Connecting B anomalies and τ observables

Claudia Cornella

University of Zurich

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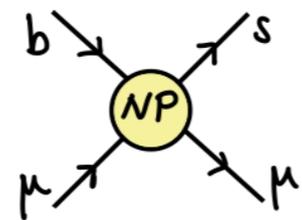
Intro

If the anomalies in $b \to sll$ and $b \to c\tau\nu$ are NP signals, **where else** can we expect to see deviations from the SM?

‣ obs. in the same partonic transition: angular obs. in $B \to D^{(*)}$

   other LFU ratios (R_ϕ , $R_{J/\psi}$...)

‣ obs. with a different quark flavor in the final state: $b \to dll$, $b \to ul\nu$

‣ obs. with τ leptons   LFU tests in τ decays,

   $b \to s\tau\tau, pp \to \tau\tau$

   τ/μ LFV $(\tau \to \mu\phi, \tau \to \mu\gamma, b \to s\tau\mu)$

‣ $B \to K\nu\bar{\nu}...$

Theory lessons from $b \to sll$ and $b \to c\tau\nu$

$b \to sll$ $\qquad$ $\mathcal{O}_{9-10} = (\bar{s}_L \gamma^\mu b_L)(\bar{\mu}_L \gamma_\mu \mu_L)$ (or $\mathcal{O}_9$)
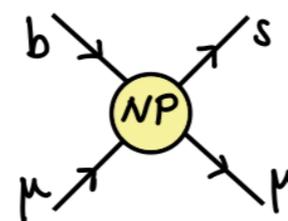
 $\sim 3 \times 10^{-5} \, G_F$ $\quad \Rightarrow \quad$ $\dfrac{g_{\mathrm{NP}}^2}{M_{\mathrm{NP}}^2} \sim \dfrac{1}{(40 \ \mathrm{TeV})^2}$

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▸ mediators: Z', leptoquarks (U_1, S_3, U_3)
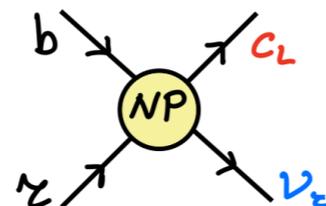
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$b \to c\tau\nu$ $\quad\quad \mathcal{O}_{V_L} = (\bar{c}_L \gamma_\mu b_L)(\bar{\tau}_L \gamma^\mu \nu_L)$ and or $\mathcal{O}_{S_R}/\mathcal{O}_{S_L}, \mathcal{O}_T$

$$ 10^{-2}\, G_F \quad\quad \Rightarrow \quad \frac{g_{\rm NP}^2}{M_{\rm NP}^2} \sim \frac{1}{(2\ \text{TeV})^2} $$

▸ large effects in both low- and high-energy observables around the corner

▸ mediators: leptoquarks (U_1, S_1, R_2)

Towards a combined explanation

To explain both anomalies, NP must be:

- ~ TeV
- coupled dominantly to 3rd family
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$$\mathscr{L} = -\frac{1}{v^2} \left(C_{lq}^{(3)} (\bar{l}_L \gamma^\mu \tau^a l_L)(\bar{q}_L \gamma^\mu \tau^a q_L) + C_{lq}^{(1)} (\bar{l}_L \gamma^\mu l_L)(\bar{q}_L \gamma^\mu q_L) \right)$$

- reproduces known EFT solutions: $C_9^\mu = -C_{10}^\mu$ for $b \to sll$, V_L for $b \to c\tau\nu$

- need $C_{\ell q}^{(3)} \approx C_{\ell q}^{(1)} \equiv C_{LL}$ to avoid too large $b \to s\nu\bar{\nu}$.

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Few **leptoquark** scenarios work:
- $U_1 \sim (3,1)_{2/3}$ (+ UV completion)
- $S_1 \sim (\bar{3},1)_{1/3} + S_3 \sim (\bar{3},3)_{1/3}$
- $R_2 \sim (\bar{3},2)_{7/6} + S_3$ *

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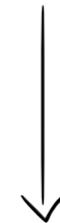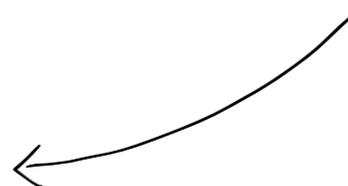No reason to expect large discrepancies in τs based on $b \to sll$ alone.

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If $R_{D^{(*)}}$ is due to NP, the same NP has to show up somewhere in τs:

