

Simulation of low energy nuclear recoils using Geant4 [heavy ions, $T < 1 \text{ keV/amu}$]

Ricardo Pinho
ricardo@lipc.fis.uc.pt

12th Geant4
Collaboration
Workshop



LIP – Coimbra

Hebden Bridge
14 Set 2007

Outline

1. a review of Geant4's models of stopping power
2. nuclear recoils full cascade
3. nuclear quenching estimates
4. summary and present work

Versions used: geant4.8.2 and geant4.8.2.p01

Stopping power

ion = projectile particle Z_1

atom = target particle Z_2

Stopping power

ion = projectile particle Z_1

G4hLowEnergyIonisation

target particle Z_2

G4hBetheBlochModel
G4hParametrisedLossModel
G4hIonEffChargeSquare

$$-\frac{dE}{dx} = S_e + S_n$$

(assuming no correlations)

G4hNuclearStoppingModel

Stopping power

ion = projectile particle Z_1

G4hLowEnergyIonisation

target particle Z_2

G4hBetheBlochModel
G4hParametrisedLossModel
G4hIonEffChargeSquare

$$-\frac{dE}{dx} = S_e + S_n$$

(assuming no correlations)

G4hNuclearStoppingModel

- Low energy extension
- Physics list: only
G4hLowEnergyIonisation
- Very small step sizes (~1% of the particle range)
- Continuous slowing down approximation:
- Cut energy very high: no delta-rays
- Actually no secondaries what so ever! Just follow our primary till it stops (default MinKineticEnergy = 10eV)
- No fluctuations

Stopping power

ion = projectile particle Z_1

G4hLowEnergyIonisation

target particle Z_2

G4hBetheBlochModel
G4hParametrisedLossModel
G4hIonEffChargeSquare

$$-\frac{dE}{dx} = S_e + S_n$$

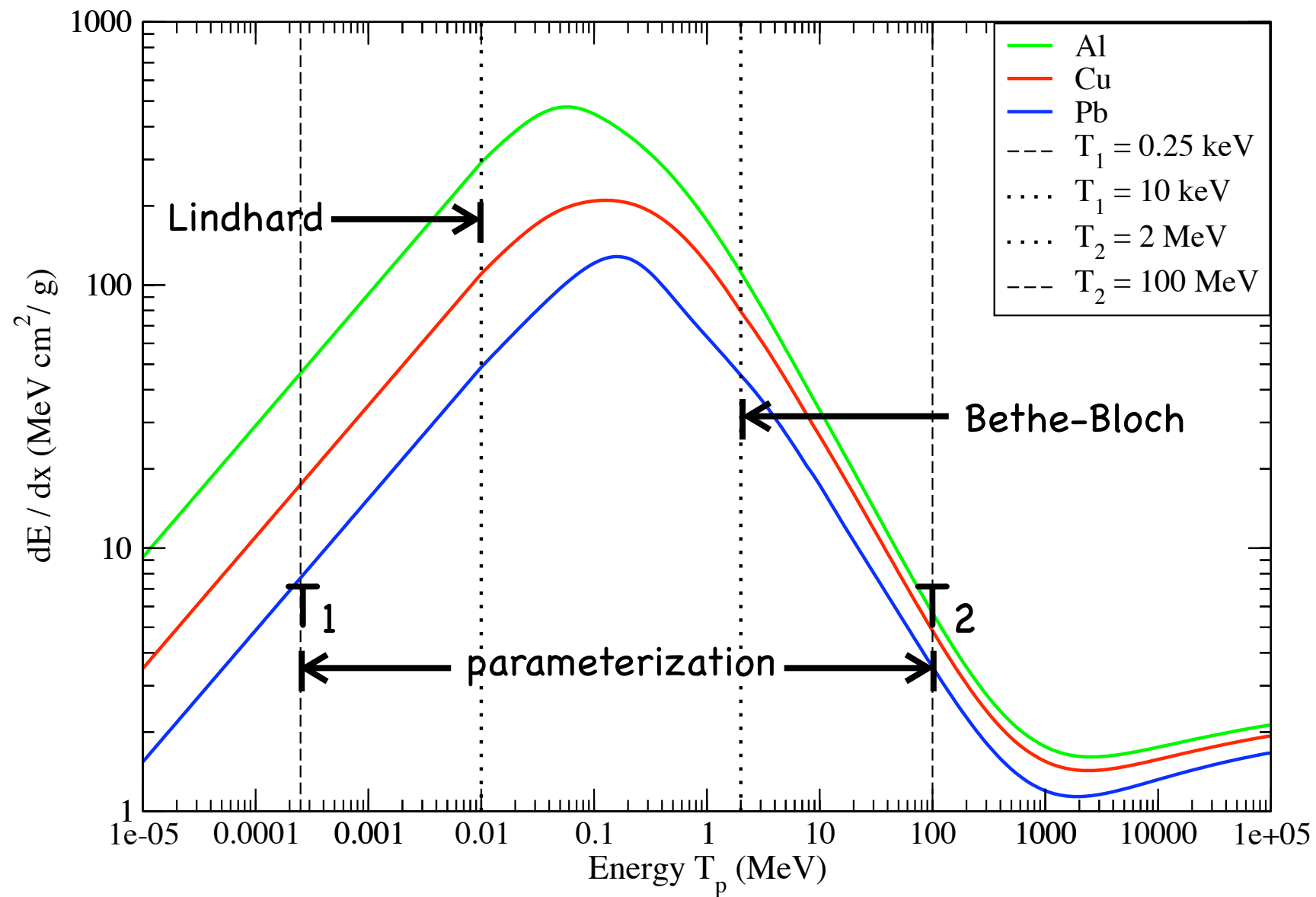
(assuming no correlations)

G4hNuclearStoppingModel

- Low energy extension
- Physics list: only
G4hLowEnergyIonisation
- Very small step sizes (~1% of the particle range)
- Continuous slowing down approximation:
- Geometry:
 - just a square box, filled with pure material
- Cut energy very high: no delta-rays
- Actually no secondaries what so ever! Just follow our primary till it stops (default MinKineticEnergy = 10eV)
- No fluctuations
- shoot primary from the center of the box, random direction

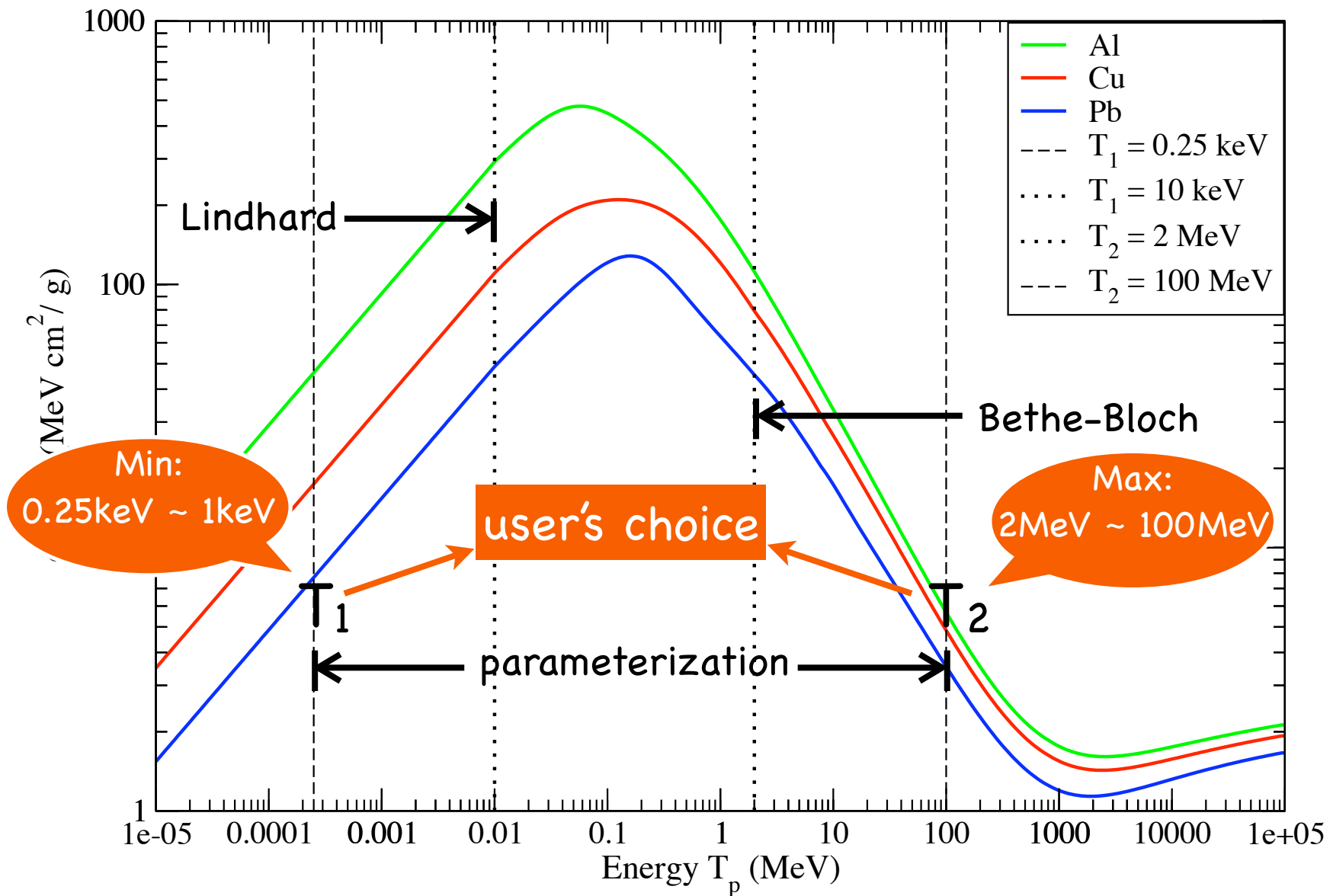
S_e : proton \rightarrow Al, Cu, Pb

GEANT4 (Lindhard+ICRU_R49p+Bethe-Bloch)



