

# Neutron Total Cross Section Measurement on Plastic Scintillator Using 3D Projection Scintillator Tracker Prototype

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# Joint T2K-DUNE 3D Scintillator R&D Group Institutions

CERN



University of Geneva, Switzerland



Louisiana State University, USA



Imperial college, UK



University of Pittsburgh, USA



University of Rochester, USA



Stony Brook University, USA



University of Tokyo, Japan



ETH Zurich, Switzerland



Chung-Ang University, South Korea



University of Pennsylvania, USA



High Energy Accelerator Research Organization (KEK), Japan



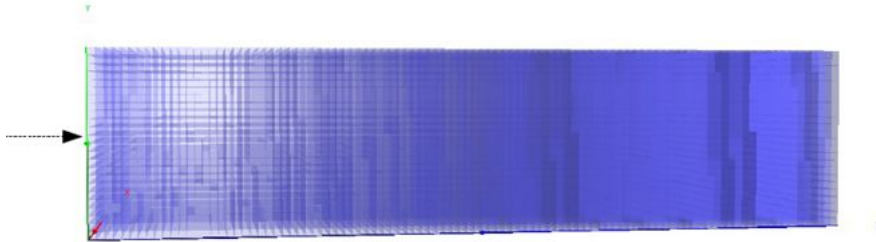
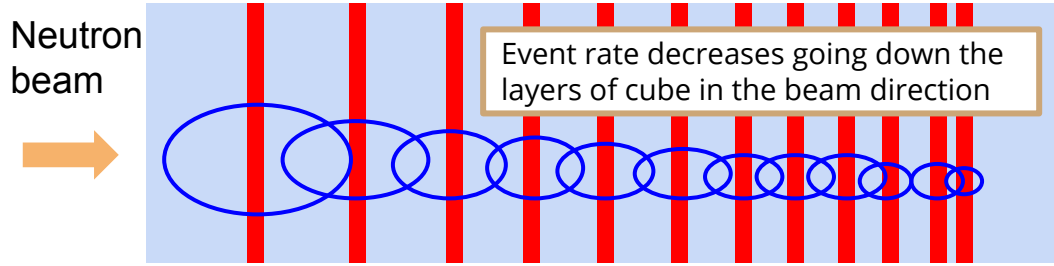
South Dakota School of Mines and Technology, USA



# Introduction

- Following the talk from Sitraka (South Dakota School of Mines and Technology), large amount of neutron beam data has been taken with two 3D projection scintillator tracker prototypes in 2019 and 2020 both in Los Alamos National Laboratory
- The main goal of the beam test is to characterize the neutron detection capability for SuperFGD and 3DST (3D projection scintillator tracker)  
=> a number of measurements can be done including neutron cross section, detection efficiency, threshold, double scattering signature etc

# Neutron Total Cross Section Measurement



$$N(l) = N_0 \cdot \exp(-T \cdot \sigma_{\text{total}} \cdot l)$$

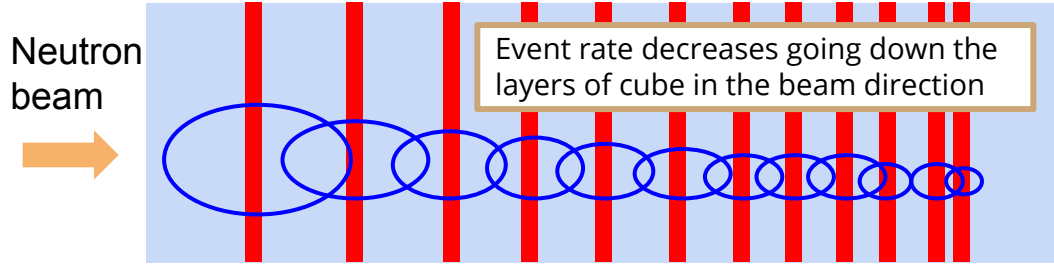
Nuclear density

total xsec

depth along the beam, i.e. layer

Measurement of event rate at each layer indicates a total cross section

# Neutron Total Cross Section Measurement



**Chosen event topology for the measurement: Single track events**

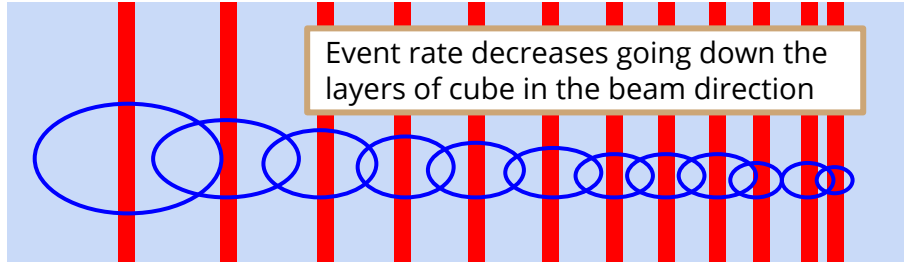
$$\begin{matrix} \text{Energy} & \text{Layer} \\ \nearrow & \nearrow \\ \mathbf{N}_{e,l} = \sum & \left( \begin{matrix} \mathbf{N}_{\text{single-track},e,l} \\ \mathbf{N}_{\text{invisible},e,l} \\ \mathbf{N}_{\text{two-track},e,l} \\ \dots \\ \mathbf{N}_{\text{100-track},e,l} \end{matrix} \right) = \sum \left( \begin{matrix} \mathbf{N}_{\text{single-track},e,l} \\ \mathbf{N}_{\text{single-track},e,l} \times \epsilon_e \end{matrix} \right)
 \end{matrix}$$

$\epsilon$  is the cross section  
 Ratio between "other-than single track" and single-track, it only depends on energy, regardless of layer

$$\begin{matrix} \text{Layer l} & \text{Layer m} \\ \nearrow & \nearrow \\ \mathbf{N}_{e,l} / \mathbf{N}_{e,m} = \mathbf{N}_{\text{single-track},e,l} / \mathbf{N}_{\text{single-track},e,m}
 \end{matrix}
 \quad \rightarrow \quad \text{Single track attenuation indicates a total cross-section}$$

# Neutron Total Cross Section Measurement

Neutron beam

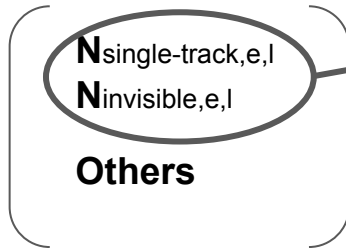


Event rate decreases going down the layers of cube in the beam direction

Presence of **invisible scattering of neutrons** in the detector:  
*Elastic and inelastic scattering of neutrons without visible energy deposit in the detector*

$$N_{e,l} = \sum$$

Energy Layer



They are always combined due to threshold and..  
 $N_{invisible} / N_{single-track}$  is **Layer-dependent !**

**A correction for the invisible scattering is needed and it should be data-based.**

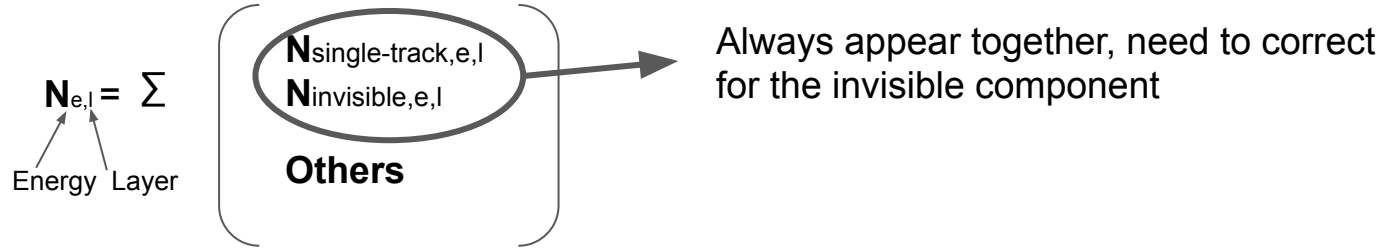
# Total Cross Section Measurement Systematics

