New Physics from Oscillations at DUNE Near Detector

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Minimal model: Seesaw Model

• Simplest extension of SM able to account for neutrino masses. Consists in the addition of heavy fermion singlets (N_R) to the SM field content:

Minkowski 77; Gell-Mann, Ramond, Slansky 79 Yanagida 79; Mohapatra, Senjanovic 80

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The New Physics Scale



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The New Physics Scale



 $0\nu\beta\beta$ decay, CLFV, Colliders, direct searches...

Are Long Baseline Neutrino Oscillation experiments sensitive to New Physics beyond 3ν framework





Neutrino Oscillations vs NP scale



Both limits can be studied in a unified & model independent way

Blennow, Coloma, Fernandez-Martinez, Hernandez-Garcia, JLP 1609.08637 Coloma, JLP, Rosauro-Alcaraz, **Urrea** 2105.11466.

Model Independent Approach

$U = \left(\begin{array}{cc} N & \Theta \\ R & S \end{array}\right)$

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$$V_i - \nu_{\alpha} \text{ mixing}$$
$$U = \left(\begin{array}{cc} N & \Theta \\ R & S \end{array}\right)$$

Model Independent Approach $N_i - \nu_{\alpha}$ mixing $U = \left(\begin{array}{c} N & \Theta \\ R & S \end{array} \right)$

Deviation from unitarity of the PMNS matrix

Langacker, London 1988 Antusch, Biggio, Fernandez-Martinez, Gavela, JLP 2006

General Parameterizations

Triangular parameterization

$$N = (I + T)U$$

Deviation from unitarity

$$T = \begin{pmatrix} \alpha_{ee} & 0 & 0\\ \alpha_{\mu e} & \alpha_{\mu \mu} & 0\\ \alpha_{\tau e} & \alpha_{\tau \mu} & \alpha_{\tau \tau} \end{pmatrix}$$

Unitary matrix (standard unitary PMNS matrix up to small corrections)

Z.-z. Xing 2008, 2012 Escrihuela, Forero, Miranda, Tortola 2015

Far Detector vs Near detector



$$N_{\nu_{\alpha} \to \nu_{\beta}} \sim \frac{\Phi_{\alpha}(E)}{L^2} P_{\alpha\beta}(L/E) \sigma_{\beta}(E) \epsilon_{\beta}(E)$$

- Sources of systematics
- Cross sections
- Neutrino flux

- Near detector measurements reduce far detector systematic uncertainties
- New Physics at near detector (strongly affected by systematic uncertainties)

Far Detector

Far Detector

• What is measured in neutrino oscillation experiments





(1) Non-Unitary Mixing

• What is measured in neutrino oscillation experiments

$$\mathcal{P}_{\alpha\beta} = \frac{\left| (N \exp(-iHL)N^{\dagger})_{\beta\alpha} \right|^2}{\left[(NN^{\dagger})_{\alpha\alpha} \right]^2}.$$

• When $NN^{\dagger} = I \implies \mathcal{P}_{\alpha\beta} = P_{\alpha\beta}$ (SM limit recovered)



1. The light-heavy oscillations averaged out at the near detector. Identical to the heavy non-unitarity case

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2. The light-heavy oscillations have not yet developed at the near detector. No normalization factor

DUNE: $0.1 \,\mathrm{eV}^2 \lesssim \Delta m^2 \lesssim 1 \,\mathrm{eV}^2$

1. The light-heavy oscillations averaged out at the near detector. Identical to the heavy non-unitarity case

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The oscillation frequency dictated by the light-heavy frequency matches the near detector distance.
 Oscillations could be observed at the near detector

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> Low Scale Non-Unitarity

Present Bounds

	High-scale Non-Unitarity	$-\nu_i$
	$(m > \mathrm{EW})$	$Z \sim (N^{\dagger}N)_{ii}$
α_{ee}	$1.3 \cdot 10^{-3}$	$\begin{bmatrix} (1 & 1 &)ij \\ \nu_j \end{bmatrix}$
$lpha_{\mu\mu}$	$2.2\cdot 10^{-4}$	$\frac{l_{\alpha}}{l}$
$\alpha_{ au au}$	$2.8\cdot 10^{-3}$	$\sim N_{\alpha i}$
$ lpha_{\mu e} $	$6.8 \cdot 10^{-4} \ (2.4 \cdot 10^{-5})$	EW & CLFV ν_i
$ \alpha_{ au e} $	$2.7\cdot 10^{-3}$	data
$ lpha_{ au\mu} $	$1.2\cdot 10^{-3}$	

