

Magnetic charge and Massive Photon*

APS DPF Meeting

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Maxwell's Equations for Massive Photons

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- In this talk we consider two changes to Maxwell's equations.
- First for a massive photon one has

$$\begin{aligned}\nabla \cdot \mathbf{E} + m^2\phi &= 4\pi\rho_{(e)} \quad ; \quad \nabla \times \mathbf{E} + \frac{\partial \mathbf{B}}{\partial t} = 0 \\ \nabla \cdot \mathbf{B} &= 0 \quad ; \quad \nabla \times \mathbf{B} - \frac{\partial \mathbf{E}}{\partial t} + m^2\mathbf{A} = 4\pi\mathbf{J}_{(e)}\end{aligned}$$

- Above equations not gauge invariant, unlike standard Maxwell
- Above equations have a reduced (gauge) symmetry
- For ordinary Maxwell $\mathbf{E} = -\nabla\phi - \partial_t\mathbf{A}$ & $\mathbf{B} = \nabla \times \mathbf{A}$ invariant under

$$\phi' \rightarrow \phi - \partial_t\lambda \quad ; \quad \mathbf{A}' \rightarrow \mathbf{A} + \nabla\lambda$$

Electrostatic Interaction for Massive Photons

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- If the photon is massive one finds

$$\mathbf{E}_{Yukawa} = \frac{qe^{-mr}\hat{r}}{r^2}(1 + mr) \quad ; \quad \phi_{Yukawa} = \frac{qe^{-mr}}{r}$$

- The Yukawa field, Yukawa potential, and Yukawa Force

$$\mathbf{F}_{Yukawa} = \frac{q_1q_2e^{-mr}\hat{r}}{r^2}(1 + mr)$$

- Experimentally $m < 10^{-18}$ eV (see Particle Data Handbook).

Maxwell's Equations Electric & Magnetic Charges

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- Maxwell's equations for Magnetic and Electric Charges are

$$\nabla \cdot \mathbf{E} = 4\pi\rho_{(e)} \quad ; \quad \nabla \times \mathbf{E} + \frac{\partial \mathbf{B}}{\partial t} = -4\pi\mathbf{J}_{(m)}$$

$$\nabla \cdot \mathbf{B} = 4\pi\rho_{(m)} \quad ; \quad \nabla \times \mathbf{B} - \frac{\partial \mathbf{E}}{\partial t} = 4\pi\mathbf{J}_{(e)}$$

- This makes Maxwell's equations symmetric under a dual rotation

$$\mathbf{E} = \mathbf{E}' \cos \theta + \mathbf{B}' \sin \theta \quad ; \quad \mathbf{B} = -\mathbf{E}' \sin \theta + \mathbf{B}' \cos \theta .$$

$$\rho_{(e)} = \rho'_{(e)} \cos \theta + \rho'_{(m)} \sin \theta \quad ; \quad \rho_{(m)} = -\rho'_{(e)} \sin \theta + \rho'_{(m)} \cos \theta .$$

$$\mathbf{J}_{(e)} = \mathbf{J}'_{(e)} \cos \theta + \mathbf{J}'_{(m)} \sin \theta \quad ; \quad \mathbf{J}_{(m)} = -\mathbf{J}'_{(e)} \sin \theta + \mathbf{J}'_{(m)} \cos \theta .$$

Magnetic Charges

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- For electric charge q , $\phi = \frac{q}{r}$ gives $\mathbf{E} = -\nabla\phi = \frac{q\hat{r}}{r^2}$
- For magnetic charge g one naively wants $\mathbf{B} = \nabla \times \mathbf{A} = \frac{g\hat{r}}{r^2}$
- For electric charge $\nabla \cdot \mathbf{E} = q \left[\nabla \cdot \left(\frac{\hat{r}}{r^2} \right) \right] = 4\pi q \delta^3(\mathbf{r}) = 4\pi \rho_{(e)}$
- For magnetic charge $\nabla \cdot \mathbf{B} = g \left[\nabla \cdot \left(\frac{\hat{r}}{r^2} \right) \right] = 4\pi g \delta^3(\mathbf{r}) = 4\pi \rho_{(m)}$
- But canonically $\nabla \cdot (\nabla \times \mathbf{A}) = 0$
- Bullet point 4 contradicts bullet point 5. Problem?