S_e : proton \rightarrow Al, Cu, Pb

GEANT4 (Lindhard+ICRU_R49p+Bethe-Bloch)



Geant4

Electronic stopping

- low energy stopping power is parameterized

Geant4

Electronic stopping

- low energy stopping power is parameterized
- parameterization is for protons and alphas only

Geant4

Electronic stopping

- ➔ low energy stopping power is parameterized
- ➔ parameterization is for protons and alphas only
- ➔ for other heavy ions, the stopping is scaled by:

Geant4

Electronic stopping

- low energy stopping power is parameterized
- parameterization is for protons and alphas only
- for other heavy ions, the stopping is scaled by:

$$S_{ei}(T) = Z_{eff}^2 \cdot S_{ep}(T_p), \quad T_p = \frac{M_p}{M} T$$

Geant4

Electronic stopping

- ➔ low energy stopping power is parameterized
- ➔ parameterization is for protons and alphas only
- ➔ for other heavy ions, the stopping is scaled by:

$$S_{ei}(T) = Z_{eff}^2 \cdot S_{ep}(T_p), \quad T_p = \frac{M_p}{M} T$$

proton stopping
parameterizations



- (default) • ICRU_R49p
- Ziegler1977p
- Ziegler1985p
- SRIM2000p

Geant4

Electronic stopping

- ➔ low energy stopping power is parameterized
- ➔ parameterization is for protons and alphas only
- ➔ for other heavy ions, the stopping is scaled by:

$$S_{ei}(T) = Z_{eff}^2 \cdot S_{ep}(T_p), \quad T_p = \frac{M_p}{M} T$$

proton stopping
parameterizations

alpha stopping
parameterizations

- (default) • ICRU_R49p
- Ziegler1977p
 - Ziegler1985p
 - SRIM2000p

- ICRU_R49He
- Ziegler1977He

Geant4

Electronic stopping

- ➔ low energy stopping power is parameterized
- ➔ parameterization is for protons and alphas only
- ➔ for other heavy ions, the stopping is scaled by:

$$S_{ei}(T) = Z_{eff}^2 \cdot S_{ep}(T_p), \quad T_p = \frac{M_p}{M} T$$

proton stopping
parameterizations

alpha stopping
parameterizations

effective charge
parameterization

$$Z_{eff}(v, v_F)$$

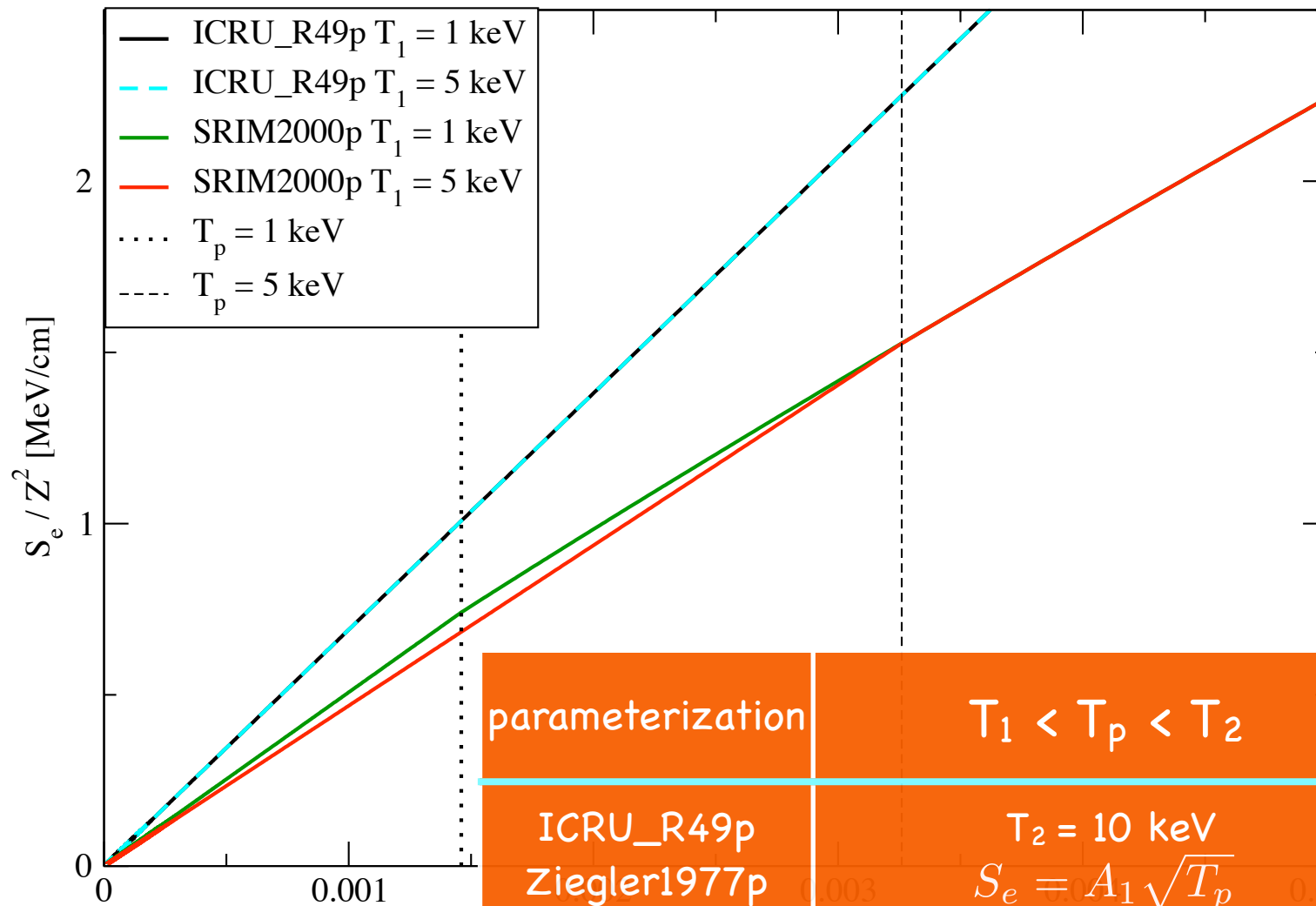
- (default) • ICRU_R49p
- Ziegler1977p
 - Ziegler1985p
 - SRIM2000p

- ICRU_R49He
- Ziegler1977He

- Ziegler, 1985,
Brandt and
Kitagawa, 1982

$S_e: \text{Xe} \rightarrow \text{LXe}$

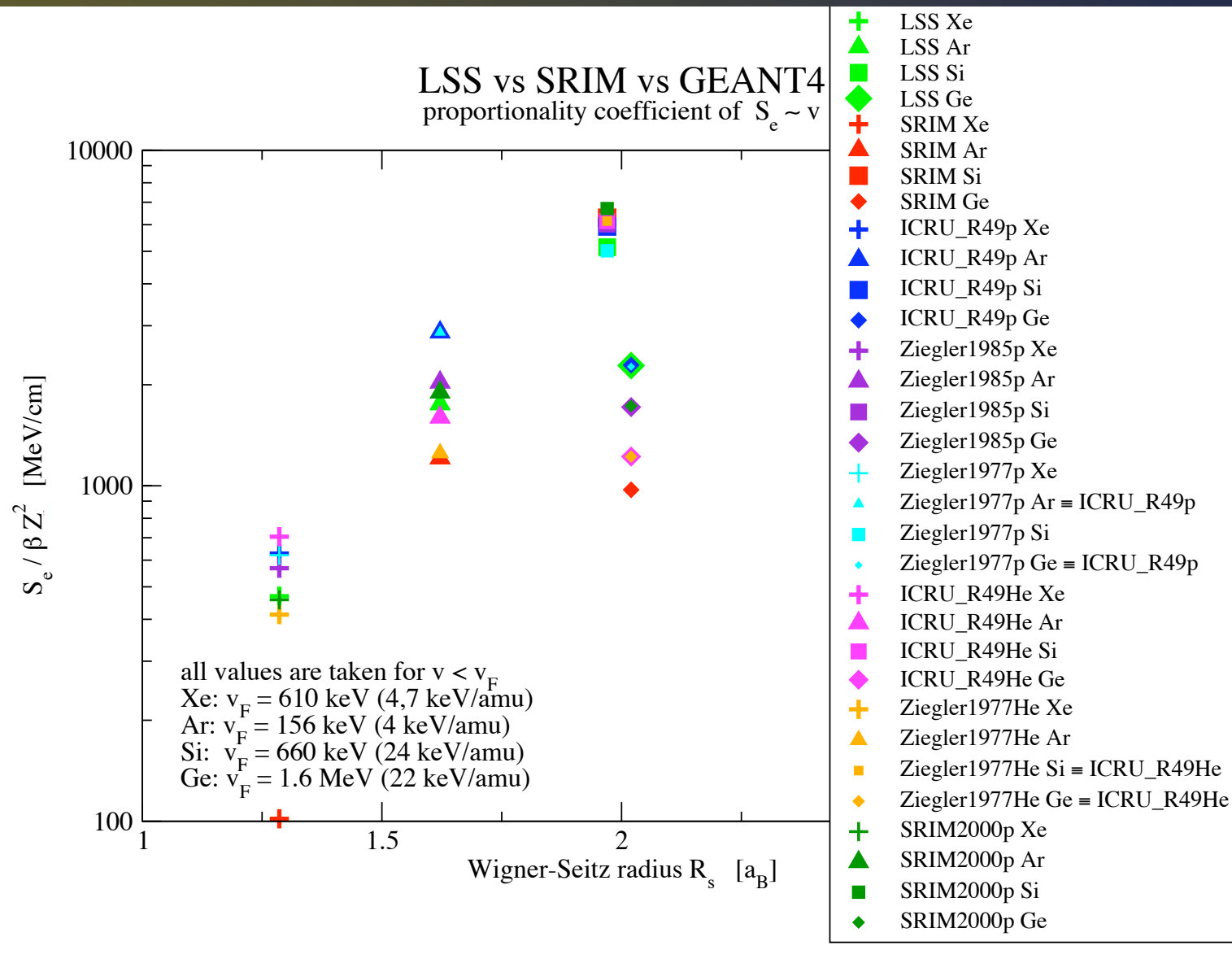
Geant4 low energy



parameterization	$T_1 < T_p < T_2$	$T_p < T_1$ (default = 1 keV)
ICRU_R49p Ziegler1977p	$T_2 = 10 \text{ keV}$ $S_e = A_1 \sqrt{T_p}$	$S_e = A_1 \sqrt{T_p}$
SRIM2000p Ziegler1985p	$S_e = A_1 T_p^{0.25}, Z_2 < 6$ $S_e = A_1 T_p^{0.375}, Z_2 = 6, 14, 32$ $S_e = A_1 T_p^{0.45}, \text{ other}$	$S_e = A \sqrt{T_p}$