LFUV in τ decays

LFV in τ and B decays
$(\tau \to \mu\phi, \tau \to \mu\gamma, b \to s\tau\mu)$

$B \to K\tau\tau, B_s \to \tau\tau$

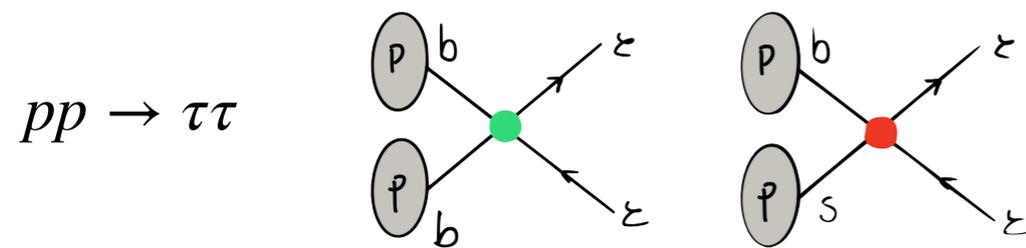$pp \to \tau\tau$
$pp \to \tau\nu$

Implications of $R_{D^{(*)}}$

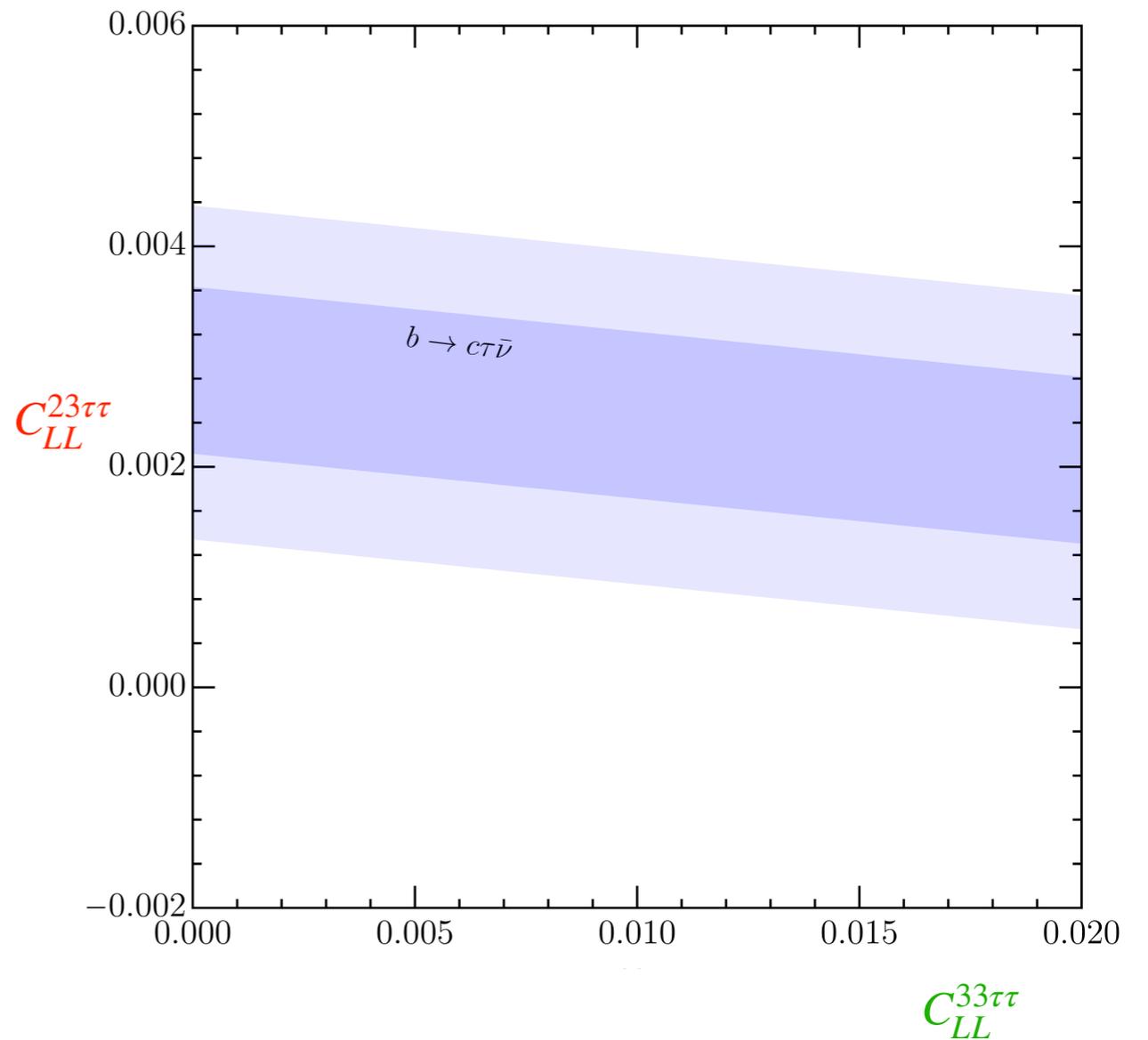$$\mathscr{L} = -\frac{2}{v^2} C_{LL}^{ij\alpha\beta} \left(\bar{q}_L^i \gamma^\mu l_L^\alpha \right)(\bar{l}_L^\beta \gamma_\mu q_L^j)$$

$$\delta R_{D^{(*)}} \equiv \frac{R_{D^{(*)}}}{R_{D^{(*)}}^{\text{SM}}} - 1 = 2\text{Re}\left(C_{LL}^{33\tau\tau} + \frac{V_{cs}}{V_{cb}} C_{LL}^{23\tau\tau} \right)$$

$$\sim 0.01 \qquad \sim 0.002$$

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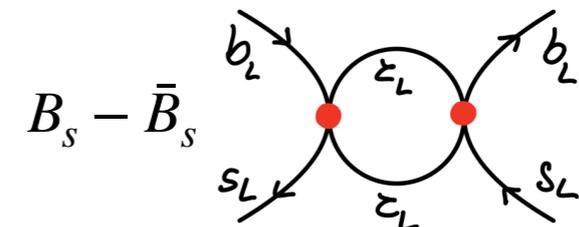$$\mathscr{L} = -\frac{2}{v^2} C_{LL}^{ij\alpha\beta} \left(\bar{q}_L^i \gamma^\mu l_L^\alpha\right)\left(\bar{l}_L^\beta \gamma_\mu q_L^j\right)$$
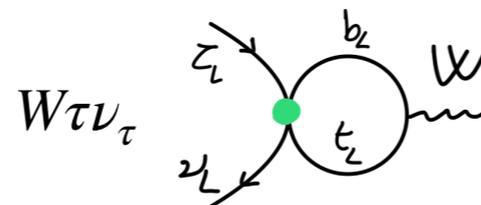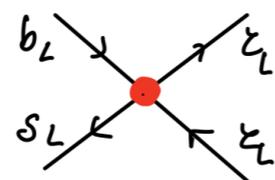
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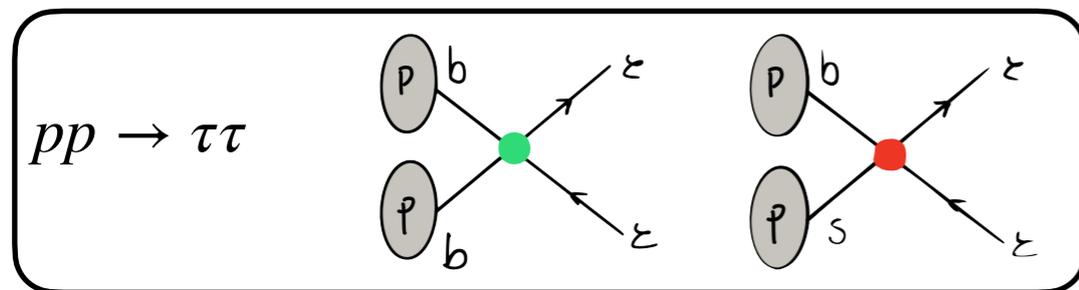$W\tau\nu_\tau$

