Exploration of the phase diagram within a transport approach

Olga Soloveva

P. Moreau, L. Oliva, T. Song, I. Grishmanovskii, V. Voronyuk, V. Kireyeu, E. Bratkovskaya

In collaboration with D. Fuseau, J. Aichelin

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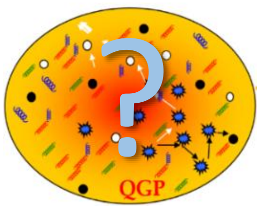
Busan, Republic of Korea

Properties of QGP: transport coefficients

"Numerical simulations are now essential to make contact between theory and experiment" - from talk J. Kapusta

! One has to specify transport and microscopic properties as well as EoS for theoretical simulations of HICs (hydro / transport approaches)



EoS(ε,n)
$\sigma(\sqrt{s}, m_q, m_q, T, \mu_B)$
$m(T, \mu_B)$

On practice: effective models for QGP

Transport simulations with QGP phase:

Catania transport – QuasiParticle Model

F. Scardina, S. K. Das, V. Minissale, S. Plumari, and V. Greco, PRC 96, 044905 (2017). See talk M. L. Sambataro

PHSD – Dynamical QPM for partonic phase

W. Cassing, E. Bratkovskaya, PRC 78 (2008) 034919
P. Moreau, O. S , L. Oliva, T. Song, W. Cassing, E. Bratkovskaya, PRC 100 (2019) , 014911;
O. S, P. Moreau, L. Oliva, V. Voronyuk, V. Kireyeu, T. Song, E. Bratkovskaya, Particles 3 (2020), 178-192

AMPT – PNJL EoS (Mean field potentials)

K.J. Sun, C. M. Ko, and Z.-W. Lin, PRC 103(2021)

Hybrid simulations with QGP:

vHLLE/Music+UrQMD/SMASH

Iu.A. Karpenko, P. Huovinen, H. Petersen and M. Bleicher PRC 91 (2015), 064901. See poster Iu. Karpenko BLK15
S. Ryu, J.F.Paquet, C. Shen, G.S. Denicol, B. Schenke PRL 115 (2015), 132301

Today:

Transport coefficients at finite **T** and μ_B

1.) crossover, CEP and 1st order phase transition (Nf =3 PNJL model)
2.) crossover + CEP (Nf = 3 DQPM)

Properties of QGP: transport coefficients

Hydrodynamics

$$\begin{cases} \partial_\mu T^{\mu\nu} = 0 & T^{\mu\nu} = -Pg^{\mu\nu} + wu^\mu u^\nu + \Delta T^{\mu\nu} \\ \partial_\mu J_B^\mu = 0 & J_B^\mu = n_B u^\mu + \Delta J_B^\mu \end{cases}$$

$$\boxed{\eta}\left(D^\mu u^\nu + D^\nu u^\mu + \tfrac{2}{3}\Delta^{\mu\nu}\partial_\rho u^\rho\right) - \boxed{\zeta}\Delta^{\mu\nu}\partial_\rho u^\rho$$

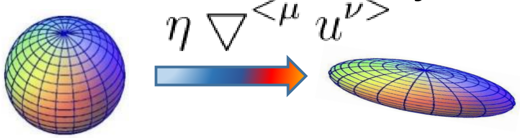
$$\Delta J_B^\mu = \boxed{\kappa_B} D^\mu\left(\frac{\mu_B}{T}\right)$$

input for hydro

Transport coefficients:

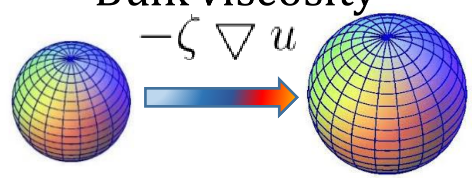
Shear viscosity

$\eta \nabla^{<\mu} u^{\nu>}$



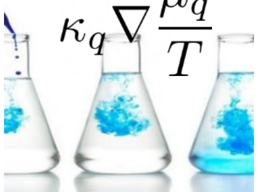
Bulk viscosity

$-\zeta \nabla u$



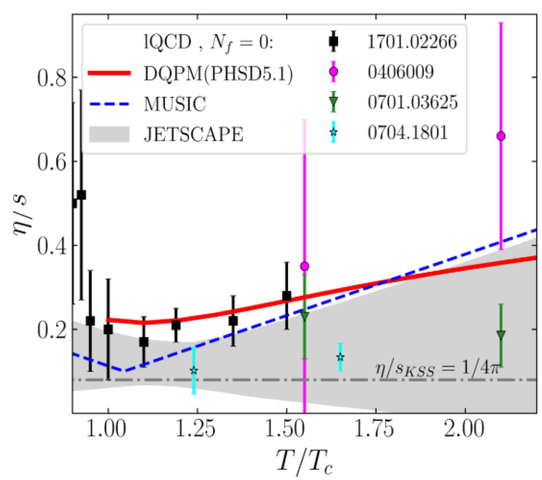
(B, Q, S) diffusion coefficients

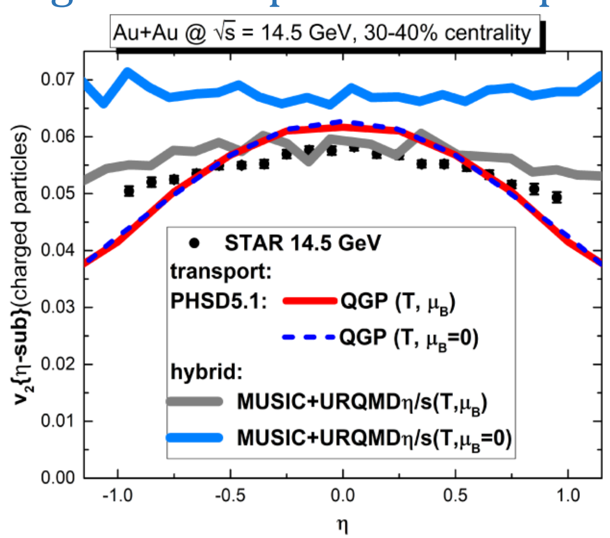
$\kappa_q \nabla \frac{\mu_q}{T}$



Model predictions for QGP: **!** same EoS but different transport coefficients

Transport coefficients can serve a bridge for comparison transport and hydro

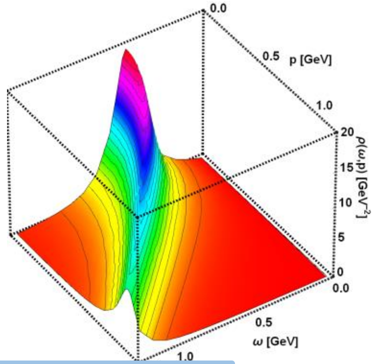




MUSIC: C. Shen, S.Alzhrani, PRC 102 (2020) 1, 014909

Dynamical Quasi-Particle Model

The QGP phase is described in terms of strongly-interacting quasiparticles - quarks and gluons with Lorentzian spectral functions:



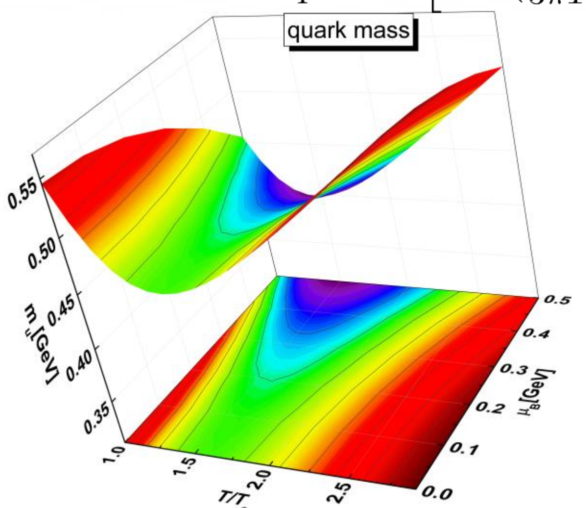
$$\rho_j(\omega,\mathbf{p}) = \frac{\gamma_j}{\tilde{E}_j}\left(\frac{1}{(\omega-\tilde{E}_j)^2+\gamma_j^2} - \frac{1}{(\omega+\tilde{E}_j)^2+\gamma_j^2}\right)$$

$$\equiv \frac{4\omega\gamma_j}{\left(\omega^2-\mathbf{p}^2-M_j^2\right)^2+4\gamma_j^2\omega^2}$$

resummed propagators: $\Delta_i(\omega,\mathbf{p}) = \frac{1}{\omega^2-\mathbf{p}^2-\Pi_i}$ & self-energies: $\Pi_i = m_i^2 - 2i\gamma_i\omega$

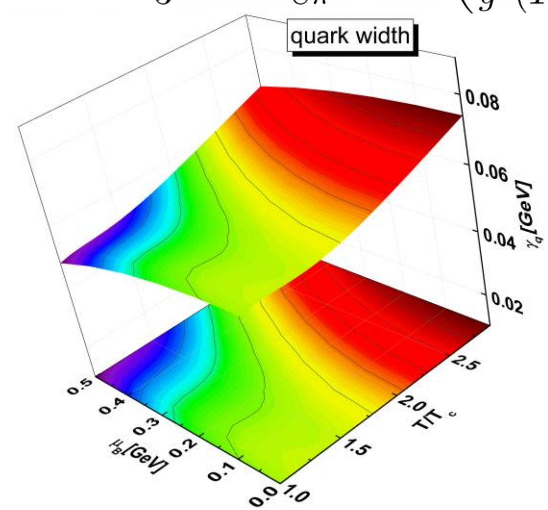
Re $\mathbf{\Pi_i}$: thermal mass (M_g, M_q)

$$m^2_{q(\bar{q})}(T,\mu_\mathrm{B}) = C_q \frac{g^2(T,\mu_\mathrm{B})}{4} T^2 \left[1+\left(\frac{\mu_B}{3\pi T}\right)^2\right]$$



Im $\mathbf{\Pi_i}$: interaction width (γ_g, γ_q)

$$\gamma_j(T,\mu_\mathrm{B}) = \frac{1}{3} C_j \frac{g^2(T,\mu_\mathrm{B})T}{8\pi} \ln\left(\frac{2c_m}{g^2(T,\mu_\mathrm{B})}+1\right)$$



Peshier, Cassing, PRL 94 (2005) 172301; Cassing, NPA 791 (2007) 365; H. Berrehrah, E. Bratkovskaya, T. Steinert, W. Cassing, Int. J. Mod. Phys. E 25 (2016), 164200; P. Moreau, O. Soloveva , L. Oliva, T. Song, W. Cassing, E. Bratkovskaya, PRC 100 (2019) , 014911;

DQPM: EoS

Entropy and baryon density in the quasiparticle limit
(G. Baym 1998, Blaizot et al. 2001):

$$s^{dqp} = -\int \frac{d\omega}{2\pi}\frac{d^3p}{(2\pi)^3}\Big[d_g \frac{\partial n_B}{\partial T}\left(\mathrm{Im}(\ln -\Delta^{-1}) + \mathrm{Im}\,\Pi\,\mathrm{Re}\,\Delta\right) + \sum_{q=u,d,s} d_q \frac{\partial n_F(\omega-\mu_q)}{\partial T}\left(\mathrm{Im}(\ln -S_q^{-1}) + \mathrm{Im}\,\Sigma_q\,\mathrm{Re}\,S_q\right) + \sum_{\bar q=\bar u,\bar d,\bar s} d_{\bar q} \frac{\partial n_F(\omega+\mu_q)}{\partial T}\left(\mathrm{Im}(\ln -S_{\bar q}^{-1}) + \mathrm{Im}\,\Sigma_{\bar q}\,\mathrm{Re}\,S_{\bar q}\right)\Big]$$

$$n^{dqp} = -\int \frac{d\omega}{2\pi}\frac{d^3p}{(2\pi)^3}\Big[\sum_{q=u,d,s} d_q \frac{\partial n_F(\omega-\mu_q)}{\partial \mu_q}\left(\mathrm{Im}(\ln -S_q^{-1}) + \mathrm{Im}\,\Sigma_q\,\mathrm{Re}\,S_q\right) + \sum_{\bar q=\bar u,\bar d,\bar s} d_{\bar q} \frac{\partial n_F(\omega+\mu_q)}{\partial \mu_q}\left(\mathrm{Im}(\ln -S_{\bar q}^{-1}) + \mathrm{Im}\,\Sigma_{\bar q}\,\mathrm{Re}\,S_{\bar q}\right)\Big]$$

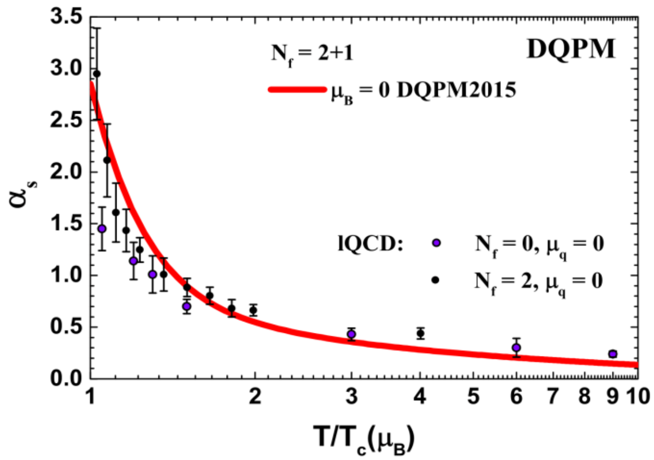
- Input: entropy density as a f(T , $\boldsymbol{\mu_B = 0}$)

$$g^2(s/s_{SB}) = d\left((s/s_{SB})^e - 1\right)^f$$
$$s^{DQPM}(\Pi, \Delta, S_q, \Sigma) = s^{lattice}$$