Fernandez-Martinez, Hernandez-Garcia, JLP 1605.08774 Blennow, Coloma, Fernandez-Martinez, Hernandez-Garcia, JLP 1609.08637

Present Bounds

	High-scale Non-Unitarity	Low-scale Non-Unitarity			
	$(m > \mathrm{EW})$	$\Delta m^2 \gtrsim 100 \ {\rm eV}^2 \Delta m^2 \sim 0.1 - 1 \ {\rm eV}^2$			
α_{ee}	$1.3 \cdot 10^{-3}$	$2.4 \cdot 10^{-2}$ bugey $1.0 \cdot 10^{-2}$ bugey			
$lpha_{\mu\mu}$	$2.2\cdot 10^{-4}$	$2.2 \cdot 10^{-2}$ SK $1.4 \cdot 10^{-2}$ MINOS			
$\alpha_{\tau\tau}$	$2.8\cdot 10^{-3}$	$1.0 \cdot 10^{-1}$ SK $1.0 \cdot 10^{-1}$ SK			
$ lpha_{\mu e} $	$6.8 \cdot 10^{-4} (2.4 \cdot 10^{-5})$	$2.5 \cdot 10^{-2}$ Nomad $1.7 \cdot 10^{-2}$			
$ \alpha_{ au e} $	$2.7\cdot 10^{-3}$	$6.9 \cdot 10^{-2}$ $4.5 \cdot 10^{-2}$			
$ lpha_{ au\mu} $	$1.2\cdot 10^{-3}$	$1.2 \cdot 10^{-2}$ NOMAD $5.3 \cdot 10^{-2}$			
Fernande 1605.087 Blennow, Hernande 1609.086	ez-Martinez, Hernandez-Garcia, JLP 74 Coloma, Fernandez-Martinez, ez-Garcia, JLP 637	$\alpha_{\alpha\beta} \le 2\sqrt{\alpha_{\alpha\alpha}\alpha_{\beta\beta}}$			

Present Bounds

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	$(m > \mathrm{EW})$	$\Delta m^2 \gtrsim 100 \ {\rm eV}^2 \Delta m^2 \sim 0.1 - 1 \ {\rm eV}^2$				
α_{ee}	$1.3\cdot 10^{-3}$	$2.4 \cdot 10^{-2}$ $1.0 \cdot 10^{-2}$				
$lpha_{\mu\mu}$	0.1-0.001%	$2.2 \cdot 10^{-2}$ 10-1% ²				
$\alpha_{ au au}$	eve	$1.0 \cdot 10^{-1}$ ·1				
$ lpha_{\mu e} $		$2.5 \cdot 10^{-2}$ 0^{-2}				
$ \alpha_{ au e} $	$2.7\cdot 10^{-3}$	$6.9 \cdot 10^{-2}$ $4.5 \cdot 10^{-2}$				
$ lpha_{ au\mu} $	$1.2\cdot 10^{-3}$	$1.2 \cdot 10^{-2}$ $5.3 \cdot 10^{-2}$				

Fernandez-Martinez, Hernandez-Garcia, JLP 1605.08774 Blennow, Coloma, Fernandez-Martinez, Hernandez-Garcia, JLP 1609.08637

See also Park, Ross-Lonergan 1508.05095 Ellis, Kelly, Weishi Li 2004.13719 Ellis, Kelly, Weishi Li 2008.01088







Blennow, Coloma, Fernandez-Martinez, Hernandez-Garcia, JLP 1609.08637. DUNE CDR configuration 1606.09550

Near Detector

Coloma, JLP, Rosauro-Alcaraz, Urrea 2105.11466.

See also Escrihuela, Forero, Miranda, Tortola, Valle arXiv:1503.08879 for other Near Detector configurations (without including tau detection).