- Interaction : invisible scattering
- Flux : energy scale uncertainty
- Detection : detector anisotropy due to geometry asymmetry and readout nonuniformity
- Reconstruction : selection-induced uncertainty

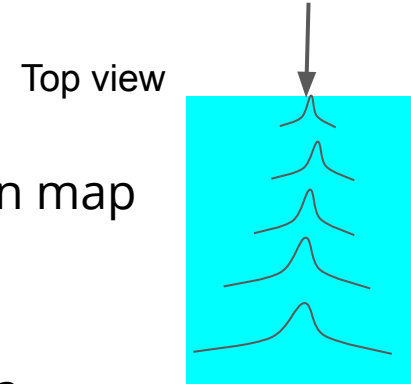


- Detector and reconstruction systematics will be combined
- A detector geometry correction will be applied

# Interaction: Invisible scattering

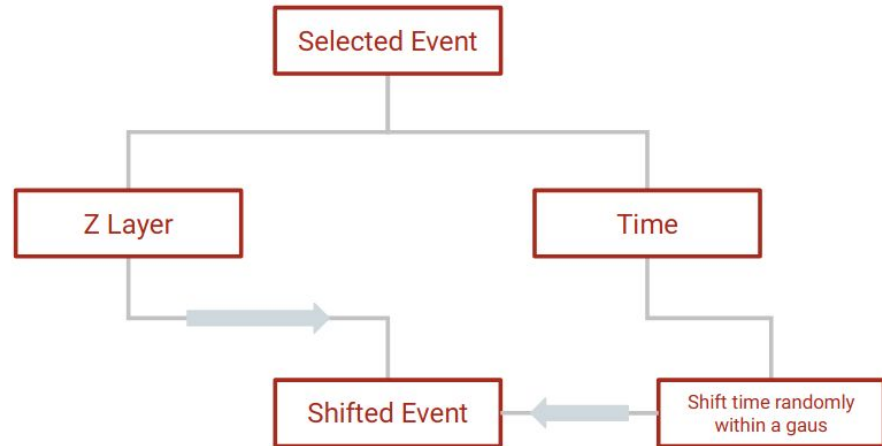


- Mainly due to threshold of the detector
- MC (Geant4 -> Bertini model) can provide a correction map to correct for the invisible component  
=> MC needs to be tuned to data
- Use transverse spread of events to tune MC with data



# Flux: Energy scale uncertainty

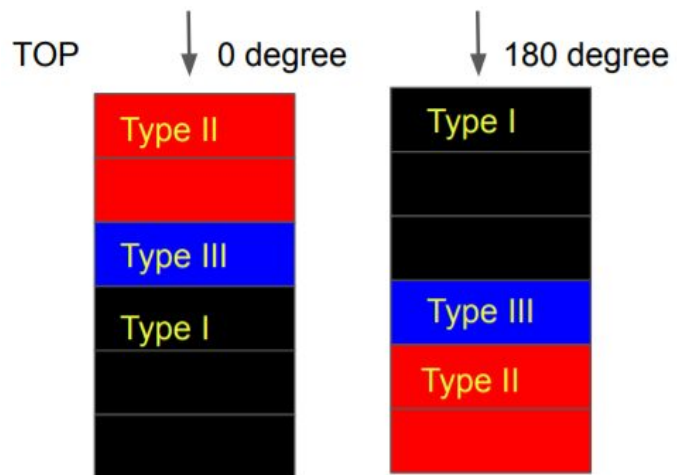
- Due to timing resolution ( $\sim 1.1$  ns [[arXiv:2008.08861](#)])
- Randomly shift the measured timing based on the time resolution, the resulting Z layer distribution for each energy range will have an uncertainty band



# Readout nonuniformity, geometry asymmetry and reconstruction

-> Work in progress

- Selection cuts optimized by minimizing the difference between 0 degree and 180 degree runs (rotated by Y)
- A geometry correction based on the proton MC will be applied:  
our detector is not geometrically symmetric
- We will use 0 degree configuration data for our first result (180 degree configuration data used for systematic evaluation)



# Cross Section Fitter: Workflow

- Fill an histogram (number of events vs z-layer) for every energy bin
- Correct this distribution for the invisible scattering
- Propagate the systematic in the following way:
  - Fit an exponential ( $N_0 e^{-\lambda z}$ ) for every variation of a systematic uncertainty
  - Fill an histogram with every value of the cross-section computed as  $\lambda/N_{\text{target}}$
  - The systematic uncertainty will be the RMS of this distribution
- Statistical uncertainty: Fit the nominal distribution and vary the number of events using a Poisson distribution

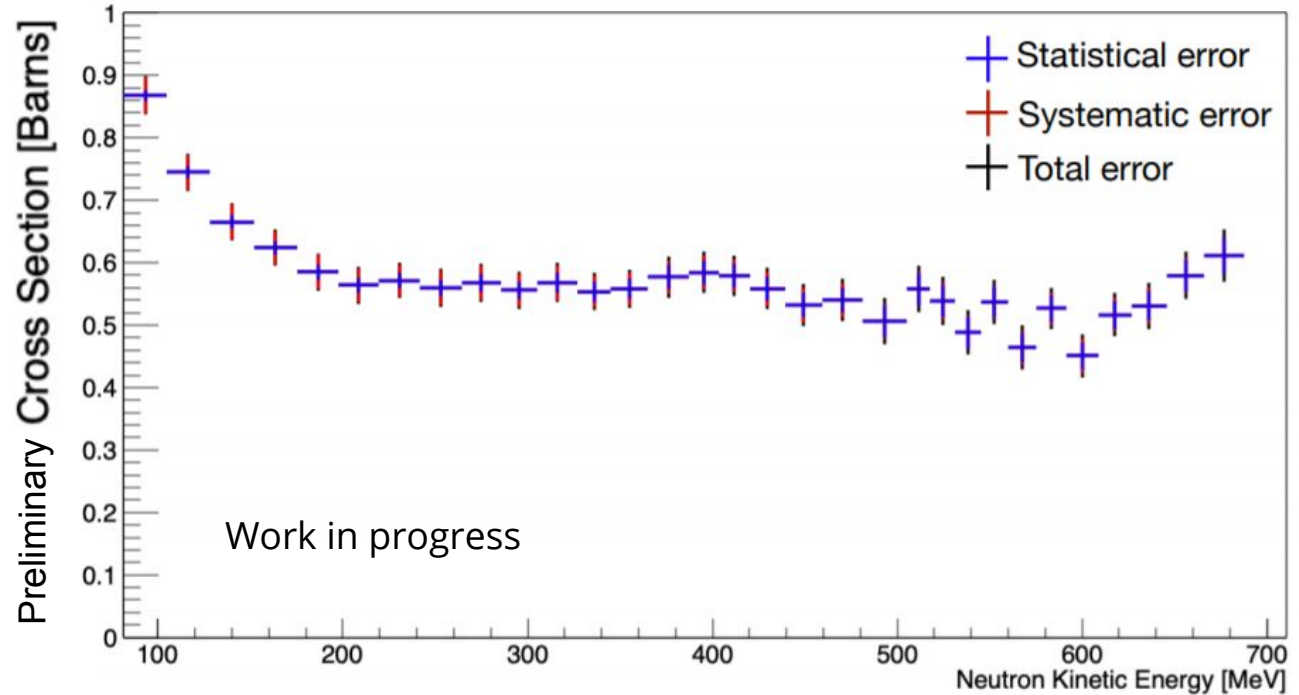
# Cross Section Fitter: Workflow

- The source of systematics uncertainties are:
  - Time resolution:  $\sim 1.1$  ns
  - Invisible scattering: 100%
  - Detector inhomogeneity: overall 10% normalization (preliminary)
- Time resolution will change the energy distributions
- Invisible scattering and detector inhomogeneity will vary the event normalization in each layer

