

# New-physics prospects in hyperon decays

**J. Martin Camalich**



New Physics Searches at Kaon and Hyperon Factories

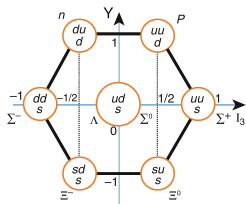
May 11 2021

# Outline

- 1 Presenting the hyperons
- 2 Heavy New Physics
- 3 Two-body decays
  - The SN bound
- 4 Dark baryons

# The hyperons

## ● Hyperons are the strange siblings of the proton and neutron



▶ **Half lives:**  $\tau_Y \sim 10^{-10}$  s

★ Sensitivity loss  $\sim 10^3$  w.r.t. to  $K^+$ ,  $K_L$

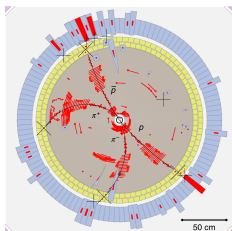
▶ **Rich phenomenology:**

★ Spin  $\rightarrow$  sensitivity to various NP structures

★  $SU(3)$ -relations to nucleon-structure

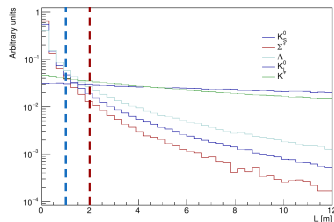
## ● Recent experimental “revolution” after 40<sup>+</sup> yrs ...

▶ **Polarized-hyperon factories (BESIII&SCTF)**



Nature Physics 15, 631-634(2019)

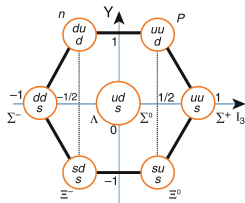
▶ **LHCb:**  $10^{2-3}$  more hyps than  $B$ 's



Alves Junior *et al.* JHEP 05 (2019) 048

# The hyperons

- **Hyperons are the strange siblings of the proton and neutron**



- ▶ **Half lifes:**  $\tau_Y \sim 10^{-10}$  s

- ★ Sensitivity loss  $\sim 10^3$  w.r.t. to  $K^+$ ,  $K_L$

- ▶ **Rich phenomenology:**

- ★ Spin  $\rightarrow$  sensitivity to various NP structures

- ★  $SU(3)$ -relations to nucleon-structure

- **We will use and recast Hai-Bo Li predictions for  $B_1 \rightarrow B_2 \nu \bar{\nu}$**

Front. Phys. 12(5), 121301 (2017)

DOI 10.1007/s11467-017-0691-9

## PERSPECTIVE

### Prospects for rare and forbidden hyperon decays at BESIII

Hai-Bo Li<sup>1,2,†</sup>

<sup>1</sup>Institute of High Energy Physics, Beijing 100049, China

<sup>2</sup>University of Chinese Academy of Sciences, Beijing 100049, China

Corresponding author. E-mail: <sup>†</sup>lihb@ihep.ac.cn

Received April 17, 2017; accepted May 8, 2017

# Form factors: Lattice, $SU(3)_F$ and ChPT

- **Predictions require strong form factors** (e.g. LQCD)

$$\langle B_2(\rho_2) | J_\mu^V | B_1(\rho_1) \rangle = \bar{u}_2(\rho_2) \left[ f_1(q^2) \gamma_\mu + \frac{f_2(q^2)}{M_1} \sigma_{\mu\nu} q^\nu + \frac{f_3(q^2)}{M_1} q_\mu \right] u_1(\rho_1),$$

$$\langle B_2(\rho_2) | J_\mu^A | B_1(\rho_1) \rangle = \bar{u}_2(\rho_2) \left[ g_1(q^2) \gamma_\mu + \frac{g_2(q^2)}{M_1} \sigma_{\mu\nu} q^\nu + \frac{g_3(q^2)}{M_1} q_\mu \right] \gamma_5 u_1(\rho_1)$$

$$\langle B_2(\rho') | J_{\mu\nu}^T | B_1(\rho) \rangle_{q^2=0} = g_T \bar{u}_2(\rho') \sigma_{\mu\nu} u_1(\rho)$$

