physics.

SMEFT and Dim-6 EFT Operators relevant to Higgs physics

If New Physics degrees of freedom can be integrated out, Higgs is SM-like and New Physics can manifest itself through higher-dimension effective interactions between SM fields; Scale-dependent Wilson coefficients $C_i(Q^2)$ encode New Physics deviations from SM and can be probed by experimental measurements: $\mathcal{L}_{Eff} = \mathcal{L}_{SM} + \sum_{i} \frac{C_i^{(6)} O_i^{(6)}}{\Lambda^2} + \mathcal{O}(\Lambda^{-4})$

Coefficient	Operator	Example process
c_{uG}	$(\bar{q}_p \sigma^{\mu\nu} T^A u_r) \widetilde{H} G^A_{\mu\nu}$	$ \begin{array}{cccccccccccccccccccccccccccccccccccc$
c_{uW}	$(\bar{q}_p \sigma^{\mu\nu} u_r) \tau^I \widetilde{H} W^I_{\mu\nu}$	q Z t t
c_{uB}	$(\bar{q}_p \sigma^{\mu\nu} u_r) \widetilde{H} B_{\mu\nu}$	$q \nearrow \overline{t}$ Π
$c_{qq}^{{\scriptscriptstyle (1)}}$	$(\bar{q}_p\gamma_\mu q_t)(\bar{q}_r\gamma^\mu q_s)$	
$c_{qq}^{{\scriptscriptstyle (3)}}$	$(\bar{q}_p \gamma_\mu \tau^I q_r) (\bar{q}_s \gamma^\mu \tau^I q_t)$	
c_{qq}	$(\bar{q}_p\gamma_\mu q_t)(\bar{q}_r\gamma^\mu q_s)$	
$c_{qq}^{\scriptscriptstyle (31)}$	$(\bar{q}_p \gamma_\mu \tau^I q_t)(\bar{q}_r \gamma^\mu \tau^I q_s)$	
c_{uu}	$(\bar{u}_p \gamma_\mu u_r)(\bar{u}_s \gamma^\mu u_t)$	q t t
$c_{uu}^{\scriptscriptstyle (1)}$	$(\bar{u}_p \gamma_\mu u_t)(\bar{u}_r \gamma^\mu u_s)$	$q \longrightarrow t^{-t}$
$c_{qu}^{{\scriptscriptstyle (1)}}$	$(\bar{q}_p \gamma_\mu q_t)(\bar{u}_r \gamma^\mu u_s)$	
$c_{ud}^{^{(8)}}$	$(\bar{u}_p \gamma_\mu T^A u_r)(\bar{d}_s \gamma^\mu T^A d_t)$	
$c_{qu}^{_{(8)}}$	$(\bar{q}_p \gamma_\mu T^A q_r)(\bar{u}_s \gamma^\mu T^A u_t)$	
$c_{qd}^{^{(8)}}$	$(\bar{q}_p \gamma_\mu T^A q_r)(\bar{d}_s \gamma^\mu T^A d_t)$	
c_G	$f^{ABC}G^{A u}_{\mu}G^{B ho}_{ u}G^{C\mu}_{ ho}$	$ \begin{array}{cccccccccccccccccccccccccccccccccccc$

EFT with SM fields including SM Higgs doublet defines commonly used SMEFT

(but see caveats in Tim Cohen's talk);

SM dim-2 and dim-4 interactions + higher dimensional interactions between allowed combinations of SM fields;

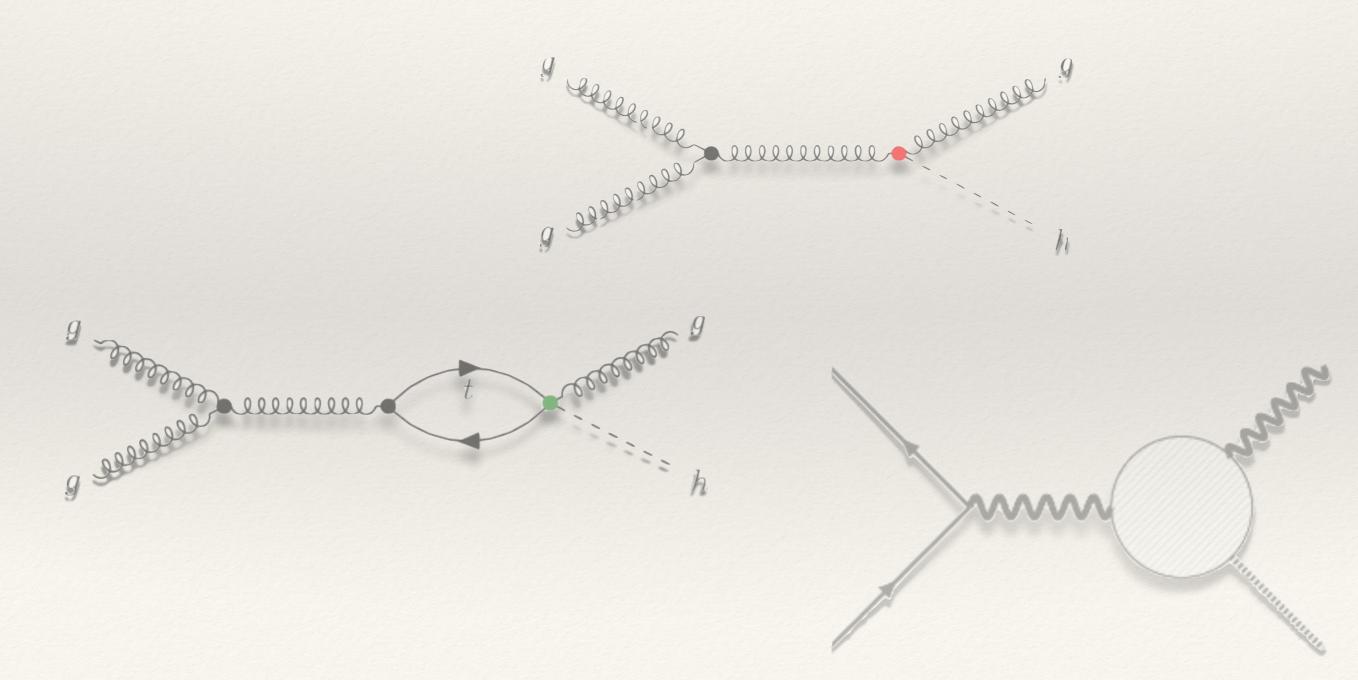
Non-redundant basis of operators defined by Warsaw basis generally used in extraction of experimental constraints by ATLAS and CMS:

Buchmuller, Wyler Nucl.Phys. B268 (1986) 621 Grzadkowski et al arXiv:1008.4884 Brivio, Trott, arXiv:1706.08945

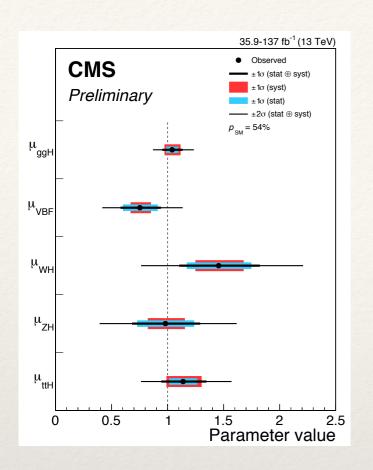
Coefficient	Operator	Example process
с _{НDD} (Т)	$\left(H^{\dagger}D^{\mu}H\right)^{*}\left(H^{\dagger}D_{\mu}H\right)$	$ \begin{array}{cccc} q & & & q \\ & & & & & q \\ Z & & & & & H \\ q & & & & & & q \end{array} $
c_{HG}	$H^\dagger H G^A_{\mu u} G^{A\mu u}$	g g
$c_{H\!B}$	$H^\dagger H B_{\mu u} B^{\mu u}$	$ \begin{array}{cccc} q & & & q \\ & & & Z \\ q & & & & q \end{array} $
c_{HW}	$H^\dagger HW^I_{\mu u}W^{I\mu u}$	$ \begin{array}{cccccccccccccccccccccccccccccccccccc$
с _{нwв}	$H^\dagger au^I H W^I_{\mu u} B^{\mu u}$	$ \begin{array}{cccc} q & & & q \\ & & & & q \\ & & & & & & q \\ q & & & & & & q \end{array} $
c_{eH}	$(H^{\dagger}H)(\bar{l}_{p}e_{r}H)$	$H \longrightarrow \ell$
$c_{Hl}^{{\scriptscriptstyle (1)}}$	$(H^{\dagger}i\overleftrightarrow{D}_{\mu}H)(\bar{l}_{p}\gamma^{\mu}l_{r})$	$ \begin{array}{c} q \\ \downarrow \qquad \qquad \downarrow \qquad \qquad \ell \\ q \\ \downarrow \qquad \qquad \qquad \downarrow \qquad \qquad \ell \\ H \end{array} $
$c_{Hl}^{^{(3)}}$	$(H^{\dagger}i\overleftrightarrow{D}_{\mu}^{I}H)(\bar{l}_{p}\tau^{I}\gamma^{\mu}l_{r})$	$q \longrightarrow W \stackrel{\nu}{\underset{\ell}{\swarrow}} \stackrel{\nu}{\underset{H}{\swarrow}}$
c_{He}	$(H^{\dagger}i\overleftrightarrow{D}_{\mu}H)(\bar{e}_{p}\gamma^{\mu}e_{r})$	q Z e e e H
$c_{Hq}^{{\scriptscriptstyle (1)}}$	$(H^{\dagger}i\overleftrightarrow{D}_{\mu}H)(\bar{q}_{p}\gamma^{\mu}q_{r})$	q Z ℓ ℓ H
$c_{Hq}^{^{(3)}}$	$(H^{\dagger}i\overleftrightarrow{D}_{\mu}^{I}H)(\bar{q}_{p}\tau^{I}\gamma^{\mu}q_{r})$	q V
c_{Hu}	$(H^{\dagger}i\overleftrightarrow{D}_{\mu}H)(\bar{u}_{p}\gamma^{\mu}u_{r})$	u U
c_{Hd}	$(H^{\dagger}i\overleftrightarrow{D}_{\mu}H)(\bar{d}_{p}\gamma^{\mu}d_{r})$	$d \sum_{\ell} Z_{\ell} \ell$