Geant4

Electronic stopping



Nuclear Stopping S_n

ion-atom interaction potencial

Nuclear Stopping S_n

ion-atom interaction potencial

$$V(r) = \frac{Z_1 Z_2 e^2}{r} \Phi \quad \left\{ \begin{array}{ll} \lim_{r \rightarrow 0} & \Phi = 1 \\ \lim_{r \rightarrow \infty} & \Phi = 0 \end{array} \right.$$


Nuclear Stopping S_n

ion-atom interaction potencial

$$V(r) = \frac{Z_1 Z_2 e^2}{r} \Phi \quad \begin{cases} \lim_{r \rightarrow 0} \Phi = 1 \\ \lim_{r \rightarrow \infty} \Phi = 0 \end{cases}$$

Classical (or statistical) treatment: the Thomas-Fermi atom

Approximations:

- elastic
- static (velocity independent)
- universal $\Phi = \Phi(r/a)$  screening length
- no shell structure (as opposed to modern Hartree-Fock solid state models)


Nuclear Stopping S_n

ion-atom interaction potencial

$$V(r) = \frac{Z_1 Z_2 e^2}{r} \Phi \quad \begin{cases} \lim_{r \rightarrow 0} \Phi = 1 \\ \lim_{r \rightarrow \infty} \Phi = 0 \end{cases}$$

Classical (or statistical) treatment: the Thomas-Fermi atom

Approximations:

- elastic
- static (velocity independent)
- universal $\Phi = \Phi(r/a)$  screening length
- no shell structure (as opposed to modern Hartree-Fock solid state models)

$$V(r) \xrightarrow[\text{pertubation treatment}]{\text{classical scattering}} d\sigma \longrightarrow S_n = \int T' d\sigma$$

T' = energy transfered to an atom at rest (target)

Universal screening functions proposed in the literature:

- Bohr
- Thomas-Fermi → Lindhard (LSS)
- Lenz-Jensen
- Moliere
- Ziegler (ZBL) → SRIM
(a universal fit to 522 (!) solid state interatomic potentials)

Universal screening functions proposed in the literature:

- Bohr
- Thomas-Fermi → Lindhard (LSS)
- Lenz-Jensen
- Moliere
- Ziegler (ZBL) → SRIM
(a universal fit to 522 (!) solid state interatomic potentials)

Geant4

Nuclear stopping

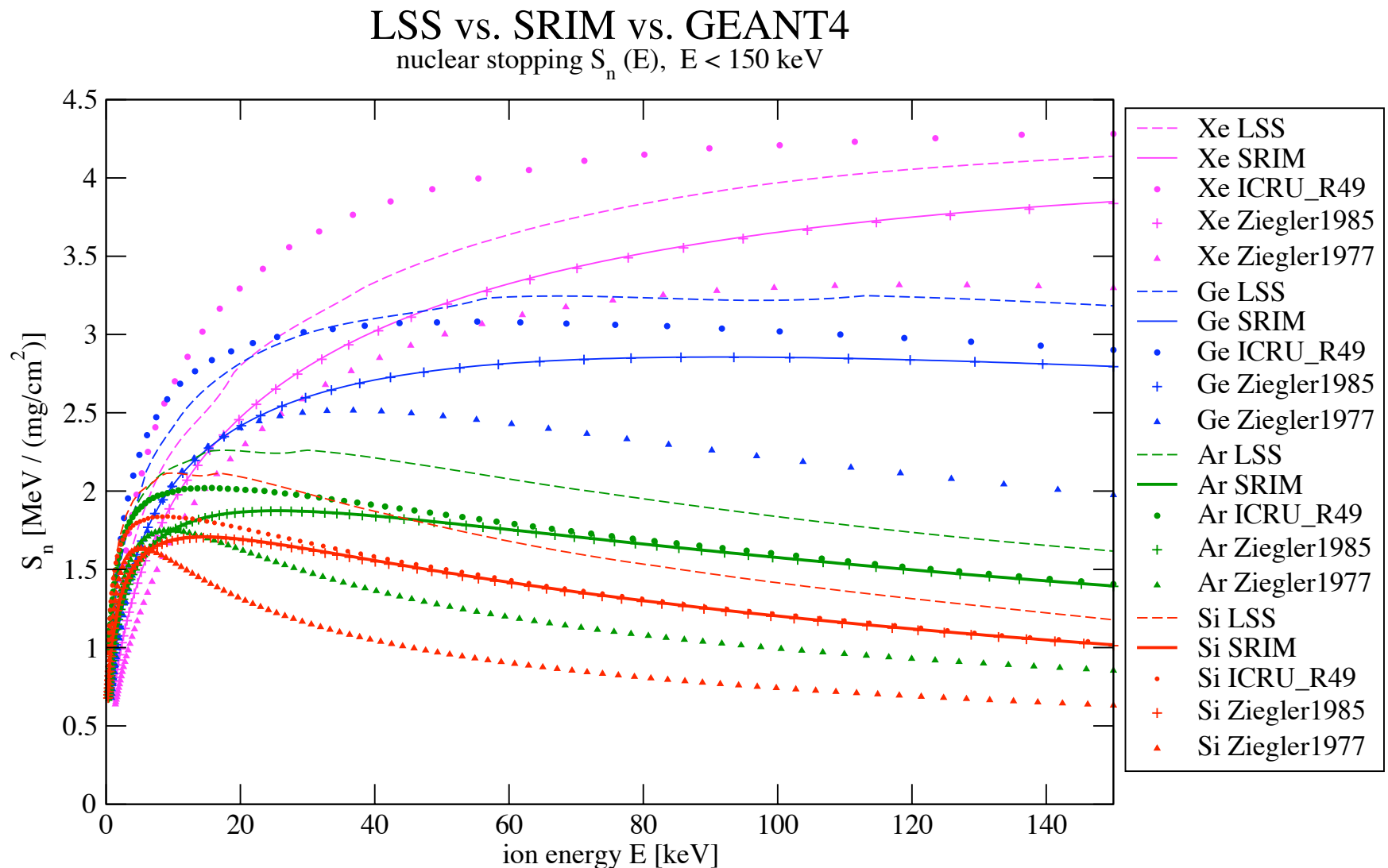
Universal stopping in reduced units scales back to ion-atom dependent stopping in physical units. Available choices:

(default)

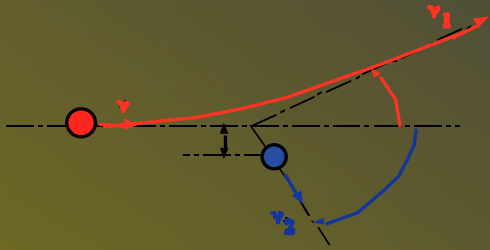
- ICRU_R49 - parameterization of Moliere's screening function
 - Ziegler1977
 - Ziegler1985
- } the ZBL universal potencial

Universal screening functions proposed in the literature:

- Bohr
- Thomas-Fermi → Lindhard
- Ziegler (ZBL) → SRIM

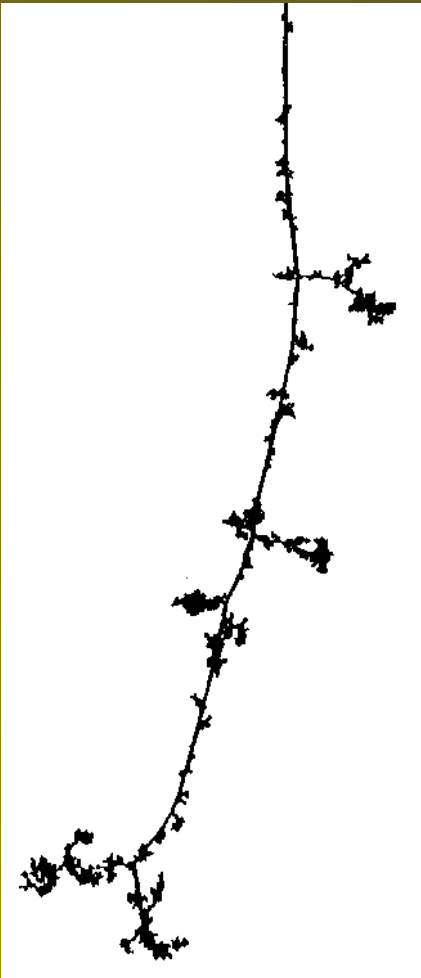


Nuclear recoil full cascade

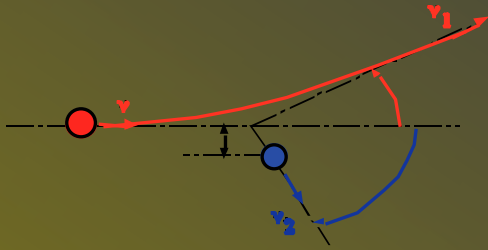


Note: the primary particle is already a recoil, doesn't matter if it originated from WIMPs or neutron scattering

recoil CUT energy (default MinKineticEnergy = 10eV)