Dirac vector-potential for massless photon

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- Bullet point 4 from above realized by the Dirac String*

$$\mathbf{A}_{\pm} = \frac{g}{r} \left(\frac{\pm 1 - \cos \theta}{\sin \theta} \right) \hat{\varphi} \implies \mathbf{B} = \nabla \times \mathbf{A}_{\pm} = \frac{g \hat{r}}{r^2}$$

- The result is a Coulomb \mathbf{B} -field.
- \mathbf{A}_{+} singular at $(\theta = \pi)$; \mathbf{A}_{-} singular at $(\theta = 0)$.



* *P.A.M. Dirac, Proc. Roy. Soc. A 133, 60-72 (1931)*

String vector-potential for massless photon

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- Consistency with $\nabla \cdot (\nabla \times \mathbf{A}) = 0$ requires string contribution*

$$\mathbf{A}_{s\pm} = \pm \frac{2g\Theta(\rho)\Theta(\mp z)}{\rho} \hat{\varphi} \implies \mathbf{B}_{s\pm} = \frac{2g\delta(\rho)\Theta(\mp z)}{\rho} \hat{z}$$

- $\mathbf{A}_{s\pm}$ is solenoid carrying flux $4\pi g$
- Adding $\mathbf{A}_{s\pm}$ gives $\mathbf{B}_{\pm} = \nabla \times (\mathbf{A}_{\pm} + \mathbf{A}_{s\pm})$ and $\nabla \cdot \mathbf{B}_{\pm} = 0$.
- Dirac condition makes this string “invisible”.

* *R. Heras, Contemp. Phys.* **59**, 331-355 (2018).

Doubly Modified Maxwell's Equations

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- Doubly modified Maxwell \rightarrow photon mass + magnetic charge

$$\nabla \cdot \mathbf{E} + m^2 \phi = 4\pi \rho_{(e)}$$

$$\nabla \times \mathbf{B} - \frac{\partial \mathbf{E}}{\partial t} + m^2 \mathbf{A} = 4\pi \mathbf{J}_{(e)}$$

$$\nabla \cdot \mathbf{B} = 4\pi \rho_{(m)}$$

$$\nabla \times \mathbf{E} + \frac{\partial \mathbf{B}}{\partial t} = -4\pi \mathbf{J}_{(m)}$$

- Gauge symmetry spoiled; dual symmetry added

Magnetic charge plus photon mass I

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- Allow a photon mass **and** a magnetic charge.
- Accomplished tentatively by multiplying **A** by the e^{-mr}
- So, $\mathbf{A}_{\pm} = g \frac{e^{-mr}}{r} \left(\frac{\pm 1 - \cos \theta}{\sin \theta} \right) \hat{\varphi}^*$
and $\mathbf{B}_{\pm} = \nabla \times \mathbf{A}_{\pm} = g \frac{e^{-mr}}{r^2} \left(\hat{r} + mr \left[\frac{\pm 1 - \cos \theta}{\sin \theta} \right] \hat{\theta} \right)^*$

* T. Evans and D. Singleton, *Int. J. Mod. Phys. A* **33**, 1850064 (2018) &

M. Dunia, T. Evans, and D. Singleton, *Phys.Rev.D* **103**, 128501 (2021)

Magnetic charge plus photon mass II

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- The above solution $\mathbf{A}_{\pm}, \mathbf{B}_{\pm}$ corrects an earlier wrong solution*
- Our solution assumes $\mathbf{A} = e^{-mr} \mathbf{A}_{\text{Dirac}}$; earlier work assumed $\mathbf{A} = \mathbf{A}_{\text{Dirac}} + \mathbf{A}_{\text{mass}}$
- String singularity is real due to its presence in the magnetic field

* A. Y. Ignatiev and G. C. Joshi, *Phys. Rev. D* **53**, 984 (1996)

Magnetic charge plus photon mass III

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- However, as in the massless case, \mathbf{A}_{\pm} , \mathbf{B}_{\pm} modified so that $\nabla \cdot \mathbf{B}_{\pm} = 0$ & $\nabla \times \mathbf{B}_{\pm} + m^2 \mathbf{A}_{\pm} \neq 0$
- Introduce vector potential regularized by ϵ

$$\mathbf{A}_{\pm} = g \frac{\Theta(\rho - \epsilon) e^{-m\sqrt{\rho^2 + z^2}}}{\rho} \left(\pm 1 - \frac{z}{\sqrt{\rho^2 + z^2}} \right) \hat{\phi}$$