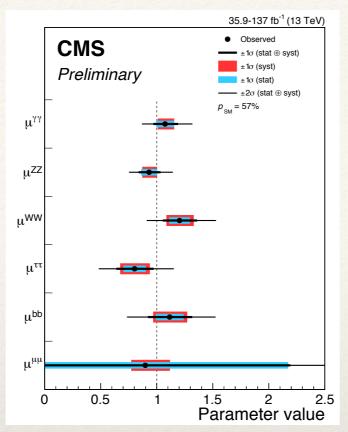
(see Ilaria Brivio's talk)

Higgs Analyses



Combined CMS Higgs boson production and decay measurements

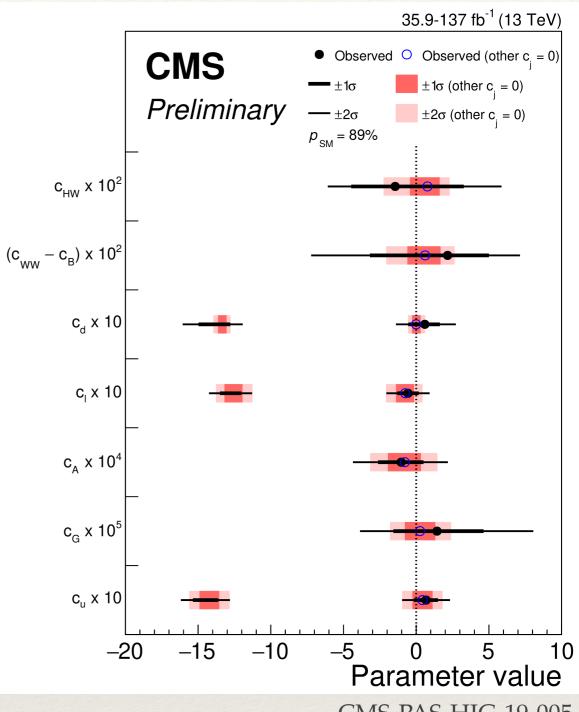




Relation between Higgs coupling modifiers in the k-framework and EFT Wilson coefficients computed from the sum of SM contribution, BSM $(1/\Lambda^4)$ and interference $(1/\Lambda^2)$ terms:

$$\mu_i(c_j) = \frac{\sigma_i^{\text{EFT}}}{\sigma_i^{\text{SM}}} = 1 + \sum_j A_j c_j + \sum_{jk} B_{jk} c_j c_k$$

Acceptance modifications and higher order missing term effects are ignored;



CMS-PAS-HIG-19-005

Simultaneous ML fit to eight coefficients leaving others fixed to SM, (C_{WW}-C_B combination in fit since C_{WW}+C_B constrained by EW data);

CMS Constraints on anomalous Higgs boson couplings to vector bosons and fermions in production and decay using four-\ell final state

Comprehensive study of CP-violation, anomalous couplings, and tensor structure of Higgs interactions with gauge bosons, gluons, and fermions, using all accessible production mechanisms and $H\rightarrow 4\ell$ decay;

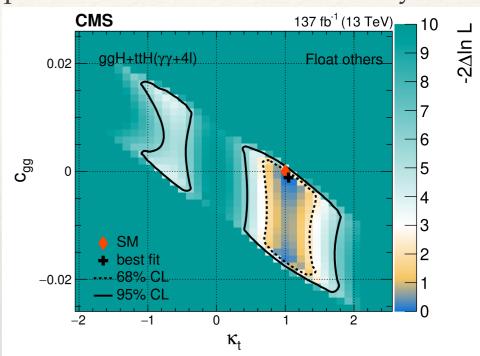
Detailed analysis of kinematics of particles associated with Higgs production in addition to decay

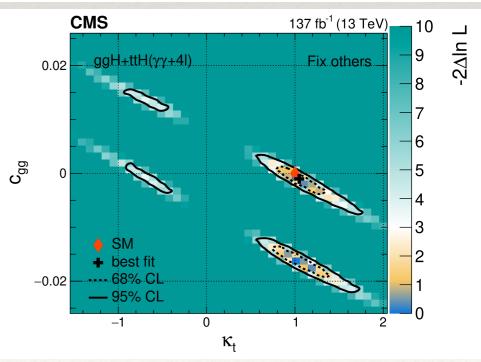
kinematic distributions

Parameterization of H boson production and decay based on scattering amplitude then connected to SMEFT formulation;

Tensor structure of HVV coupling tested with VBF and VH in $H\rightarrow VV\rightarrow 4\ell$ decay. Impose SU(2)×U(1) symmetry to relate parameters to SMEFT. Operator basis chosen as couplings of mass eigenstates: translation of SMEFT results to bosonic dim-6 operators in Warsaw basis for comparison with other results:

$\begin{array}{cccccccccccccccccccccccccccccccccccc$		Channels	Coupling	Observed	Expected
$c_{ ext{H ilde{W}B}} egin{array}{cccc} -0.25^{+0.56}_{-0.57} & 0.00^{-1.11}_{-1.21} \ c_{ ext{H ilde{B}}} & -0.06^{+0.15}_{-0.16} & 0.00^{+0.33}_{-0.33} \ \end{array}$	VBF		c _{H□} c _{HD} c _{HW} c _{HWB} c _{HW} c _{HW} c _{HW}	$0.04_{-0.45}^{+0.43} \\ -0.73_{-0.45}^{+0.97} \\ -0.73_{-4.21}^{+0.18} \\ 0.01_{-0.17}^{+0.18} \\ 0.01_{-0.18}^{+0.20} \\ 0.00_{-0.05}^{+0.05} \\ -0.23_{-0.52}^{+0.51} \\ -0.25_{-0.57}^{+0.56}$	$0.00^{+0.75}_{-0.93}$ $0.00^{+1.06}_{-4.60}$ $0.00^{+0.39}_{-0.28}$ $0.00^{+0.42}_{-0.31}$ $0.00^{+0.03}_{-0.08}$ $0.00^{+1.11}_{-1.11}$ $0.00^{+1.21}_{-1.21}$





Interpretations of ATLAS combined measurement of Higgs boson production and decay



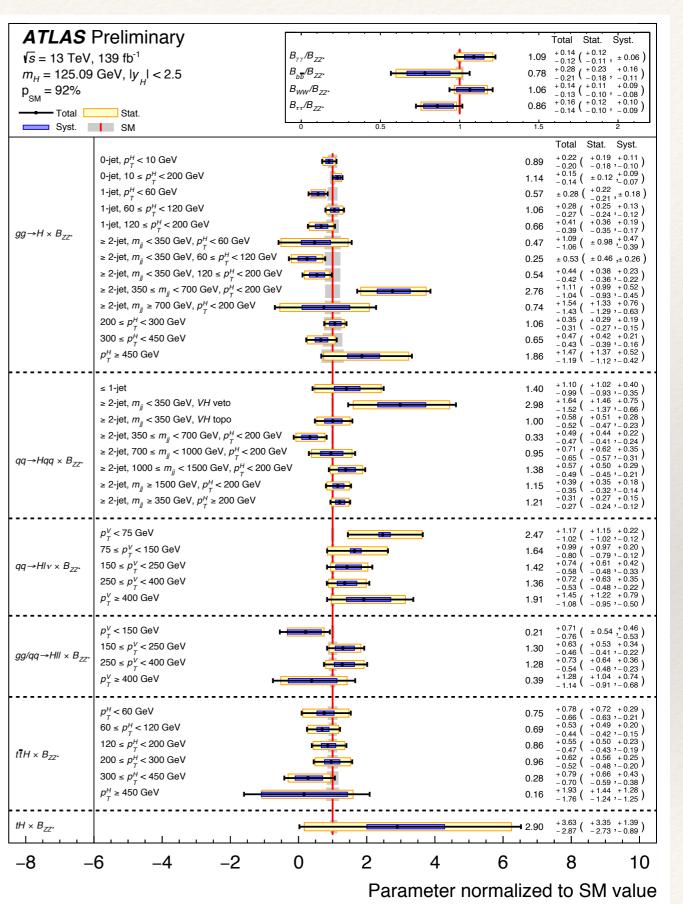
New combination of measurements of Higgs production and decay with full Run 2 stat;

Extended p_T^H reach of analyses

Measurements separated by production mode and kinematic template category (STXS) for different decay channels re-parametrised for SMEFT operators;

Measured signal strength for each STXS category used in EFT analysis to extract constraints on (combinations of) Wilson coefficients.

Probability to obtain observed data under SM-only hypothesis = 92%



Interpretations of ATLAS combined measurement of Higgs boson production and decay



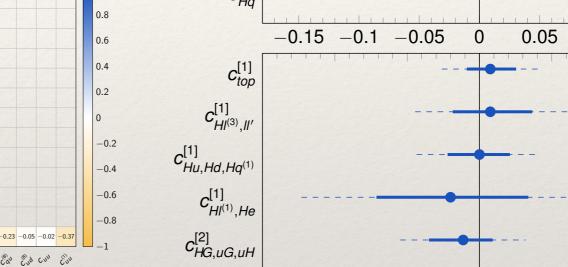
New eigenvector decomposition in sub-spaces constructed with guidance from sensitivity of individual Wilson coefficients on inputs from available measurements;

All parameters fitted simultaneously:

CHG.UG.UH (×10)

All parameters fitted simultaneously;





ATLAS Preliminary

 $m_H = 125.09 \text{ GeV}, |y_H| < 2.5$

 \sqrt{s} =13 TeV, 139 fb⁻¹

SMEFT $\Lambda = 1$ TeV

 $p_{SM} = 59\%$

Compared to previous analysis:

 Correlations between Wilson coeffs. reduced due to additional input from: H→ττ; VBF, H→bb and ttH, H→bb;

- first constraints on some combinations of four fermion top operators (c_{eH} , c_{dH} and c_{top});
- Sensitivity improved up to 70% compared to 2020 study;

 $c^{[2]}_{HW,HB,HWB,HDD,uW,uB,W}$ -3 -2 -1 0 1 2 $c^{[2]}_{Hu,Hd,Hq,Hq^{(1)}}$ $c^{[3]}_{HW,HB,HWB,HDD,uW,uB,W}$ c_{eH} c_{dH} c_{dH} c_{dH} c_{dH} c_{dH} Parameter Value

Measured parameters consistent with SM expectation within uncertainties.