- $SU(3)_F$  **symmetry** (lacking **direct** LQCD computations)

$$\langle B_a | J_b | B_c \rangle = f_{bac} F_J(q^2) + d_{bac} D_J(q^2)$$

Georgi, Lie Algebras in Particle Physics

- ▶ **Reduced matrix elements** from **data** (axial/vector) or **lattice** (tensor)

$$SU(3)_F \text{ breaking: } \delta \equiv (M_1 - M_2)/M \sim 10 - 20\%$$

- ▶ **Kinematic expansion:**  $q/M_1 \sim \mathcal{O}(\delta)$
- $SU(3)_F$ -**symmetry is embedded as the LO in ChPT**

# Heavy New Physics: $B_1 \rightarrow B_2 \bar{\nu} \nu$ Geng, JMC, Xiang, to appear

- **New physics scale of  $s \rightarrow d\nu\nu$  transitions  $\Lambda \gg 1$  GeV**

$$\mathcal{H}_{\text{eff}}^{\text{SM}} = \frac{G_F}{\sqrt{2}} \lambda_t \left( \sum_{i=1}^{10} C_i O_i + \alpha \sum_{\ell} C_{\nu\ell}^L O_{\nu\ell}^L \right)$$

- **Add RH NP**

$$\mathcal{H}_{\text{eff}}^{\text{NP}} = \frac{G_F \alpha}{\sqrt{2}} \lambda_t \sum_{\ell} C_{\nu\ell}^R (\bar{d}s)_{V+A} (\bar{\nu}_{\ell} \nu_{\ell})_{V-A}$$

- **Total  $B_1 \rightarrow B_2 \bar{\nu} \nu$  rate**

$$\Gamma = \sum_{\ell} \frac{\alpha^2 G_F^2 |\lambda_t|^2 f_1(0)^2 M_1^5 \delta^5}{60\pi^3} \cdot \left[ \left(1 - \frac{3}{2}\delta\right) |C_{\nu\ell}^L + C_{\nu\ell}^R|^2 + 3\left(1 - \frac{3}{2}\delta\right) \frac{g_1(0)^2}{f_1(0)^2} |C_{\nu\ell}^L - C_{\nu\ell}^R|^2 \right] + \mathcal{O}(\delta^2)$$

- ▶ Kinematic expansion  $\Rightarrow$  Only **2 charges**  $f_1(0)$  and  $g_1(0)$  contribute
- ▶  $f_1(0)$  and  $g_1(0)$  determined from **data** ( $B_1 \rightarrow B_2 e \nu$ )
- ▶ Baryon decays sensitive to **axial current** combination

- **SM predictions including NLO corrections in ChPT**

	$\Lambda n$	$\Sigma^+ p$	$\Xi^0 \Sigma^0$	$\Xi^0 \Lambda$
SM-NLO	$6.0(2) \times 10^{-13}$	$3.8(2) \times 10^{-13}$	$8.4(5) \times 10^{-14}$	$5.7(3) \times 10^{-13}$

# Heavy New Physics: Interplay with kaons

- $K \rightarrow \pi \nu \bar{\nu}$  is only sensitive to **vectorial** combination

$$\text{BR}(K^+ \rightarrow \pi^+ \nu \bar{\nu}) = \frac{2\alpha^2 |\lambda_t|^2 \text{BR}(K^+ \rightarrow \pi^0 e^+ \nu_e)}{|V_{us}|^2} \sum_{\ell=e,\mu,\tau} |C_{\nu\ell}^L + C_{\nu\ell}^R|^2$$