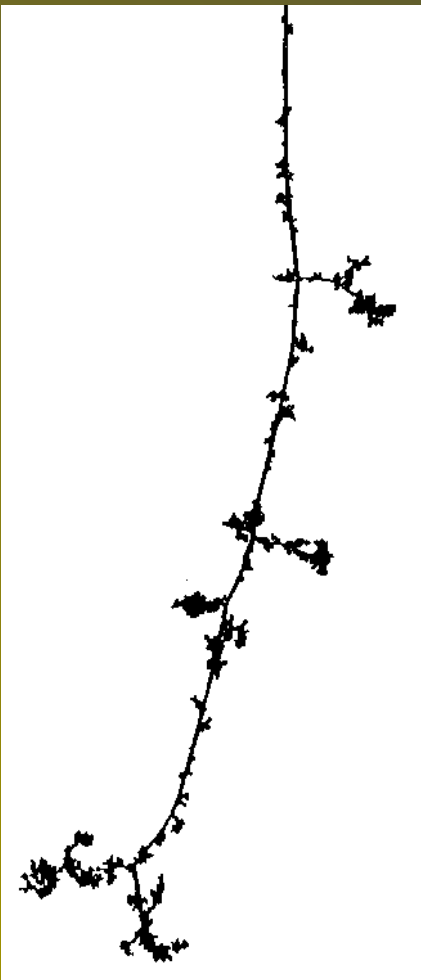
Nuclear recoil full cascade



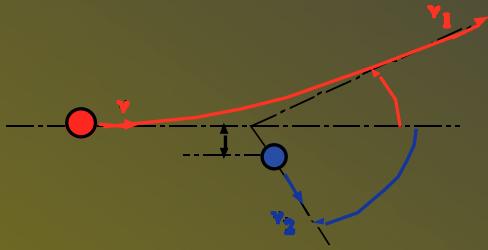
Note: the primary particle is already a recoil, doesn't matter if it originated from WIMPs or neutron scattering

recoil CUT energy (default MinKineticEnergy = 10eV)

Physics list:



Nuclear recoil full cascade

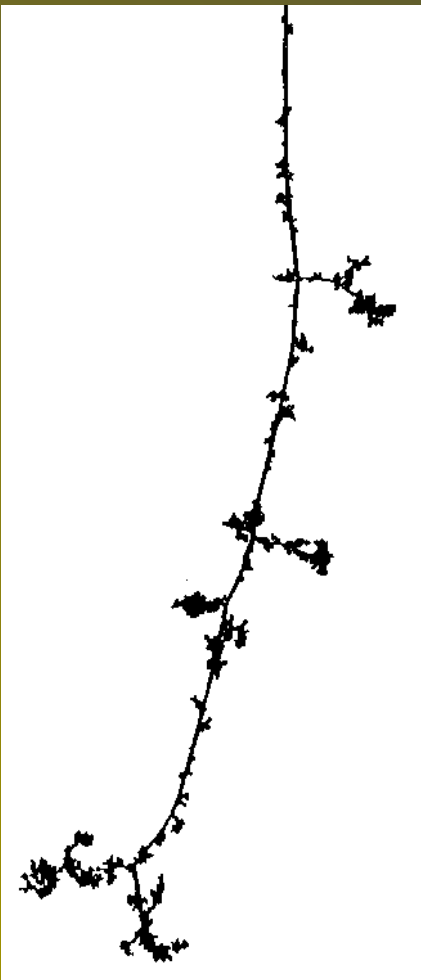


Note: the primary particle is already a recoil, doesn't matter if it originated from WIMPs or neutron scattering

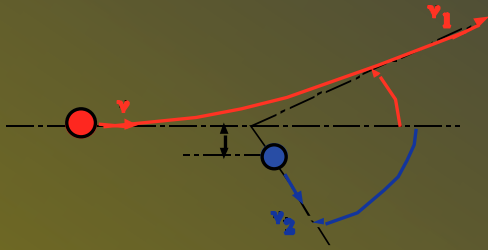
recoil CUT energy (default MinKineticEnergy = 10eV)

Physics list:

- Introduced a new physical process for the recoils



Nuclear recoil full cascade

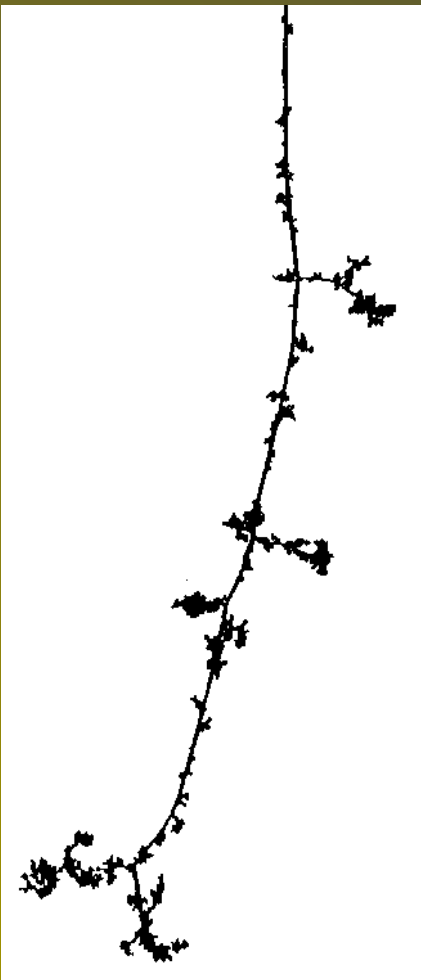


Note: the primary particle is already a recoil, doesn't matter if it originated from WIMPs or neutron scattering

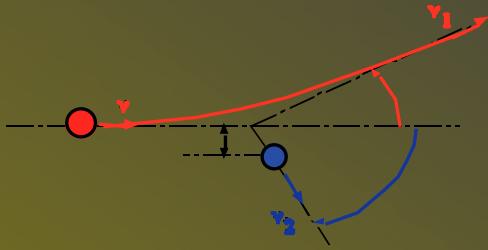
recoil CUT energy (default MinKineticEnergy = 10eV)

Physics list:

- Introduced a new physical process for the recoils
- Had to make some changes to G4hLowEnergyIonisation:



Nuclear recoil full cascade

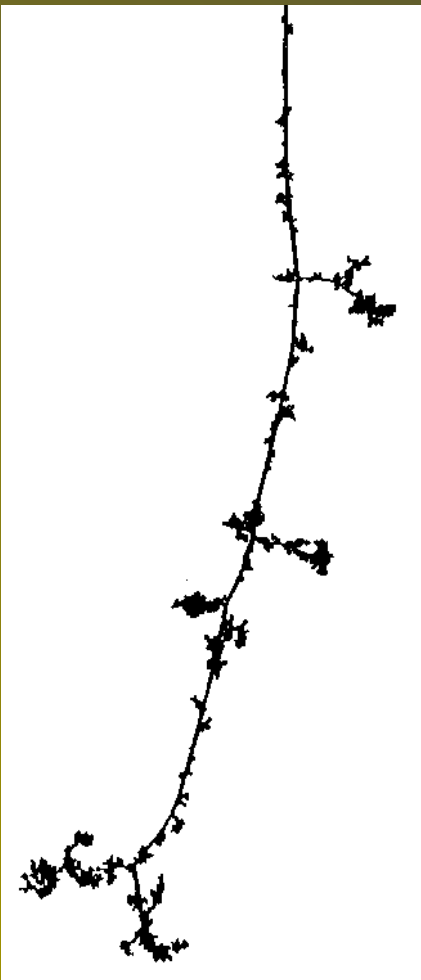


Note: the primary particle is already a recoil, doesn't matter if it originated from WIMPs or neutron scattering

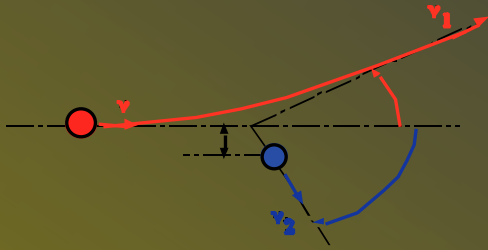
recoil CUT energy (default MinKineticEnergy = 10eV)

Physics list:

- Introduced a new physical process for the recoils
- Had to make some changes to G4hLowEnergyIonisation:
 - ★ don't deposit T when $T < \text{CUT}$



Nuclear recoil full cascade

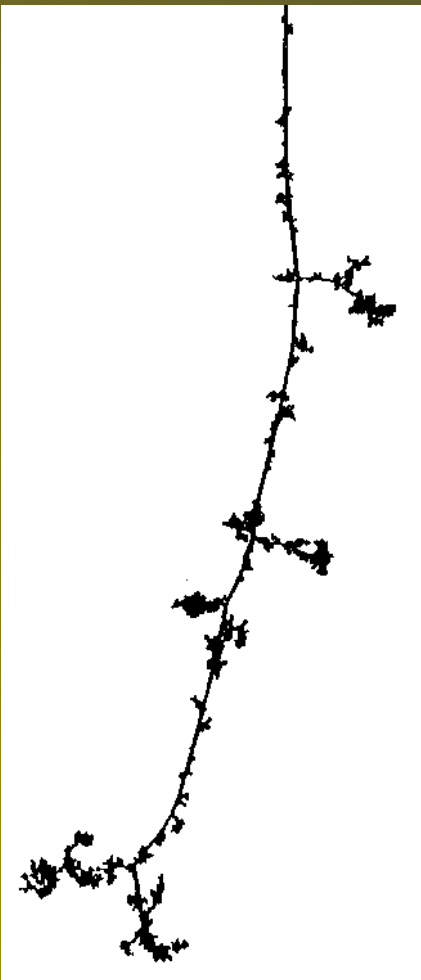


Note: the primary particle is already a recoil, doesn't matter if it originated from WIMPs or neutron scattering