# Cross Section Fitter: Workflow Test

Caveat:

Only 16 minutes  
of data used



# Summary

- Expected neutron detection potential is observed from preliminary analysis on the neutron data taken
- A number of results can be expected in the next few years
- A neutron total cross section measurement on plastic scintillator publication is expected by this summer

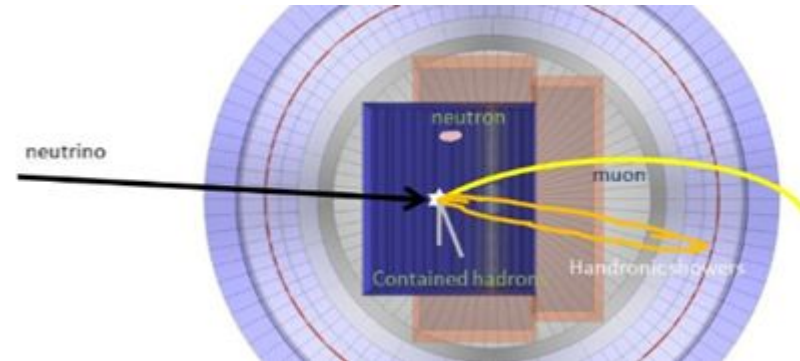


# Backup Slides



# Introduction

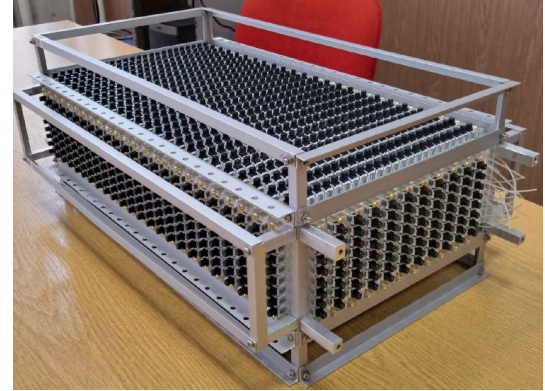
- Missing neutron energy is one of the dominant systematic uncertainties in long-baseline neutrino oscillation analyses.
- Neutron information desired in the near detectors of the long-baseline experiments in the precision era.
- Neutrons kinetic energy can be measured in 3DST (3D Projection Scintillator Tracker) with time-of-flight technique.
- Neutron detection capability, efficiency, backgrounds etc. can be studied with analysis on the neutron beam test data taken in LANL.



# Prototype Detectors

- SuperFGD prototype being used for the charged particle beam test at CERN and both 2019 + 2020 neutron beam tests (24 x 8 x 48)
- US-Japan prototype being fully used for the 2020 beam test (8 x 8 x 32)

superFGD

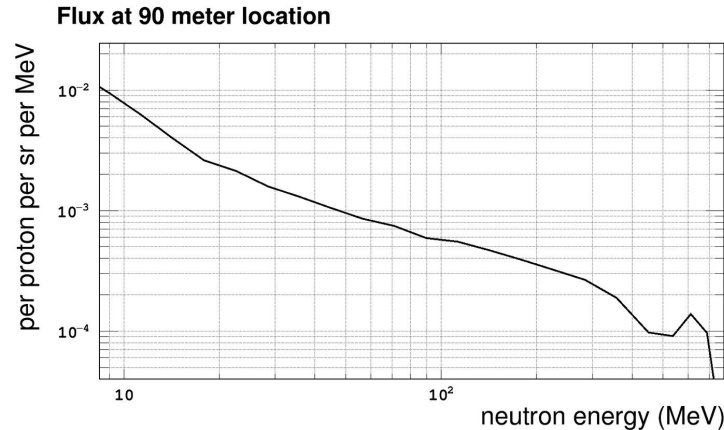
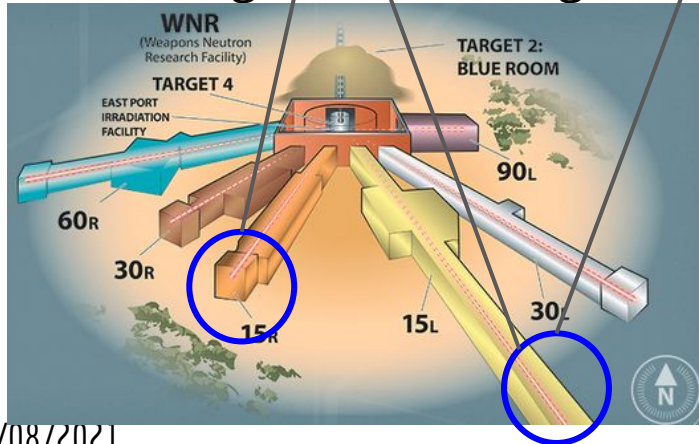


US-Japan



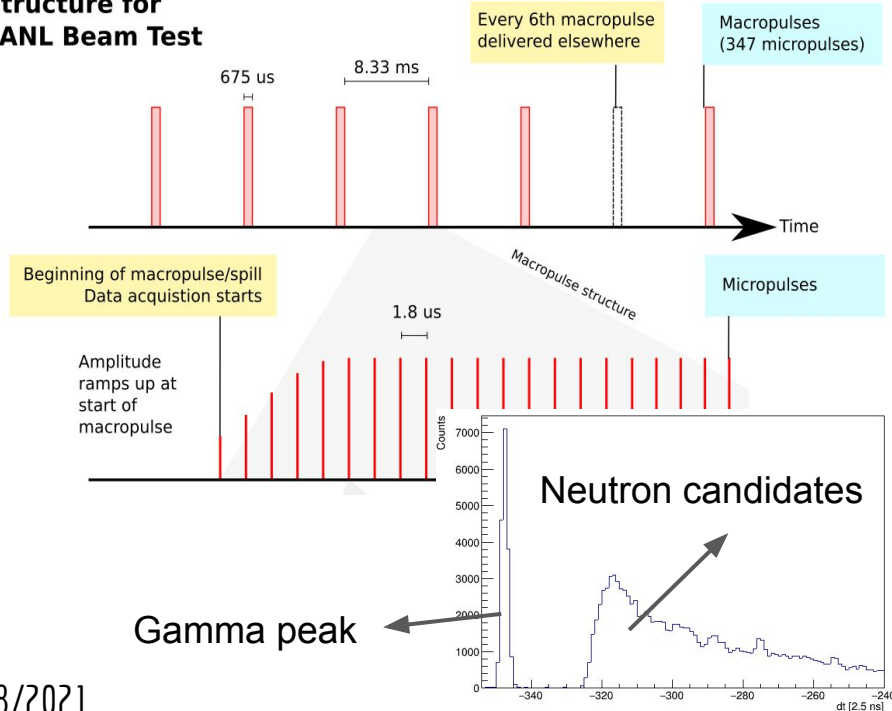
# Neutron Beam Test Facility

- Los Alamos National Lab LANCSE facility provides neutron beam ranged from 0 - 800 MeV.
- 2019: 15R 20 m 3 days + 15L 90 m 2 weeks
- 2020: 15L 90 m 2 weeks (various collimator, pulse spacing, detector configuration settings.)