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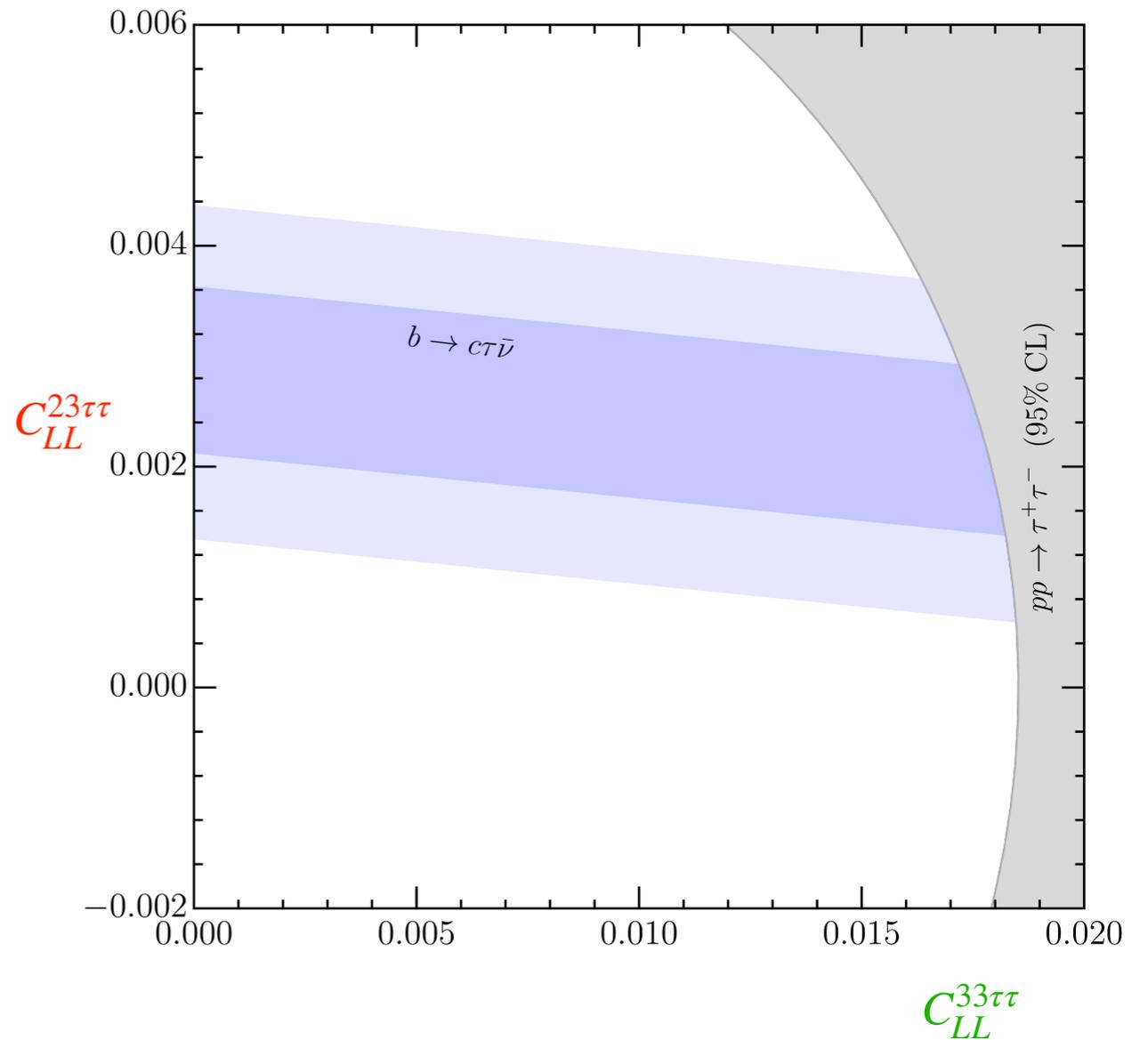$$\mathcal{L} = -\frac{2}{v^2} C_{LL}^{ij\alpha\beta} \left(\bar{q}_L^i \gamma^\mu l_L^\alpha\right)\left(\bar{l}_L^\beta \gamma_\mu q_L^j\right)$$

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$pp \to \tau^+\tau^-$ (95% CL)