68 % CL

95 % CL

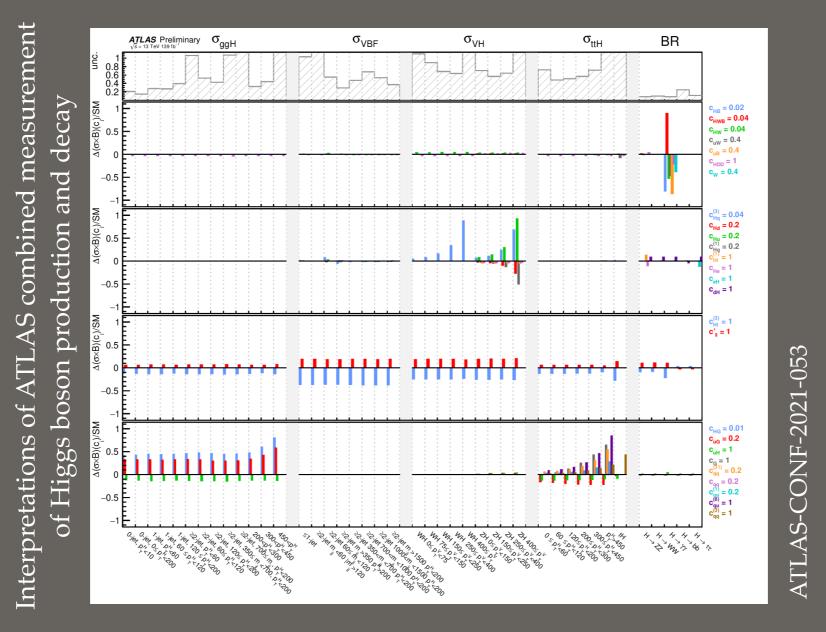
Best Fit

0.1

0.15

The Case for High p_T Higgs Measurements

While sensitivity to new physics scale Λ from static Higgs properties scales as $\sim v/\Lambda$, Higgs bosons produced at large p_T^H access regions where effects may be enhanced by powers of p_T^H/Λ .



Measurements with extended energy reach may outperform lower-energy measurements of higher precision in their BSM sensitivity;

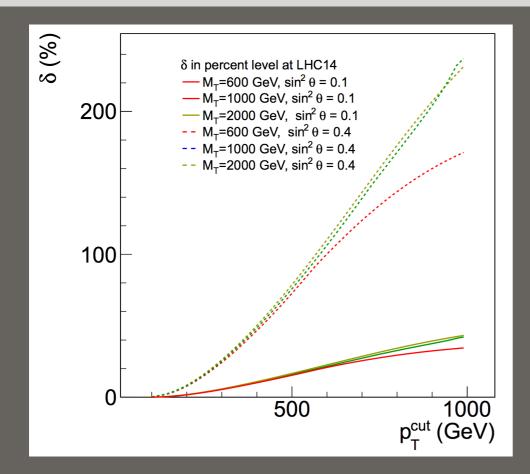
Furthermore S/B for Higgs processes typically increases with energy.

EFT and Explicit BSM Models

If new particle sufficiently heavy (M>> M_h and M> $p_T^H_{max}$) contribution can be matched to local EFT operators, scales with p_T x coupling; Increasing to $p_T^H > M$ resolves contribution from new particles yielding terms of p_T/M but EFT approximation breaks down;

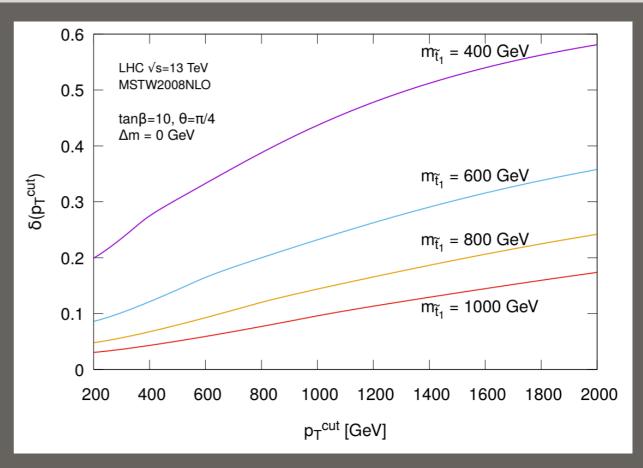
Heavy Top Partner

$$\left| \mathcal{M}_{+++} \right|_{m_t \ll p_T \ll M_T} \propto \frac{m_t^2 \,\kappa_t}{p_T} \left(A_{t,0} + A_{t,1} \, \ln \left(\frac{p_T^2}{m_t^2} \right) + A_{t,2} \, \ln^2 \left(\frac{p_T^2}{m_t^2} \right) \right) + \kappa_T \, p_T$$



SUSY Scalar Top Quarks

$$\propto \frac{m_t^2 \, \kappa_t}{p_T} \Big(A_{t,0} + A_{t,1} \, \ln \Big(\frac{p_T^2}{m_t^2} \Big) + A_{t,2} \, \ln^2 \Big(\frac{p_T^2}{m_t^2} \Big) \Big) + \kappa_T \, p_T \\ \mathcal{M}_{+++}|_{m_t \ll p_T \ll m_{\tilde{t}_1}, m_{\tilde{t}_2}} \, \frac{m_t^2}{p_T} \left(A_0 + A_1 \ln \Big(\frac{p_T^2}{m_t^2} \Big) + A_2 \ln^2 \Big(\frac{p_T^2}{m_t^2} \Big) \right) + 8 \frac{\Delta}{p_T^2} \, F_g \, R_{t,1}$$



Grojean, Salvioni, Schaffer, Weiler, JHEP 05 (2014) 022 Banfi, Martin, Sanz, JHEP 08 (2014) 053 Banfi, Bond, Martin, Sanz, JHEP 11 (2018) 171



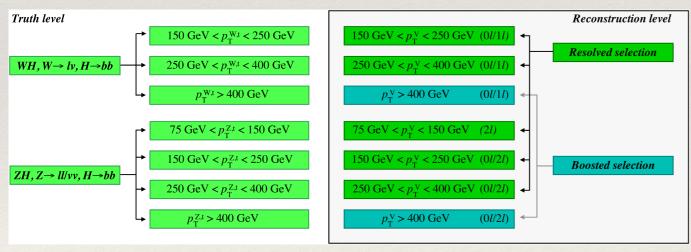
Combination of ATLAS measurements of Higgs boson production in association with W or Z boson in bb decay channel

First combination of ATLAS measurements of associated VH production with $H\rightarrow bb$; W, $Z\rightarrow \ell s$ in resolved and boosted topologies optimised for different p_T^V regimes:

- Resolved VH, H→bb with p_TV<400 GeV with Higgs boson candidate reconstructed as two resolved jets
- Boosted VH, H \rightarrow bb with $p_T^V>400$ GeV with Higgs boson candidate reconstructed as single large-radius jet.

with $Z \rightarrow \ell\ell$ and $W \rightarrow \ell\nu$

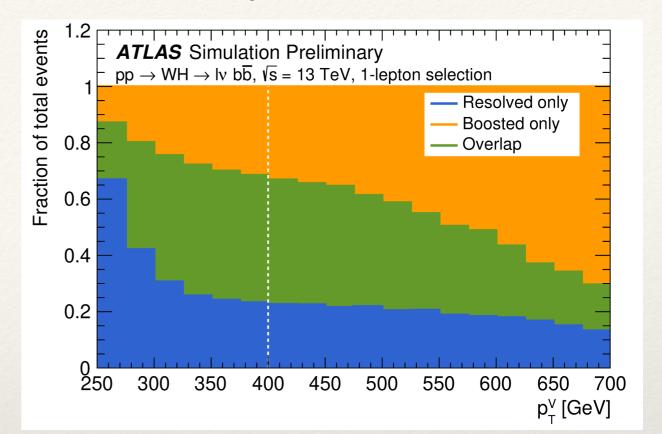
Cross-sections measured in fiducial STXS volumes:

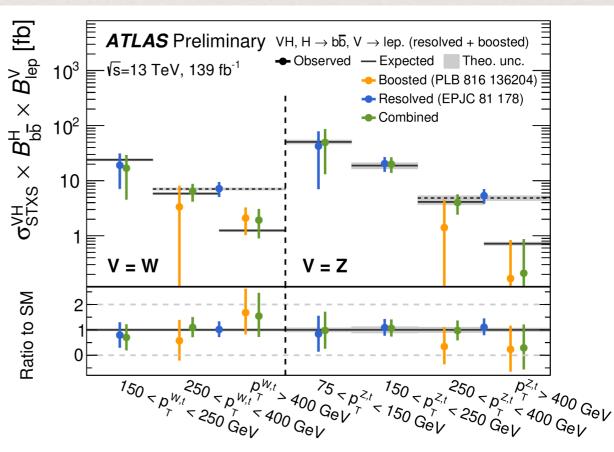


Results compatible with SM within uncertainties of 30% to 300%:

Input analyses:

Eur. Phys. J. C 81 (2021) 178, Phys. Lett. B 816 (2021) 136204





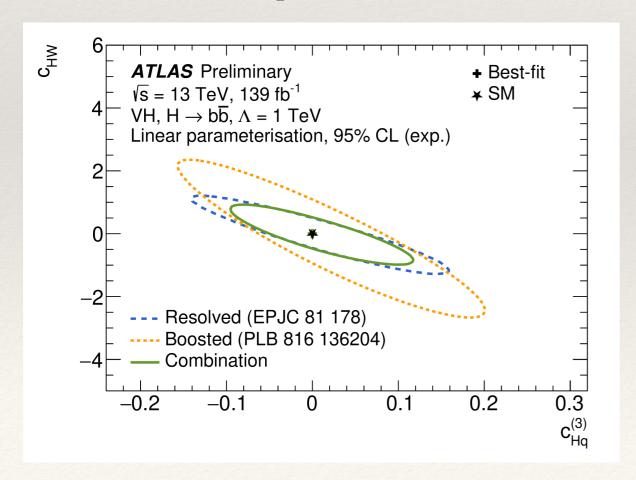


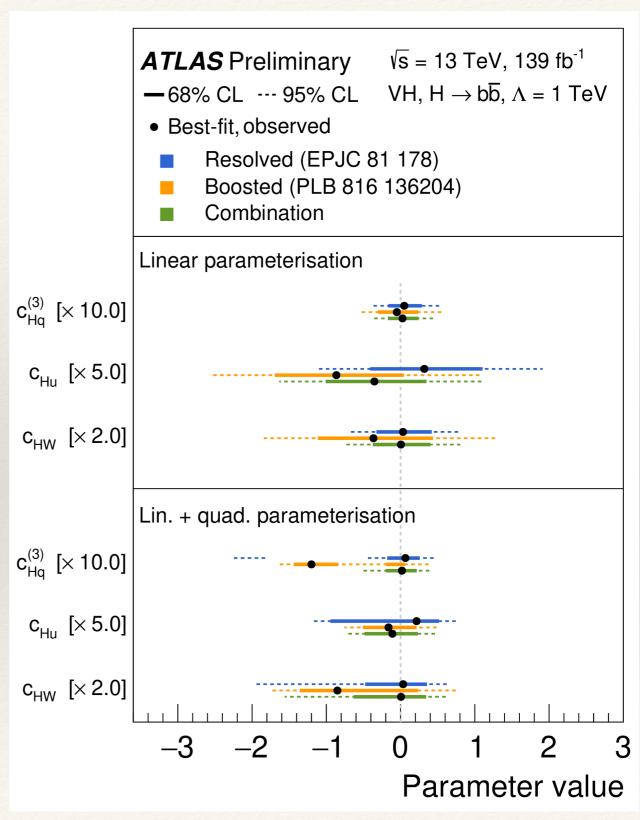
Combination of ATLAS measurements of Higgs boson production in association with W or Z boson in bb decay channel

Combination yields more stringent limits than individual analyses on EFT parameters strongly dependent on p_T^H .