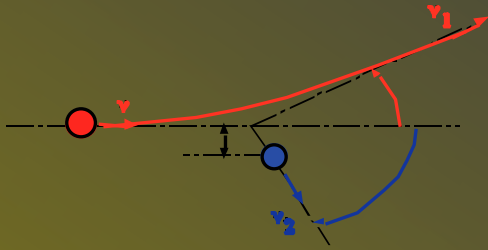
recoil CUT energy (default MinKineticEnergy = 10eV)

Physics list:

- Introduced a new physical process for the recoils
- Had to make some changes to G4hLowEnergyIonisation:
 - ★ don't deposit T when $T < \text{CUT}$
 - ★ changed CUT from T_p to T



Nuclear recoil full cascade

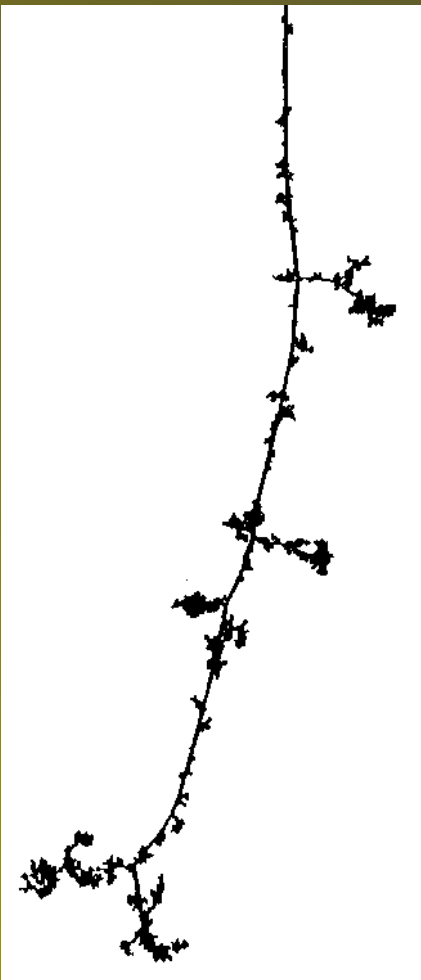


Note: the primary particle is already a recoil, doesn't matter if it originated from WIMPs or neutron scattering

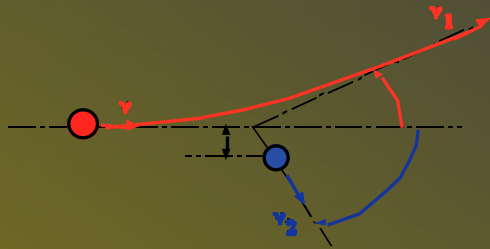
recoil CUT energy (default `MinKineticEnergy` = 10eV)

Physics list:

- Introduced a new physical process for the recoils
- Had to make some changes to `G4hLowEnergyIonisation`:
 - ★ don't deposit T when $T < \text{CUT}$
 - ★ changed CUT from T_p to T
 - ★ `SetNuclearStoppingOff()` is mandatory



Nuclear recoil full cascade



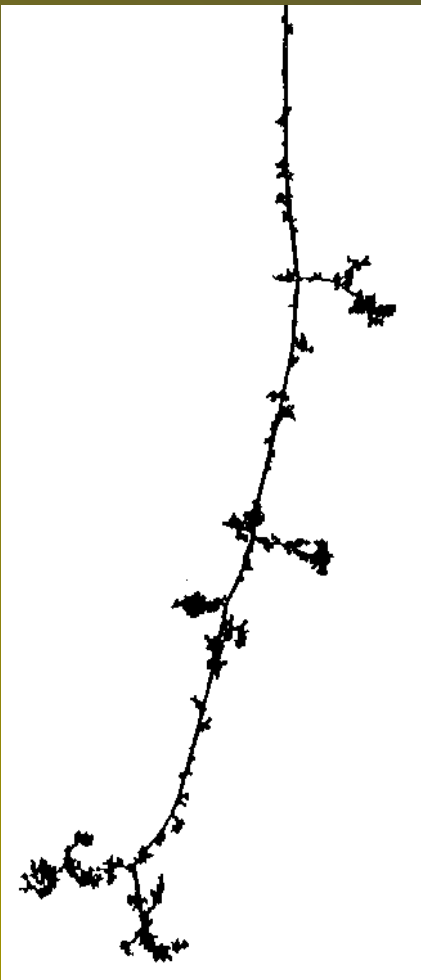
Note: the primary particle is already a recoil, doesn't matter if it originated from

code structure based on δ -electrons
emission, fluorescence and atom
deexcitation of G4hLowEnergyIonisation
and on the continuous photon generation
in G4Cerenkov

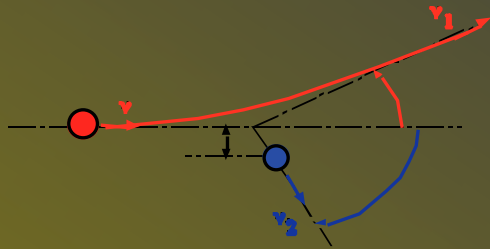
kineticEnergy = 10eV)

Phys.

- Introduced a new physical process for the recoils
- Had to make some changes to G4hLowEnergyIonisation:
 - ★ don't deposit T when $T < \text{CUT}$
 - ★ changed CUT from T_p to T
 - ★ SetNuclearStoppingOff() is mandatory



Nuclear recoil full cascade



Note: the primary particle is already a recoil, doesn't matter if it originated from

code structure based on δ -electrons
emission, fluorescence and atom
deexcitation of G4hLowEnergyIonisation
and on the continuous photon generation
in G4Cerenkov

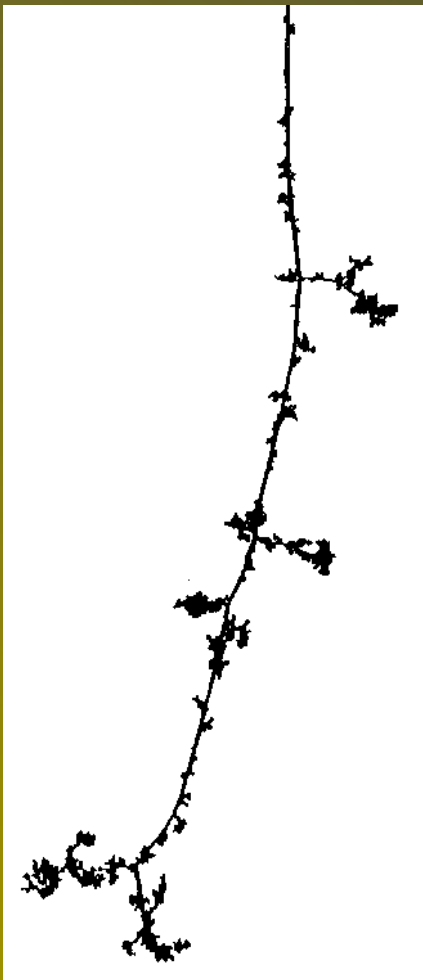
kineticEnergy = 10eV)

Phys.

- Introduced a new physical process for the recoils
- Had to make some changes to G4hLowEnergyIonisation:
 - ★ don't deposit T when $T < \text{CUT}$
 - ★ changed CUT from T_p to T
 - ★ SetNuclearQuenchingOff() is mandatory

nuclear quenching

→ competition between total energy ultimately given to electrons and energy permanently lost to recoils



Geant4

Recoil cascade

- Geant4 has no nuclear recoils (at least heavy ion)
- It also doesn't have the respective cross sections (neither differential nor integrated)

Geant4

Recoil cascade

- Geant4 has no nuclear recoils (at least heavy ion)
 - It also doesn't have the respective cross sections (neither differential nor integrated)
- ➔ 1st approximation, using only what Geant4 already has:
- ➔ the recoil atom kinetic energy is taken from the parameterized nuclear stopping power

$$T' = \left(-\frac{dE}{dx} \right)_n \Delta x = S_n \times \text{step size}$$

Geant4

Recoil cascade

- Geant4 has no nuclear recoils (at least heavy ion)
- It also doesn't have the respective cross sections (neither differential nor integrated)

➡ 1st approximation, using only what Geant4 already has:

➡ the recoil atom kinetic energy is taken from the parameterized nuclear stopping power

$$T' = \left(-\frac{dE}{dx} \right)_n \Delta x = S_n \times \text{step size}$$

- ➡ the so called continuous slowing down approximation (in a loosely sense) => a continuous process
- ➡ major disadvantage: very strong dependence on step size
=> an external parameter in need of fine tuning

Geant4

Recoil cascade

- Geant4 has no nuclear recoils (at least heavy ion)
- It also doesn't have the respective cross sections (neither differential nor integrated)

➔ 1st approximation, using only what Geant4 already has:

➔ the recoil atom kinetic energy is taken from the parameterized nuclear stopping power

$$T' = \left(-\frac{dE}{dx} \right)_n \Delta x$$

actually, we shouldn't even try to approximate nuclear collisions as a continuous process (as opposed to electronic stopping) as you're not supposed to get a recoil every step, just every now and then