High Scale Non-Unitarity

• What is measured in Near Detector

$$\mathcal{P}_{\alpha\beta} = \left| (NN^{\dagger})_{\beta\alpha} \right|^2 = |\alpha_{\alpha\beta}|^2 \qquad \begin{array}{c} \text{zero} \\ \text{distance} \\ \text{effect:} \end{array}$$

$$\mathcal{P}_{\alpha\alpha} = \left| (NN^{\dagger})_{\alpha\alpha} \right|^2 = 1 - 4 \,\alpha_{\alpha\alpha}$$

Sterile Neutrinos: 3+1

• What is measured in Near Detector

$$\mathcal{P}_{\alpha\beta} = 4|\boldsymbol{U}_{\alpha4}||\boldsymbol{U}_{\beta4}|\sin^2\frac{\Delta m_{41}^2 L}{4E}$$

$$\mathcal{P}_{\alpha\alpha} = 1 - 4|\boldsymbol{U}_{\alpha4}|^2 \sin^2 \frac{\Delta m_{41}^2 L}{4E}$$

Averaged-out regime

• What is measured in Near Detector $\Delta m^2_{41}\gtrsim 100\,{\rm eV^2}$

$$\langle \mathcal{P}_{\alpha\beta} \rangle = 2 |U_{\alpha4}| |U_{\beta4}|$$

zero distance effect:

$$\langle \mathcal{P}_{\alpha\alpha} \rangle = 1 - 2 |U_{\alpha4}|^2$$

Averaged-out regime

• What is measured in Near Detector $\Delta m^2_{41}\gtrsim 100\,{\rm eV^2}$

$$\langle \mathcal{P}_{\alpha\beta} \rangle = 2 |\alpha_{\alpha\beta}|^2$$

zero distance effect:

$$\langle \mathcal{P}_{\alpha\alpha} \rangle = 1 - 4 |\alpha_{\alpha\alpha}|^2$$

Low Scale Non-Unitarity

Role of shape uncertainty



- · Sensitivity driven by spectral information.
- Marginal impact of global normalization error.

ν_{τ} appearance channel

• Energy threshold of au production 3.2 GeV.



• $\mathcal{V}_{\mathcal{T}}$ detection: we follow de Gouvêa, Kelly, Stenico, Pasquini 1904.07265

Low Scale Non-Unitarity







3+1 Sterile Neutrinos: $P_{\mu\mu} + P_{\mu e} + P_{ee}$



Conclusions

 Near detectors can play a relevant role in testing the robustness of the 3-neutrino picture
 Low scale Non-Unitarity, sterile neutrino oscillations, NSI (Non-Unitarity results can be easily mapped to NSI framework, see 2105.11466)

• Keeping under control shape uncertainties is a key issue.

- Sensitivity beyond present bounds only when systematics below 5%
- Joint experimental and theoretical effort required to reduce systematics.
- Independent measurements of the cross sections would give very relevant information regarding the energy dependence. (see for instance $\nu STORM$ proposal)
- Neutrino tau detection opens a complementary New Physics channel.



ν_{τ} appearance channel

$\mathcal{V}_{\mathcal{T}}$ detection:

- Energy threshold of au production 3.2 GeV.
- Short lifetime of τ , indirect measurement via hadronic decays (~ 65% branching ratio).
- NC background. We have considered a sample in which 30% of the hadronic events are identified keeping 0.5% of NC background.

de Gouvêa, Kelly, Stenico, Pasquini 1904.07265

DUNE set up

Globes files

DUNE Collaboration, arXiv:2103.04797 [hep] 8 Mar 2021.