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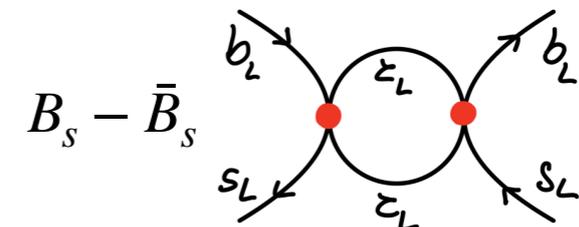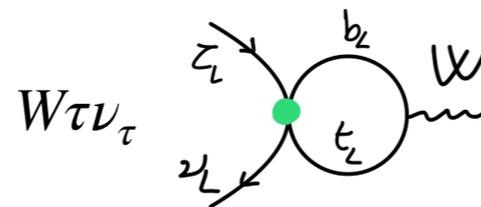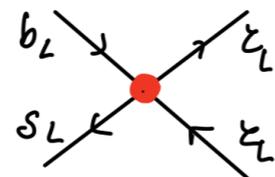
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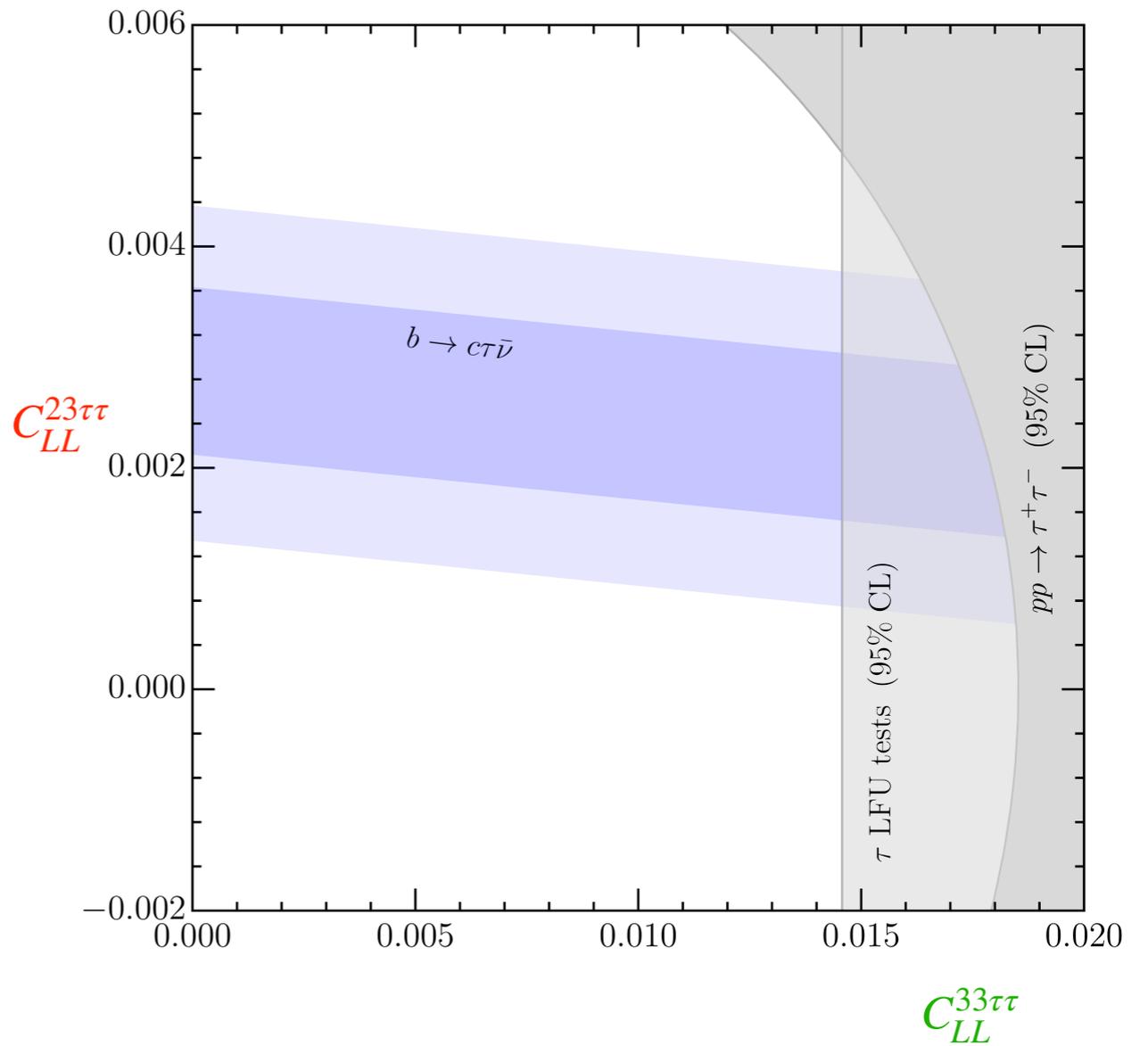
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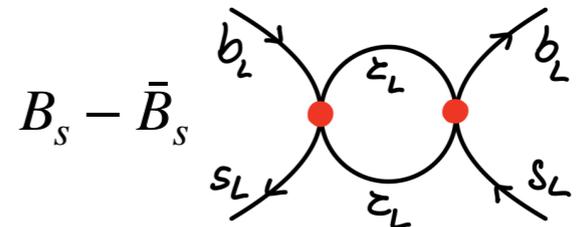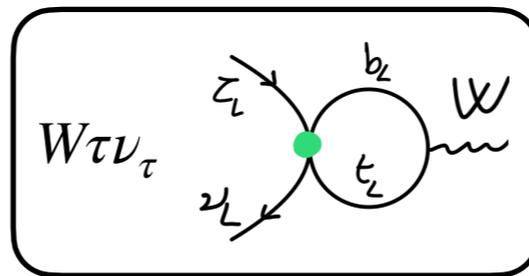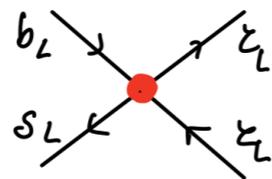
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LFUV in τ decays

▸ If NP couples only to LH fermions, $\delta R_{D^{(*)}}$ unavoidably implies per mil LFUV in τs:

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▸ If $\delta R_{D^{(*)}}$ is addressed only via scalar/tensor (e.g. R_2), no τ LFUV.

$b \to s\tau\tau$ at low energy

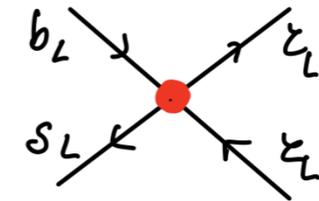Another "consequence" of $\delta R_{D^{(*)}}$ is a large enhancement of $b \to s\tau\tau$ transitions.

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▸ LH case:

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$$\frac{\mathscr{B}(B_s \to \tau\tau)}{\mathscr{B}(B_s \to \tau\tau)_{\text{SM}}} \approx \frac{\mathscr{B}(B \to K\tau\tau)}{\mathscr{B}(B \to K\tau\tau)_{\text{SM}}} \approx \frac{|\Delta C_9^\tau|^2}{|C_{9,\text{SM}}|^2} \approx 1 \cdot 10^2 \quad \text{i.e. } \text{BR} \sim \text{few} \cdot 10^{-5}$$

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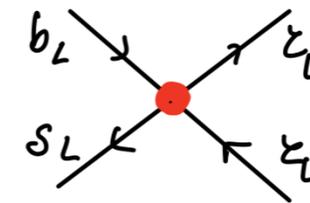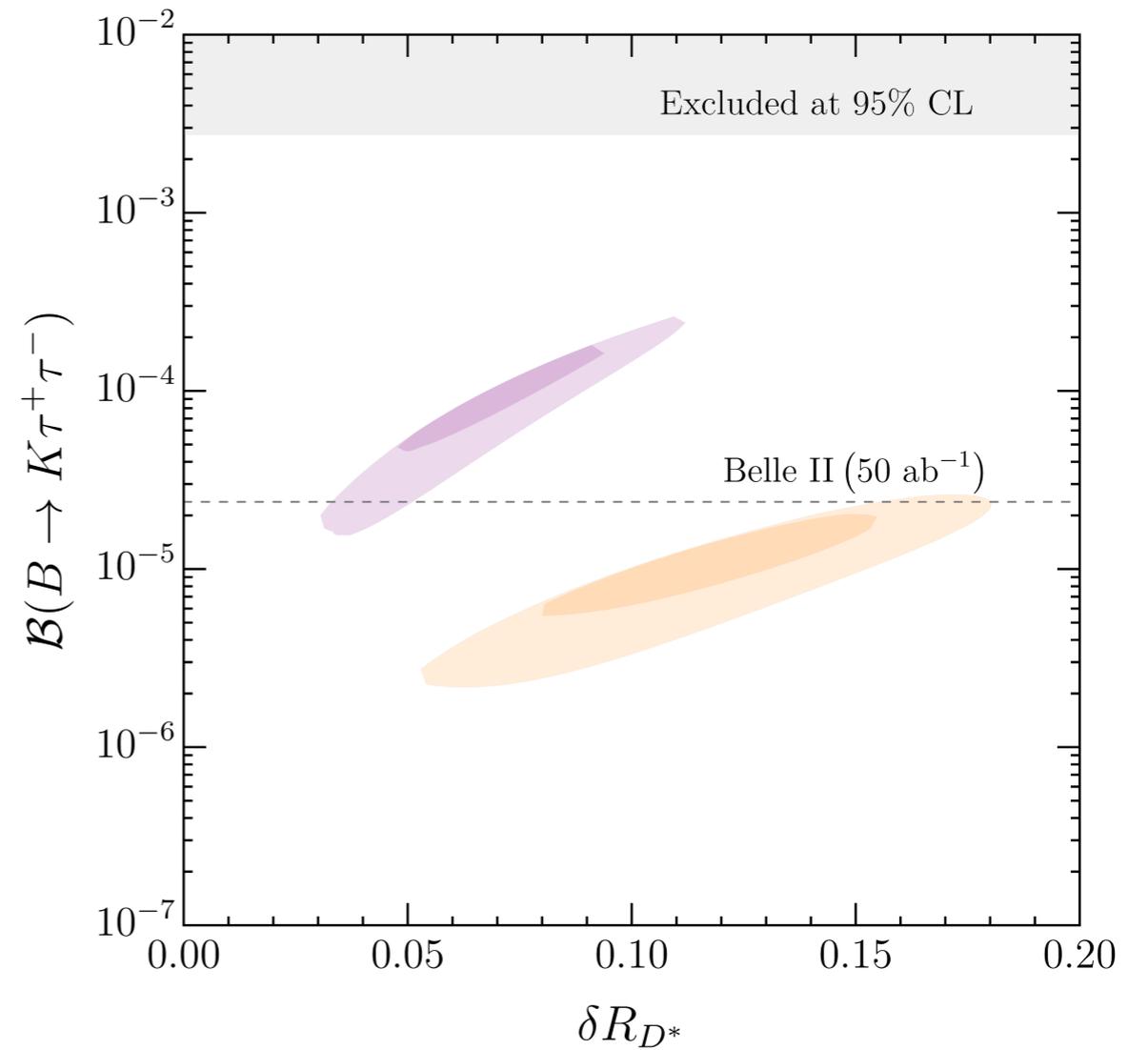