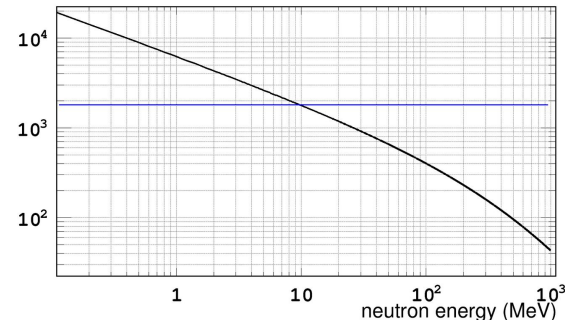
# Neutron Beam Time Structure

## Structure for LANL Beam Test

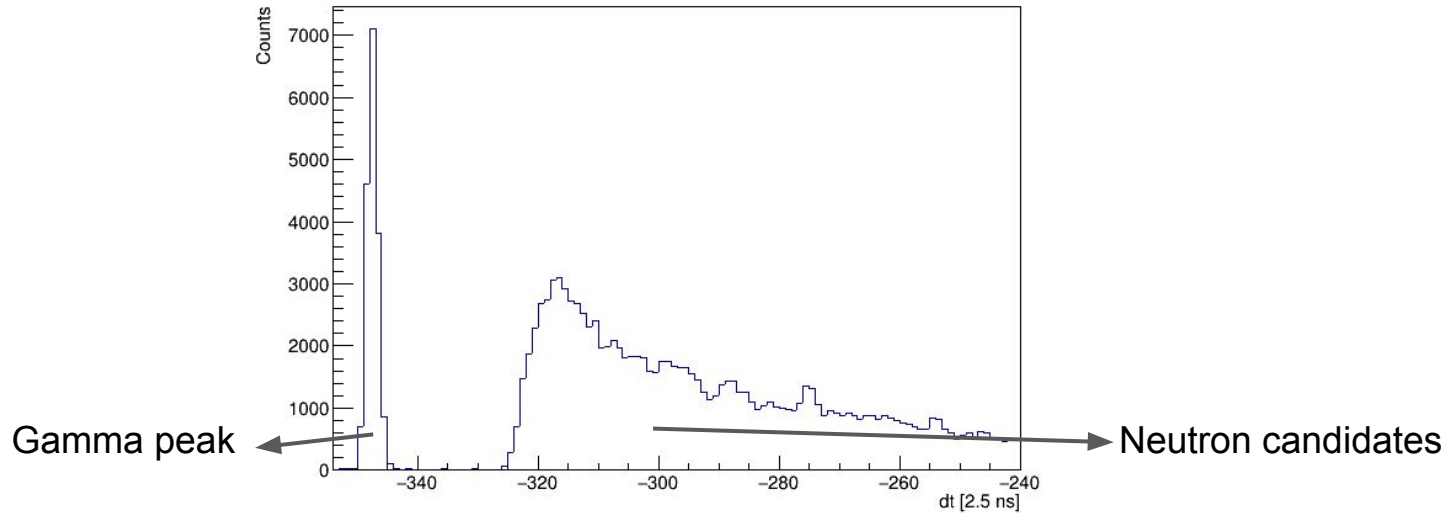


- Beam pulse is narrow enough to provide single neutron energy by time-of-flight
- We have 675  $\mu$ s trigger window to cover each macropulse
- Gamma flash and t0 are available for micropulses

neutron energy vs time diff. at 90 m location

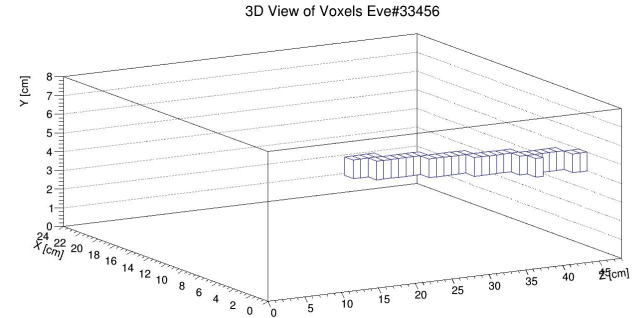
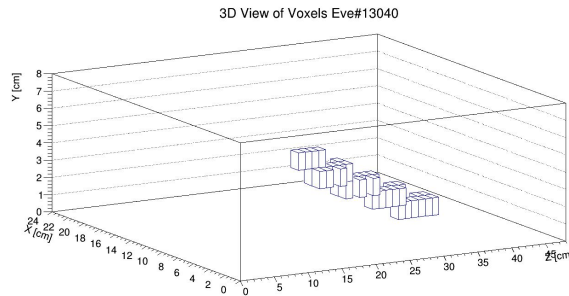
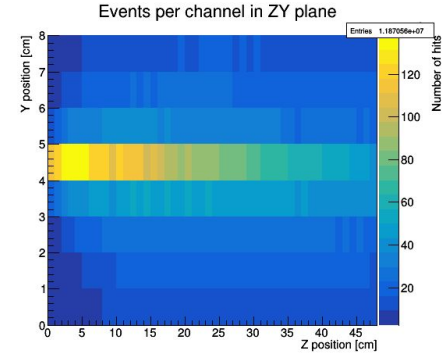
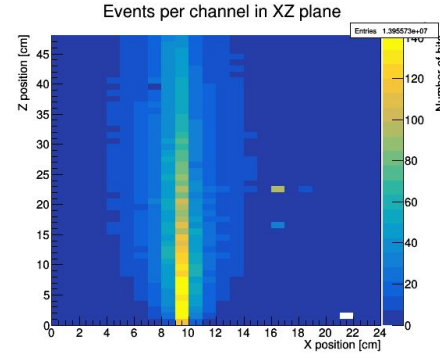
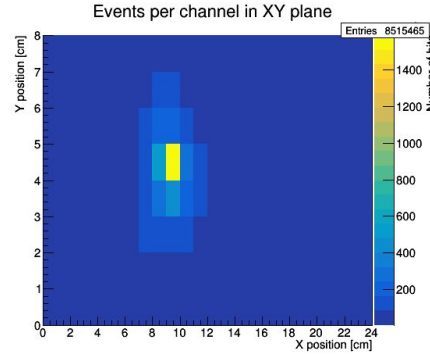


# Neutron Beam Data



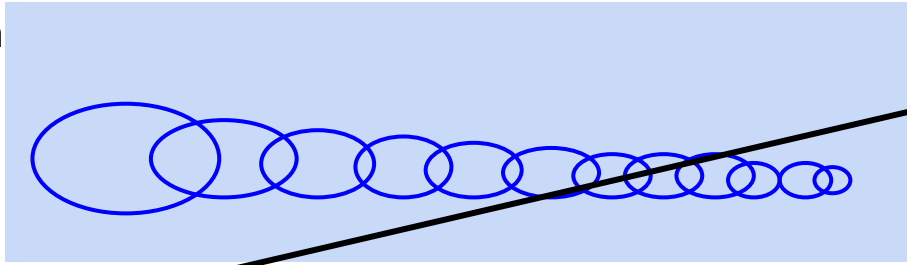
# Neutron Beam Data

- Neutron beam observed in SuperFDG prototype detector
- Neutron events in SuperFDG prototype detector