Flux configuration								
	Beam configuration	Power	E_p	PoT/yr	t_{ν} (yr)	$t_{\bar{\nu}} (\mathrm{yr})$	$M_{ m det}$	
	Nominal	$1.2 \ \mathrm{MW}$	$120 \mathrm{GeV}$	1.1×10^{21}	3.5	3.5	67.2 tons	
	High-Energy	$1.2 \ \mathrm{MW}$	$120~{\rm GeV}$	$1.1 imes 10^{21}$	3.5	—	67.2 tons	

Running mode	Sample	Contribution	Event rates $(\times 10^5)$	$E_{\rm obs}^{\rm max}$ (GeV)	
		Intrinsic cont.	20.18		
	ν_e -like	Flavor mis-ID	4.61	7.125	
		NC	6.77		
ν mode (nominal)	1:1	$\nu_{\mu}, \bar{\nu}_{\mu} \ CC \ (P_{\mu\mu} = 1) $ 2,235.72		7 195	
	$ u_{\mu}$ -nke	NC 17.35		1.120	
	ν_{τ} -like	$\nu_{\tau}, \bar{\nu}_{\tau} \ CC \ (P_{\mu\tau} = 1)$ 39.33		19	
		NC	3.23	10	
		Intrinsic cont.	11.18		
	$\bar{\nu}_e$ -like	Flavor mis-ID	1.07	7.125	
		NC	NC 3.89		
$\bar{\nu}$ mode (nominal)	$\bar{\nu}_{\mu}$ -like	$\nu_{\mu}, \bar{\nu}_{\mu} \text{ CC } (P_{\mu\mu} = 1)$ 1,013.42		7.195	
		NC	VC 9.76		
	$\bar{\nu}$. Blo	$ u_{ au}, \bar{\nu}_{ au} \ \mathrm{CC} \ (P_{\mu au} = 1) $	27.75	18	
	ν_{τ} -nke	NC	C 1.80		
	ν_e -like	Intrinsic cont.	38.10		
		Flavor mis-ID	12.98	18	
		\mathbf{NC}	30.51		
ν mode (HE)	1:1.0	$\nu_{\mu}, \bar{\nu}_{\mu} \text{ CC } (P_{\mu\mu} = 1)$	5,784.30	10	
	$ u_{\mu}$ -нке	NC 72.15		10	
	ν_{τ} -like	$\nu_{\tau}, \bar{\nu}_{\tau} \text{ CC } (P_{\mu\tau} = 1)$	259.67	19	
		NC 9.42		10	

Event cample	Contribution	Benchmark 1		Benchmark 2		Benchmark 3	
Event sample	Contribution	σ_{norm}	σ_{shape}	σ_{norm}	σ_{shape}	σ_{norm}	σ_{shape}
	Signal	5%	_	5%	_	5%	_
	Intrinsic cont.	10%	_	10%	2%	10%	5%
ν_e -fike	Flavor mis-ID	5%		5%	2%	5%	5%
	\mathbf{NC}	10%		10%	2%	10%	5%
	$\nu_{\mu}, \bar{\nu}_{\mu} \text{ CC (signal)}$	10%		10%	2%	10%	5%
ν_{μ} -fike	NC	10%		10%	2%	10%	5%
u liko	Signal	20%		20%		20%	
ν_{τ} -nke	NC	10%	_	10%	2%	10%	5%



• Mapping: $2|\alpha_{\beta\gamma}|^2 = |\epsilon^d_{\beta\gamma}|^2 + |\epsilon^s_{\beta\gamma}|^2 + 2|\epsilon^d_{\beta\gamma}||\epsilon^s_{\beta\gamma}|\cos(\Phi^s_{\beta\gamma} - \Phi^d_{\beta\gamma})$



Blennow, Coloma, Fernandez-Martinez, Hernandez-Garcia, JLP 1609.08637. DUNE CDR configuration 1606.09550

Far Detector



• Including ν_{τ} appearance channel does not change the picture.

de Gouvêa, Kelly, Steni Pasquini 1904.07265

3+1 Sterile Neutrinos: $P_{\mu\mu}$





3+1 Sterile Neutrinos: $P_{\mu e} + P_{ee}$



Systematics: Disappearance

$$N_{\nu_{\alpha} \to \nu_{\beta}} \sim \frac{\Phi_{\alpha}(E)}{L^2} P_{\alpha\beta}(L/E) \sigma_{\beta}(E) \epsilon_{\beta}(E)$$

• Using near detectors is a very effective way of reducing systematics in disappearance experiments (K2K, MINOS, reactors...).