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▸ with RH couplings one gets scalar/pseudo-scalar operators, whose contribution is chirally enhanced

E.g. U_1 coupled to LH & RH 3rd fam.

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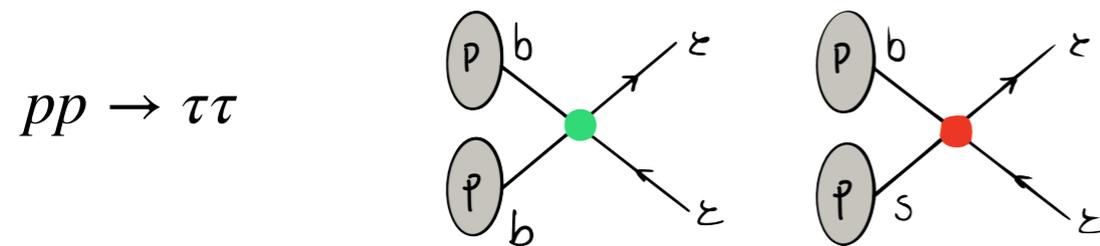
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$b \rightarrow s\tau\tau$ at low energy for the U_1

$bb(s) \to \tau\tau$ at high pT

The same interaction can be probed in **di-tau tails** at the LHC.

Generally **stronger** than low-E bounds!

$pp \to \tau\tau$

U_1 sol. will be completely probed at HL-LHC, same for $R_2 + S_3$, still space left for $S_1 + S_3$.

Models for $R_D^{(*)}$ only yield similar enhancements in $B \to K\tau\tau$, $B_s \to \tau\tau$ and $pp \to \tau\tau$.

U_1 (LH)

[CC, Fuentes-Martin, Faroughy, Isidori, Neubert, 2101.11626]

g_U

M_U [TeV]

$pp \to U_1 U_1^*$ (13 TeV 137 fb^{-1})

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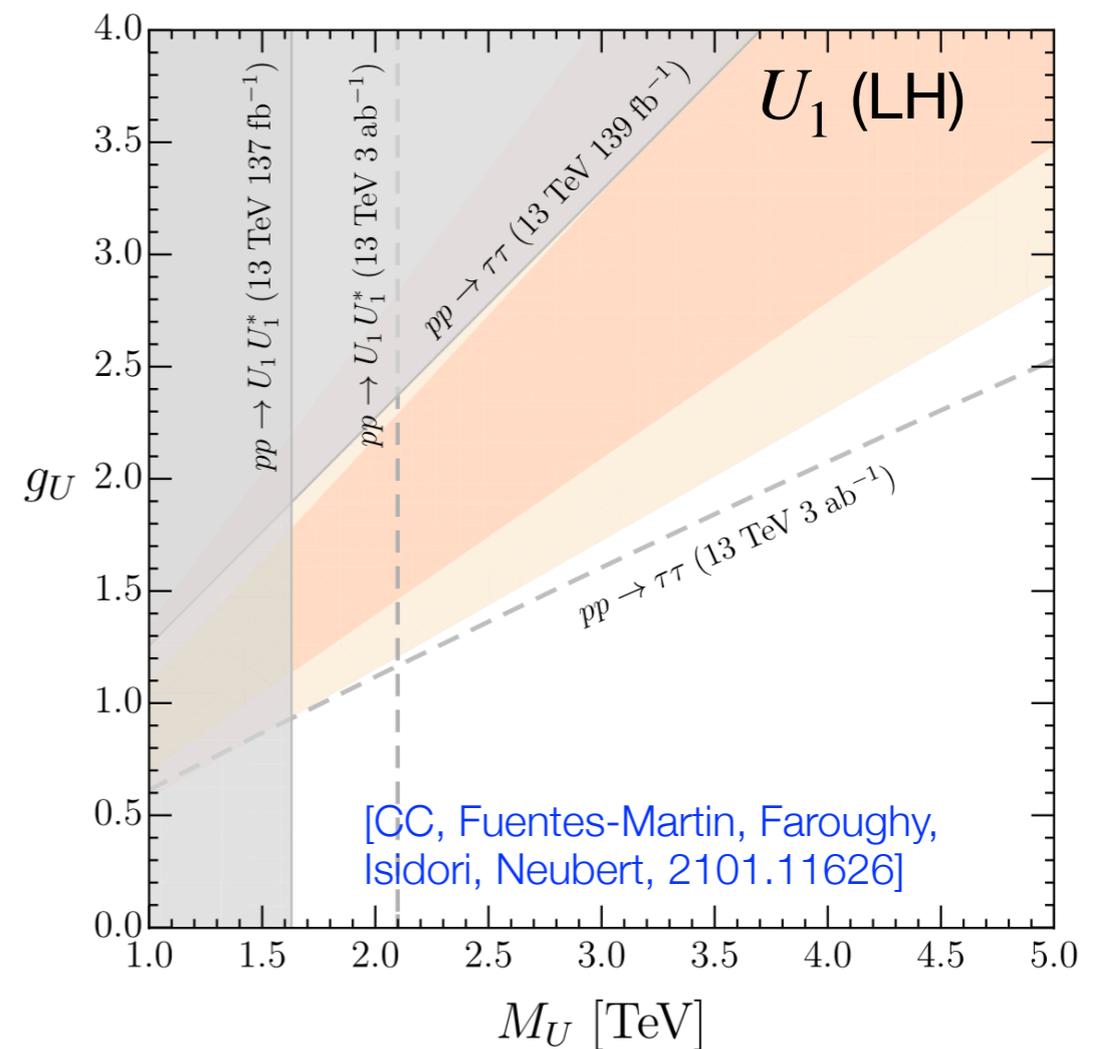
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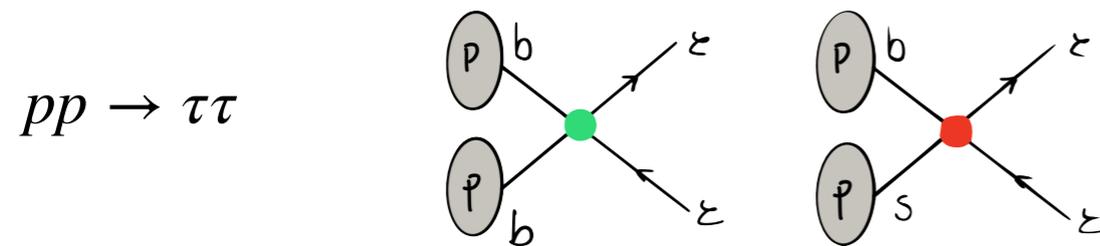
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LFV in $b \rightarrow s\tau\mu$ and τ decays