fix the parameters

- Scaling hypothesis for the crossover region at finite $\boldsymbol{\mu_B}$

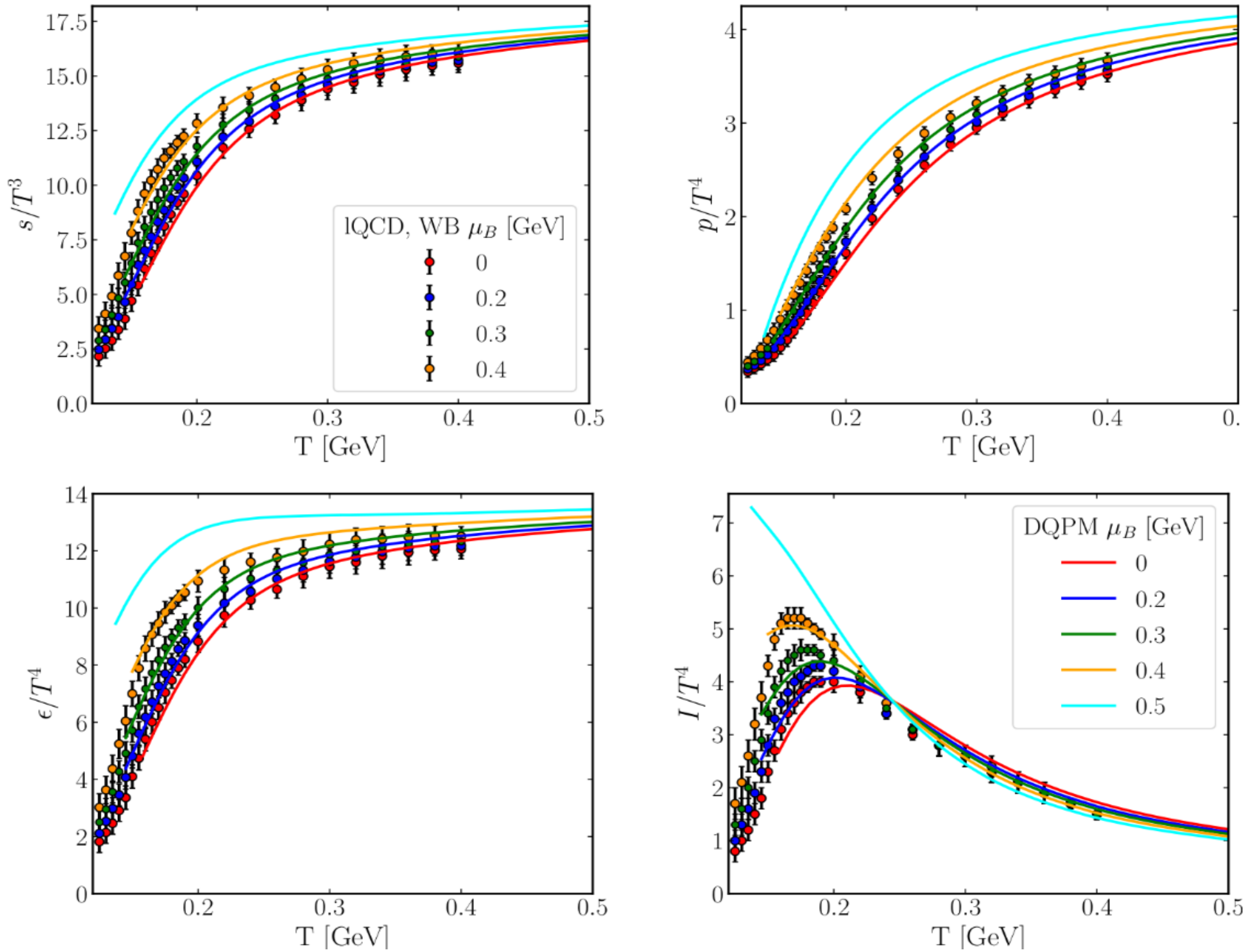
$$g^2(T/T_c, \mu_B) = g^2\left(\frac{T^*}{T_c(\mu_B)}, \mu_B = 0\right) \quad \text{with} \quad T^* = \sqrt{T^2 + \mu_q^2/\pi^2}$$



DQPM: EoS

Input:
lattice EoS
$\mu_B = 0$
(red dots)

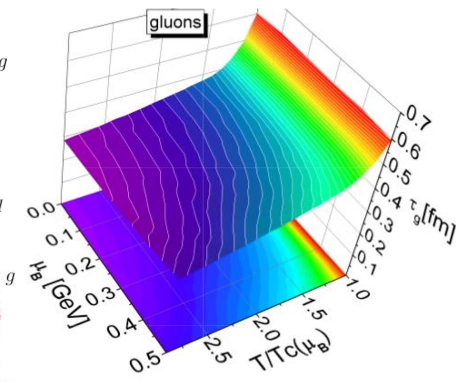
Output:
DQPM EoS
$\mu_B \geq 0$
(lines)



Transport coefficients at finite μ_B

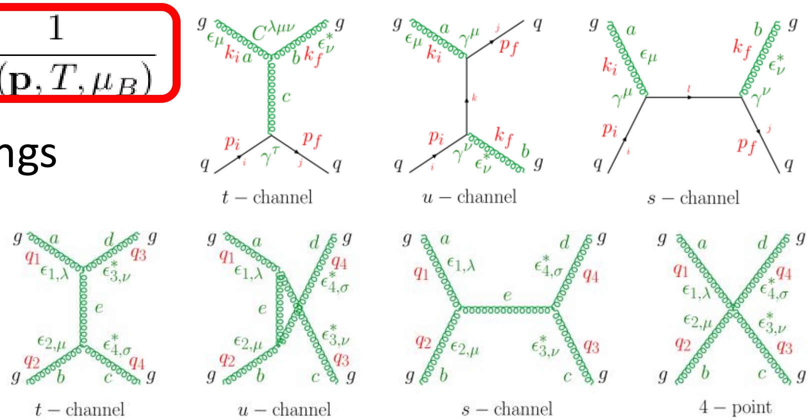
$$\eta^{\mathrm{RTA}}(T,\mu_B) = \frac{1}{15T} \sum_{i=q,\bar{q},g} \int \frac{d^3p}{(2\pi)^3} \frac{\mathbf{p}^4}{E_i^2} \tau_i(\mathbf{p},T,\mu_B)\, d_i(1 \pm f_i) f_i$$