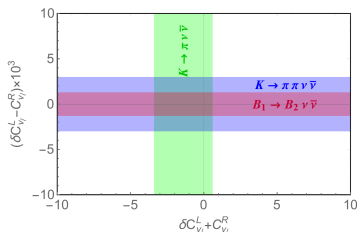
- $K \rightarrow \pi \pi \nu \bar{\nu}$  is sensitive to **axial** combination

- ▶ Form factors predicted using **isospin symmetry** and semileptonic  $K \rightarrow \pi \pi \nu e$  decays

Geng, JMC, Xiang, to appear; see also Littenber&Valencia, PLB385,379 (1996), Chiang, Gilman, PRD62,094026(2000)

$K^+ \pi^+ \pi^0$	$K_L \pi^0 \pi^0$
$7.90(34) \times 10^{-15}$	$3.11(6) \times 10^{-13}$

- Compare **current kaon constraints** with **BESIII projections** (Hai-bo's paper)



Geng, JMC, Xiang, to appear

Can we do much better than this with hyperons or  $K \rightarrow \pi \pi \nu \bar{\nu}$ ?

## Two body decays with a final dark boson $X^0$

- **Axions:**  $\mathcal{L}_a = \frac{\partial_{\mu a}}{2f_a} \bar{\psi}_i \gamma^\mu (c_{ij}^V + c_{ij}^A \gamma_5) \psi_j + \frac{1}{2} m_a^2 a^2$

$$\Gamma_a = \frac{\bar{p}\bar{\omega}^2}{8\pi f_a^2} (f_1^2 |c_{ds}^V|^2 + (1 - \frac{x_a^2}{\delta^2}) g_1^2 |c_{ds}^A|^2) + \mathcal{O}(\delta^2)$$

- ▶ Only **2 hadronic charges**  $f_1$  and  $g_1$ :  $SU(3)_F$  + **data**
- ▶  $B_1 \rightarrow B_2 a$  sensitive to **axial** current  $c_{ds}^A$  (worse for larger  $x_a = m_a/\delta$ )
- **Massless dark photon:**  $\mathcal{L}_{\gamma'} = \frac{1}{\Lambda_{UV}} \bar{\psi}_i \sigma^{\mu\nu} (c^{ij} + c_5^j \gamma_5) \psi_j F'_{\mu\nu}$

$$\Gamma_{\gamma'} = \frac{4\bar{\omega}^3 g_T^2}{\pi \Lambda_{UV}^2} (|c_{ds}^{ds}|^2 + |c_5^{ds}|^2)$$

- ▶ Only **1 hadronic charge**  $g_T$ :  $SU(3)_F$  + **LQCD** Gupta et al. PRD98,091501
- ▶ Best lab limit on  $s \rightarrow d\gamma'$  stems from  $\Xi^0 \rightarrow \Sigma^0 \gamma'$  and invisible BR of  $\Xi^0$

$$\text{BR}_{\text{inv}}(\Xi^0) \leq 2.3 \times 10^{-4} \quad \Rightarrow \quad \Lambda_{UV} \leq 3.8 \times 10^7 \text{ (95\% C.L.)}$$



## Two body decays with a final dark boson $X^0$

- **Axions:**  $\mathcal{L}_a = \frac{\partial_{\mu a}}{2f_a} \bar{\psi}_i \gamma^\mu (c_{ij}^V + c_{ij}^A \gamma_5) \psi_j + \frac{1}{2} m_a^2 a^2$

$$\Gamma_a = \frac{\bar{p}\bar{\omega}^2}{8\pi f_a^2} (f_1^2 |c_{ds}^V|^2 + (1 - \frac{x_a^2}{\delta^2}) g_1^2 |c_{ds}^A|^2) + \mathcal{O}(\delta^2)$$