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$$\delta R_{D^{(*)}} \equiv \frac{R_{D^{(*)}}}{R_{D^{(*)}}^{\mathrm{SM}}} - 1 = 2\mathrm{Re}\left(C_{LL}^{33\tau\tau} + \frac{V_{cs}}{V_{cb}} C_{LL}^{23\tau\tau} \right) \qquad \Delta C_9^\mu = -\frac{2\pi}{\alpha V_{ts}^* V_{tb}} C_{LL}^{23\mu\mu}$$

LFV in $b \to s\tau\mu$ and τ decays

$b \to sll$ **and** $b \to c\tau\nu$ $\longleftrightarrow$ τ/μ **LFV** in $b \to s\tau\mu$ and τ decays

$$\delta R_{D^{(*)}} \equiv \frac{R_{D^{(*)}}}{R_{D^{(*)}}^{\text{SM}}} - 1 = 2\text{Re}\left({\color{green}C_{LL}^{33\tau\tau}} + \frac{V_{cs}}{V_{cb}}{\color{red}C_{LL}^{23\tau\tau}}\right) \qquad \Delta C_9^\mu = -\frac{2\pi}{\alpha V_{ts}^* V_{tb}}{\color{blue}C_{LL}^{23\mu\mu}}$$

$$\mathscr{L} \supset \frac{g_U}{\sqrt{2}} U_1^\mu \beta_L^{i\alpha}(\bar{q}_L^i \gamma_\mu \ell_L^\alpha) + \text{h.c.} \qquad \beta_L = \begin{bmatrix} 0 & 0 & 0 \\ 0 & \beta_{s\mu} & \beta_{s\tau} \\ 0 & \beta_{b\mu} & \beta_{b\tau} \end{bmatrix}$$

LFV in $b \to s\tau\mu$ and τ decays

$b \to sll$ **and** $b \to c\tau\nu$ $\longleftrightarrow$ τ/μ **LFV** in $b \to s\tau\mu$ and τ decays

$$\delta R_{D^{(*)}} \equiv \frac{R_{D^{(*)}}}{R_{D^{(*)}}^{\mathrm{SM}}} - 1 = 2\mathrm{Re}\left(C_{LL}^{33\tau\tau} + \frac{V_{cs}}{V_{cb}} C_{LL}^{23\tau\tau} \right) \qquad \Delta C_9^\mu = -\frac{2\pi}{\alpha V_{ts}^* V_{tb}} C_{LL}^{23\mu\mu}$$

$$\mathscr{L} \supset \frac{g_U}{\sqrt{2}} U_1^\mu \beta_L^{i\alpha} (\bar{q}_L^i \gamma_\mu \ell_L^\alpha) + \mathrm{h.c.} \qquad \beta_L = \begin{bmatrix} 0 & 0 & 0 \\ 0 & \beta_{s\mu} & \beta_{s\tau} \\ 0 & \beta_{b\mu} & \beta_{b\tau} \end{bmatrix}$$

LFV in $b \to s\tau\mu$ and τ decays

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$$\mathscr{L} \supset \frac{g_U}{\sqrt{2}} U_1^\mu \beta_L^{i\alpha} (\bar{q}_L^i \gamma_\mu \ell_L^\alpha) + \text{h.c.} \qquad \beta_L = \begin{bmatrix} 0 & 0 & 0 \\ 0 & \beta_{s\mu} & \beta_{s\tau} \\ 0 & \beta_{b\mu} & \beta_{b\tau} \end{bmatrix}$$

$$B_s \to \tau\mu, B \to K\tau\mu$$

LFV in $b \rightarrow s\tau\mu$ and τ decays

$b \rightarrow sll$ **and** $b \rightarrow c\tau\nu$ $\longleftrightarrow$ τ/μ **LFV** in $b \rightarrow s\tau\mu$ and τ decays

$$\delta R_{D^{(*)}} \equiv \frac{R_{D^{(*)}}}{R_{D^{(*)}}^{\text{SM}}} - 1 = 2\text{Re}\left(C_{LL}^{33\tau\tau} + \frac{V_{cs}}{V_{cb}}C_{LL}^{23\tau\tau}\right) \qquad \Delta C_9^\mu = -\frac{2\pi}{\alpha V_{ts}^* V_{tb}}C_{LL}^{23\mu\mu}$$

