

# RPV Decays of Chargino and Neutralino LSPs in the B-L MSSM

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- The  $B - L$  extension of the MSSM arises naturally in the context of heterotic  $E_8 \times E_8$  string theory; [[Ambroso, Ovrut, arXiv:1005.5392](#)].
- $SO(10)$  unification at  $M_U \approx 3 \times 10^{16} \text{ GeV}$ ;
- Two Wilson lines break  $SO(10)$  to  $SU(3)_C \otimes SU(2)_L \otimes U(1)_{T_{3R}} \otimes U(1)_{B-L}$  gauge group below the Calabi-Yau scale;
- Introduces 3 species of **right-handed neutrinos** for an anomaly-free theory;
- $B - L$  broken by a right-handed sneutrino VEV.

$$U(1)_{3R} \otimes U(1)_{B-L} \rightarrow U(1)_Y,$$

and the  $Z_R$  boson is produced (constraints lift  $B - L$  breaking scale  $> 4.1 \text{ TeV}$ ).

The spectrum is exactly that of the MSSM with three extra right-handed neutrino chiral multiplets (one per family).

- The **superpotential** of the  $B - L$  MSSM is given by

$$W = Y_u QH_u u^c - Y_d QH_d d^c - Y_e LH_d e^c + Y_\nu LH_u \nu^c + \mu H_u H_d .$$

- The **soft supersymmetry breaking** Lagrangian is given by

$$\begin{aligned} -\mathcal{L}_{\text{soft}} = & \left( \frac{1}{2} M_3 \tilde{g}^2 + \frac{1}{2} M_2 \tilde{W}^2 + \frac{1}{2} M_R \tilde{W}_R^2 + \frac{1}{2} M_{BL} \tilde{B}^2 \right. \\ & \left. + a_u \tilde{Q} H_u \tilde{u}^c - a_d \tilde{Q} H_d \tilde{d}^c - a_e \tilde{L} H_d \tilde{e}^c + a_\nu \tilde{L} H_u \tilde{\nu}^c + b H_u H_d + h.c. \right) \\ & + m_{\tilde{Q}}^2 |\tilde{Q}|^2 + m_{\tilde{u}^c}^2 |\tilde{u}^c|^2 + m_{\tilde{d}^c}^2 |\tilde{d}^c|^2 + m_{\tilde{L}}^2 |\tilde{L}|^2 + m_{\tilde{\nu}^c}^2 |\tilde{\nu}^c|^2 \\ & + m_{\tilde{e}^c}^2 |\tilde{e}^c|^2 + m_{H_u}^2 |H_u|^2 + m_{H_d}^2 |H_d|^2 . \end{aligned}$$

No CP-violation  $\longrightarrow$  real parameters

Problem: Most general MSSM superpotential allows for baryon and lepton number violating terms at tree level and, therefore, **rapid proton decay**.

- Typical solution: add an *ad-hoc* R-parity defined as

$$R = (-1)^{3(B-L)+2s}.$$

- However, the lightest supersymmetric particle (LSP) is stable and therefore must be neutral to avoid a disallowed density of charged relics  $\Rightarrow$  R-parity severely narrows the SUSY phenomenological landscape.
- Our  $B - L$  MSSM model:  $B - L$  arises naturally from string theory; a third generation sneutrino acquires a VEV and hence  $B - L$  is broken at low energy ( $U(1)_{3R} \otimes U(1)_{B-L} \rightarrow U(1)_Y$ ), with small RPV couplings.

After the right handed sneutrino acquires a VEV  $v_R$ , there are two sources of RPV couplings in the Lagrangian:

- The RPV operators in the superpotential:

$$W \supset \epsilon_i e_i H_u^+ - \frac{1}{\sqrt{2}} Y_{ei} v_{Li} H_d^- e_i^c ,$$

where

$$\epsilon_i \equiv \frac{1}{\sqrt{2}} Y_{\nu i3} v_R, \quad v_{Li} = \frac{\frac{v_R}{\sqrt{2}} (Y_{\nu i3}^* \mu v_d - a_{\nu i3}^* v_u)}{m_{L_i}^2 - \frac{g_2^2}{8} (v_u^2 - v_d^2) - \frac{g_{BL}^2}{8} v_R^2} .$$

- The RPV operators in the super-covariant derivatives:

$$\begin{aligned} \mathcal{L} \supset & -\frac{1}{2} v_{Li}^* \left[ g_2 \left( \sqrt{2} e_i \tilde{W}^+ + \nu_{L_i} \tilde{W}^0 \right) - g_{BL} \nu_{L_i} \tilde{B}' \right] \\ & - \frac{1}{2} v_R \left[ -g_R \nu_3^c \tilde{W}_R + g_{BL} \nu_3^c \tilde{B}' \right] + \text{h.c.} \end{aligned}$$