- ▶ Only **2 hadronic charges**  $f_1$  and  $g_1$ :  $SU(3)_F$  + **data**
- ▶  $B_1 \rightarrow B_2 a$  sensitive to **axial** current  $c_{ds}^A$  (worse for larger  $x_a = m_a/\delta$ )
- **Massless dark photon:**  $\mathcal{L}_{\gamma'} = \frac{1}{\Lambda_{UV}} \bar{\psi}_i \sigma^{\mu\nu} (C^{ij} + C_5^{ij} \gamma_5) \psi_j F'_{\mu\nu}$

$$\Gamma_{\gamma'} = \frac{4\bar{\omega}^3 g_T^2}{\pi \Lambda_{UV}^2} (|C^{ds}|^2 + |C_5^{ds}|^2)$$

- ▶ Only **1 hadronic charge**  $g_T$ :  $SU(3)_F$  + **LQCD** Gupta et al. PRD98,091501
- ▶ Best lab limit on  $s \rightarrow d\gamma'$  stems from  $\Xi^0 \rightarrow \Sigma^0 \gamma'$  and invisible BR of  $\Xi^0$

$$\text{BR}_{\text{inv}}(\Xi^0) \leq 2.3 \times 10^{-4} \quad \Rightarrow \quad \Lambda_{UV} \leq 3.8 \times 10^7 \text{ (95\% C.L.)}$$

# Contribution of $\Lambda \rightarrow nX^0$ to SN's dark luminosity

JMC, Pospelov, Vuong, Ziegler, Zupan PRD 102 (2020) 1, 015023, JMC, Terol-Calvo, Tolos & Ziegler, arXiv: 2012.11632 [hep-ph]

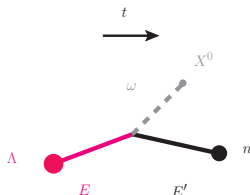
## • There are hyperons ( $\Lambda$ 's) in proto-neutron stars!

- ▶ Abundancies sustained by weak processes

$$pe^- \leftrightarrow \Lambda \nu_e, \quad \Lambda \rightarrow pe^- \bar{\nu}, \dots$$

- ▶ High Temperatures and supranuclear densities

**Thermal effects:**  $n_\Lambda \simeq n_n \exp\left(-\frac{m_\Lambda - m_n}{T}\right) \simeq 0.01 \times n_n$



- ▶ **Spectrum of dark cooling rate per unit volume**

$$\frac{dQ}{d\omega} = \frac{m_\Lambda^2 \Gamma \omega}{2\pi^2 \bar{\omega}} \int_{E_0}^{\infty} dE f_\Lambda (1 - f_n)$$

- ★  $\Gamma$ :  $\Lambda \rightarrow nX^0$  decay rate in vacuum
- ★  $f_B$  Fermi-Dirac distributions

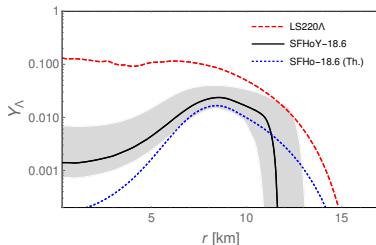
# The SN 1987A bound on $\Lambda \rightarrow nX^0$

“Raffelt’s Criterion” G. Raffelt’95

$$L_d \lesssim L_\nu = 3 \times 10^{52} \text{ erg s}^{-1}$$

- Recent SN simulations for SN 1987A + EoS from [compOSE database](#)

Bollig *et al.* PRL125(2020)051104



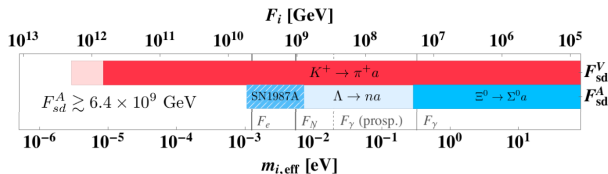
- Most conservative limit, including reabsorption/trapping, medium effects ...