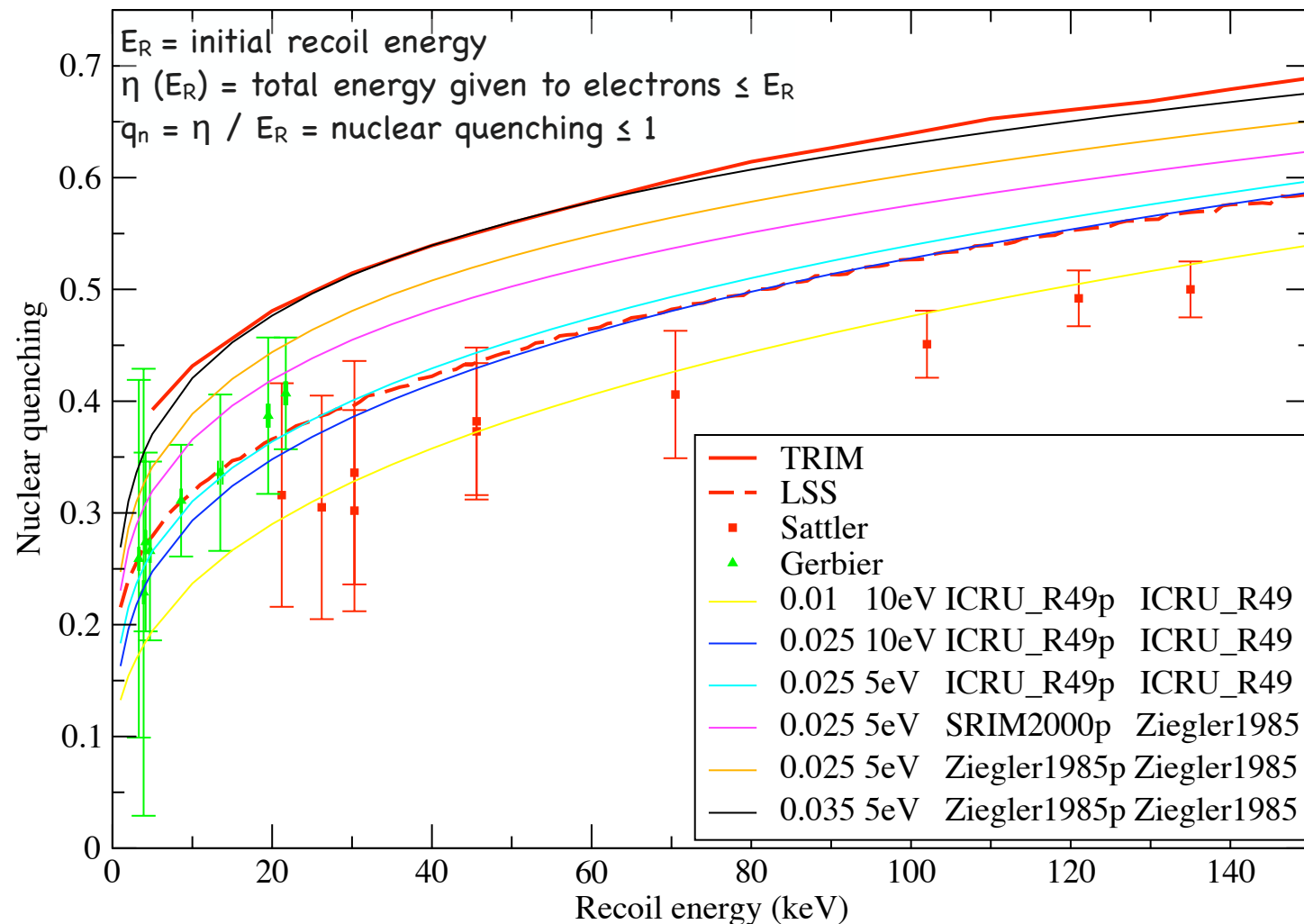
- ➔ the so called continuous slowing down approximation (in a loosely sense) => a continuous process
- ➔ major disadvantage: very strong dependence on step size
=> an external parameter in need of fine tuning

Geant4

Nuclear quenching: Si

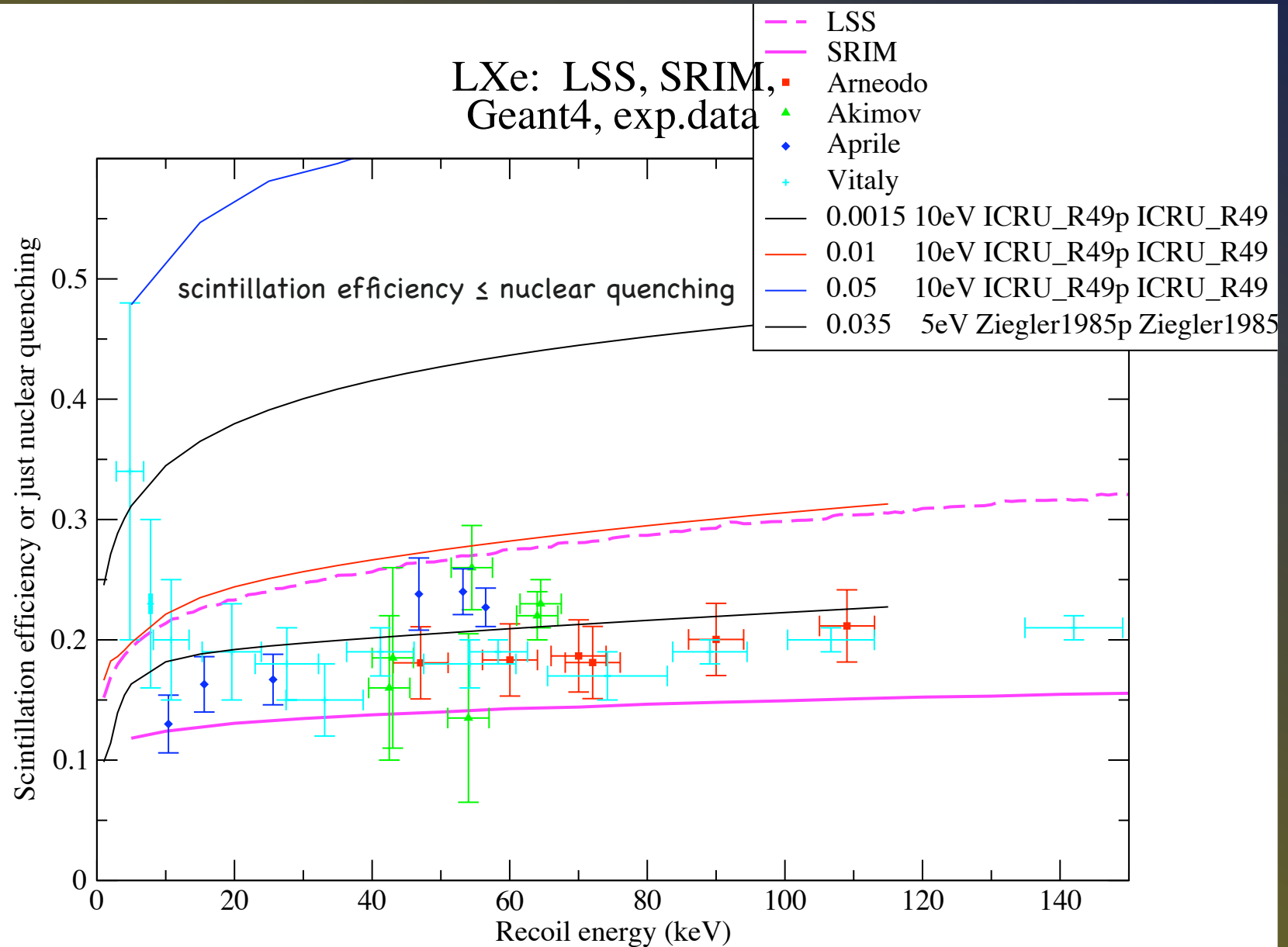
Si: LSS vs. SRIM vs. GEANT4 vs. exp. data

total fraction of energy transferred to electrons



Geant4

Nuclear quenching: LXe



Conclusions

1. Geant4 provides a good parameterization of both electronic and nuclear stopping powers (not shown here, is our validation of Geant4's stopping by comparing with NIST databases as well as experimental data, in different target materials, besides SRIM and LSS)

Conclusions

1. Geant4 provides a good parameterization of both electronic and nuclear stopping powers (not shown here, is our validation of Geant4's stopping by comparing with NIST databases as well as experimental data, in different target materials, besides SRIM and LSS)
2. The user has a total of 3x6 combinations of stopping parameterizations at his/her disposal

Conclusions

1. Geant4 provides a good parameterization of both electronic and nuclear stopping powers (not shown here, is our validation of Geant4's stopping by comparing with NIST databases as well as experimental data, in different target materials, besides SRIM and LSS)
2. The user has a total of 3x6 combinations of stopping parameterizations at his/her disposal
3. The continuous approximation provides a good prediction of the nuclear quenching of the recoils cascade, but also provides to much flexibility, by fine-tuning the total of 4 parameters that control the simulation:

electronic S_e :

- the parameterization
- step size

nuclear recoils :

- the parameterization
- energy cut

Conclusions

1. Geant4 provides a good parameterization of both electronic and nuclear stopping powers (not shown here, is our validation of Geant4's stopping by comparing with NIST databases as well as experimental data, in different target materials, besides SRIM and LSS)
2. The user has a total of 3x6 combinations of stopping parameterizations at his/her disposal
3. The continuous approximation provides a good prediction of the nuclear quenching of the recoils cascade, but also provides to much flexibility, by fine-tuning the total of 4 parameters that control the simulation:

electronic S_e :

- the parameterization
- ~~step size~~

 mean free path $\sim \frac{1}{\sigma}$

nuclear recoils :

- the parameterization
- energy cut

Conclusions

1. Geant4 provides a good parameterization of both electronic and nuclear stopping powers (not shown here, is our validation of Geant4's stopping by comparing with NIST databases as well as experimental data, in different target materials, besides SRIM and LSS)
2. The user has a total of 3x6 combinations of stopping parameterizations at his/her disposal
3. The continuous approximation provides a good prediction of the nuclear quenching of the recoils cascade, but also provides to much flexibility, by fine-tuning the total of 4 parameters that control the simulation:

electronic S_e :

- the parameterization
- ~~step size~~

mean free path $\sim \frac{1}{\sigma}$

+ $d\sigma \longrightarrow$ a full discrete process (for the recoils)

nuclear recoils :

- screening function
- energy cut

Conclusions

1. Geant4 provides a good parameterization of both electronic and nuclear stopping powers (not shown here, is our validation of Geant4's stopping by comparing with NIST databases as well as experimental data, in different target materials, besides SRIM and LSS)
2. The user has a total of 3x6 combinations of stopping parameterizations at his/her disposal
3. The continuous approximation provides a good prediction of the nuclear quenching of the recoils cascade, but also provides to much flexibility, by fine-tuning the total of 4 parameters that control the simulation:

electronic S_e :

- the parameterization
- ~~step size~~

mean free path $\sim \frac{1}{\sigma}$

+ $d\sigma \longrightarrow$ a full discrete process (for the recoils)

nuclear recoils :

- screening function
- energy cut

work in progress

Present work

$$d\sigma = \Pi a^2 \frac{f(t^{1/2})}{2t^{3/2}}$$

Present work

a_F, a_L, a_U

$$d\sigma = \Pi a^2 \frac{f(t^{1/2})}{2t^{3/2}}$$

Thomas-Fermi, Lenz-Jensen, Moliere, ZBL, ect.