VH-sensitive EFT operators pin down tensor structure of HVV couplings;

Study constraints on Wilson coefficients using linear and linear + quadratic EFT terms:



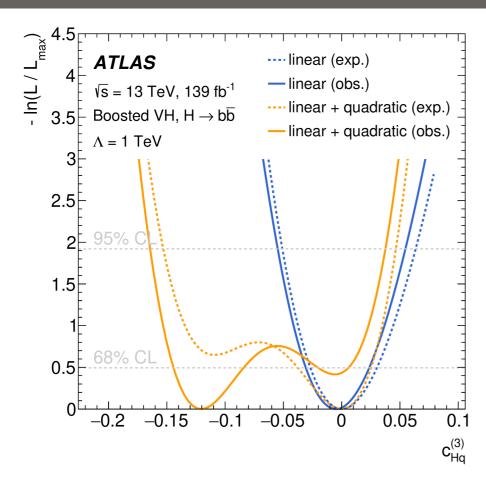


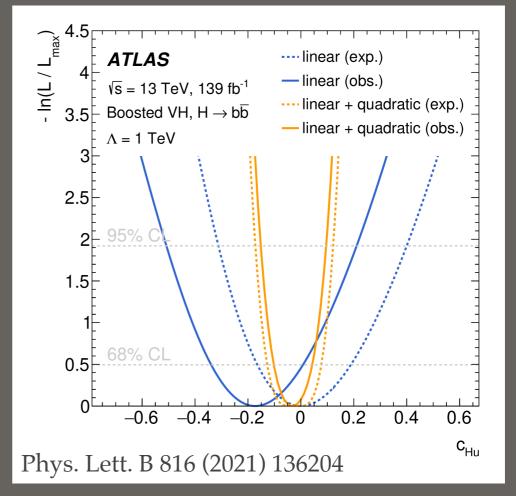
Linear and Quadratic EFT terms

Whether quadratic terms should be included or not in dim-6 SMEFT analysis an open question: in SMEFT expansion, Λ^4 term not complete with dim-6 operators: difference of linear vs linear +quadratic may offer us way to estimate impact of dim-8 operators, whose precise contributions is unknown but has been estimated to be small;

Observe that inclusion can change fit results opening additional minima or shrinking allowed region for Wilson coefficients of some operators;

ATLAS Measurement of the associated production of a Higgs boson decaying into b-quarks with a vector boson at high transverse momentum





Answer may also be related to decoupling behaviour of underlying UV model as in case of non-decoupling models (e.g. heavy top partner) and decoupling models (e.g. MSSM with light scalar top and bottom quarks). 13

CMS Inclusive Search for Highly Boosted H → bb

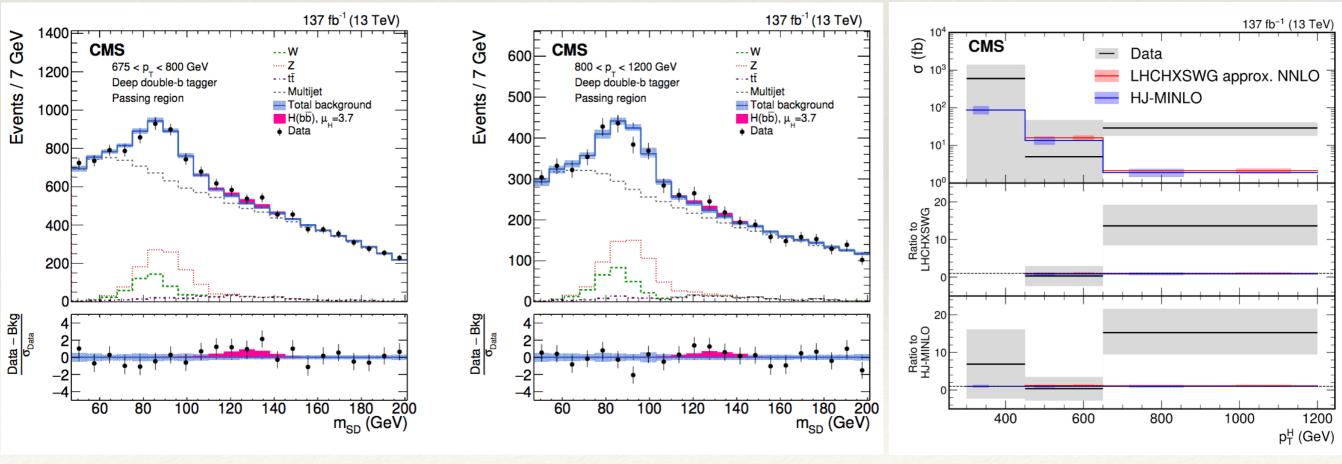
Pioneering study of boosted H \rightarrow bb reconstructed as single large-radius jets with p_T>450 GeV:

- Signal selected with jet substructure and dedicated b-tagging technique based on deep NN.
- Method validated using $Z \rightarrow bb$ decays.

Excess of events above the background assuming no 125 GeV Higgs boson production observed with local significance of 2.5 σ (expectation 0.7):

• Signal strength $\mu_{\rm H}$ (p_T>450 GeV) = 3.7 ± 1.2(stat) +0.8/-0.7 (syst) +0.8/-0.5 (theo).

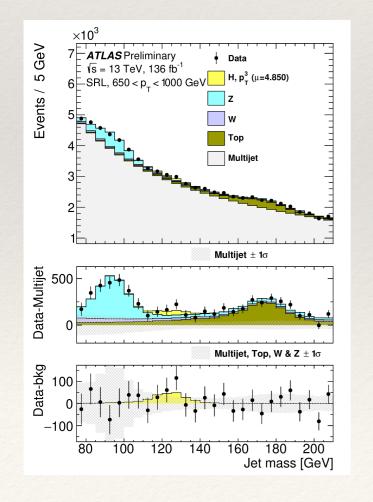
Unfolded differential ggF cross section in p_T^H bins assuming other production modes at SM rates.

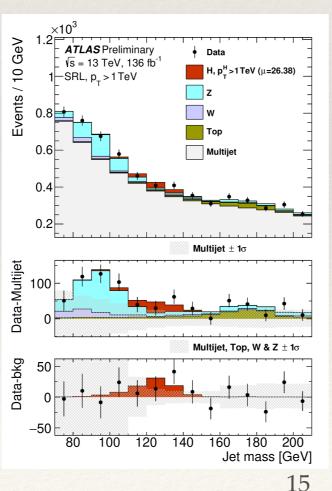


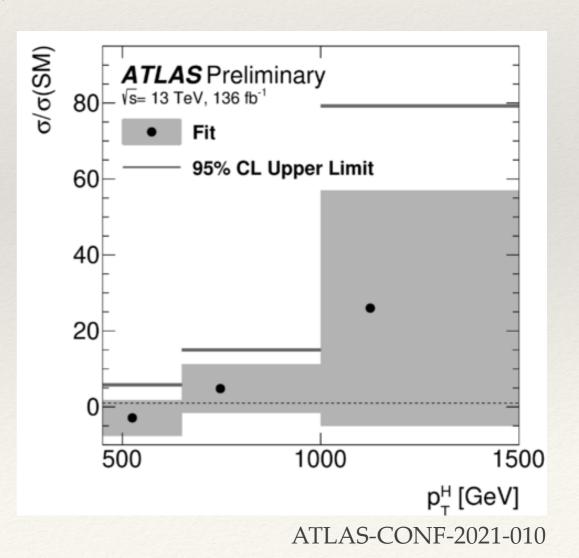
ATLAS Study of Higgs-boson production with large transverse momentum using H→bb

First ATLAS study of Higgs bosons production above 1 TeV in all-hadronic bb final state;

- $H \rightarrow bb$ reconstructed as single large-radius jets, with double b-tag of VR jets+recoil hadronic system.
- Experimental techniques validated in same kinematic regime using anti-tagged Validation Region and double b-tagged $Z \rightarrow bb$;
- Top bkg constrained in dedicated Control Region, W/Z mass resolution measured vs $p_{T;}$ Inclusive analysis with no restrictions to select particular Higgs boson production mode aside from requiring energetic hadronic recoil system.
- Simultaneous fit of V+jets, QCD yield+mass spectrum shape and inclusive $H \rightarrow bb$ yield in multiple STXS-like p_T^H bins:
- Signal strength $\mu_{\rm H}$ (p_T>450 GeV) = 0.7 ± 2.8 (stat) ± 1.7 (syst)