Relaxation times

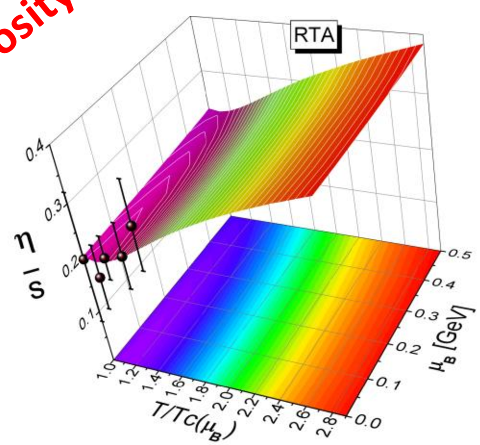


$$\boxed{\tau_i(\mathbf{p},T,\mu_B) = \frac{1}{\Gamma_i(\mathbf{p},T,\mu_B)}}$$

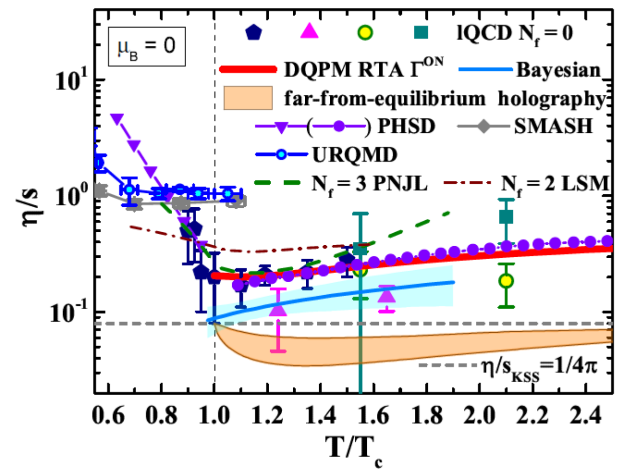
2 ↔ 2 scatterings



Specific shear viscosity



μ_B

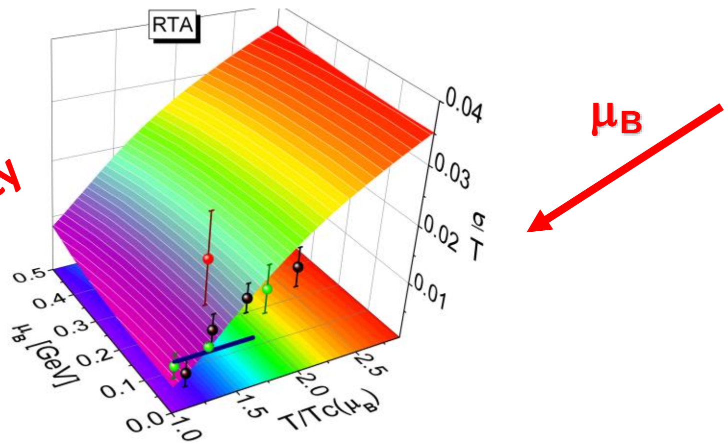


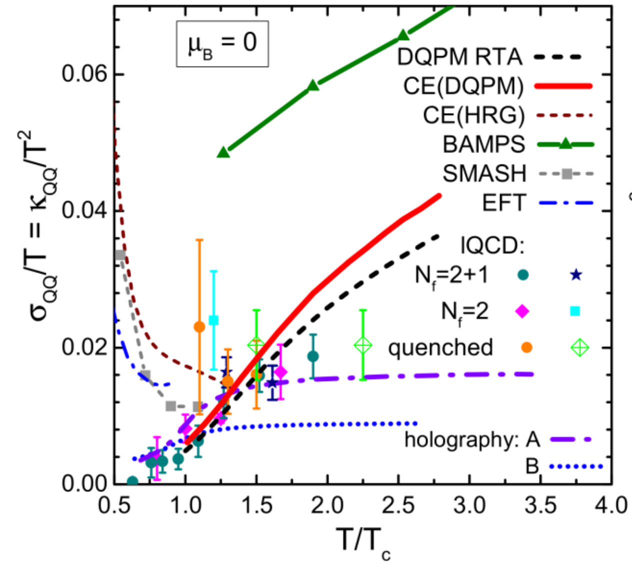
O. Soloveva, P. Moreau and E. Bratkovskaya, PRC 101 (2020), 045203

- Good agreement with lQCD predictions and Bayesian estimates
- Light increase with μ_B in the crossover region for viscosities and electric conductivity

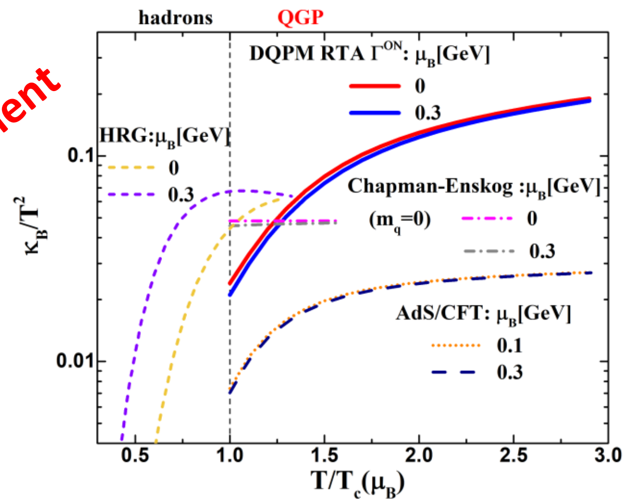
Transport coefficients at finite μ_B

Electric conductivity





Baryon diffusion coefficient



+ Full diffusion coefficient matrix has been evaluated $\sigma_q = \kappa_q/T$

$$\begin{pmatrix} j_B^\mu \\ j_Q^\mu \\ j_S^\mu \end{pmatrix} = \begin{pmatrix} \kappa_{BB} & \kappa_{BQ} & \kappa_{BS} \\ \kappa_{QB} & \kappa_{QQ} & \kappa_{QS} \\ \kappa_{SB} & \kappa_{SQ} & \kappa_{SS} \end{pmatrix} \cdot \begin{pmatrix} \nabla^\mu \alpha_B \\ \nabla^\mu \alpha_Q \\ \nabla^\mu \alpha_S \end{pmatrix}$$

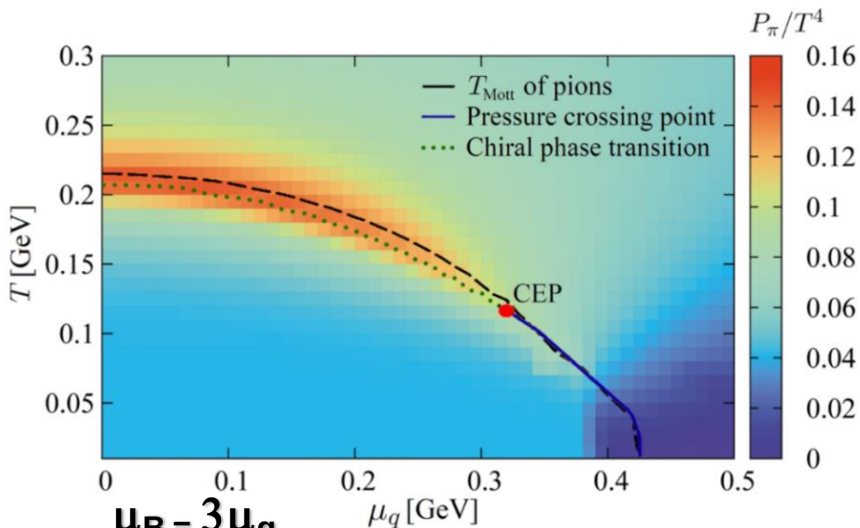
J. A. Fotakis, O. Soloveva, C. Greiner, O. Kaczmarek and E. Bratkovskaya PRD 104 (2021) , 034014

- Light increase with μ_B in the crossover region for shear and bulk viscosities and electric conductivity
- Baryon diffusion coefficients decrease with μ_B

O. Soloveva, P. Moreau and E. Bratkovskaya, PRC 101 (2020), 045203

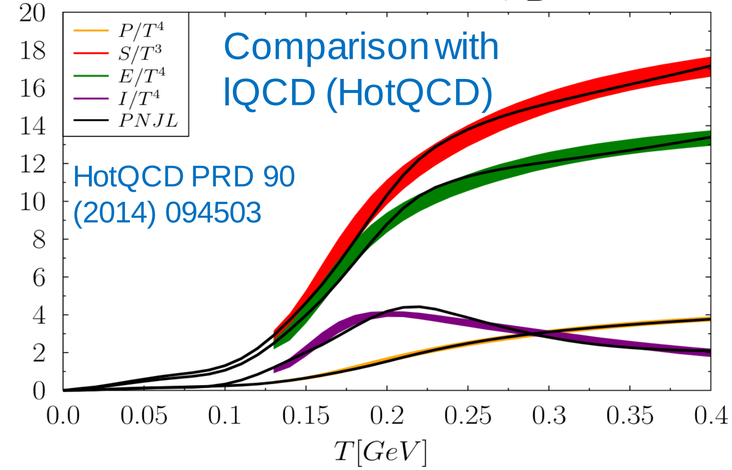
QGP in the Polyakov extended NJL model

- PNJL allows for prediction of macroscopic properties of QGP at finite **T** and μ_B
- & QGP transport coefficients for $0 \le \mu_B \le 1.2$ GeV



$\mu_B = 3\mu_q$

➢ Parameters fixed, EoS at $\mu_B = 0$:



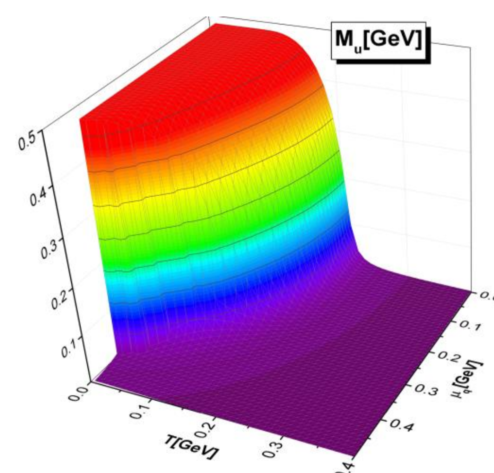
➢ CEP: $(T, \mu_B) = (110{,}960)$ MeV, $\mu_B/T = 8.73$

