Origin of Spectral Hardening of Secondary Cosmic Ray Nuclei

(NK & Lee 2021, ApJ, 917, 61)



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XXVIII Cracow EPIPHANY Conference on Recent Advances in Astroparticle Physics 11/1/2022

Galactic Cosmic-rays (p, He, Li-Be-B, C,...)

(probably) produced via diffusive shock acceleration at SNRs

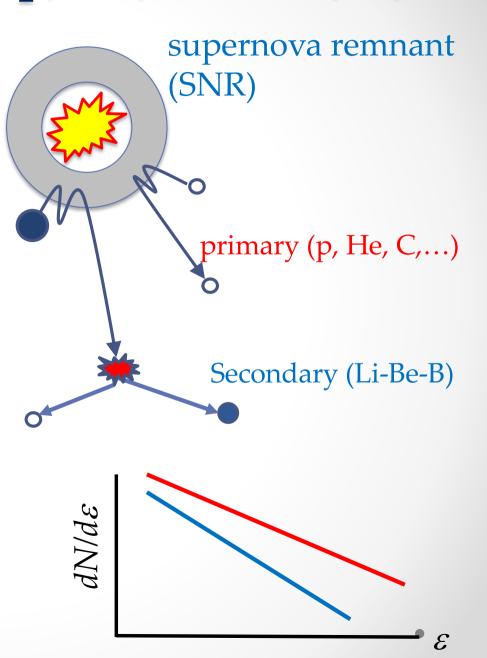
Conventional picture

proton, He, C, O, etc.:

primarily produced at

SNRs, power-law spectrum

Li-Be-B: secondarily produced via spallation of heavier nuclei during propagation, having steeper spectrum than primary CRs



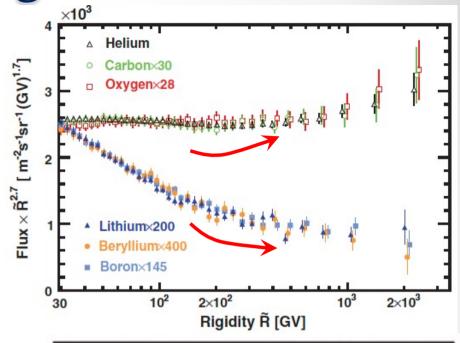
Spectral hardening of CR nuclei

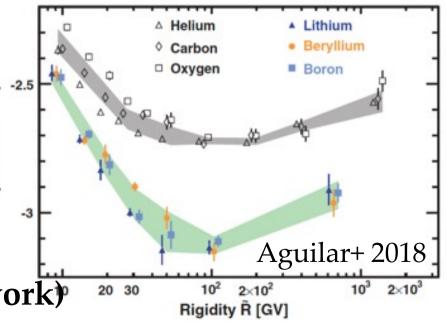
AMS-02:

- (1) The spectra of p, He, C, O, and Li-Be-B harden above ≥ 200GV
- (2) Li, Be, and B harden more than He, C and O
- propagation effect? (Thoudam & Horandel 14; Blasi+ 12 etc.)
- reacceleration? (Bresci+ 19 etc.)
- superposition of different kinds of sources? (Niu+ 20 etc.)
 Only Li hardens? primary Li
- source?

(NK & Yanagita 2018; Boschini+ 2020)

Primary Li-Be-B source? (This work)





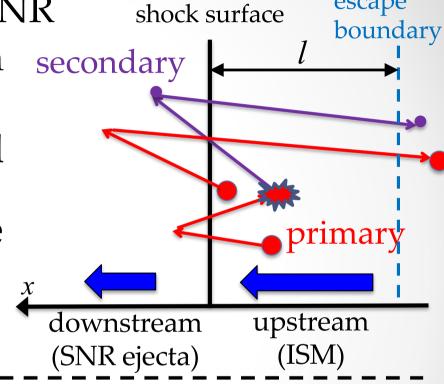
Production, Acceleration, and Escape of CRs

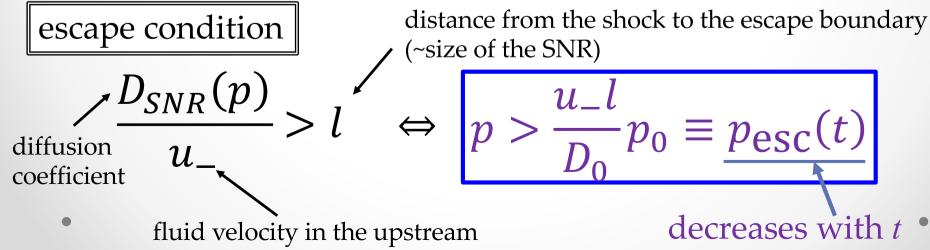
escape

Pri. CRs are accelerated at the SNR

- → interact w/ ambient medium secondary
- → sec. CRs production ***
- → sec. CRs are also accelerated (Mertsch & Sarkar 2009 etc.)
- → High energy CRs escape the SNR into the ISM earlier

(Gabici+ 2007, 2009; Ohira+ 2010)