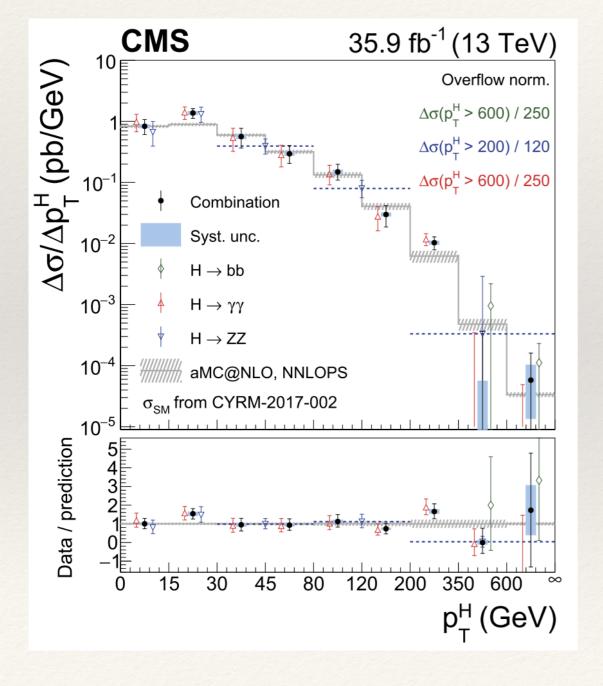


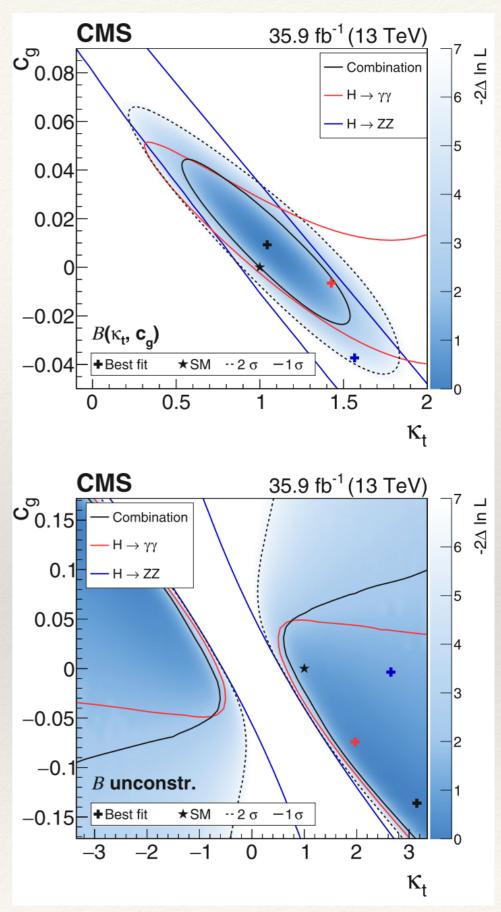


Interpretation of CMS differential cross sections for Higgs boson production and cg

Combined $H \rightarrow \gamma \gamma$, $H \rightarrow ZZ$, and $H \rightarrow$ bb (resolved and preliminary first boosted analysis) decay channels;

Higgs p_T spectrum (including high pT bin from boosted $H\rightarrow bb$) used to set limits on couplings to top, bottom, and charm quarks, as well as c_g coupling to gluon field.



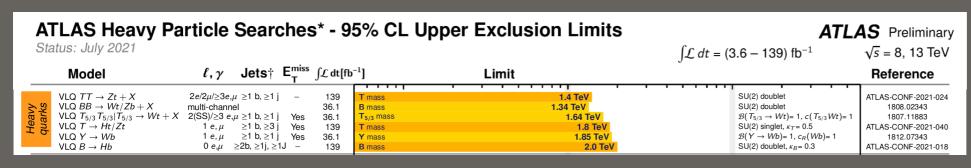


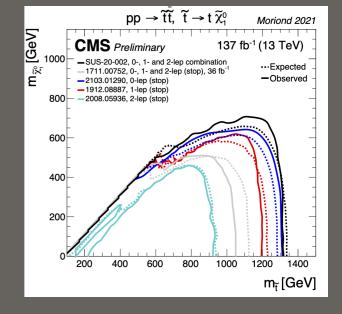
Phys. Lett. B 792 (2019) 369

Experimental Sensitivity and the Limit of EFT validity

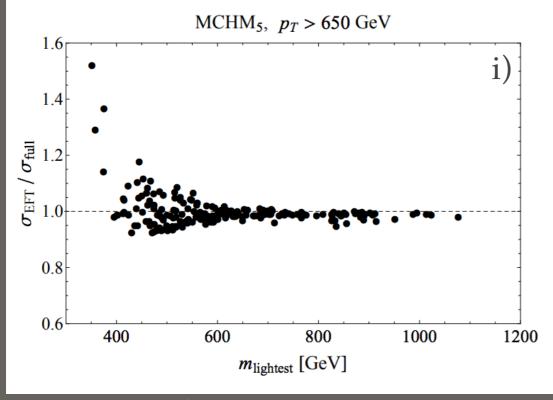
EFT explicitly assumes that BSM degrees of freedom can be integrated out;

Validity of this assumption may become questionable as ATLAS and CMS extend sensitivity at (and beyond) 1 TeV, at scale comparable to (or higher than) direct search limits for relevant new particles:

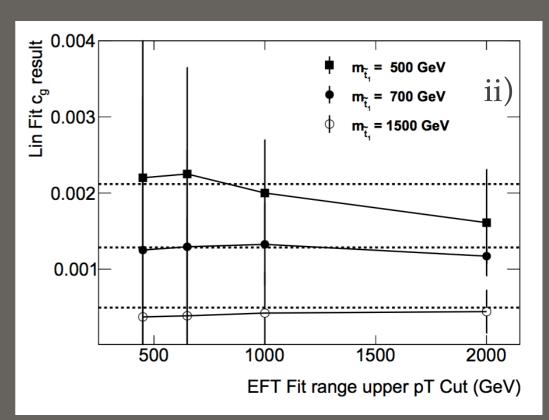




EFT break down observed from mismatch between i) Higgs cross section in EFT matching approximation and full theory and ii) Wilson coefficients extracted by EFT matching and theory:



Top Partner



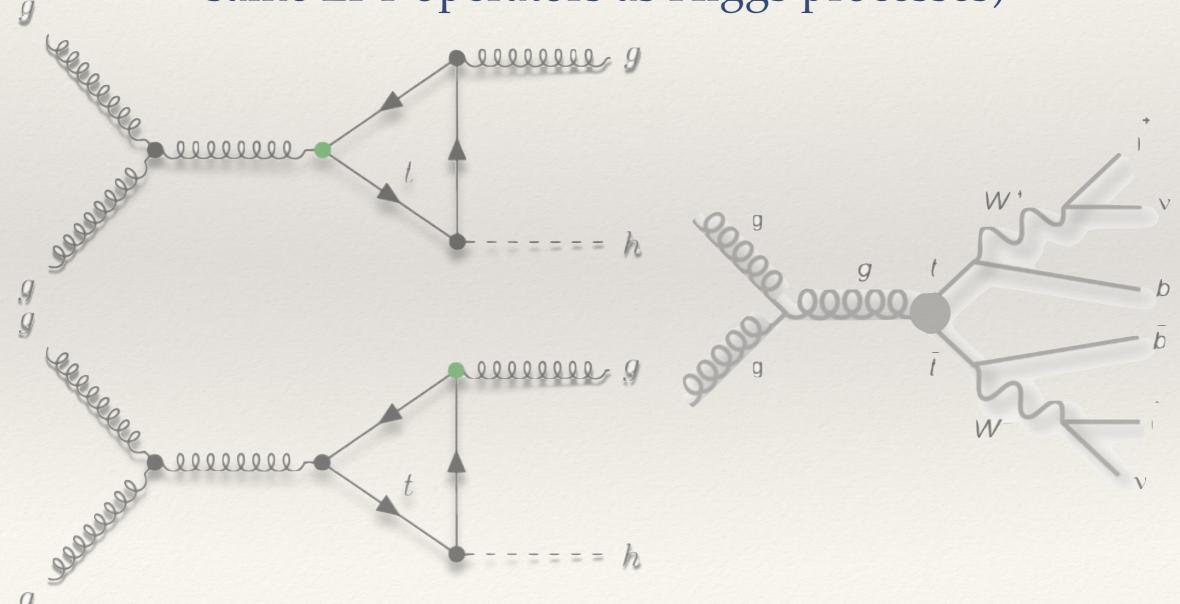
SUSY Scalar Top Quarks

Grojean, Salvioni, Schaffer, Weiler, JHEP 05 (2014) 022

MB, Grazzini, Spira, Wiesemann, arXiv:2109.02987

... and beyond

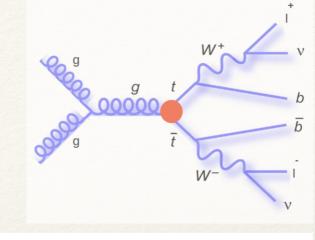
(some analysis of Top and EW physics sensitive to same EFT operators as Higgs processes)

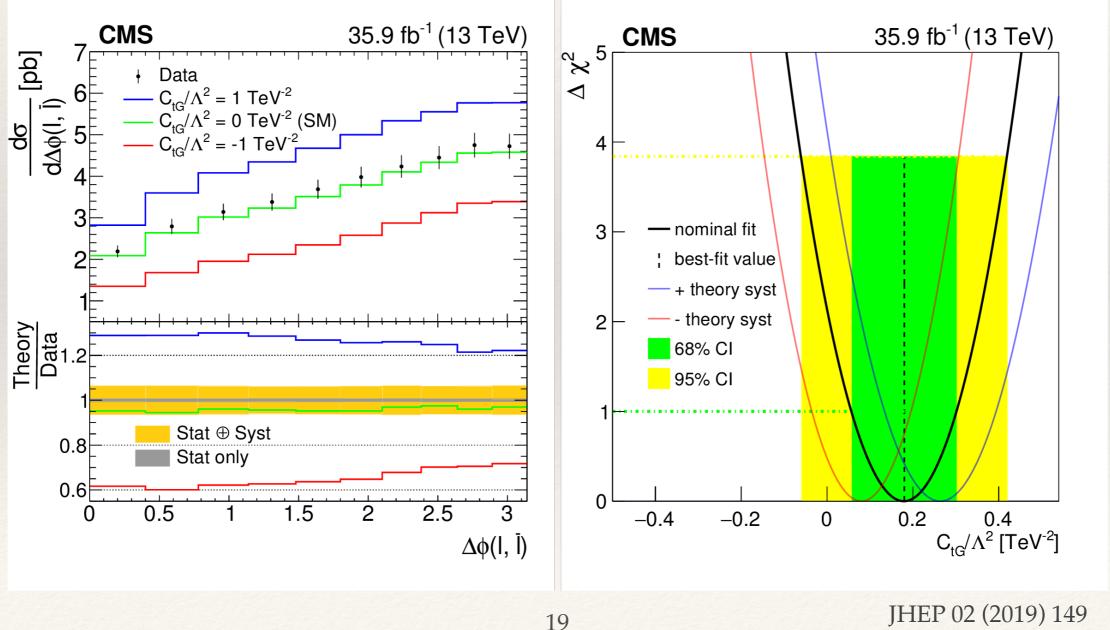


CMS Extractions of Chromo-magnetic operator Wilson coefficient constraints from measurements of ttbar differential cross sections using events containing two leptons