$$\text{BR}(\Lambda \rightarrow nX^0) \lesssim 8.0 \times 10^{-9}$$

# Application of the SN bound

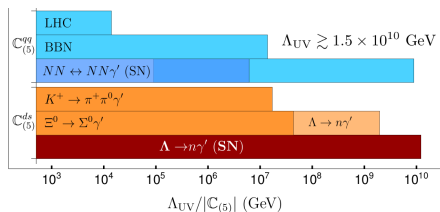
## ● Axions

- ▶ Define  $F_{sd}^{V,A} = 2f_a/c_{sd}^{V,A}$



- ▶ Best limit on axion models with **tuned axial couplings**
- ▶ Small change for ALP – **significant less tuning** for the  $K \rightarrow 2\pi$  region
- ▶ **BESIII projections** (Hai-bo's) lie within 1 order of magnitude from SN bound

## ● Massless dark photon



- ▶ Strongest absolute limit on the quark couplings of the massless dark-photon

# How well *should* we do experimentally?

- If we take the SN bound as a benchmark

Kaon and Hyperon decays

● SN:  $\Lambda/d_M^{sd} = 1.2 \times 10^7$  TeV

Decay	Branching Ratio	Decay	Branching Ratio
$K^+ \rightarrow \pi^0 \pi^+ A'$	$7.8 \times 10^{-11}$	$K_L \rightarrow \pi^0 \gamma A'$	$3.4 \times 10^{-11}$
$K^+ \rightarrow \pi^+ \gamma A'$	$1.9 \times 10^{-11}$	$K_L \rightarrow \pi^+ \pi^- A'$	$3.2 \times 10^{-10}$
$K_L \rightarrow \gamma A'$	$4.0 \times 10^{-8}$		
$\Sigma^+ \rightarrow p A'$	$4.9 \times 10^{-10}$	$\Lambda \rightarrow n A'$	$8.0 \times 10^{-9}$
$\Xi^0 \rightarrow \Sigma^0 A'$	$3.02 \times 10^{-9}$	$\Xi^- \rightarrow \Sigma^- A'$	$3.6 \times 10^{-9}$
$\Xi^0 \rightarrow \Lambda A'$	$1.4 \times 10^{-9}$	$\Omega^- \rightarrow \Sigma^- A'$	$3.6 \times 10^{-8}$

- ▶ Using current notation in white book
- ▶ Predictions for kaons and  $\Omega^-$  decay from [Su & Tandean, PRD101, 035044\(2020\)](#)

**Any projection one order of magnitude short of this would be interesting**

## How well can we do experimentally?

- Our recast of **Hai-bo's projections**  $\text{BR}(B_1 \rightarrow B_2 X^0) \simeq \text{BR}(B_1 \rightarrow B_2 \nu \bar{\nu})$  is **wrong**
  - ▶ For 2-body with massless  $X^0$  we need to subtract SM's  $B_1 \rightarrow B_2 \gamma$  (missing  $\gamma$ )
- **Idea:** Exploit polarized  $B_1$  and **experimental**  $\alpha_\gamma$  parameter

$$\frac{d\Gamma}{d\cos\theta} = \frac{\Gamma}{2}(1 + P\alpha_\gamma \cos\theta)$$

where

$$\alpha_\gamma = \frac{2\text{Re}(SP^*)}{|S|^2 + |P|^2}$$

measures interference between parity conserving and violating conts. to amplitude

Formulas for axions/Dark photons

**Useful for realistic experimental projections at BESIII & SCTF?**

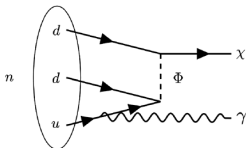
# Dark baryons Alonso-Alvarez, Elor, Escudero, Fornal, Grinstein, JMC – WIP

- **Dark particles carrying baryon number  $\Rightarrow$  Hyperon BNV signatures**

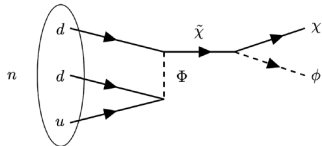
- ▶  $m_\chi > m_p$  to avoid proton decay!