CR distribution function f_i at the SNR

NK and Lee 2021

diffusion-convection equation

$$u(x)\frac{\partial f_i}{\partial x} = \frac{\partial}{\partial x} \left[D_i(p) \frac{\partial f_i}{\partial x} \right] + \frac{p}{3} \frac{du}{dx} \frac{\partial f_i}{\partial p} - \Gamma_i f_i + q_i$$

$$+ u_- Q_i \delta(x) \delta(p - p_0),$$

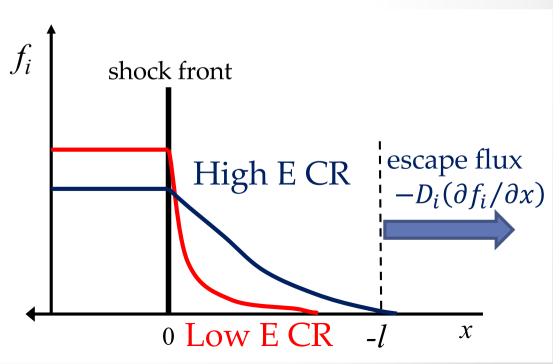
i = p, He, Li, Be, B, C, N, O

 $\Gamma_i = \sum_{i>j} \Gamma_{i\to j}$: spallation rate of nuclei *i*

 Q_i : injection rate of nuclei i $q_i \sim \sum_{i < j} \Gamma_{i \to i} f_i$: injection rate due to the spallation

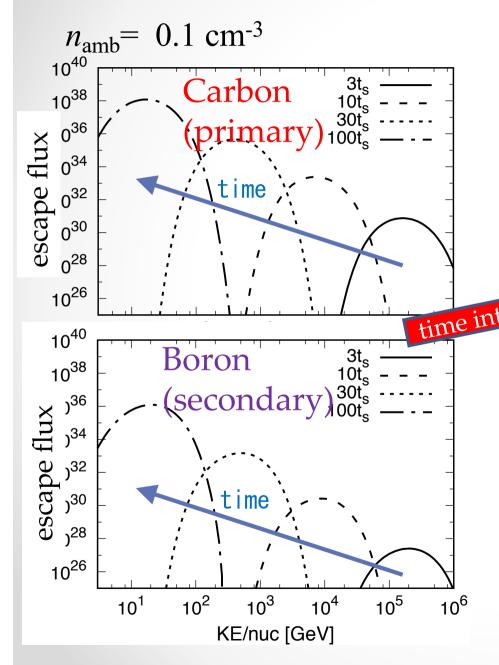
Boundary conditions

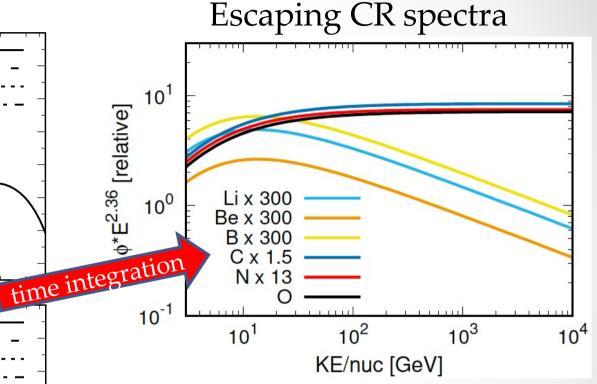
- $\lim_{x \to -0} f_i = \lim_{x \to +0} f_i,$
- $\lim_{x \to -1} f_i = 0,$ (ii)
- $\left|\lim_{x\to+\infty}f_i\right|<\infty,$ (iii)
- (iv) $\left[D_i(p)\frac{\partial f_i}{\partial x}\right]_{r=+0}^{x=-0}$ $= \frac{1}{3}(u_+ - u_-)p\frac{\partial f_{i,0}}{\partial p} + u_-Q_i\delta(p - p_0),$



Case 1: uniform ISM (n_{amb} = const.)