# Neutron Total Cross Section Measurement

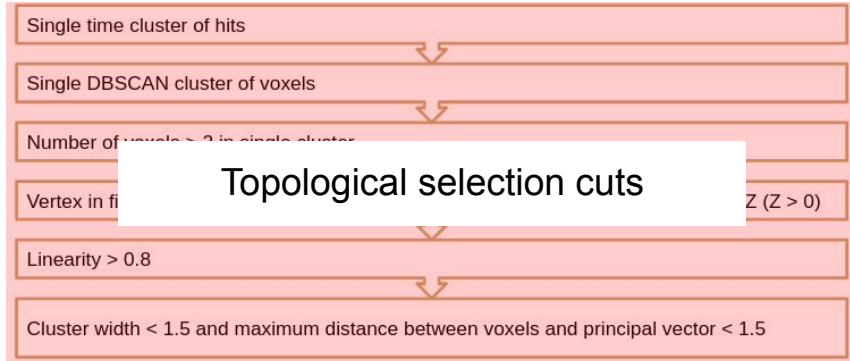
Neutron beam



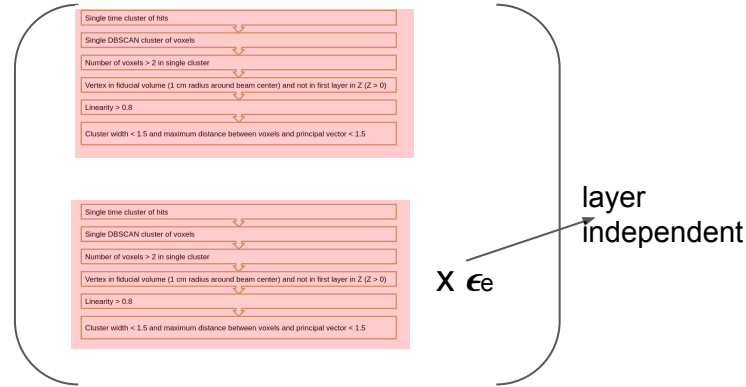
$$N_{e,l} = \sum \begin{pmatrix} N_{\text{single-track},e,l} \\ N_{\text{invisible},e,l} \\ N_{\text{two-track},e,l} \\ \dots \\ N_{\text{100-track},e,l} \end{pmatrix}$$

Energy Layer

$N_{\text{single-track},e,l}$  can be any topological selection, e.g



Blue arrow pointing to  $N_{e,l} = \sum$



# Principal Component Analysis (PCA) Overview

- Calculate the centroid for a distribution of points
- Calculate the covariance matrix with the centroid

$$[Cov]_{ij} = \frac{\sum_{i=1}^N (A_i - \bar{A}_i) \cdot (A_j - \bar{A}_j)}{N}$$

- Perform eigen decomposition on the covariance matrix to obtain the eigenvalues of the covariance matrix
- Sort the obtained eigenvalues by  $\lambda_1 \geq \lambda_2 \geq \lambda_3 \geq 0$
- Evaluate the linearity, planarity and sphericity of the distribution of points

Linearity	$(\lambda_1 - \lambda_2)/\lambda_1$
Planarity	$(\lambda_2 - \lambda_3)/\lambda_1$
Sphericity	$\lambda_3/\lambda_1$

# Cluster Width Overview

- 1D projection of voxels to the eigenvector with the second largest eigenvalue (from PCA calculation)

$$d_i = \mathbf{v}_2 \cdot (\mathbf{r}_i - \bar{\mathbf{r}})$$

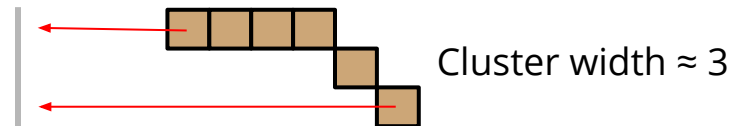
$\mathbf{v}_2$  : Eigenvector with the second largest eigenvalue

$\mathbf{r}_i$  : 3D coordinate of voxel  $i$

$\bar{\mathbf{r}}$  : Mean 3D coordinate of voxels in the same cluster

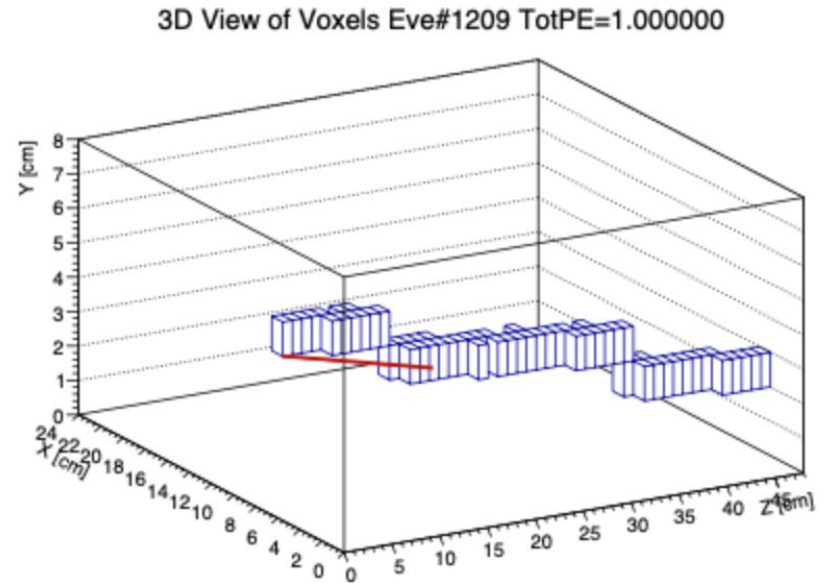
- Calculate the distance between the 2 voxels furthest away from each other in this eigenbasis (cluster width)

$$d = d_{max} - d_{min} \quad (\text{Cluster width})$$



# Max-Vox-Line Overview

- Calculate the eigenvectors for a cluster of voxels using PCA
- Shift the origin of the eigenvectors from the centroid of the cluster to the vertex of the cluster
- Obtain the main eigenvector which is the eigenvector with the largest eigenvalue (red line in the figure)
- Compute the maximum distance between the voxels and the main eigenvector (max-vox-line)



# Single Track Events Selection (MC)

- Definitions:

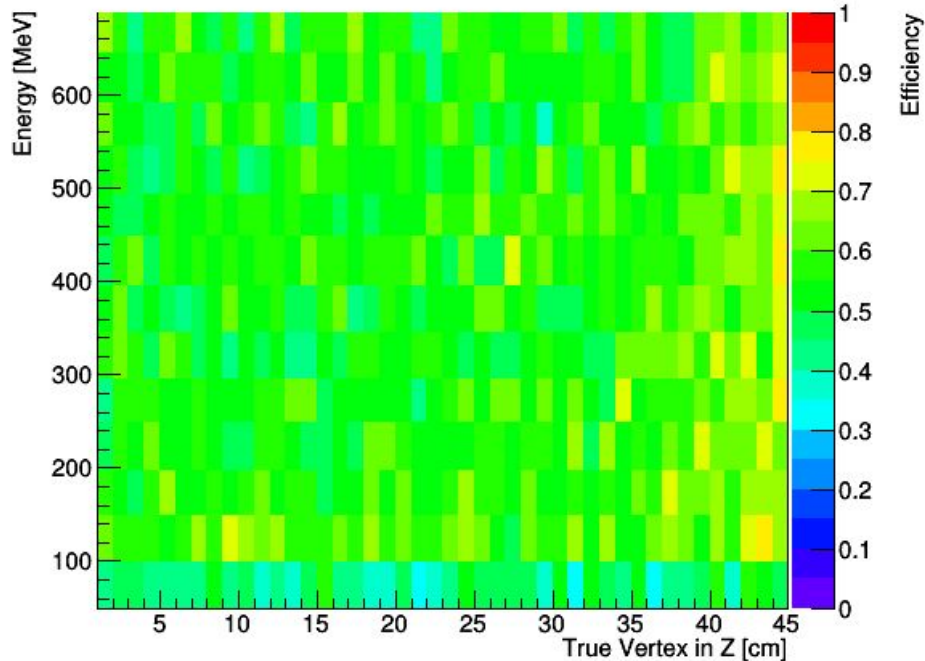
$$\text{Efficiency} = \frac{\text{Number of true single track events that got reconstructed as single track events}}{\text{Number of true single track events}}$$

$$\text{Purity} = \frac{\text{Number of reconstructed single track events that are true single track events}}{\text{Number of reconstructed single track events}}$$

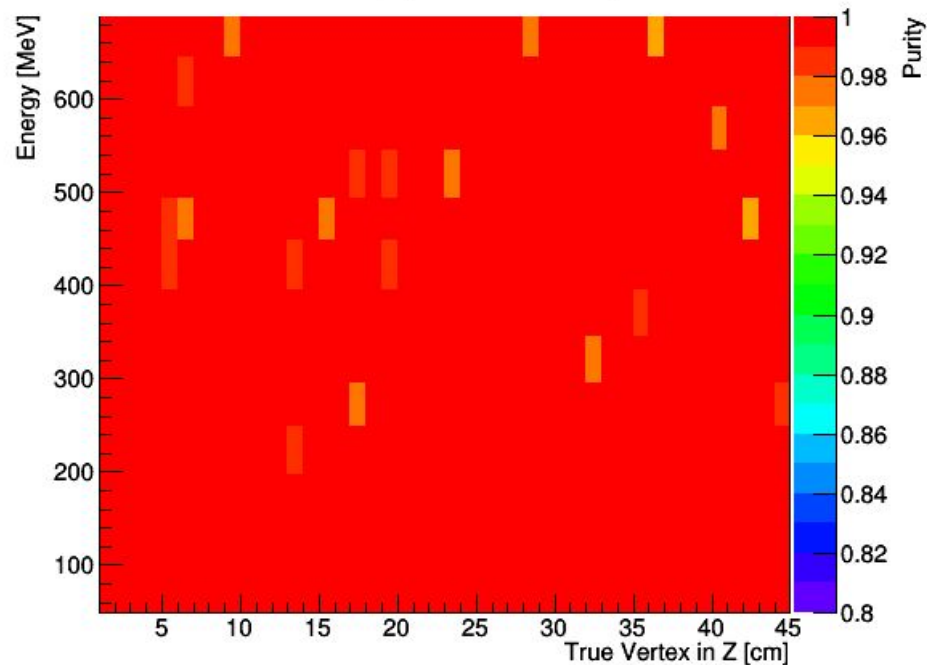
- Neutron MC sample
  - Selected a sample of events that got tagged as either true single track events or true multi track events in the MC file
  - True single track events: Only 1 proton as the final state particle ( $\geq 2$  MeV energy deposit)
  - True multi track events: At least 3 final state particles (each one  $\geq 2$  MeV energy deposit)

# Single Track Events Selection (MC)

Efficiency (Neutron MC)

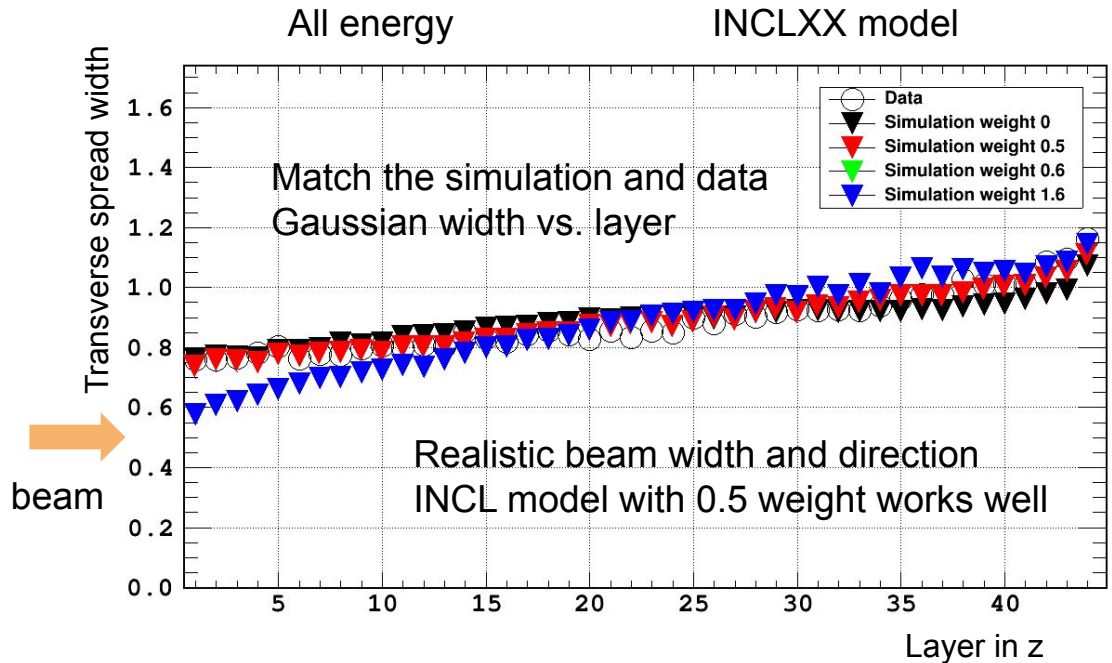


Purity (Neutron MC)