➢ 1st order PT at high μ_B

➢ same symmetries for the quarks as QCD

Chiral masses (M_l, M_s)

$$m_i = m_{0i} - 4G\langle\langle\bar{\psi}_i\psi_i\rangle\rangle + 2K\langle\langle\bar{\psi}_j\psi_j\rangle\rangle\langle\langle\bar{\psi}_k\psi_k\rangle\rangle$$

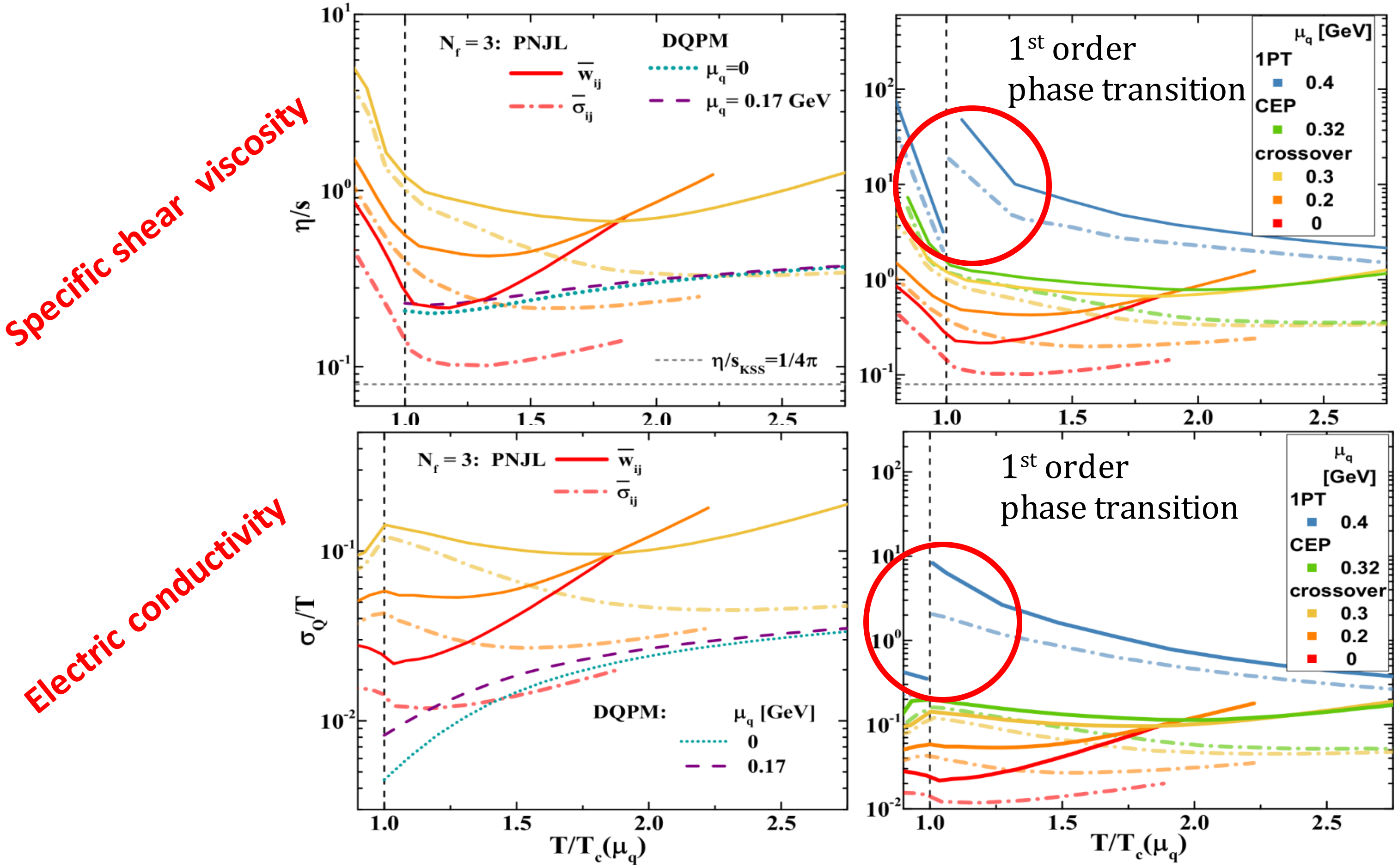


Improved thermodynamics by NNLO in Ω and Polyakov loop

J. M. Torres-Rincon, J. Aichelin PRC 96 (2017) 4 045205

D. Fuseau, T. Steinernert, J. Aichelin PRC 101 (2020) 6 065203

Specific shear viscosity at high μ_B



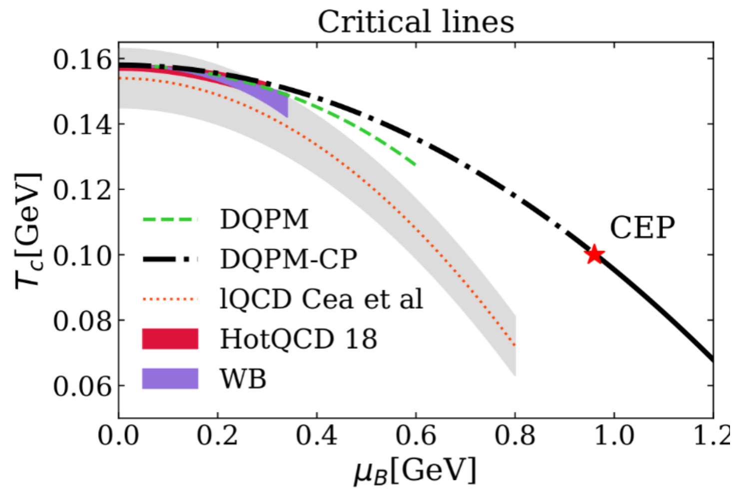
Similar jump at 1st order PT
Also seen in the holographic model: Talk J. Grefa 14.06

- Two different models have similar increase with μ_B –dependence in the crossover region
- Drastic change of T–dependence for all transport coefficients after 1st order phase transition See also Nf=2 NJL results C. Sasaki et al, NPA 832 (2010)

O. S., D. Fuseau, J. Aichelin and E. Bratkovskaya, PRC 103 (2021) no.5, 054901

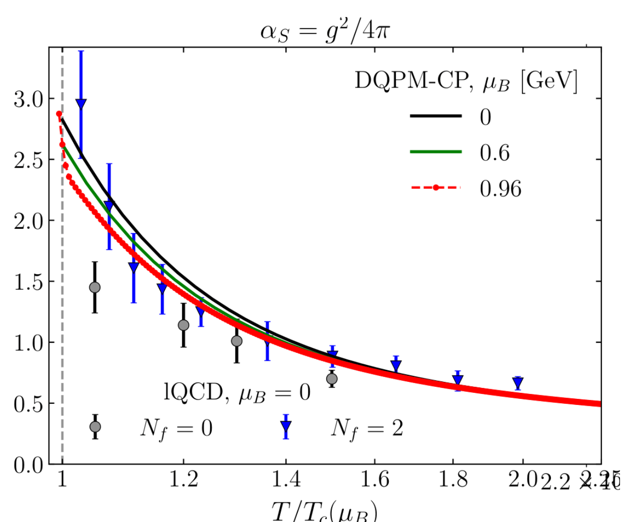
Quasiparticle model with CEP at high μ_B