# Chargino and Neutralino States

The RPV couplings  $\epsilon_i$  and  $\nu_{L_i}$  extend the chargino and neutralino gauge eigenstates:

$$\psi^+ = (\tilde{W}^+, \tilde{H}_u^+, \mathbf{e}^c_i) \quad \psi^- = (\tilde{W}^-, \tilde{H}_d^-, \mathbf{e}_i)$$
$$\psi^0 = (\tilde{W}_R, \tilde{W}_0, \tilde{H}_d^0, \tilde{H}_u^0, \tilde{B}', \nu_3^c, \nu_1, \nu_2, \nu_3)$$

These states can be combined into mass terms inside the Lagrangian of the form

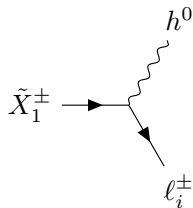
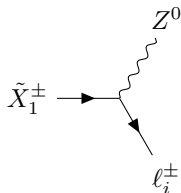
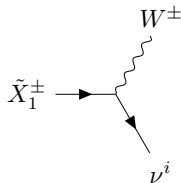
$$\mathcal{L} \supset -\frac{1}{2} (\psi^0)^T \mathcal{M}_{\tilde{\chi}^0} \psi^0 - \frac{1}{2} (\psi^+ \psi^-) \begin{pmatrix} 0 & \mathcal{M}_{\tilde{\chi}^\pm}^T \\ \mathcal{M}_{\tilde{\chi}^\pm} & 0 \end{pmatrix} \begin{pmatrix} \psi^+ \\ \psi^- \end{pmatrix} + h.c.$$

The mass eigenstates  $\chi^\pm$ ,  $\chi^0$  are related to the gauge eigenstates  $\psi^\pm$ ,  $\psi^0$  by the unitary matrices  $\mathcal{V}$ ,  $\mathcal{U}$ ,  $\mathcal{N}$ , which diagonalize the mass mixing matrices:

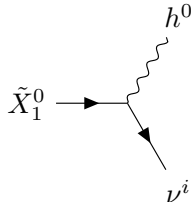
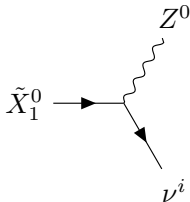
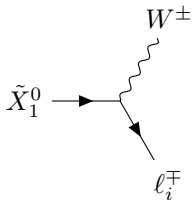
$$\mathcal{M}_{\tilde{\chi}^\pm}^D = \mathcal{U}^* \mathcal{M}_{\tilde{\chi}^\pm} \mathcal{V}^{-1} \quad \mathcal{M}_{\tilde{\chi}^0}^D = \mathcal{N}^* \mathcal{M}_{\tilde{\chi}^0} \mathcal{N}^\dagger,$$
$$\Rightarrow \tilde{\chi}^+ = \mathcal{U} \psi^+, \quad \chi^- = \mathcal{V} \psi^-, \quad \tilde{\chi}^0 = \mathcal{N} \psi^0.$$

# Chargino Decay Channels

- After diagonalization:  $\tilde{\chi}^\pm = (\tilde{\chi}_1^\pm, \tilde{\chi}_2^\pm, e_1^\pm, e_2^\pm, e_3^\pm)$ . 3 possible RPV decay modes:



- After diagonalization:  $\tilde{\chi}^0 = (\tilde{\chi}_1^0, \tilde{\chi}_2^0, \tilde{\chi}_3^0, \tilde{\chi}_4^0, \tilde{\chi}_5^0, \tilde{\chi}_6^0, \nu_1, \nu_2, \nu_3)$ . 3 possible RPV decay modes:





# Simulation - RGE run initial input

- The decay rates have complicated analytic expressions, which depend on the matrices  $\mathcal{U}$ ,  $\mathcal{V}$  and  $\mathcal{N}$ , the RPV couplings  $\epsilon_i$  and  $v_L^i$ , as well as other mass parameters and gauge couplings.
- A computer simulation [Ovrut, Purves, Spinner arXiv:1412.6103] allows us to compute the LSP decay rates numerically and analyze which decay channel dominates statistically.
- We scatter all 24 dimensionful soft supersymmetry breaking parameters in the interval, with a **log-uniform** distribution.

$$\left[ \frac{M}{f}, Mf \right] \quad \text{where} \quad M = 1.5 \text{ TeV}, \quad f = 6.7.$$

Parameter	Range	Prior
$m_{\tilde{q}_1} = m_{\tilde{q}_2}, m_{\tilde{q}_3} : \tilde{q} = \tilde{Q}, \tilde{u}^c, \tilde{d}^c$	(224GeV–10TeV)	log
$m_{\tilde{l}_1}, m_{\tilde{l}_2}, m_{\tilde{l}_3} : \tilde{l} = \tilde{L}, \tilde{e}^c, \tilde{\nu}^c$	(224GeV–10TeV)	log
$m_{H_u}, m_{H_d}$	(224GeV–10TeV)	log
$ A_f  : f = t, b, \tau$	(224GeV–10TeV)	log
$ M_a  : a = R, BL, 2, 3$	(224GeV–10TeV)	log
$\tan \beta = \frac{v_u}{v_d}$	(1.2-65)	flat
Sign of $\mu, a_f, M_a : f = t, b, \tau \quad a = R, BL, 2, 3$	[-,+]	flat

# Simulation: Experimental Constraints - RGE flow

We run an RGE analysis to simulate possible initial conditions in agreement with current experimental constraints.

