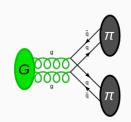
# The Scalar Glueball in Lattice QCD

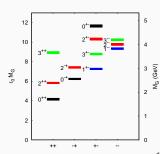
Mixing with Quarks and Mesons

Ruairí Brett (GWU) 2nd May 2021 Mini-workshop on glueball hunting Hadron Spectroscopy and Glueballs

#### HADRON SPECTROSCOPY

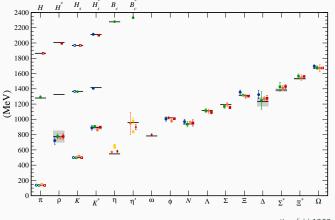
- many puzzles remain in hadron spectroscopy
  - e.g. tetra/pentaquarks, glueballs, hybrids, ...
- theoretical descriptions from QCD even more elusive
- glueballs on the lattice: early successes in pure Yang-Mills
  - next step to full QCD "challenging"
- can Lattice QCD complement experiment?
  - eg.  $\pi\pi$  hard to (directly) probe experimentally





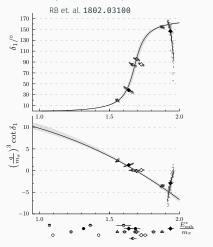
# **Historically:** (QCD) stationary states in finite-volume

- stable hadron masses  $\rightarrow$  precision: isospin breaking relevant
- no decays  $\rightarrow$  no resonance information



Historically: (QCD) stationary states in finite-volume

**Last 10 years/today:**  $2 \rightarrow 2$  scattering from finite-vol. energies



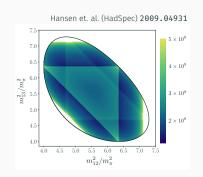
- meson-meson:  $\pi\pi$ ,  $K\pi$ ,  $\pi\eta$ ,  $D\pi$ , . . .
- meson-baryon accessible: review: Bulava 1909.13097
- on the horizon: baryon-baryon at light  $m_\pi$  e.g. Green et. al. 2103.01054

Historically: (QCD) stationary states in finite-volume

**Last 10 years/today:**  $2 \rightarrow 2$  scattering from finite-vol. energies

**Today:**  $3 \rightarrow 3$  the frontier

- formalisms for 3-body scattering rapidly developing
- first applications so far to  $3\pi^+$ ,  $3K^-$  e.g. from GWU: Culver et. al. 1911.09047 Alexandru et. al. 2009.12358 RB et al. 2101.06144



Historically: (QCD) stationary states in finite-volume

**Last 10 years/today:**  $2 \rightarrow 2$  scattering from finite-vol. energies

**Today:**  $3 \rightarrow 3$  the frontier

**Near future:** what about scattering producing glueballs?

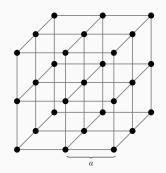
# \_\_\_\_

Extracting the finite-volume

spectrum

# LATTICE QCD

- Euclidean space-time
- finite volume L
- lattice spacing a
- (often) unphysical light quark masses
- → amenable to numerical calculations:

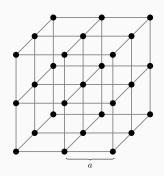


$$\langle \pmb{A} \rangle = \frac{1}{Z} \int \mathcal{D}[U, \overline{\psi}, \psi] \, \pmb{A}[U, \overline{\psi}, \psi] \, e^{-\overline{\psi} M \psi - S_G[U]}$$
 - Monte Carlo for gluon integration 
$$= \frac{1}{Z} \int \mathcal{D}[U] \, \pmb{A}[U, \pmb{M}^{-1}] \, \det M[U] \, e^{-S_G[U]} \quad \text{- Dirac inverse expensive:}$$
 
$$\dim M \sim 10^8 - 10^9$$

# LATTICE QCD

- Euclidean space-time
- finite volume L
- lattice spacing a
- (often) unphysical light quark masses
- → correlation functions:

$$C_{ij}(t) = \langle \mathcal{O}_i(t)\bar{\mathcal{O}}_j(0)\rangle$$
$$= \sum_n \langle 0|\mathcal{O}_i|n\rangle\langle n|\bar{\mathcal{O}}_j|0\rangle e^{-E_n t}$$