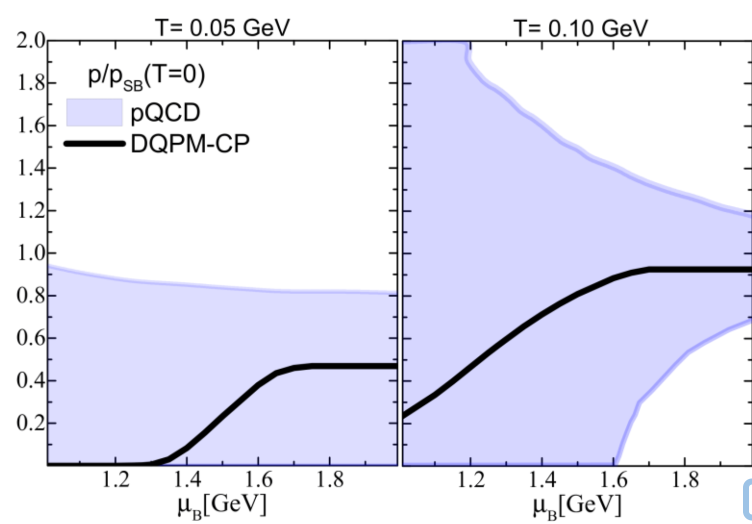
- DQPM-CP for high μ_B, including the CEP region based on the scaling properties of the entropy density from the PNJL model
- DQPM-CP interpolates EoS and microscopic properties between two asymptotics - high $T \gg Tc, \mu_B = 0$ and $T > Tc, \mu_B \gg T$
- EoS and transport coefficients of the QGP phase for the wide range of $T > Tc, \mu_B$



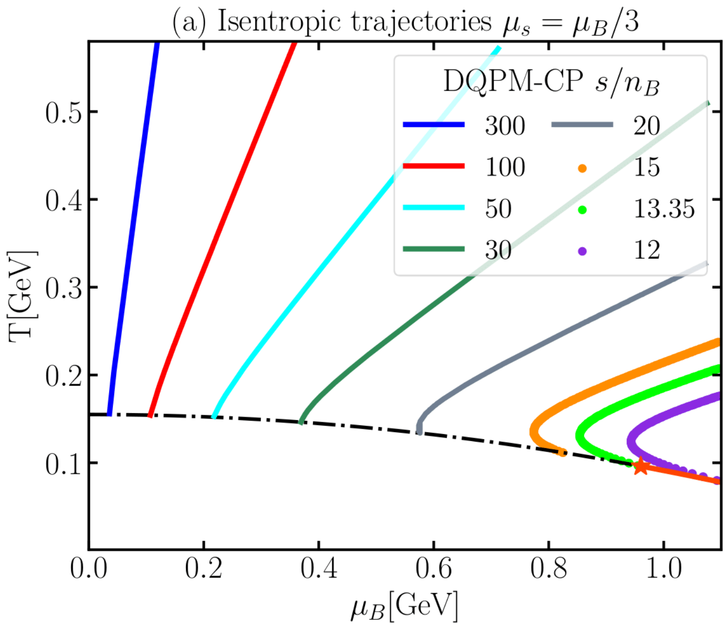
- CEP: $(T, \mu_B) = (100, 960)$ MeV , $\mu_B/T = 9.6$
- EoS : for $\mu_B/T < 2$ agreement with lQCD for $\mu_B/T > 6$ agreement with pQCD

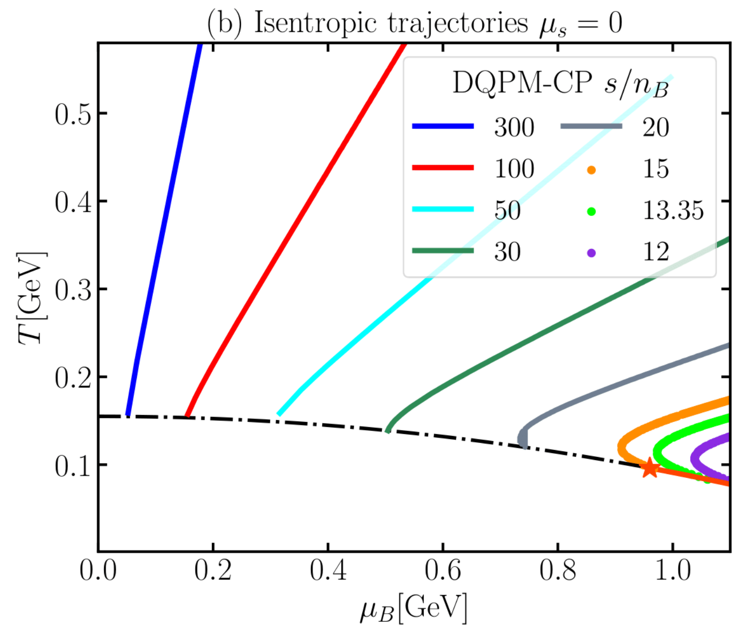
Near CEP: $g^2 = f(s^{PNJL}(T/T_c)) \rightarrow g^2(T/T_c)$