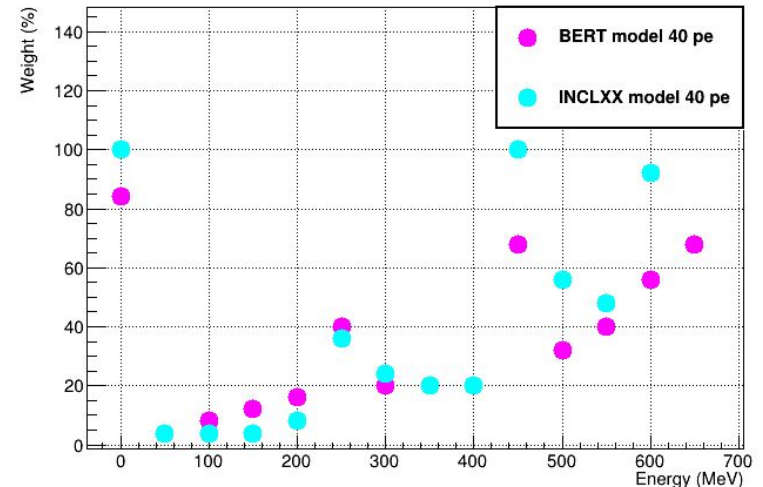
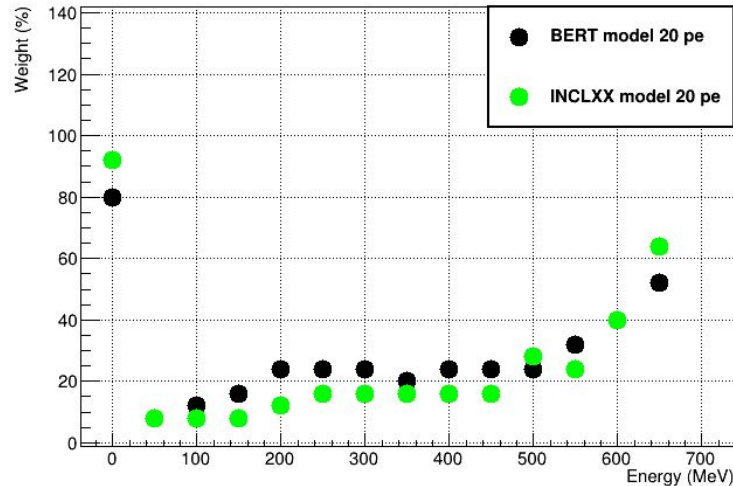
# Looking at the transverse spread => our “control sample”

- Along beamline, the interaction spread increases => assume this merely comes from the invisible scattering.
- Invisible scattering in simulation can be weighted to match data.
- Best matched weight can be found and this weight can be applied to simulation => simulation will be tuned.
- It is implemented for each energy range with multiple G4 models.
- Uncertainty comes from the correction discrepancy between two models.

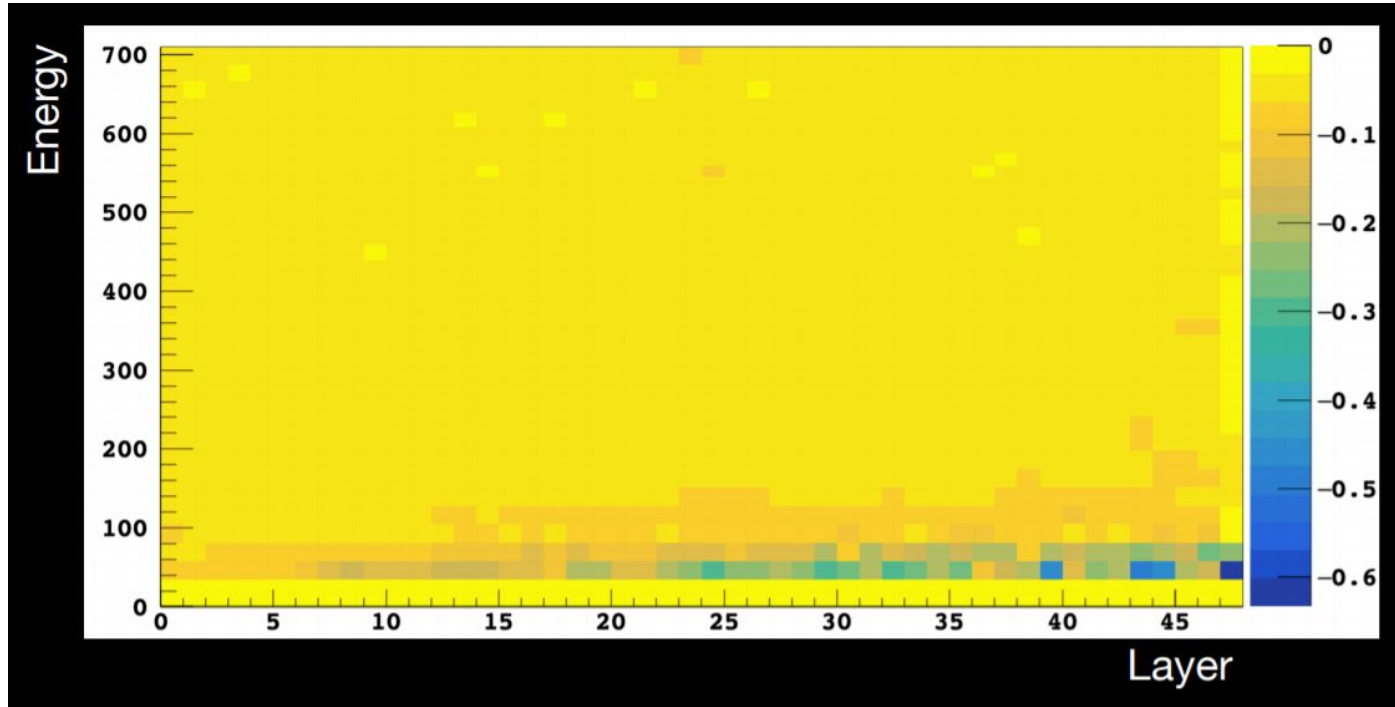


# Invisible scattering reweight in various cases

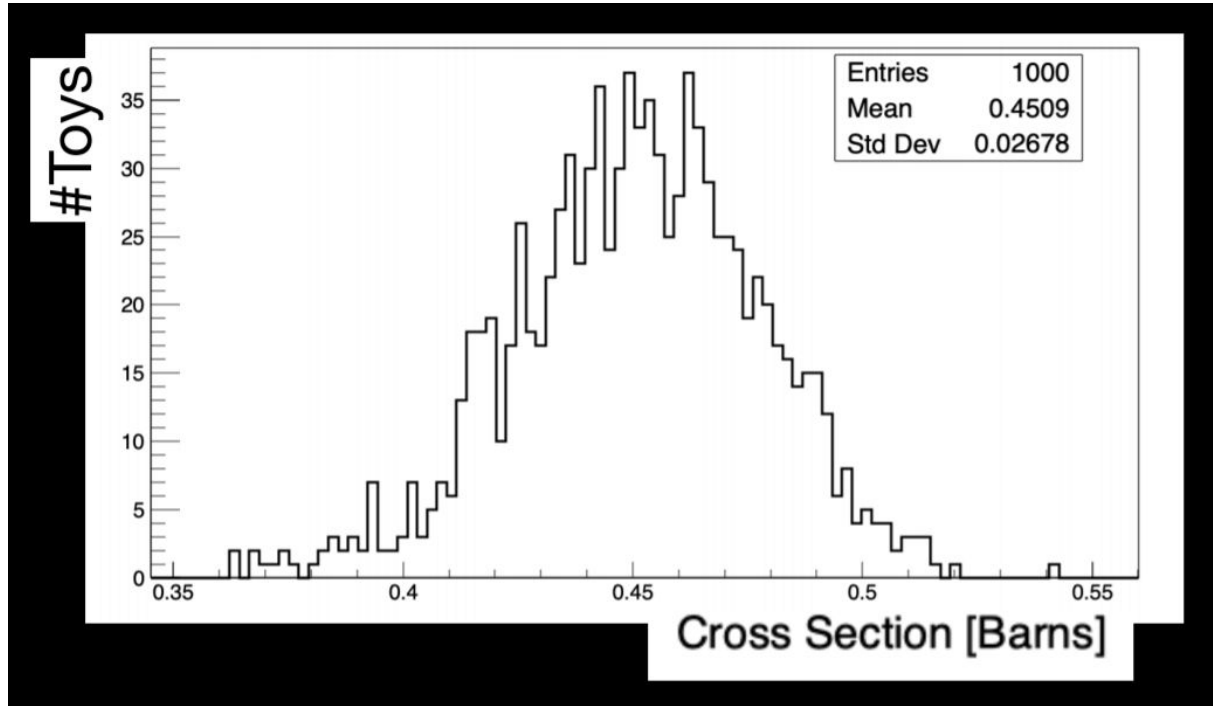
- Two models (BERT and INCLXX) were used to check the correction -> difference can be a systematic uncertainty source.
- Two pe cuts were used to check the correction stability -> systematics from pe cut will be included in the detection+reconstruction uncertainty