Present work

a_F, a_L, a_U

$$d\sigma = \Pi a^2 \frac{f(t^{1/2})}{2t^{3/2}}$$

Thomas-Fermi, Lenz-Jensen, Moliere, ZBL, ect.

numerical integration

$$\sigma(T, T_c)$$

Present work

Thomas-Fermi, Lenz-Jensen, Moliere, ZBL, ect.

a_F, a_L, a_U

$$d\sigma = \Pi a^2 \frac{f(t^{1/2})}{2t^{3/2}}$$

numerical integration

$$\sigma(T, T_c)$$

To do:

put it in tables and
load them in Geant4,
with interpolation

Present work

Thomas-Fermi, Lenz-Jensen, Moliere, ZBL, ect.

a_F, a_L, a_U

$$d\sigma = \Pi a^2 \frac{f(t^{1/2})}{2t^{3/2}}$$

numerical integration

$$\sigma(T, T_c)$$

To do:

put it in tables and
load them in Geant4,
with interpolation

write the
DiscreteProcess functions

Present work

Thomas-Fermi, Lenz-Jensen, Moliere, ZBL, ect.

a_F, a_L, a_U

$$d\sigma = \Pi a^2 \frac{f(t^{1/2})}{2t^{3/2}}$$

numerical integration

$$\sigma(T, T_c)$$

implement the
sampling

To do:

put it in tables and
load them in Geant4,
with interpolation

write the
DiscreteProcess functions

Present work

Thomas-Fermi, Lenz-Jensen, Moliere, ZBL, ect.

a_F, a_L, a_U

$$d\sigma = \Pi a^2 \frac{f(t^{1/2})}{2t^{3/2}}$$

numerical integration

$$\sigma(T, T_c)$$

implement the
sampling

$$T' = \gamma T \sin^2 \frac{\theta}{2}$$

To do:

put it in tables and
load them in Geant4,
with interpolation

write the
DiscreteProcess functions

Present work

Thomas-Fermi, Lenz-Jensen, Moliere, ZBL, ect.

a_F, a_L, a_U

$$d\sigma = \Pi a^2 \frac{f(t^{1/2})}{2t^{3/2}}$$

numerical integration

$$\sigma(T, T_c)$$

implement the
sampling

To do:

put it in tables and
load them in Geant4,
with interpolation

write the
DiscreteProcess functions

$$T' = \gamma T \sin^2 \frac{\theta}{2}$$

Opposed to the standard method of solving the scattering integral:

Present work

Thomas-Fermi, Lenz-Jensen, Moliere, ZBL, ect.

a_F, a_L, a_U

$$d\sigma = \Pi a^2 \frac{f(t^{1/2})}{2t^{3/2}}$$

numerical integration

$$\sigma(T, T_c)$$

implement the
sampling

To do:

put it in tables and
load them in Geant4,
with interpolation

write the

DiscreteProcess functions

$$T' = \gamma T \sin^2 \frac{\theta}{2}$$

Opposed to the standard method of solving the scattering integral:

- SRIM/TRIM (the standard!)

Present work

Thomas-Fermi, Lenz-Jensen, Moliere, ZBL, ect.

a_F, a_L, a_U

$$d\sigma = \Pi a^2 \frac{f(t^{1/2})}{2t^{3/2}}$$

numerical integration

$$\sigma(T, T_c)$$

implement the sampling

To do:

put it in tables and load them in Geant4, with interpolation

write the

DiscreteProcess functions

$$T' = \gamma T \sin^2 \frac{\theta}{2}$$

Opposed to the standard method of solving the scattering integral:

- SRIM/TRIM (the standard!)
- M.H. Mendenhalla and R.A. Weller, An algorithm for computing screened Coulomb scattering in GEANT4, Nuclear Instruments and Methods in Physics Research B227, 3 (2005) 420

Extras

Electronic Stopping S_e

Fermi-Teller: $S_e \propto v$ (ion velocity) if $v < v_F \sim v_0$

Lindhard: uniform free electron gas, Thomas-Fermi atom, particle-plasma interaction as a perturbation

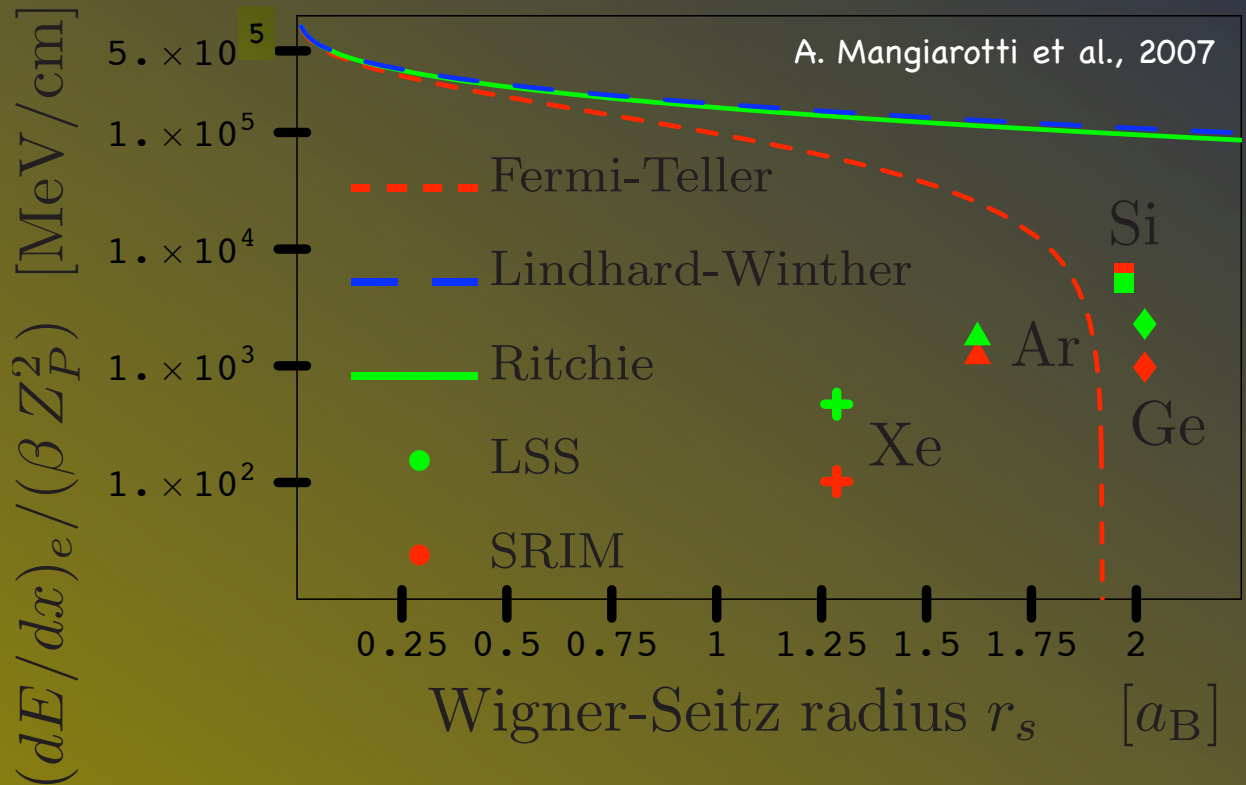
$S_e = k\epsilon^{1/2}$ if $v < \text{Bragg peak}$ ($\sim v_0 Z_1^{2/3} \approx 500 \text{ MeV}$ for Xe)

$$0.10 < k < 0.20$$

(reduced units)

$$k_{Xe} \approx 0.166$$

SRIM: local density approximation using Hartree-Fock solid state atoms; semi-empirical fit of charge state of the ion



Nuclear quenching ?!

Nuclear quenching ?!

- E = initial particle energy (e.g. $E_\gamma = 122\text{keV}$ for Co57 or E_R for recoils produced by neutrons or WIMPs)

Nuclear quenching ?!

- E = initial particle energy (e.g. $E_\gamma = 122\text{keV}$ for Co57 or E_R for recoils produced by neutrons or WIMPs)
- $\eta(E)$ = total energy ultimately given to electrons $\leq E$ ($\eta_\gamma \approx E_\gamma$)

Nuclear quenching ?!

- E = initial particle energy (e.g. $E_\gamma = 122\text{keV}$ for Co57 or E_R for recoils produced by neutrons or WIMPs)
- $\eta(E)$ = total energy ultimately given to electrons $\leq E$ ($\eta_\gamma \approx E_\gamma$)
- $q_n = \eta_R / E_R$ = nuclear quenching (of the recoils) ≤ 1

Nuclear quenching ?!