Chromo-magnetic dipole operator O_{tG} modifies gtt vertex inducing effects on ttbar production rate and $\Delta \varphi(ll)$ for azimuthal angle between ℓs from di-leptonic ttbar events;

Probes interplay between top and Higgs sectors;



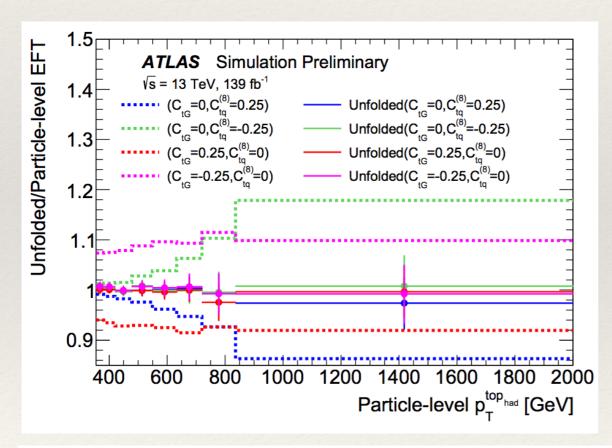


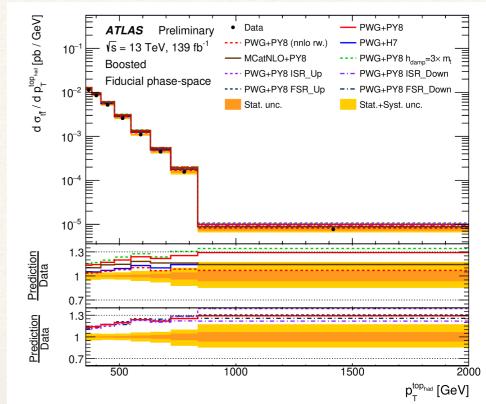
ATLAS Measurements of differential cross-sections in top-quark pair events with a high transverse momentum top quark and limits on BSM contributions to top-quark pair production

Precise cross-section measurements of ttbar events containing high-p_T hadronically decaying top:

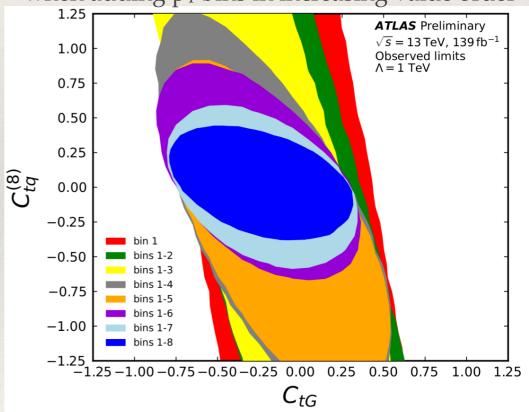
- Distributions unfolded to particle level.
- Precision significantly improved by use of invariant mass of hadronic top decay jet to reduce impact of JES uncertainties.

Probes chromo-magnetic dipole + four-fermion top operators:





Evolution of 95% C.L. Wilson coeff. region when adding p_T bins in increasing value order



Wilson coefficient	Marginalised 95% intervals		Individual 95% intervals		
wilson coemcient	Expected	Observed	Expected	Observed	Global fit [101]
C_{tG}	[-0.44, 0.44]	[-0.68, 0.21]	[-0.41, 0.42]	[-0.63, 0.20]	[0.007, 0.111]
$C_{tq}^{(8)}$	[-0.35, 0.35]	[-0.30, 0.36]	[-0.35, 0.36]	[-0.34, 0.27]	[-0.40, 0.61]

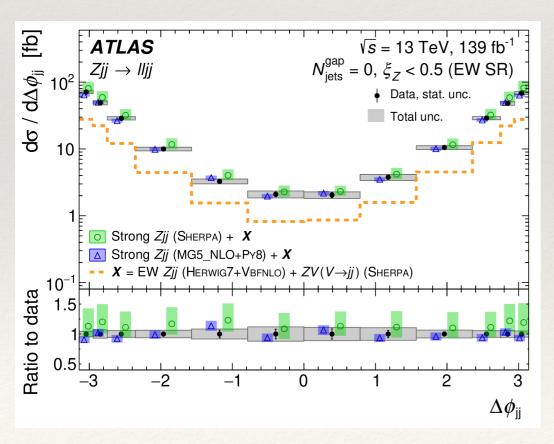
ATLAS-CONF-2021-031

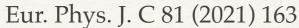
See also ATLAS-CONF-2021-050

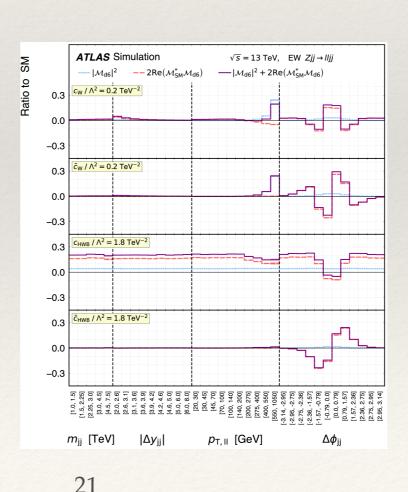
Differential ATLAS cross-section measurements for electro-weak production of dijets in association with Z boson

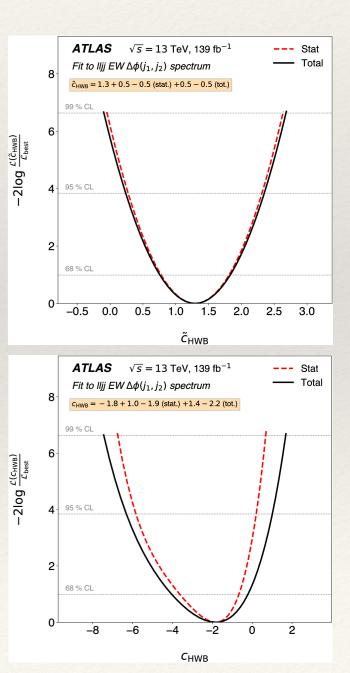
Differential cross-section measurements for EW production of two jets + Z; Measurements sensitive to vector-boson fusion production mechanism and provide a fundamental test of SM gauge structure;

Differential cross-sections used in search for anomalous weak-boson self-interactions bosonic dim-6 EFT operators. Signed azimuthal angle between two jets sensitive to interference between the SM and dim-6 scattering amplitudes.



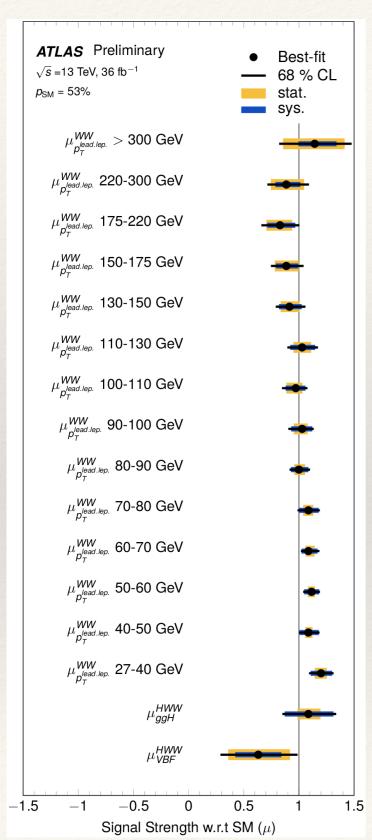






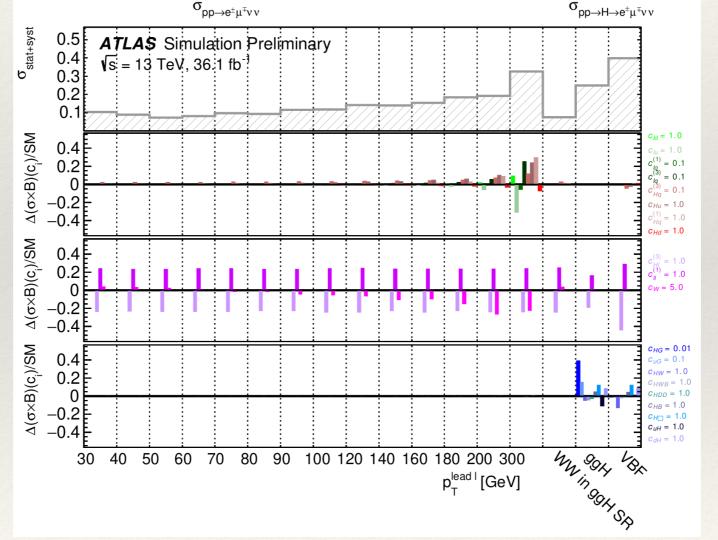
Combined effective field theory interpretation of ATLAS H→WW* and WW measurements

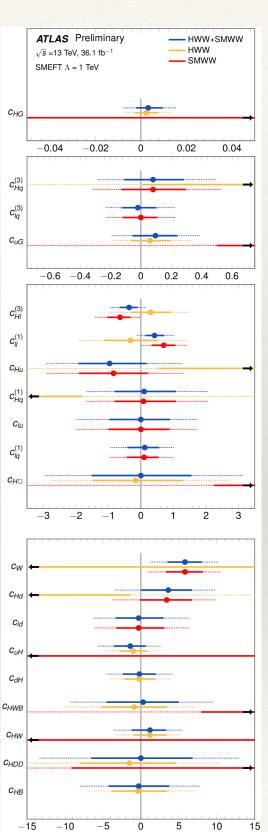
EFT constraints by combination of $H\rightarrow WW^*$ in ggF and vector fusion production with differential results on WW di-boson production in 0-jet channel.



Likelihood function from μ of Higgs measurement together with unfolded differential WW cross-sections;

Constrain 22 Wilson coeffs. of bosonic and two-fermion operators in SMEFT framework.

