

Lepton Flavor Universality

$$R_K = \frac{\Gamma(B \rightarrow K e^+ e^-)}{\Gamma(B \rightarrow K \mu^+ \mu^-)} \stackrel{SM}{=} 1 + \mathcal{O}(10^{-3})$$

[for K -inclusive rate]

- Why LFU holds within the SM?
- Why is it interesting to test LFU?
- The $b \rightarrow s \ell \ell$ system : R_K is not alone ...

I. Why LFU holds within the SM?

LFU is an approximate accidental symmetry of \mathcal{L}_{SM}

accidental symm. = symm not imposed, but emerging in the renormalizable (= low-energy) part of QFT
 (E.g. baryon number)

The diagram illustrates the energy scale. On the left, there is a vertical axis labeled "Energy" with an upward arrow. Below this axis, there is a bracket labeled "SM spectrum" with three horizontal tick marks. Above the axis, there is a red symbol resembling a Greek letter Λ with a minus sign, indicating a cutoff or scale. To the right of the axis, the SM Lagrangian is shown as $\mathcal{L}_{SM} = \mathcal{L}^{d=4} + \sum_i \frac{g_i^{[d]}}{\Lambda^{d-4}} \mathcal{O}_i^d(\gamma_{SM})$. Below the SM Lagrangian, there is a double-headed arrow labeled "d=4".

the only $d=4$ ops. we can write distinguishing
 e, μ, τ are the Yukawa interactions

$$\left\{ \begin{array}{l} g_e \approx 3 \times 10^{-6} \quad g_\mu \approx 3 \times 10^{-4} \quad g_\tau \approx 1 \times 10^{-2} \quad \ll g_{SM}^{gauge} \\ \text{Limit } g_{e,\mu,\tau} \rightarrow 0 \end{array} \right. \Rightarrow \text{LFU in } \mathcal{L}_{SM}$$

II. Why is it interesting to test LFU?

$$\mathcal{L}_{\text{SM-EFT}} = \mathcal{L}^{\text{deg}} + \sum_i \frac{c_i[\lambda]}{\lambda^{d-4}} \partial_i^d (\psi_{\text{SM}})$$

Testing accidental symm. is a very powerful probe of high-energy dynamics

LFU tested very precisely in many systems,
 $(\pi \rightarrow e\nu, K \rightarrow e\nu, \tau \rightarrow e\nu\bar{\nu}, Z \rightarrow ee)$
 but not yet in semileptonic b [3rd gen.] decays

A "semi-historical" example :

1974 : \mathcal{L}_{SM} (2 gen.)

- Renormalizable
- Excellent descr. of all known phenomena but CP in $K-\bar{K}$

CP is an
accidental symm.

[single observable!]

$$\hookrightarrow \mathcal{L}_{\text{EFT}} = \mathcal{L}_{\text{SM}} (2 \text{ gen.}) + \frac{e^{i\delta}}{\Lambda_{\text{SW}}^2} (\bar{s} \gamma^\mu d)^2$$

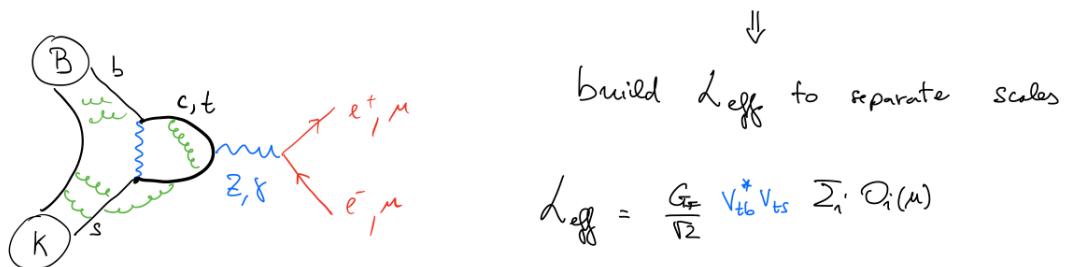
$$\Lambda_{\text{SW}} \approx 10^4 \text{ TeV} \Rightarrow \text{"Super-weak" interaction} \quad [\text{Wolfenstein}]$$

We now know that $\frac{1}{\Lambda_{\text{SW}}^2} = \frac{G_F^2 M_t^2}{4\pi^2} |V_{ts} V_{td}|^2$

- Lessons :
- Accidental symm. probe high-energy scales
 - Scales of higher-dim ops can be misleading
 - Light gen. easily insensitive to NP

III. The $b \rightarrow s \ell \ell$ system : R_K is not alone ...

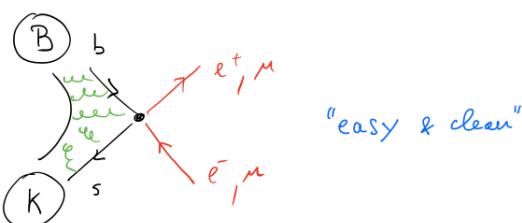
$b \rightarrow s \ell \ell$ transitions occurs to 2nd order in G_F



FCNC ops

$$Q_9^e = \bar{b}_L \gamma^\mu \xi_L \bar{\ell} \gamma_\mu e$$

$$Q_{10}^e = \bar{b}_L \gamma^\mu \xi_L \bar{\ell} \gamma_\mu \gamma_5 e$$



$$C_i^e = C_i^{\mu}$$

\xrightarrow{LFV}

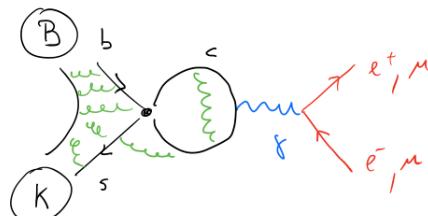
$$C_9^{SM} \approx -C_{10}^{SM}$$

(LH int.)

4-quark ops

$$Q_2^c = \bar{b}_L \gamma^\mu \xi_L \bar{\ell} \gamma^\mu e_L$$

⋮



- obscure short-distance for $q^2 \sim M_Z^2, q^4$

- small for $q^2 \ll M_Z^2$
non-pert effects $\mathcal{O}\left(\frac{\Lambda_{\text{QCD}}^2}{m_c^2}\right)$

KEY observation:

$$\begin{cases} \text{cannot induce } \cancel{LFV} \\ \text{cannot induce } \Delta C_{10} [\Rightarrow B_s \rightarrow \mu^+ \mu^-] \end{cases}$$

If th. error underestimated, we can expect a q^2 -dependent shift in C_9 universal in e & μ