NK & Lee 2021



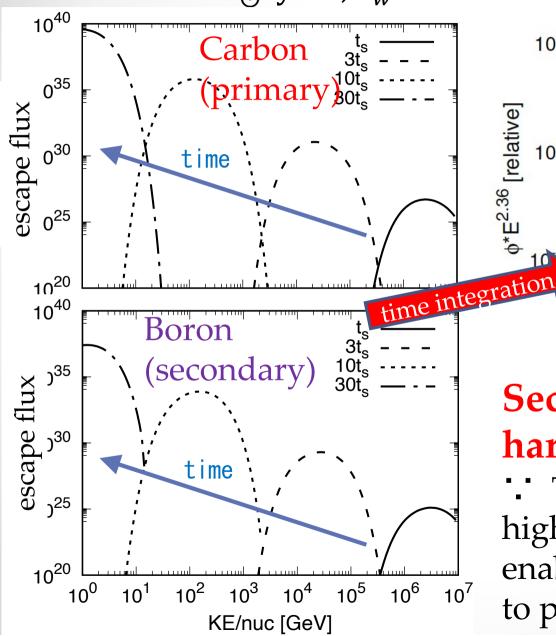


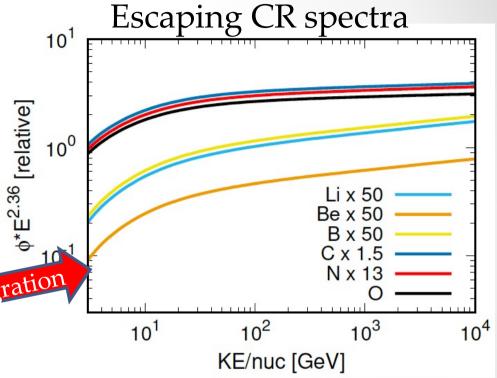
Secondaries are always softer than primaries ☺

High energy primaries escape the SNR earlier and produce less secondaries.

Case 2: SN with dense CSM $(n_{\rm amb} \propto r^{-2})$

 $\dot{M} = 1 \times 10^{-3} M_{\odot} \text{ yr}^{-1}$, $v_w = 100 \text{ km s}^{-1}$ NK and Lee 2021





Secondaries are always harder than primaries ©

The ambient density is higher in the early phase, which enables high energy primaries to produce more secondaries.

Fitting to the AMS-02 data

NK & Lee 2021

a local SNR with CSM:

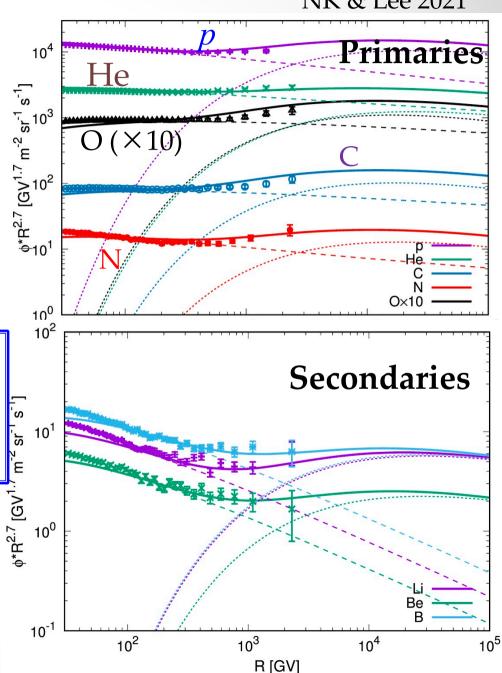
(dotted lines)

$$\dot{M} = 2.5 \times 10^{-3} M_{\odot} \text{ yr}^{-1},$$
 $v_w = 100 \text{ km s}^{-1},$
 $E_{\text{SN}} = 10^{51} \text{erg}, \eta_{\text{CR}} = 0.1$
age: $1.6 \times 10^5 \text{ yr}$
distance: 1.6 kpc

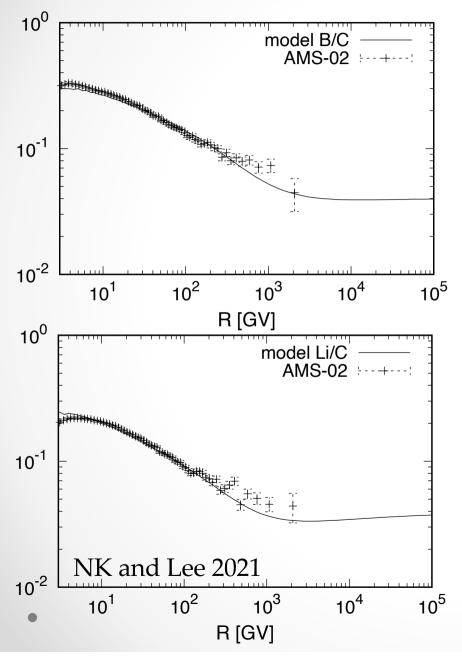
Hardening of both pri. and sec. CRs are reproduced simultaneously!

... if abundances of CNO in the CSM enhanced by a factor ~3 (metal-rich CSM)

....stripped-envelope SN (Ib/c)?



Prediction



<u>background</u> (≤ 200 GV): Secondaries are <u>softer</u> than primaries.

a local SNR contribution(≥ 200 GeV):Secondaries are <u>harder</u> than primaries



Energy dependence of secondary-to-primary ratios would flatten at higher energies

¹⁰⁵ # It may rise with energy!

Summary

- We propose a local supernova with dense circumstellar medium as the birth place of the hard CR Li-Be-B component appearing ≥ 200 GV.
- We calculate the production and acceleration of secondary CR nuclei in the SNR, as well as their escape into the ISM in a consistent way.
- The energy spectra of p, He, Li, Be, B, C, N, and O predicted from our model are consistent with the observations of AMS-02.
- Our scenario may be tested by secondary-toprimary ratios (e.g., B/C, Li/C, etc.) in ≥ TeV range.
- AMS-02, CALET, DAMPE, AMS-100 etc.