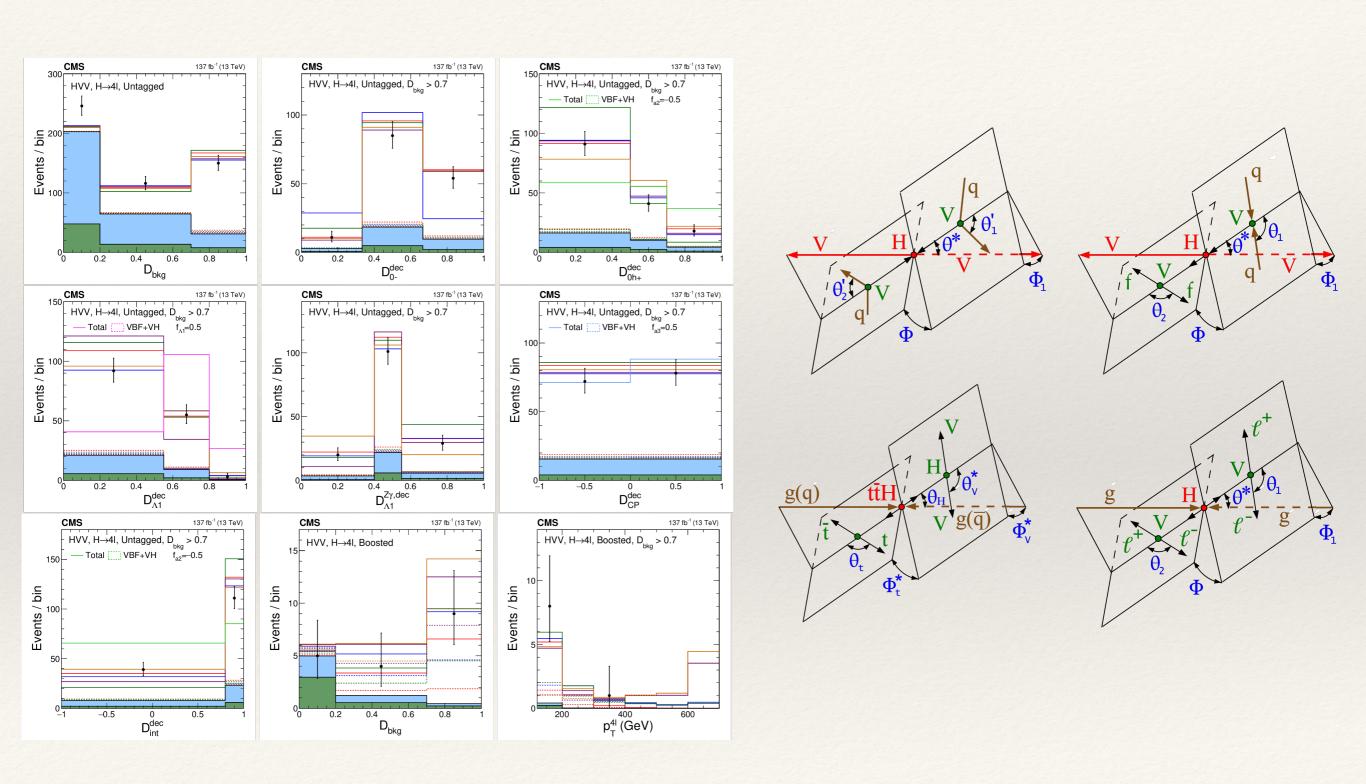






CMS Constraints on anomalous Higgs boson couplings to vector bosons and fermions in production and decay using four-\ell final state

Discriminants from matrix element calculations



CMS Constraints on anomalous Higgs boson couplings to vector bosons and fermions in production and decay using four-\ell final state

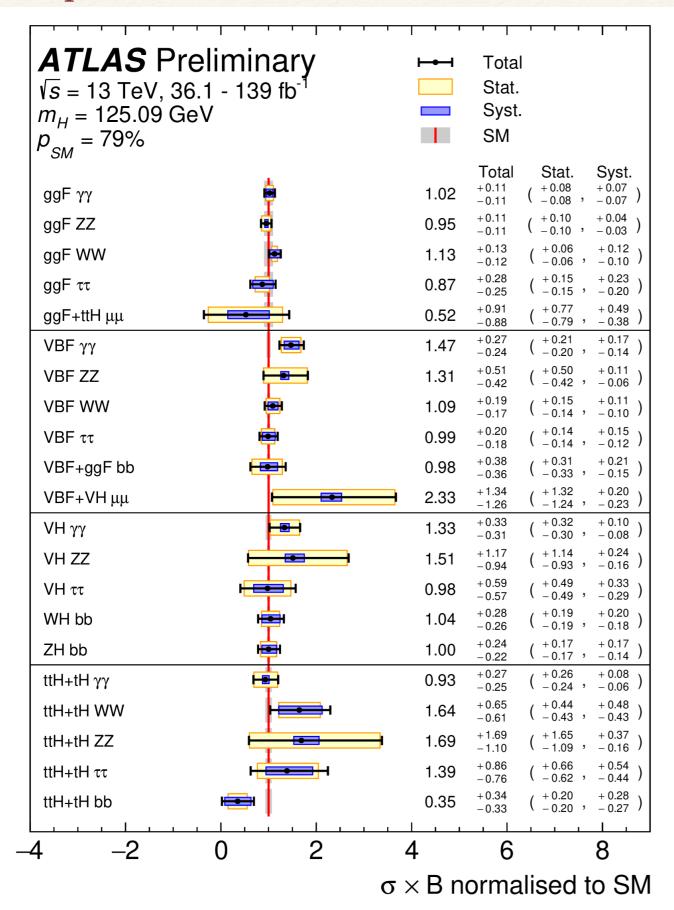
$$\begin{split} A(\mathrm{HVV}) &= \frac{1}{v} \left[a_1^{\mathrm{VV}} + \frac{\kappa_1^{\mathrm{VV}} q_{\mathrm{V1}}^2 + \kappa_2^{\mathrm{VV}} q_{\mathrm{V2}}^2}{\left(\Lambda_1^{\mathrm{VV}}\right)^2} + \frac{\kappa_3^{\mathrm{VV}} (q_{\mathrm{V1}} + q_{\mathrm{V2}})^2}{\left(\Lambda_Q^{\mathrm{VV}}\right)^2} \right] m_{\mathrm{V1}}^2 \epsilon_{\mathrm{V1}}^* \epsilon_{\mathrm{V2}}^* \\ &\quad + \frac{1}{v} a_2^{\mathrm{VV}} f_{\mu\nu}^{*(1)} f^{*(2),\mu\nu} + \frac{1}{v} a_3^{\mathrm{VV}} f_{\mu\nu}^{*(1)} \tilde{f}^{*(2),\mu\nu} \,, \end{split}$$

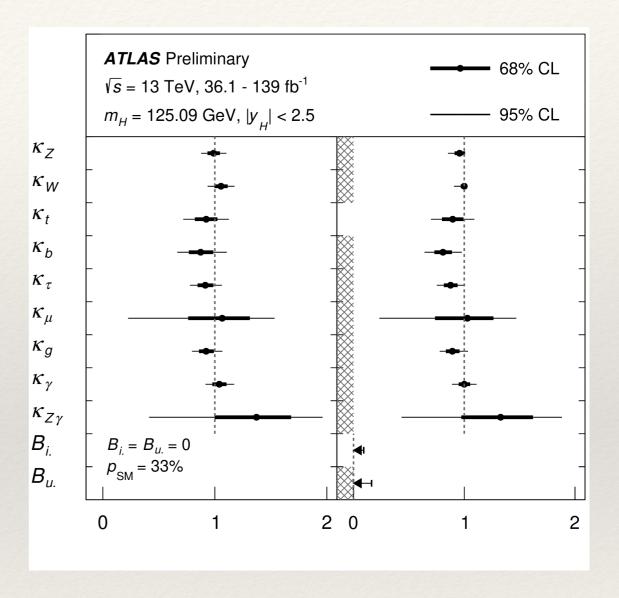
$$\begin{array}{rcl} a_1^{\rm WW} & = & a_1^{\rm ZZ} \,, \\ a_2^{\rm WW} & = & c_w^2 a_2^{\rm ZZ} + s_w^2 a_2^{\gamma\gamma} + 2 s_w c_w a_2^{\rm Z\gamma} \,, \\ a_3^{\rm WW} & = & c_w^2 a_3^{\rm ZZ} + s_w^2 a_3^{\gamma\gamma} + 2 s_w c_w a_3^{\rm Z\gamma} \,, \\ \frac{\kappa_1^{\rm WW}}{(\Lambda_1^{\rm WW})^2} (c_w^2 - s_w^2) & = & \frac{\kappa_1^{\rm ZZ}}{(\Lambda_1^{\rm ZZ})^2} + 2 s_w^2 \frac{a_2^{\gamma\gamma} - a_2^{\rm ZZ}}{m_Z^2} + 2 \frac{s_w}{c_w} (c_w^2 - s_w^2) \frac{a_2^{\rm Z\gamma}}{m_Z^2} \,, \\ \frac{\kappa_2^{\rm Z\gamma}}{(\Lambda_1^{\rm Z\gamma})^2} (c_w^2 - s_w^2) & = & 2 s_w c_w \left(\frac{\kappa_1^{\rm ZZ}}{(\Lambda_1^{\rm ZZ})^2} + \frac{a_2^{\gamma\gamma} - a_2^{\rm ZZ}}{m_Z^2} \right) + 2 (c_w^2 - s_w^2) \frac{a_2^{\rm Z\gamma}}{m_Z^2} \,, \end{array}$$

$$\delta c_z = rac{1}{2} a_1 - 1 \, , \ c_{z\square} = rac{m_Z^2 s_w^2}{e^2} rac{\kappa_1}{(\Lambda_1)^2} \, , \ c_{zz} = -rac{2 s_w^2 c_w^2}{e^2} a_2 \, , \ ilde c_{zz} = -rac{2 s_w^2 c_w^2}{e^2} a_3 \, .$$

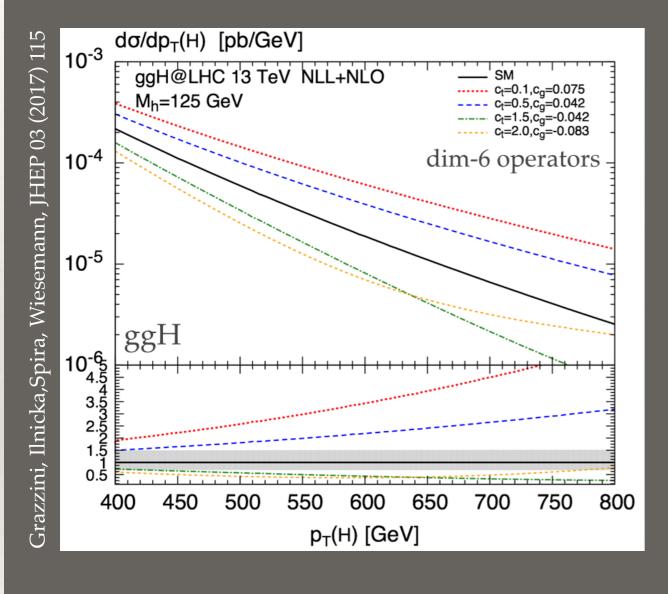
$$c_{gg}=-rac{1}{2\pilpha_s}a_2^{
m gg}$$
 , $au_{gg}=-rac{1}{2\pilpha_s}a_3^{
m gg}$.