- ① **Models addressing *neutron lifetime anomaly*** Fornal&Grinstein, PRL120,191801(2018)

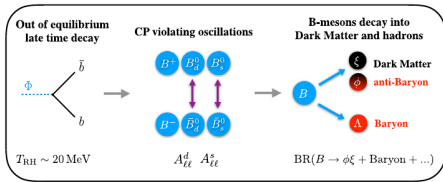
- ▶ Expt signature with one SM particle



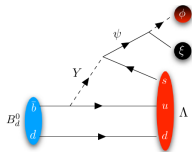
- ▶ Expt signature purely invisible



- ② **B-mesogenesis** Elor, Escudero, Nelson, PRD99(2019)3,035031, Alonso-Alvarez, Elor, Escudero, arXiv: 2101.02706



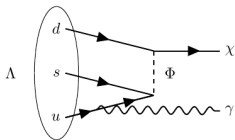
- ▶ **Expt signature: *B* BNV decays**



- **Other ideas in this direction:** McKeen&Pospelov arXiv: 2003.09865, McKeen, Pospelov&Raj arXiv: 2012.09865, ...

# The dark baryon - hyperon connection

- These models can have flavor entries triggering hyperon BNV signatures!



- However, there is generally not a one-to-one correspondence ...

## 1 Identify mediators

$$\mathcal{L}_1 \supset -y_{d_a d_b} \epsilon_{ijk} \Psi^i d_{Ra}^j d_{Rb}^k - y_{\chi u_c} \Psi_i^* \chi u_{Rc}^i + \text{h.c.}$$

$$\mathcal{L}_2 \supset -y_{u_a d_b} \epsilon_{ijk} \Phi^i u_{Ra}^j d_{Rb}^k - y_{\chi d_c} \Phi_i^* \chi d_{Rc}^i + \text{h.c.}$$

$$\mathcal{L}_3 \supset -y_{Q_a d_b} \epsilon_{ijk} \epsilon_{\alpha\beta} X_\mu^{i\alpha} Q_{La}^{j\beta} \gamma^\mu d_{Rb}^k - y_{\chi Q_c} X_\mu^{i\alpha} Q_{Lc}^{i\alpha} \gamma^\mu \chi + \text{h.c.},$$

$$\Psi = (3, 1)_{\frac{2}{3}}$$

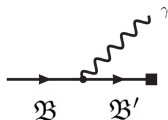
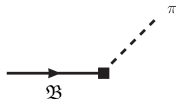
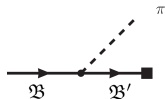
$$\Phi = (3, 1)_{-\frac{1}{3}}$$

$$X_\mu = (3, 2)_{\frac{1}{6}}$$

## 2 Match to $SU(2) \times U(1)$ and $U(1)_{\text{em}}$ EFTs

- ★ Fierz transformations, RGEs, etc [Weinberg PRL43,1566 \(1979\)](#), [Abbot&Wise PRD22,9,22 \(1980\)](#)

## 3 Match low-E EFT to chiral Lagrangian and calculate matrix elements



- ★ Leading order parameters obtained from LQCD calculations for proton decay!



# Experimental hyperon BNV signatures

## • Simplest experimental signatures

(1)  $\mathfrak{B} \rightarrow \text{inv}$  (2)  $\mathfrak{B} \rightarrow \gamma + \text{inv}$  (3)  $\mathfrak{B} \rightarrow \pi + \text{inv}$