- (n-)hadron interpolation operators
- discrete energies from2-pt. functions

## EXTRACTING THE FINITE-VOL. SPECTRUM

- temporal correlation matrix:

$$C_{ij}(t) \equiv \langle 0|\mathcal{O}_i(t+t_0)\overline{\mathcal{O}}_j(t_0)|0\rangle$$
$$= \sum_n \langle 0|\mathcal{O}_i|n\rangle\langle n|\overline{\mathcal{O}}_j|0\rangle e^{-E_n t}$$

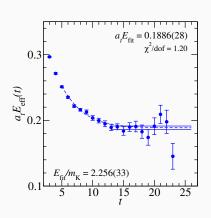
- solve generalized eigenvalue problem

$$C(\tau_0)^{-1/2} C(t) C(\tau_0)^{1/2} v_n(t, \tau_0)$$
  
=  $\lambda_n(t, \tau_0) v_n(t, \tau_0)$ 

- eigenvalues tend to lowest N energies

$$\lim_{t \to \infty} \lambda_n(t) = b_n e^{-E_n t} [1 + \mathcal{O}(e^{\Delta_n t})]$$

$$\Rightarrow E_{\text{eff}}^{n}(t) = \frac{1}{\Delta t} \ln \left( \frac{\lambda_{n}(t)}{\lambda_{n}(t + \Delta t)} \right)$$



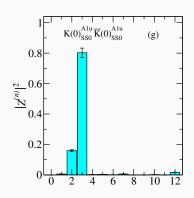
## EXTRACTING THE FINITE-VOLUME SPECTRUM

- correlation matrix:

$$C_{ij}(t) = \sum_{n} \langle 0|\mathcal{O}_{i}|n\rangle\langle n|\overline{\mathcal{O}}_{j}|0\rangle e^{-E_{n}t}$$
$$= \sum_{n} Z_{i}^{(n)} Z_{j}^{(n)*} e^{-E_{n}t}$$

- level ID inferred from *Z* overlaps with *probe* operators:

$$|\Phi_j\rangle \equiv \mathcal{O}_j|0\rangle \quad \Rightarrow \quad Z_j^{(n)} = \langle \Phi_j|n\rangle$$



ightarrow overlaps give **qualitative** measure of mixing between states

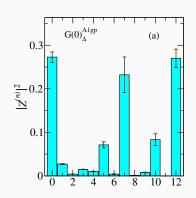
## EXTRACTING THE FINITE-VOLUME SPECTRUM

- correlation matrix:

$$C_{ij}(t) = \sum_{n} \langle 0|\mathcal{O}_{i}|n\rangle\langle n|\overline{\mathcal{O}}_{j}|0\rangle e^{-E_{n}t}$$
$$= \sum_{n} Z_{i}^{(n)} Z_{j}^{(n)*} e^{-E_{n}t}$$

 level ID inferred from Z overlaps with probe operators:

$$|\Phi_j\rangle \equiv \mathcal{O}_j|0\rangle \quad \Rightarrow \quad Z_j^{(n)} = \langle \Phi_j|n\rangle$$

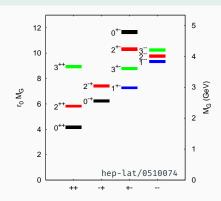


- ightarrow overlaps give **qualitative** measure of mixing between states
- ightarrow large operator bases crucial for reliably probing the spectrum

# From Pure Glue to QCD

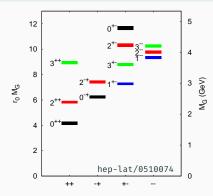
#### THE SCALAR GLUEBALL

- Glueball: bound state of gluons
- experimental evidence elusive, light scalar candidates:
  - $f_0(1370)$ ,  $f_0(1500)$ ,  $f_0(1710)$
- lattice studies to date:
  - only glueball operators
  - light scalar  $\approx 1.7~\text{GeV}$
  - most in pure  $SU_c(3)$ /quenched QCD



#### THE SCALAR GLUEBALL

- Glueball: bound state of gluons
- experimental evidence elusive, light scalar candidates:
  - $f_0(1370)$ ,  $f_0(1500)$ ,  $f_0(1710)$
- lattice studies to date:
  - only glueball operators
  - light scalar  $\approx 1.7~\text{GeV}$
  - most in pure  $SU_c(3)$ /quenched QCD



here: low-lying  $A_{1q}^+$  ( $J^{PG}=0^{++}$ ) spectrum up to  $\sim 2$  GeV

- **first study** including glueball,  $\overline{q}q$ , and two-hadron operators
- stepping stone to scattering studies