$$\frac{N_{\nu_{\alpha}}^{\rm FD}}{N_{\nu_{\alpha}}^{\rm ND}} \sim \frac{L_{\rm ND}^2}{L_{\rm FD}^2} \frac{\Phi_{\alpha} \sigma_{\alpha} \epsilon_{\alpha}}{\Phi_{\alpha} \sigma_{\alpha} \epsilon_{\alpha}} P_{\alpha \alpha}$$

Systematics: Appearance (CP violation)

$$N_{\nu_{\alpha} \to \nu_{\beta}} \sim \frac{\Phi_{\alpha}(E)}{L^2} P_{\alpha\beta}(L/E) \sigma_{\beta}(E) \epsilon_{\beta}(E)$$

• For appearance experiments the situation is more complicated

$$\frac{N_{\nu_e}^{\rm FD}}{N_{\nu_{\mu}}^{\rm ND}} \sim \frac{L_{\rm ND}^2}{L_{\rm FD}^2} \frac{\sigma_e \epsilon_e}{\sigma_{\mu} \epsilon_{\mu}} P_{\mu e}$$

• CP violation requires comparison between neutrino and anti-netrino signals.

$$\frac{N_{\nu_e}^{\text{Far}}}{N_{\bar{\nu}_e}^{\text{Far}}} \sim \frac{N_{\nu_{\mu}}^{\text{ND}}}{N_{\bar{\nu}_{\mu}}^{\text{ND}}} \frac{\sigma_e \epsilon_e}{\sigma_{\mu} \epsilon_{\mu}} \frac{\sigma_{\bar{\mu}} \epsilon_{\bar{\mu}}}{\sigma_{\bar{e}} \epsilon_{\bar{e}}} \frac{P_{\mu e}}{P_{\bar{\mu}\bar{e}}}$$

Huber, Mezzetto, Schwetz, 0711.2950 Coloma, Huber, Kopp, Winter, 1209.5973

Nuclear Cross sections

 Neutrino-nucleus cross section missmodeling could lead to unacceptably large systematic uncertainties or biased measurements, even after the inclusion of a near detector.



$$\begin{split} \chi^2_{\min}(\{\Theta\}) &= \min_{\{\xi,\zeta\}} \left[\chi^2_{\text{stat}}(\{\Theta,\xi,\zeta\}) + \sum_s \left(\frac{\zeta_s}{\sigma_{\text{norm},s}}\right)^2 + \sum_b \left(\frac{\zeta_b}{\sigma_{\text{norm},b}}\right)^2 \\ &+ \sum_i \left(\frac{\xi_i^{\text{sig}}}{\sigma_{\text{shape,sig}}}\right)^2 + \sum_i \left(\frac{\xi_i^{\text{bg}}}{\sigma_{\text{shape,bg}}}\right)^2 \right] \,, \end{split}$$

$$\chi^2_{\text{stat}}(\{\Theta,\xi,\zeta\}) = \sum_i 2\left(N_i(\{\Theta,\xi,\zeta\}) - O_i + O_i \ln \frac{O_i}{N_i(\{\Theta,\xi,\zeta\})}\right)$$
$$N_i(\{\Theta,\xi,\zeta\}) = \sum_s (1+\xi_i^{\text{sig}}+\zeta_s) \, s_i(\{\Theta\}) + \sum_b (1+\xi_i^{\text{bg}}+\zeta_b) \, b_i(\{\Theta\})$$

General Parameterizations

Hermitian parameterization

$$N = (I - \eta)U'$$

Deviation from unitarity

$$\eta = \begin{pmatrix} \eta_{ee} & \eta_{e\mu} & \eta_{e\tau} \\ \eta_{e\mu}^* & \eta_{\mu\mu} & \eta_{\mu\tau} \\ \eta_{e\tau}^* & \eta_{\mu\tau}^* & \eta_{e\tau} \end{pmatrix} = \frac{\Theta\Theta^{\dagger}}{2}$$

Unitary matrix (standard unitary PMNS matrix up to small corrections)

Broncano, Gavela, Jenkins 2003 Fernandez-Martinez, Gavela, JLP, Yasuda 2007