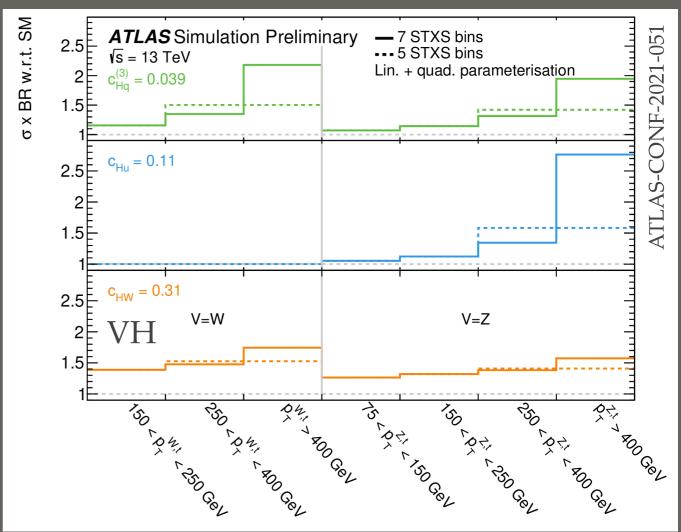
Interpretations of ATLAS combined measurement of Higgs boson production and decay



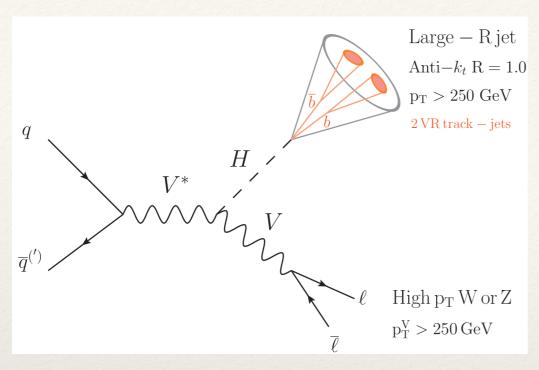


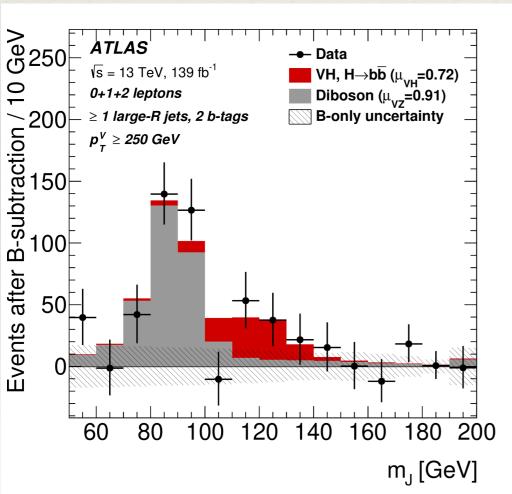
The Case for High p_T Higgs Measurements

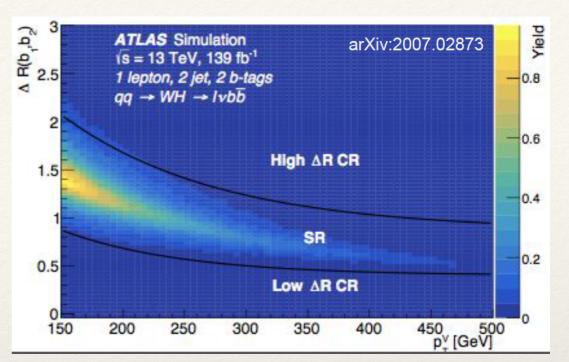


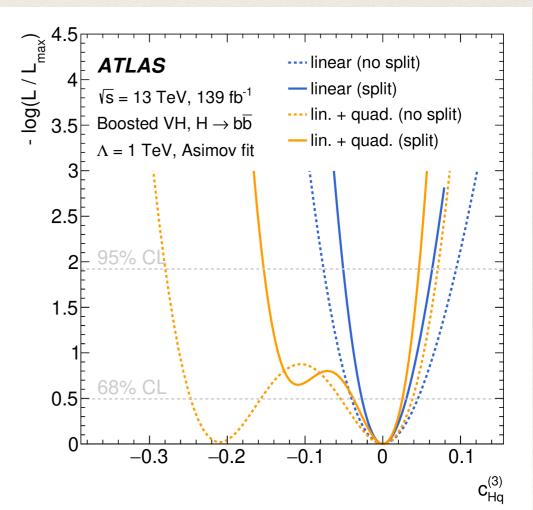


ATLAS Measurement of the associated production of a Higgs boson decaying to b quarks with a vector boson at high transverse momentum

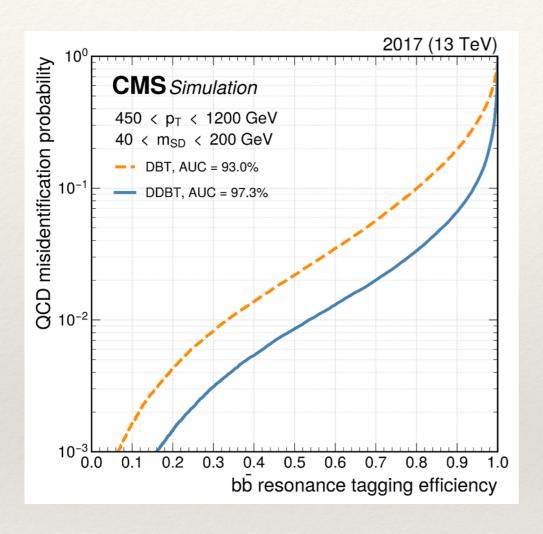


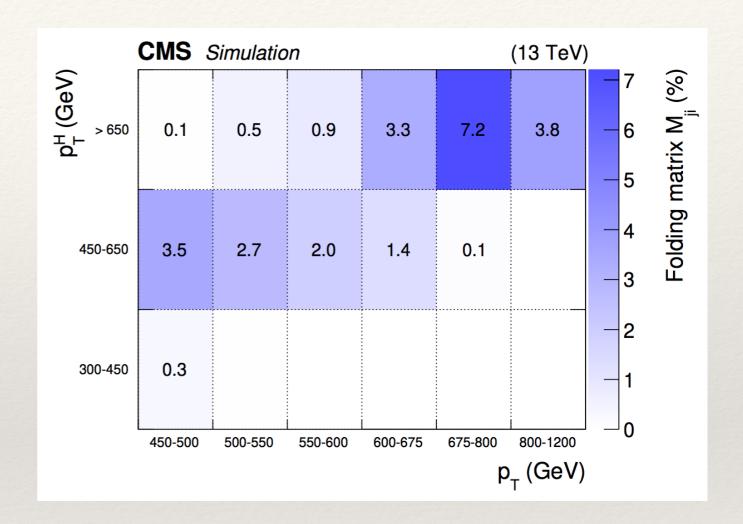




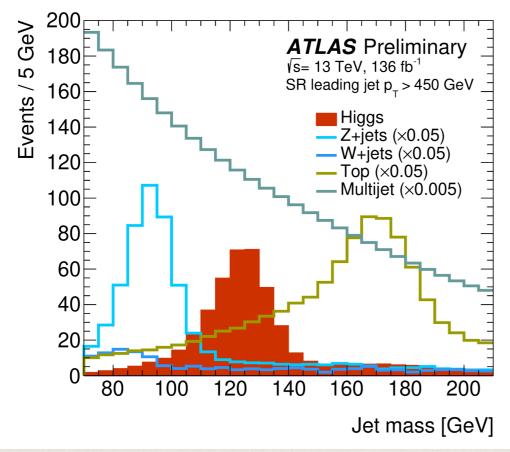


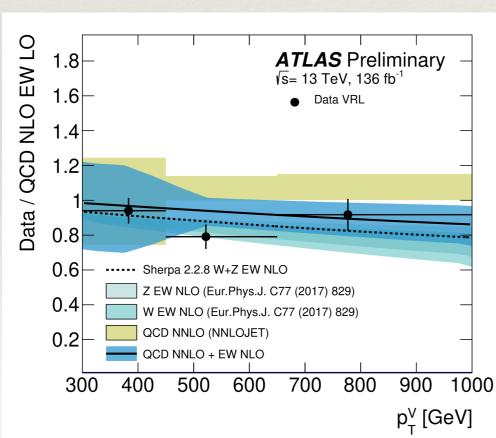
CMS Inclusive Search for Highly Boosted H → bb

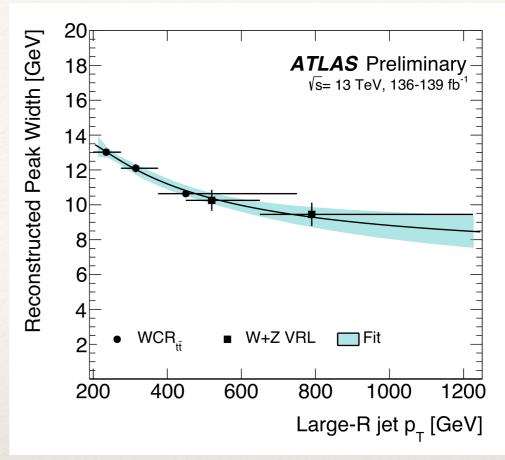


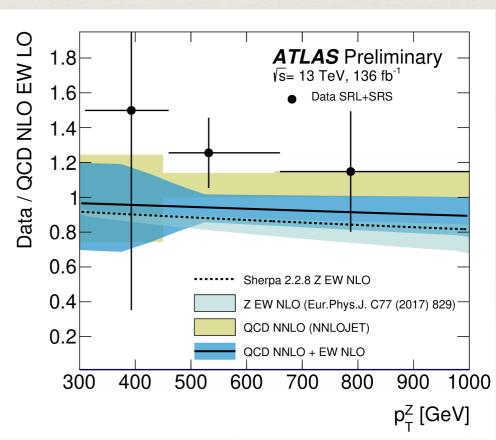


ATLAS Study of Higgs-boson production with large transverse momentum using H→bb









Combined effective field theory interpretation of ATLAS H→WW* and WW measurements