# (MULTI) MESON OPERATORS

0

single-site



singly-displaced



doubly-displaced-L



triply-displaced-U



triply-displaced-O

q's = smeared, displaced quark fields

$$\overline{\Phi}^{AB}_{\alpha\beta;ij}(\boldsymbol{p},t) = \sum_{\boldsymbol{x}} e^{i\boldsymbol{p}\cdot(\boldsymbol{x}+\frac{1}{2}(d_{\alpha}+d_{\beta}))} \delta_{ab} \, \overline{q}^{B}_{b\beta j}(\boldsymbol{x},t) \, q^{A}_{a\alpha i}(\boldsymbol{x},t)$$

group-theoretic projections for definite "spin", parity, etc.

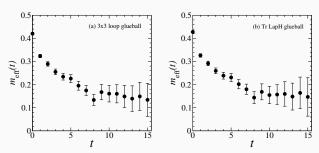
$$\overline{M}_{l}(\boldsymbol{p},t) = c_{\alpha\beta}^{(l)*} \overline{\Phi}_{\alpha\beta}^{AB}(\boldsymbol{p},t)$$

multi-hadron operators from single-hadron building blocks i.e.

$${\mathcal O}_{\pi(p_1)\pi(p_2)} \sim \sum_{i,j} c_{ij} \pi(p_i) \pi(p_j)$$

Morningstar et. al. 1303.6816

# A DIFFERENT GLUEBALL OPERATOR



any gauge invariant, purely gluonic object with the correct transformation properties is valid as a glueball operator<sup>1</sup>

scalar  ${\rm Tr}\,{\rm LapH}$  operator:

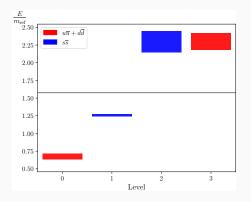
$$G_\Delta=-\operatorname{Tr}[\Theta(\sigma_s^2+ ilde\Delta) ilde\Delta],$$
  $ilde\Delta\equiv$  covariant laplacian (function of gauge links/gluons)

Morningstar et. al. 1303.6816

# $A_{1g}^+$ spectrum - $\overline{q}\,q$ only

 $N_f=2+1$  Wilson-clover, anisotropic lattice,  $m_\pi pprox 390$  MeV,  $m_K pprox 550$  MeV:

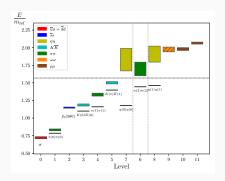
- first:  $\overline{q}q$  operators only
- only 2 states below  $\sim 2~\text{GeV}$
- $\sigma$  and  $f_0(980)$
- no sign of  $\overline{q}q$ -dominated states  $\sim 1.5-1.8~{\rm GeV^*}$



$$m_{\rm ref} = 1.82 m_K \approx 1 \text{ GeV}$$

# $A_{1g}^{+}$ SPECTRUM

 $N_f=2+1$  Wilson-clover, anisotropic lattice,  $m_\pi pprox 390$  MeV,  $m_K pprox 550$  MeV:

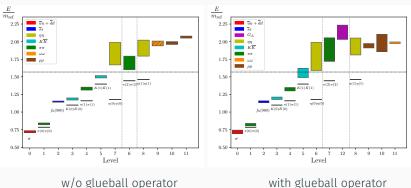


w/o glueball operator

Hatched boxes: significant overlap  $(|\langle n|\overline{\mathcal{O}}_{\alpha}|0\rangle|^2)$  with multiple operators  $m_{ref}=1.82m_K\approx 1~{\rm GeV}$ 

# $A_{1g}^+$ SPECTRUM

 $N_f=2+1$  Wilson-clover, anisotropic lattice,  $m_\pi\approx 390$  MeV,  $m_K\approx 550$  MeV:

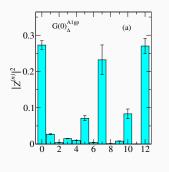


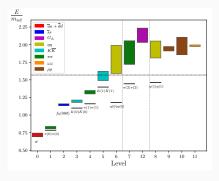
with glueball operator

Hatched boxes: significant overlap  $(|\langle n|\overline{\mathcal{O}}_{\alpha}|0\rangle|^2)$  with multiple operators  $m_{\rm ref} = 1.82 m_K \approx 1 \text{ GeV}$ 