- ✓  $B - L$  gauge symmetry breaking

$$M_{Z_R} > 4.1 \text{ TeV}$$

- ✓ EW symmetry breaking

$$M_{Z^0} = 91.1876 \pm 0.0021 \text{ GeV}, \quad M_{W^\pm} = 80.379 \pm 0.012 \text{ GeV}$$

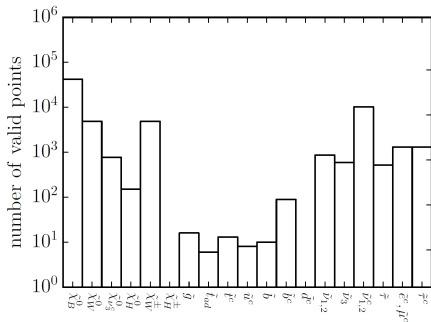
- ✓ Sparticle masses

SUSY Particle	Lower Bound
Left-handed sneutrinos	45.6 GeV
Charginos, sleptons	100 GeV
Squarks, except stop or bottom LSP	1000 GeV
Stop LSP (admixture)	550 GeV
Stop LSP (right-handed)	400 GeV
Sbottom LSP	500 GeV
Gluino	1300 GeV

- ✓ Higgs boson mass

$$M_{h^0} = 124.97 \pm 0.72 \text{ GeV}$$

# Final Particle Spectra



A histogram of the LSPs associated with a random scan of 100 million initial data points.

stop/ sbottom LSP

<100 points

[Marshall,Ovrut,Purves, arXiv:1401.7989]

Wino Chargino LSP

4,869 points

Wino Neutralino LSP

4,858 points

[Dumitru,Ovrut,Purves, arXiv:1811.05581]

Bino Neutralino LSP

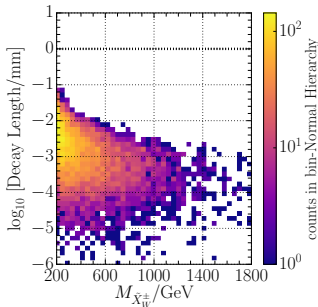
42,039 points

[Dumitru,Herwig,Ovrut, arXiv:1906.03174]

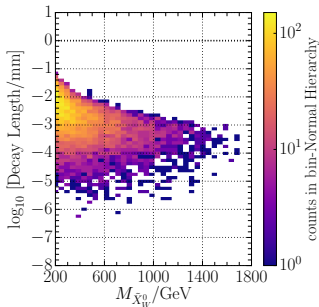
# LSP RPV Decays - Promptness

A decay process is **prompt** if its decay length

$$L = c \times \frac{1}{\Gamma} < 1 \text{ mm} .$$



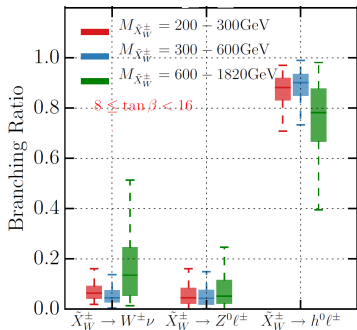
(a) Wino chargino LSP



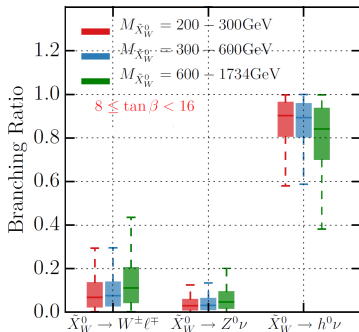
(b) Wino neutralino LSP

Scans of the Wino chargino and Wino neutralino LSP decay lengths in millimeters, for the normal hierarchy.

# Branching Ratios



Branching ratios for the three possible decay channels of the Wino charged LSP.



Branching ratios for the three possible decay channels of the Wino neutralino LSP.

# Conclusions

- In the B-L MSSM, the third generation sneutrino VEV breaks R-parity at TeV scale  $\rightarrow$  the LSP can decay into SM particles.
- RGE flow simulation - SUSY can exist just above the TeV scale, in agreement with most current experimental bounds. [Ovrut, Purves, Spinner [arXiv:1412.6103](#)]
- The Wino chargino  $\tilde{X}_W^\pm$  has three RPV decay channels:  $W^\pm \nu$ ,  $Z^0 \ell^\pm$ ,  $h^0 \ell^\pm$ . The Wino neutralino  $\tilde{X}_W^0$  has three RPV decay channels:  $W^\pm \ell^\mp$ ,  $Z^0 \nu$ ,  $h^0 \ell^\nu$ . [Dumitru, Ovrut, Purves, [arXiv:1810.11035](#)]
- The Wino chargino and Wino neutralino can have masses between 200-1800 GeV and can decay promptly at LHC, into SM particles  $\rightarrow$  possible detection in the near future. [Dumitru, Ovrut, Purves, [arXiv:1811.05581](#)]