# Invisible Scattering Correction

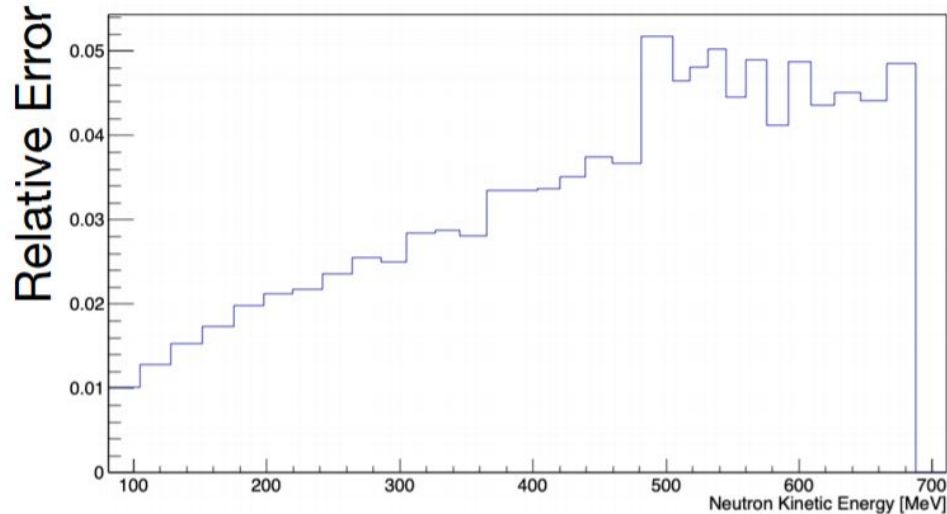


# Details about Uncertainty Propagation

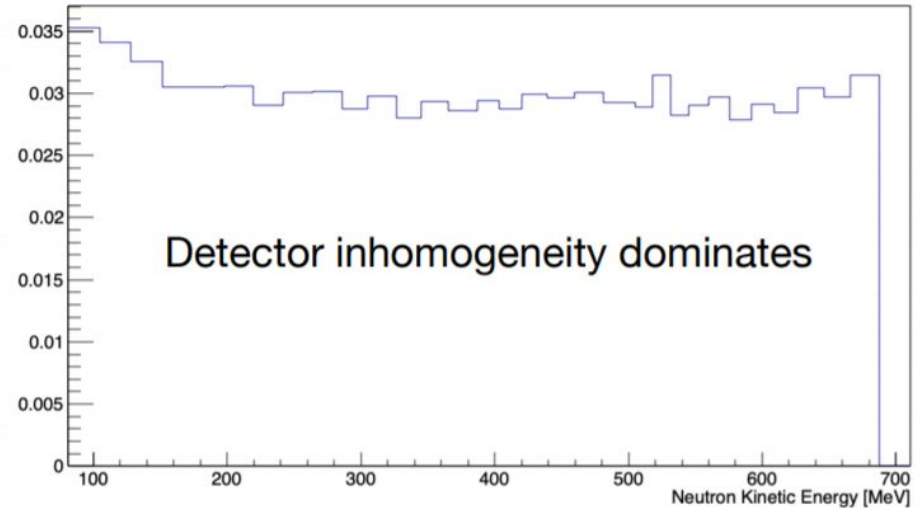


# Cross Section Fitter: Initial Result

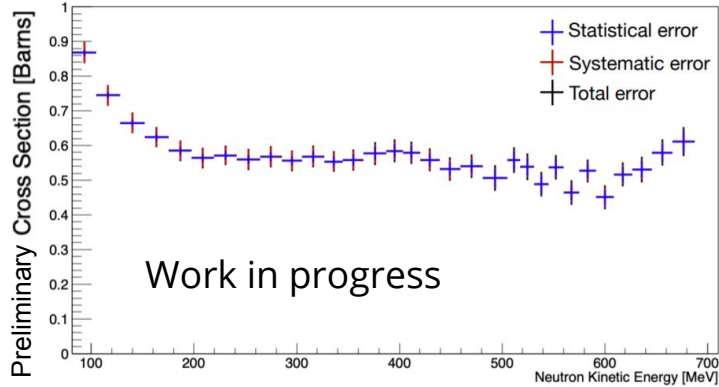
Statistical uncertainty



Systematic uncertainty



Caveat: Only 16 minutes of data used



DOI: 10.1103/PhysRevC.47.237

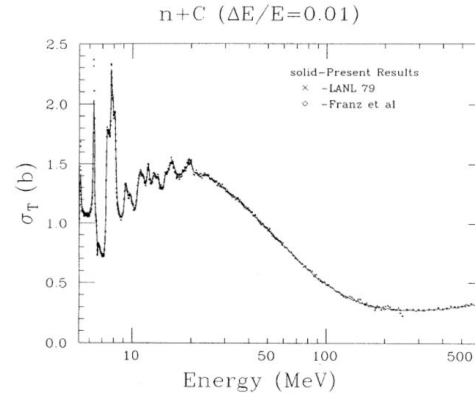


FIG. 8. Present results for carbon (solid line) compared with earlier work from Refs. [3] and [23]. The seemingly better resolution of the 1979 data is a consequence of the choice of binning function for the present work. The resolution of the unbinned raw data is about a factor of 2 better than the 1979 work.

DOI: 10.1103/PhysRevC.63.044608

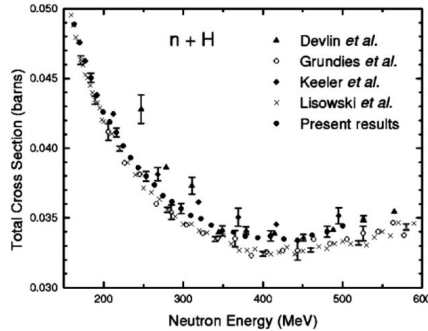


FIG. 8. Results are shown for the H cross section extracted from the two  $CH_2$  samples combined with those from the  $C_8H_{18}$  sample.

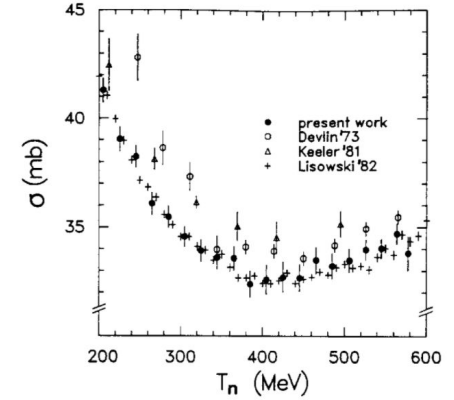


Fig. 2. Total neutron–proton cross sections between 200 and 600 MeV.

DOI:  
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0370-269  
3(85)907  
29-4