- E = initial particle energy (e.g. $E_Y = 122\text{keV}$ for Co57 or E_R for recoils produced by neutrons or WIMPs)
- $\eta(E)$ = total energy ultimately given to electrons $\leq E$ ($\eta_Y \approx E_Y$)
- $q_n = \eta_R / E_R$ = nuclear quenching (of the recoils) ≤ 1
 - ➔ But what do we actually measure in the lab ?!

Nuclear quenching ?!

- E = initial particle energy (e.g. $E_\gamma = 122\text{keV}$ for Co57 or E_R for recoils produced by neutrons or WIMPs)
- $\eta(E)$ = total energy ultimately given to electrons $\leq E$ ($\eta_\gamma \approx E_\gamma$)
- $q_n = \eta_R / E_R$ = nuclear quenching (of the recoils) ≤ 1

➡ But what do we actually measure in the lab ?!

- relative scintillation efficiency $\approx \frac{\eta_R(E_R)}{E_R} \times \frac{W_S(E_\gamma)}{W_S(E_R)}$

Nuclear quenching ?!

- E = initial particle energy (e.g. $E_\gamma = 122\text{keV}$ for Co57 or E_R for recoils produced by neutrons or WIMPs)
- $\eta(E)$ = total energy ultimately given to electrons $\leq E$ ($\eta_\gamma \approx E_\gamma$)
- $q_n = \eta_R / E_R$ = nuclear quenching (of the recoils) ≤ 1

➡ But what do we actually measure in the lab ?!

- relative scintillation efficiency $\approx \frac{\eta_R(E_R)}{E_R} \times \frac{W_S(E_\gamma)}{W_S(E_R)}$

where we redefine:

$$\left. \begin{array}{l} \gamma : W_S = E_\gamma / N_{ph} \\ R : W_S = \eta_R / N'_{ph} \end{array} \right\} W_S \left\{ \begin{array}{l} \text{now they refer to the energy} \\ \text{transferred to electrons only!} \\ \text{and are both just functions of LET} \end{array} \right.$$

Nuclear quenching ?!

- E = initial particle energy (e.g. $E_\gamma = 122\text{keV}$ for Co57 or E_R for recoils produced by neutrons or WIMPs)
- $\eta(E)$ = total energy ultimately given to electrons $\leq E$ ($\eta_\gamma \approx E_\gamma$)
- $q_n = \eta_R / E_R = \text{nuclear quenching (of the recoils)} \leq 1$

➔ But what do we actually measure in the lab ?!

- relative scintillation efficiency $\approx \frac{\eta_R(E_R)}{E_R} \times \frac{W_S(E_\gamma)}{W_S(E_R)}$

where we redefine:

$$\left. \begin{array}{l} \gamma : W_S = E_\gamma / N_{ph} \\ R : W_S = \eta_R / N'_{ph} \end{array} \right\} W_S \left\{ \begin{array}{l} \text{now they refer to the energy} \\ \text{transferred to electrons only!} \\ \text{and are both just functions of LET} \end{array} \right.$$

Nuclear quenching ?!

- E = initial particle energy (e.g. $E_\gamma = 122\text{keV}$ for Co57 or E_R for recoils produced by neutrons or WIMPs)
- $\eta(E)$ = total energy ultimately given to electrons $\leq E$ ($\eta_\gamma \approx E_\gamma$)
- $q_n = \eta_R / E_R = \text{nuclear quenching (of the recoils)} \leq 1$

➡ But what do we actually measure in the lab ?!

- relative scintillation efficiency $\approx \frac{\eta_R(E_R)}{E_R} \times \underbrace{\frac{W_S(E_\gamma)}{W_S(E_R)}}_{\text{Hitachi gives constant } < 1}$

where we redefine:

$$\left. \begin{array}{l} \gamma : W_S = E_\gamma / N_{ph} \\ R : W_S = \eta_R / N'_{ph} \end{array} \right\} W_S \left\{ \begin{array}{l} \text{now they refer to the energy} \\ \text{transferred to electrons only!} \\ \text{and are both just functions of LET} \end{array} \right.$$

(relative) Scintillation efficiency

$$q_n \times \frac{W_S(E_\gamma)}{W_S(E_R)}$$

Is everybody clear about this definition ?!

derivation from the traditional/experimental formula

comparison with Hitachi definition

traditional definition of $W_S = \frac{E}{N_{ph}}$ for γ , Recoils, whatever!

back

traditional definition of $W_S = \frac{E}{N_{ph}}$ for γ , Recoils, whatever!

traditional/experimental definition of relative scintillation efficiency

$$= \frac{W_S(E_\gamma)}{W_S(E_R)}$$

< 1 because it
takes more energy
to produce a
photon in RC than
in γ

2 "quenching"
effects all
mixed up!

back

traditional definition of $W_S = \frac{E}{N_{ph}}$ for γ , Recoils, whatever!

traditional/experimental definition of relative scintillation efficiency

$$= \frac{W_S(E_\gamma)}{W_S(E_R)} = \frac{\frac{E_\gamma}{N_{ph}}}{\frac{E_R}{N'_{ph}}}$$

back

traditional definition of $W_S = \frac{E}{N_{ph}}$ for γ , Recoils, whatever!

traditional/experimental definition of relative scintillation efficiency

$$= \frac{W_S(E_\gamma)}{W_S(E_R)} = \frac{\frac{E_\gamma}{N_{ph}}}{\frac{E_R}{N'_{ph}}} = \frac{\frac{E_\gamma}{N_{ph}}}{\frac{\eta_R / q_n}{N'_{ph}}}$$

remember $\eta_\gamma \approx E_\gamma$

remember our definition of $q_n = \eta_R / E_R$

back

traditional definition of $W_S = \frac{E}{N_{ph}}$ for γ , Recoils, whatever!

traditional/experimental definition of relative scintillation efficiency

$$= \frac{W_S(E_\gamma)}{W_S(E_R)} = \frac{\frac{E_\gamma}{N_{ph}}}{\frac{E_R}{N'_{ph}}} = \frac{\frac{E_\gamma}{N_{ph}}}{\frac{\eta_R/q_n}{N'_{ph}}} = q_n \frac{W_S(E_\gamma)}{W_S(E_R)}$$

we redefine:

$$\begin{aligned} \gamma : W_S &= E_\gamma / N_{ph} \\ R : W_S &= \eta_R / N'_{ph} \end{aligned}$$

back

traditional definition of $W_S = \frac{E}{N_{ph}}$ for γ , Recoils, whatever!

traditional/experimental definition of relative scintillation efficiency

$$= \frac{W_S(E_\gamma)}{W_S(E_R)} = \frac{\frac{E_\gamma}{N_{ph}}}{\frac{E_R}{N'_{ph}}} = \frac{\frac{E_\gamma}{N_{ph}}}{\frac{\eta_R/q_n}{N'_{ph}}} = q_n \frac{W_S(E_\gamma)}{W_S(E_R)}$$

we redefine:

$$\begin{aligned} \gamma : W_S &= E_\gamma / N_{ph} \\ R : W_S &= \eta_R / N'_{ph} \end{aligned}$$

back

Hitachi

what we measure $\frac{RC}{\gamma} = \frac{q_n^{(R)} q_{el}^{(R)}}{q_n^{(\gamma)} q_{el}^{(\gamma)}}$

back

Hitachi

what we measure $\frac{RC}{\gamma} = \frac{q_n^{(R)} q_{el}^{(R)}}{q_n^{(\gamma)} q_{el}^{(\gamma)}}$

$q \equiv$ quenching or
scintillation efficiency
of the recoils

Hitachi

what we measure

$$\frac{RC}{\gamma} = \frac{q_n^{(R)} q_{el}^{(R)}}{1 \cancel{q_n^{(\gamma)}} q_{el}^{(\gamma)}}$$

$q \equiv$ quenching or
scintillation efficiency
of the recoils

Hitachi

what we measure

$$\frac{RC}{\gamma} = \frac{q_n^{(R)} q_{el}^{(R)}}{q_n^{(\gamma)} q_{el}^{(\gamma)}}$$

$q \equiv$ quenching or scintillation efficiency of the recoils

1

$$= q_n^{(R)} \frac{q_{el}^{(R)}}{q_{el}^{(\gamma)}}$$

scintillation only

Hitachi

what we measure

$$\frac{RC}{\gamma} = \frac{q_n^{(R)} q_{el}^{(R)}}{q_n^{(\gamma)} q_{el}^{(\gamma)}}$$

$q \equiv$ quenching or scintillation efficiency of the recoils

1

$$= q_n^{(R)} \frac{q_{el}^{(R)}}{q_{el}^{(\gamma)}}$$

scintillation only

competition between recombination processes

Hitachi

what we measure $\frac{RC}{\gamma} = \frac{q_n^{(R)} q_{el}^{(R)}}{q_n^{(\gamma)} q_{el}^{(\gamma)}}$

$q \equiv$ quenching or scintillation efficiency of the recoils

1 $\frac{q_n^{(R)} q_{el}^{(R)}}{q_n^{(\gamma)} q_{el}^{(\gamma)}}$ scintillation only

$= q_n^{(R)} \frac{q_{el}^{(R)}}{q_{el}^{(\gamma)}}$

competition between recombination processes

electrons escaping recombination

Hitachi

what we measure

$$\frac{RC}{\gamma} = \frac{q_n^{(R)} q_{el}^{(R)}}{q_n^{(\gamma)} q_{el}^{(\gamma)}}$$

$q \equiv$ quenching or scintillation efficiency of the recoils

1

$$= q_n^{(R)} \frac{q_{el}^{(R)}}{q_{el}^{(\gamma)}}$$

scintillation only

competition between recombination processes

electrons escaping recombination