$$\mathscr{L} \supset \frac{g_U}{\sqrt{2}} U_1^\mu \beta_L^{i\alpha}(\bar{q}_L^i \gamma_\mu \ell_L^\alpha) + \text{h.c.} \qquad \beta_L = \begin{bmatrix} 0 & 0 & 0 \\ 0 & \beta_{s\mu} & \beta_{s\tau} \\ 0 & \beta_{b\mu} & \beta_{b\tau} \end{bmatrix} \tau \rightarrow \mu\phi$$

$$B_s \rightarrow \tau\mu, B \rightarrow K\tau\mu$$

LFV in $b \to s\tau\mu$ and τ decays

$$b \to sll \text{ and } b \to c\tau\nu \qquad \longleftrightarrow \qquad \tau/\mu \text{ LFV in } b \to s\tau\mu \text{ and } \tau \text{ decays}$$

$$\delta R_{D^{(*)}} \equiv \frac{R_{D^{(*)}}}{R_{D^{(*)}}^{\text{SM}}} - 1 = 2\text{Re}\left(C_{LL}^{33\tau\tau} + \frac{V_{cs}}{V_{cb}} C_{LL}^{23\tau\tau} \right) \qquad \Delta C_9^\mu = -\frac{2\pi}{\alpha V_{ts}^* V_{tb}} C_{LL}^{23\mu\mu}$$

$$\mathcal{L} \supset \frac{g_U}{\sqrt{2}} U_1^\mu \beta_L^{i\alpha} (\bar{q}_L^i \gamma_\mu \ell_L^\alpha) + \text{h.c.} \qquad \beta_L = \begin{bmatrix} 0 & 0 & 0 \\ 0 & \beta_{s\mu} & \beta_{s\tau} \\ 0 & \beta_{b\mu} & \beta_{b\tau} \end{bmatrix} \begin{array}{l} \\ \tau \to \mu\phi \\ \\ \end{array}$$

$$\tau \to \mu\gamma$$

$$B_s \to \tau\mu \,, B \to K\tau\mu$$

These predictions are subject to large uncertainties!
(B anomalies fix product of couplings, not individual coefficients…)

LFV in $b \to s\tau\mu$ and τ decays

▸ LH case

$$\mathscr{B}(B \to K\tau\mu) \approx \mathscr{B}(B_s \to \tau\mu) \approx 10^{-7} - 10^{-5}$$

$$\mathscr{B}(\tau \to \mu\phi) \approx 10^{-11} - 10^{-8}$$

▸ With RH currents (U_1)

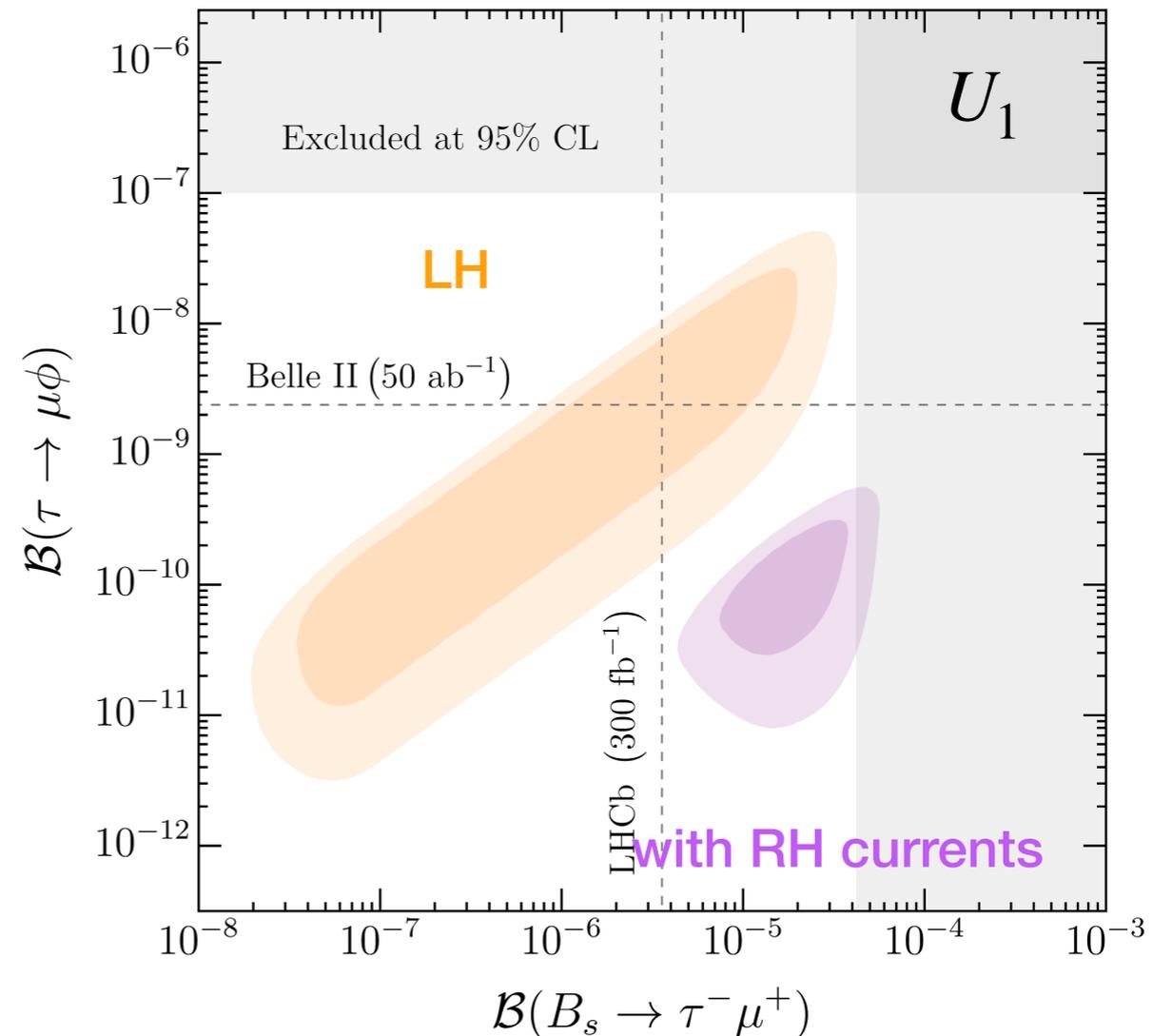$$\mathscr{B}(B_s \to \tau\mu) \approx 1 \cdot 10^{-5}$$