- Taking the curl and $\epsilon \rightarrow 0$ yields (with $\nabla \cdot \mathbf{B}_{\pm} = 0$)

$$\mathbf{B}_{\pm} = \frac{ge^{-mr}}{r^2} \left[\hat{r} + mr \left(\frac{\pm 1 - \cos \theta}{\sin \theta} \right) \hat{\theta} \right] \pm 4\pi g e^{-m|z|} \Theta(\mp z) \delta(x) \delta(y) \hat{z}$$

Dirac Quantization Condition

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- Amazing feature of magnetic charge – Dirac Condition

$$qg = n\frac{\hbar}{2}$$

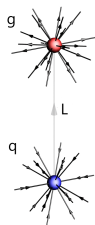
- Many Derivations (AB-Phase, Wu-Yang Fiber Bundle)
- Here we use the field angular momentum quantization method

Field Angular Momentum and Dirac Condition

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- An at rest $q + g$ system carries an E&M field angular momentum
- For massless photon, $\mathbf{L}_{EM} = \frac{1}{4\pi} \int \mathbf{r} \times (\mathbf{E} \times \mathbf{B}_{\text{Dirac}}) d^3x = qg\hat{\mathbf{r}}_0$
- Quantizing $|\mathbf{L}_{EM}| = n\frac{\hbar}{2}$ yields the Dirac condition $qg = \frac{n\hbar}{2}$
- \mathbf{L}_{EM} not dependent on r_0 [magnitude of \mathbf{r}_0]



Modified Field Angular Momentum

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- With a photon mass, field angular momentum becomes

$$\mathbf{L}_{EM} = \frac{1}{4\pi} \int \mathbf{r} \times \left[(\mathbf{E} \times \mathbf{B}) + m^2 \mathbf{A} \right] d^3x$$

- When $r_0 = 0$, this field angular momentum gives

$$\mathbf{L}_{EM} = qg\hat{z}$$

- String along z-axis; same quantization condition as Dirac!
- Here \mathbf{L}_{EM} comes from string, not the monopole itself.

Modified Dirac condition

- Now take q + string on z -axis and $r_0 \neq 0$
- Detailed calculation now gives E&M field angular momentum:
 - *String along $\mp z$ -axis, electric charge at $r = \mp r_0 \hat{z}$:*
$$L_{EM} = L_{EM}^{point} + L_{EM}^{\mathbf{B}^\theta + \mathbf{A}} = \pm 2qge^{-\frac{mr_0}{\hbar}} \hat{z}$$
 - *String along $\mp z$ -axis, electric charge at $r = \pm r_0 \hat{z}$:*
$$L_{EM} = L_{EM}^{point} + L_{EM}^{\mathbf{B}^\theta + \mathbf{A}} = 0$$
- Quantizing this angular momentum, give:
$$2qg \exp\left[-\frac{mr_0}{\hbar}\right] = \frac{n\hbar}{2} \text{ for all } n = 1, 2, 3, \dots$$
- Implies not only quantization of q and g but m and r_0 .

Summary and Conclusions

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- Combining magnetic charge plus photon mass leads to a Yukawa-Dirac String potential and magnetic field.
- Dirac string becomes a real object *i.e.* appears in **A** and **B**.
- Dirac quantization condition changes. Includes m and r_0 .

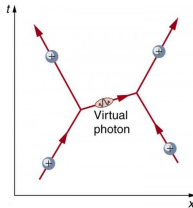


Quantum Picture for Interactions (additional material)

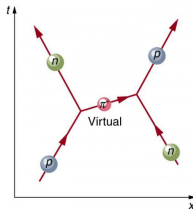
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- The Coulomb field described in QM via Feynman diagram.
- Exchanged force particle is massless.



- The Yukawa field described in QM via Feynman diagram.
- Exchanged force particle is massive.



Field Angular Momentum and Dirac Condition (Additional Material)

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- One can write the total angular momentum as the following

$$\mathbf{L}_{Tot} = \mathbf{L}_{Mech} + \mathbf{L}_{EM}, \text{ where } \mathbf{L}_{Mech} = \mathbf{r} \times (\mathbf{p} - q\mathbf{A})$$

- Yang* showed that the total angular momentum still satisfies

$$[\mathbf{L}_{Tot} \ i, \mathbf{L}_{Tot} \ j] = i\hbar\epsilon_{ijk}\mathbf{L}_{Tot} \ k$$

- Angular momentum commutator not satisfied by \mathbf{L}_{Mech} alone

* C.N. Yang, *Annals New York Academy of Sciences* **294**, 86 - 97 (1977)