Isentropic trajectories



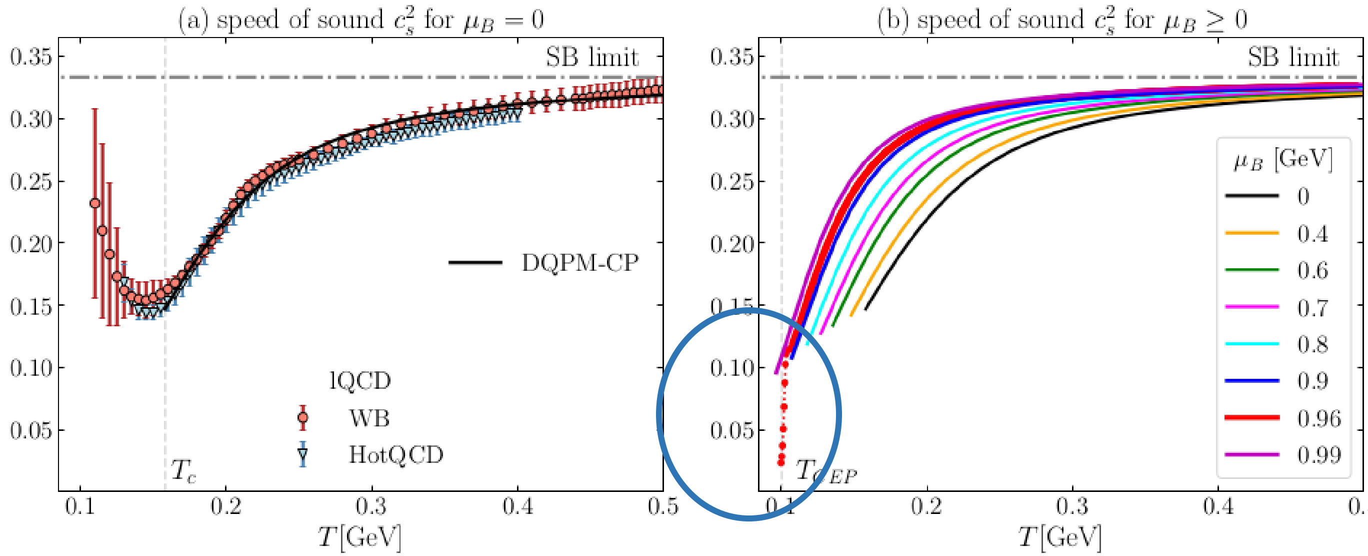


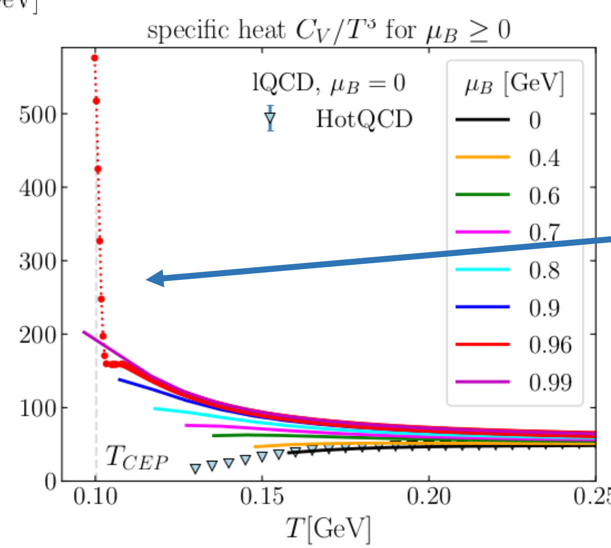
- CEP acts as an attractor of isentropic trajectories (Chiho Nonaka and Masayuki Asakawa PRC 71 (2005), 044904)
- Trajectories of s/n_B =const for $<n_S>$ =0 are shifted towards higher μ_B

O. S., J. Aichelin and E. Bratkovskaya, PRD 105 (2022) 054011

Speed of sound

- EoS : for μ_B/T <2 agreement with lQCD for μ_B/T >6 agreement with pQCD





Near CEP critical scaling can be seen:

$$\ln(C_V) = -\alpha \cdot \ln(T - T_{CEP}) + const$$

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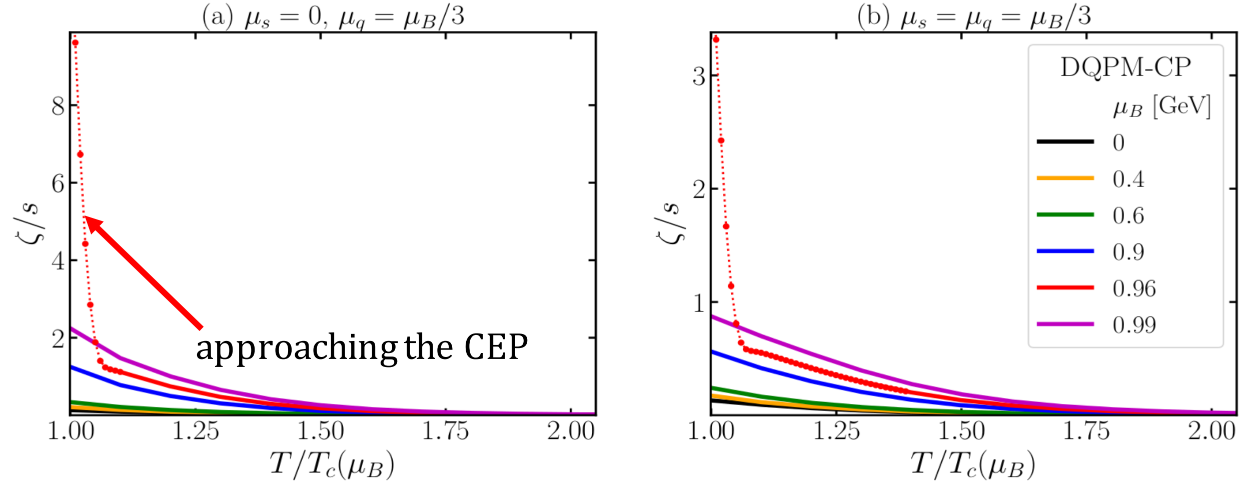
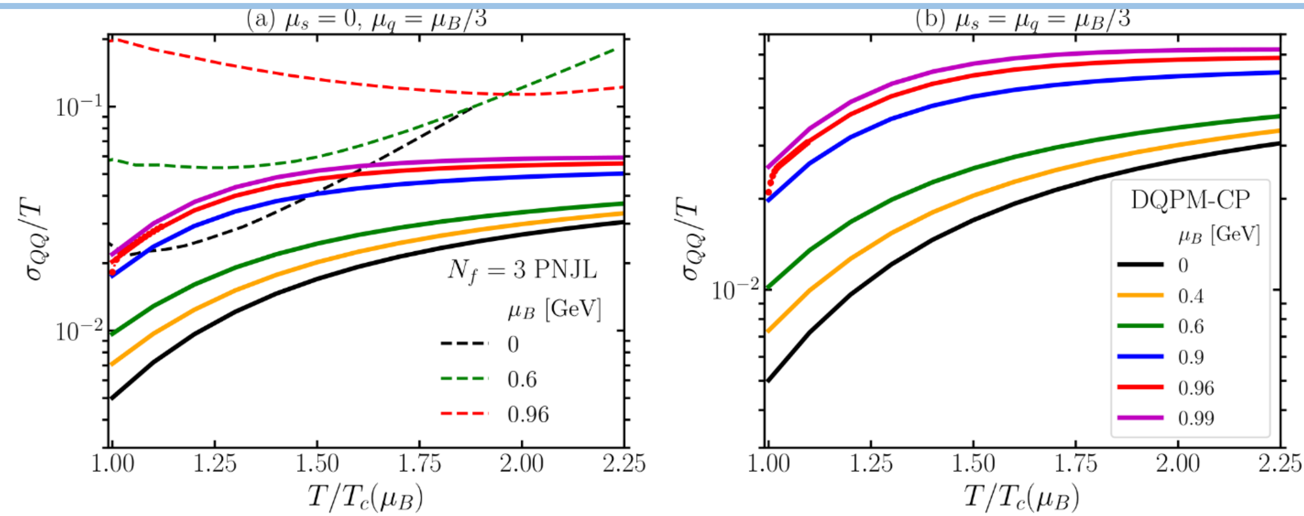
Shear and bulk iscosities near the CEP



- Sudden rise of specific bulk viscosity approaching the CEP



- B,Q,S diffusion coefficients have pronounced $\boldsymbol{\mu_B}, \boldsymbol{\mu_S}$-dependence
- Only small increase approaching the CEP

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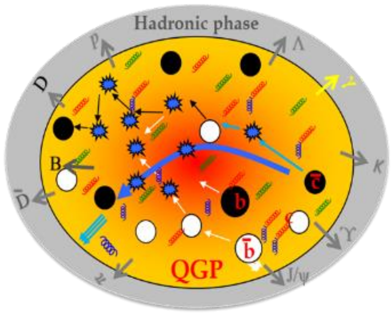
Modelling HICs: PHSD

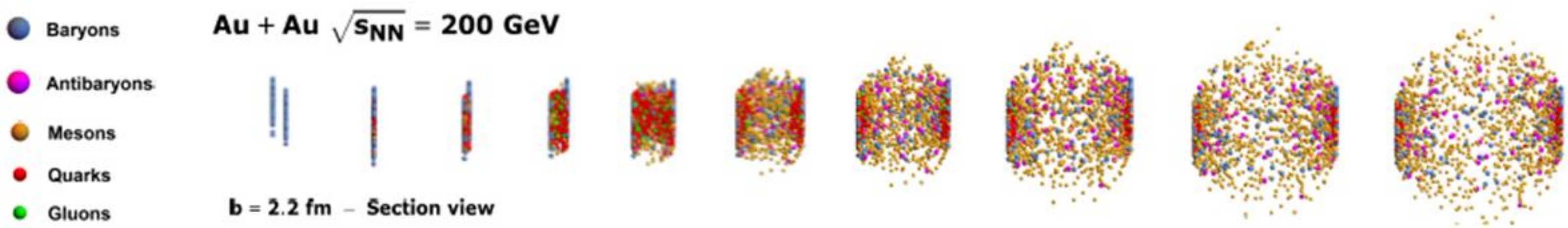


QGP out-of equilibrium ←→ HIC

Parton-Hadron-String-Dynamics (PHSD)