$$\mathscr{B}(B \to K\tau\mu) \approx 1 \cdot 10^{-6}$$

$$\mathscr{B}(\tau \to \mu\phi) \approx 10^{-11} - 10^{-10}$$

sensitive to the vector contribution only

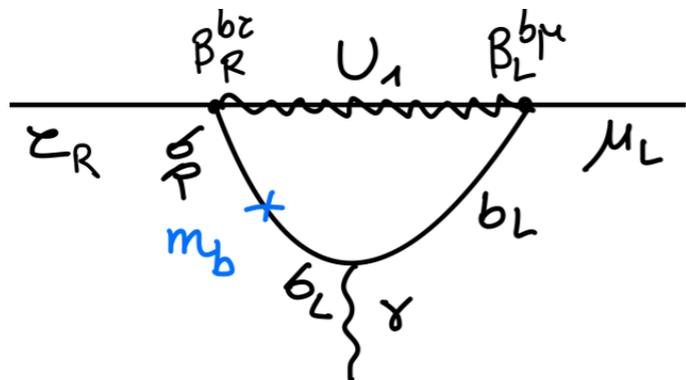sensitive to (chirally enhanced) scalar contribution, hence larger than in the LH case
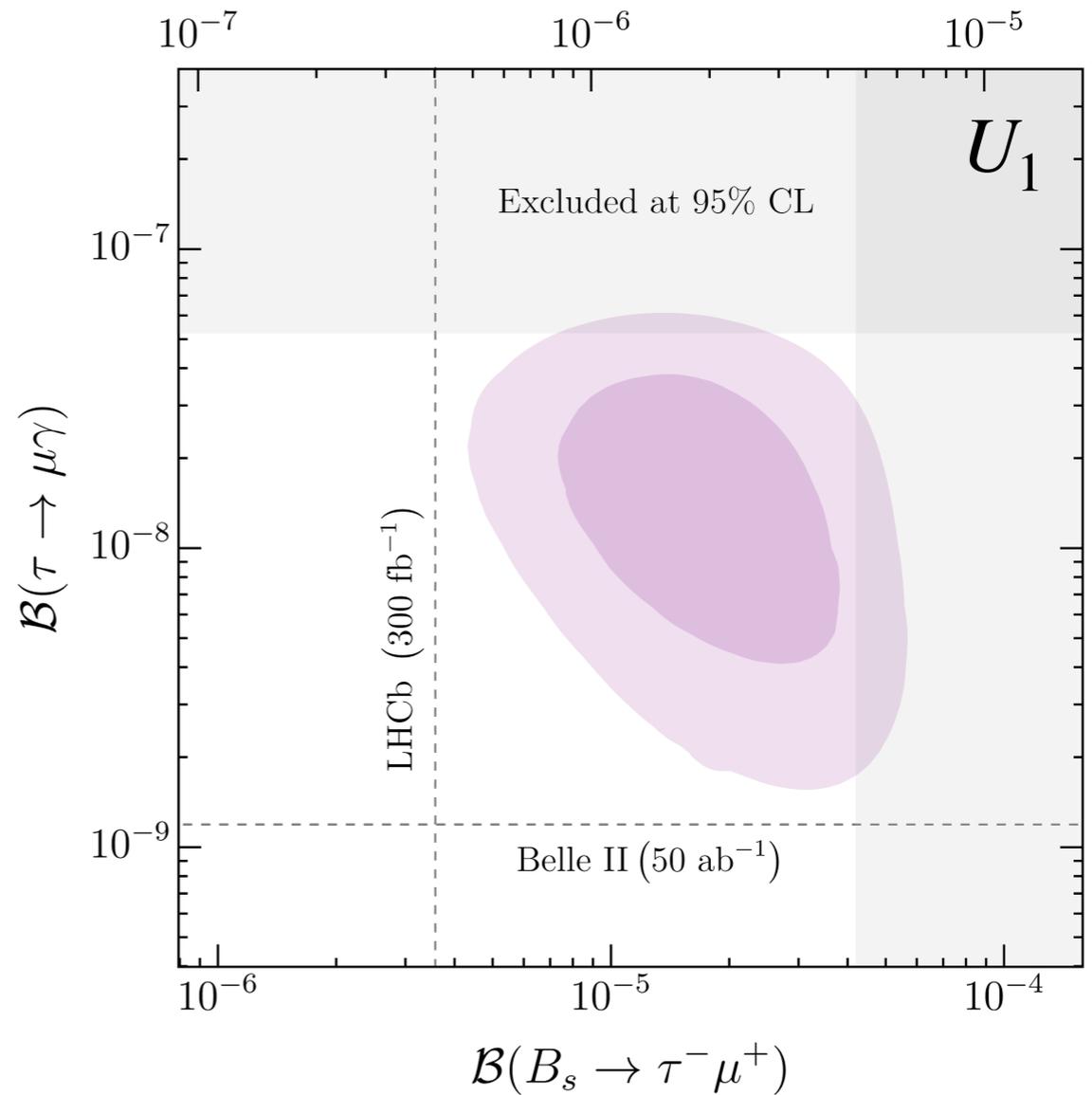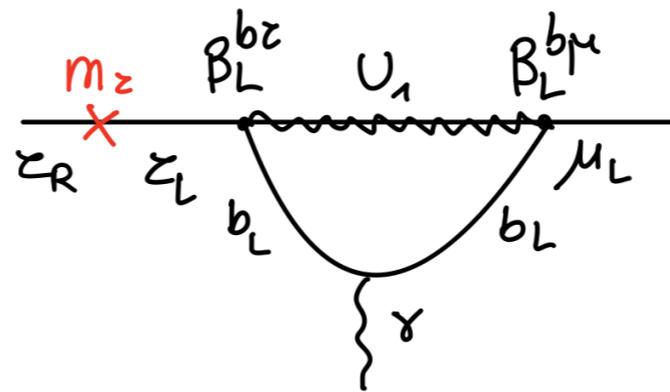
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$\tau - \mu\gamma$

▸ NP couples only to LH fermions

chiral suppression!

▸ RH couplings can lift it:

$$\mathscr{B}(\tau \to \mu\gamma) \approx 10^{-8}$$

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Conclusions

Present data in $b \to c$ imply large effects in many obs. at low and high energy, putting a fair amount of pressure on the parameter space of all* NP models for $R_{D^{(*)}}$.

No way out on the model building side: need experimental corroboration to guide us!

Obs. with τ leptons provide a strong consistency test of the $b \to c$ anomaly and help us characterize / rule out possible NP effects.

Ranked by constraining power:
- τ LFUV , $pp \to \tau\tau$
- τ/μ LFV in $b \to s\tau\mu$ and τ decays