# $A_{1g}^{+}\; {\rm SPECTRUM}$

 $N_f=2+1$  Wilson-clover, anisotropic lattice,  $m_\pi pprox 390$  MeV,  $m_K pprox 550$  MeV:





with glueball operator

Hatched boxes: significant overlap  $\left(|\langle n|\overline{\mathcal{O}}_{\alpha}|0\rangle|^2\right)$  with multiple operators  $m_{\mathrm{ref}}=1.82m_K\approx 1~\mathrm{GeV}$ 

**Looking Forward** 

#### INTERPRETING THESE RESULTS

# we find only two " $\bar{q}q$ dominated", and no new "glueball dominated" states below 2 GeV

- does this tell us there isn't a "pure glueball" state below 2 GeV? (at  $m_\pi \approx 390$  MeV, etc.)
- suggestion that the  $f_0$ 's below 2 GeV might not be  $\bar{q}q$  either?
  - → molecular states?
- more importantly, "is it even in principle possible to know?"
- either way, getting a handle on the glueball-meson mixing will be crucial
- what matrix elements will be most insightful (and calculable)?
- $4\pi$  looms overhead, but experience with  $3\pi$  makes it less intimidating

# **CONCLUSIONS**

- first study in full QCD to include the mixing between glueball,  $\bar{q}q$ , and meson-meson operators
  - relative dearth of new lattice studies (compared to other hadron scattering) indicates large hurdles that glueballs present
- including/excluding a scalar glueball operator introduces no new states
  - $\,\,
    ightarrow\,$  but plenty of noise!
- conservatively: a precursor to full blown scattering studies
  - ightarrow not so far away
- more speculative: no scalar "pure glueball" state below 2 GeV when quark dynamics fully factored in

# ANGULAR MOMENTUM ON THE LATTICE

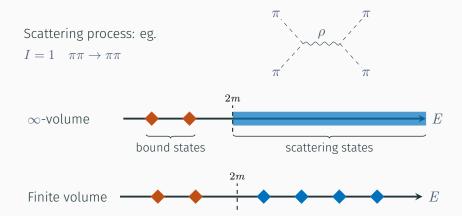
- periodic B.C. in cubic box:
  - $\Rightarrow J$  no longer a good quantum number
- label states with irreps of cubic group  ${\cal O}_{\hbar}^{D}$

- continuum spin ID:

J	$A_1$	$A_2$	E	$T_1$	$T_2$	J	$G_1$	$G_2$	Н
0	1	0	0	0	0	$\frac{1}{2}$	1	0	0
1	0	0	0	1	0	$\frac{3}{2}$	0	0	1
2	0	0	1	0	1	$\frac{5}{2}$	0	1	1
3	0	1	0	1	1	$\frac{7}{2}$	1	1	1
4	1	0	1	1	1	$\frac{9}{2}$	1	0	2

boosted frames introduce further mixing (but allow more sampling of scattering region)

# FINITE VOLUME SPECTRA

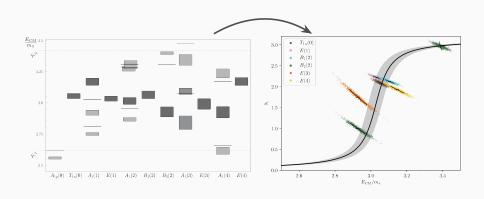


 $\rightarrow$  how to access  $\infty\text{-vol.}$  physics?

# TWO-BODY QUANTIZATION CONDITION

- quantisation condition (Lüscher formula):

$$\det[\widetilde{K}^{-1}(\underline{E}_{cm}) - B(E_{cm}, L)] = 0$$



# TWO-BODY QUANTIZATION CONDITION

- quantisation condition (Lüscher formula):

$$\widetilde{K}_{\ell}^{-1} \sim \cot \delta_{l} \xrightarrow{\det \left[\widetilde{K}^{-1}(E_{\rm cm}) - B(E_{\rm cm}, L)\right] = 0} \uparrow$$
known, mixes partial waves

- determinant over decay channel, partial waves (truncation!)

implementation with group theory for cubic box (& boosted frames):
 github.com/cjmorningstar10/TwoHadronsInBox

Morningstar et. al. 1707.05817