

1. Department of Physics, University of the Free State, Bloemfontein, South Africa 2. Indian Institute of Technology Indore, Indore, India

Radio-loud Active Galactic Nuclei (AGN) are associated with the production of relativistic jets. The spectral energy distribution (SED) of these jets can extend from radio to very high energy γ -rays and exhibit a characteristic double bump structure [1]. The emission at the lower wavelengths (radio-optical) is dominated by synchrotron radiation from a population of non-thermal electrons. Based on the properties of their emission AGN are subdivided into different classes which depend on both the morphology (e.g. FRI and FRII dichotomy) as well as the viewing angle with respect to the relativistic jet influence the properties of the AGN, numerical simulations have been employed. In the past these simulations (e.g. [3]), which do not consider micro-scale physics (such as the evolution of the non-thermal particles), or smaller scale particle in cell simulations, which become too complete emission model of large-scale fluid dynamic simulations of AGN jets, we have employed the new fluid-particle hybrid module of the PLUTO code [5].

RMHD Simulations











Modelling the synchrotron emission of RMHD AGN Jet simulations with the PLUTO particle module

Izak van der Westhuizen,^{1,} Brian van Soelen,¹ Barghav Vaidya²

Introduction

vanderwesthuizenip@ufs.ac.za

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Particle-hybrid module

Using the PLUTO hybrid framework [5], a sample of Lagrangian particles was continuously injected at the jet nozzle, with a random spatial distribution. Each Lagrangian particle represents an ensemble of non-thermal electrons, with a finite power-law energy distribution. The parameters for the initial nonthermal distribution is listed in Table 2. As the particle propagates with the fluid the spectrum is updated considering processes such as adiabatic expansion and radiative cooling. Figure 2 shows a rendering of the Lagrangian particles in a segment of the jet.

 Table 2: Injected Non-thermal electron distribution

Ymin	Υmax	p
10 ²	10 ⁷	2.1

Emission modelling

Using the spectral information of the particles the synchrotron emissivity is calculated during runtime by [5],

$$j'_{syn}(\nu', \hat{n}'_{los}, B') = \frac{\sqrt{3e^3}}{4\pi m_e c^2} |B' \times n'_{los}| \int_{E_i}^{E_f} N'(E') F(x) dE',$$

where $F(x) = x \int_{x}^{\infty} K_{\frac{5}{3}}(z) dz$ and $x = \frac{\nu'}{\nu'_{cr}} = \frac{4\pi m_e^3 c^5 \nu'}{3eE'^2 |B' \times n'_{los}|}$

In addition, we have added the calculation of absorption coefficients to the module using [7],

$$\alpha'_{syn}(\nu',\,\widehat{n}'_{los},\,B') = \frac{\sqrt{3e^3}}{8\pi(m_e c\nu')^2} |B' \times n'_{los}| \int_{E_i}^{E_f} \frac{N'(E')}{E'} \frac{d}{dE'} \left(E'^2 F(x) \right) dE',$$

Relativistic transformations are incorporated into the calculations of the coefficients. These coefficients are integrated along a line of sight using ray tracing to produce intensity maps from the simulations at different radio frequencies. Figure 3 shows the resulting emission and absorption coefficients at 1.4 GHz as well as the resulting intensity map for a viewing angle of 45° with respect to the axis of the jet.

Discussion & Conclusion

The RMHD simulations reproduce a collimated relativistic beam surrounded by a turbulent backflow and cocoon region. For the magnetic field strength considered the jet remains kinetically dominated and, therefore, we do not obtain magnetically driven instabilities. The Lagrangian particles were injected with the jet fluid and flowed along the beam of the jet, after which they were deposited in the backflow and cocoon. The slice of the emissivity shows extended emission from the cocoon region as well as the collimated beam. The highest intensity emission occurs along the edges of the cocoon close to the shock front produced by the jet. The extended cocoon dominates the integrated intensity map and obscures the beam of the jet at the chosen viewing angle.

	References
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Figure 2: Distribution of Lagrangian (1)particles in the simulation, the colour represents the time of injection

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