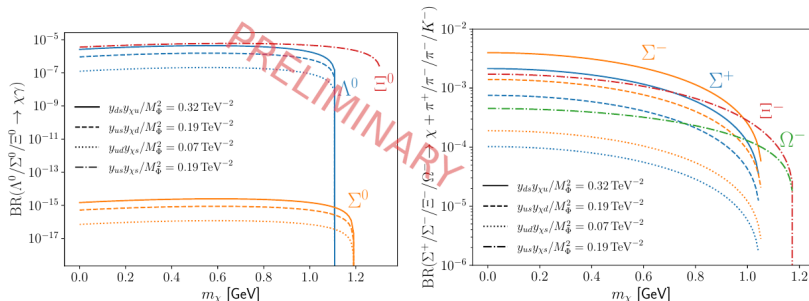
Initial State	Final State	$\Delta M$ (MeV)	Operator	$W_{RR}/\text{GeV}^3$	$\text{Br} \times \left[ \frac{M_\chi^2}{y^2} \frac{1}{5\text{TeV}^2} \right]^2$	Br sensitivity
$\Lambda^0 (uds)$	$\chi + \gamma$	1116	$\chi_{uds}$	0.01445	$7.2 \times 10^{-7}$	
$\Lambda^0 (uds)$	$\xi + \phi$	1116	$\chi_{uds}$	0.01445	$5.9 \times 10^{-3} \mu_{\xi\phi}^2$	
$\Sigma^+ (uus)$	$\chi + \pi^+ (d\bar{u})$	1050	$\chi_{uds}$	0.122	$4.7 \times 10^{-5}$	
$\Sigma^- (dds)$	$\chi + \pi^- (\bar{u}d)$	1058	$\chi_{uds}$	0.122	$9.8 \times 10^{-5}$	
$\Xi^0 (uss)$	$\chi + \gamma$	1315	$\chi_{uss}$	0.01445	$5.2 \times 10^{-6}$	
$\Xi^0 (uss)$	$\xi + \phi$	1315	$\chi_{uss}$	0.01445	$1.4 \times 10^{-3} \mu_{\xi\phi}^2$	
$\Xi^- (dss)$	$\chi + \pi^- (\bar{u}d)$	1182	$\chi_{uss}$	0.122	$3.1 \times 10^{-4}$	
$\Omega^- (sss)$	$\chi + K^- (\bar{u}s)$	1179	$\chi_{uss}$	-0.085	$1.4 \times 10^{-4}$	

- ▶ Finishing/checking calculation of matrix elements (**to happen soon!**)

## • Other signatures calculable with chiral Lagrangians

- ▶ **Two pions:** e.g.  $\Xi^- \rightarrow \pi^0 \pi^- + \text{inv}$
- ▶ **Two leptons:** e.g.  $\Xi^0 \rightarrow e^+ e^- + \text{inv}$
- ▶ **Two leptons and one pion:** e.g.  $\Xi^- \rightarrow \pi^- e^+ e^- + \text{inv}$  (LHCb!)

# Type of results to appear in the white book



- Interplay with colliders, flavor observables, SN bound will be partially covered

## Question for experimentalists

Which BNV modes could be measured and for which could we get **BESIII/SCTF** and **LHCb** projections?

## Solid progress towards completion of baryon pheno section

### 1 Heavy New Physics: Constraining axial currents

- ▶  $B_1 \rightarrow B_2 \nu \bar{\nu}$  interesting if order-of-magnitude improvement over Hai-Bo's projections
- ▶ We can provide  $K \rightarrow \pi \pi \nu \bar{\nu}$  predictions

### 2 2-body decays

- ▶ Formulae and form factors for baryonic decays for **axions** and **massless dark photon**
- ▶ Axial current for **axion** and (possibly) **best limits** for **dark photon**
- ▶ **SN bound** (should be taken as reference)
- ▶ Decays of polarized baryons to help beating backgrounds?
- ▶ @Marco Specific proposals to try converge and finish **dark photon section**

### 3 Dark baryons

- ▶ **Only accessible with hyperon decays:** Fullfledged theoretical&experimental program
- ▶ **To be written soon!!**

### What is missing?

- 1 Selection of experimental projections
- 2 Other models in interplay with kaons?