Non-equilibrium microscopic transport approach for the description of strongly-interacting hadronic and partonic matter created in heavy-ion collisions

Dynamics: based on the solution of generalized off-shell transport equations derived from Kadanoff-Baym many-body theory



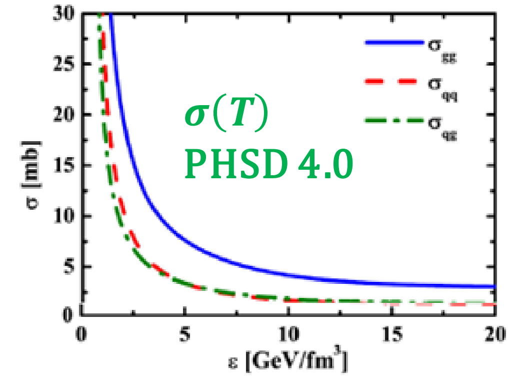
W. Cassing, E. Bratkovskaya, PRC 78 (2008) 034919; NPA831 (2009) 215; W. Cassing, EPJ ST 168 (2009) 3; …….;
P. Moreau, O. Soloveva , L. Oliva, T. Song, W. Cassing, E. Bratkovskaya, PRC 100 (2019) , 014911;
O. Soloveva, P. Moreau, L. Oliva, V. Voronyuk, V. Kireyeu, T. Song, E. Bratkovskaya, Particles 3 (2020), 178-192;….

PHSD

- PHSD 4.0 : only isotropic $\sigma(T)$ and ρ(T)

parton cross sections

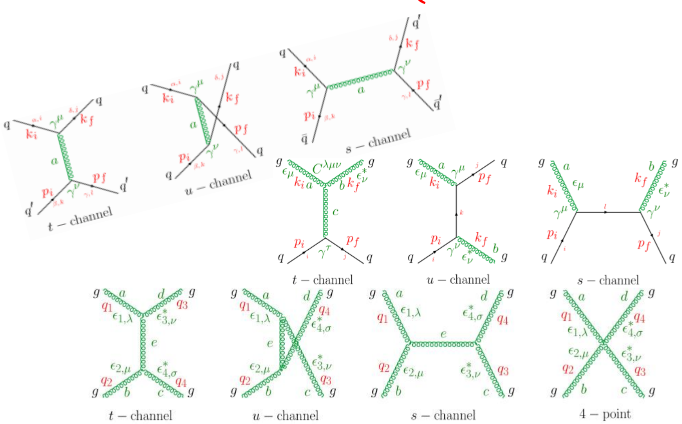
parton spectral function (masses and widths)

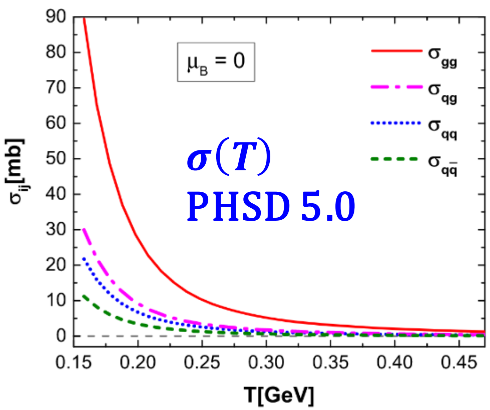


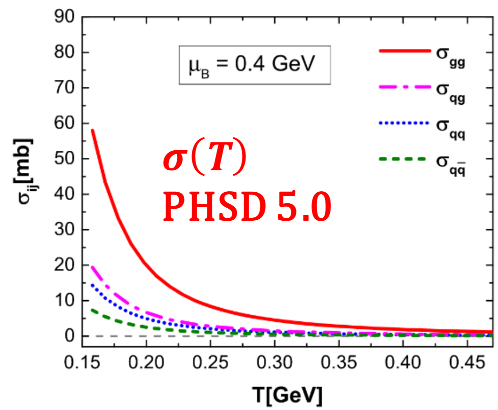
new PHSD 5 : angular dependence of dσ/d cosθ

- PHSD 5.0 : with $\sigma\left(\sqrt{s}, m_1, m_2, T, \mu_B = 0\right)$ and ρ$(T, \mu_B = 0)$

- PHSD 5.0 : with $\sigma\left(\sqrt{s}, m_1, m_2, T, \mu_B\right)$ and ρ(T, μ_B)







P. Moreau, O. Soloveva , L. Oliva, T. Song, W. Cassing, E. Bratkovskaya, PRC 100 (2019) , 014911;
O. Soloveva, P. Moreau, L. Oliva, V. Voronyuk, V. Kireyeu, T. Song, E. Bratkovskaya, Particles 3 (2020), 178-192

PHSD: QGP evolution in HICs

Input:

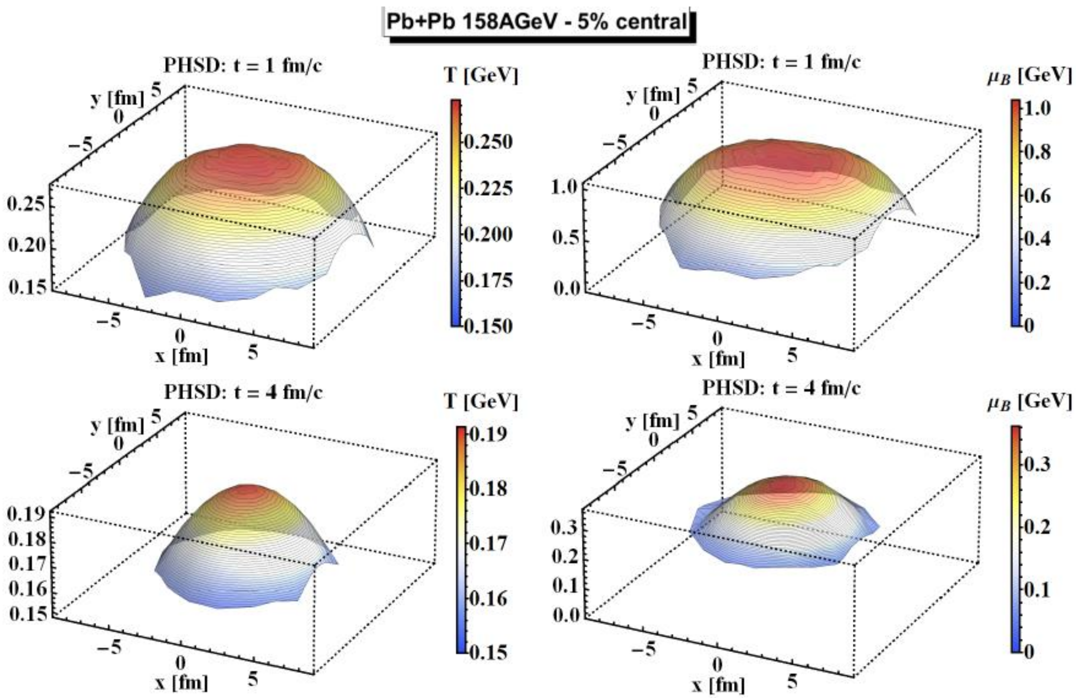
ε^{PHSD} and n_B^{PHSD}

Output:

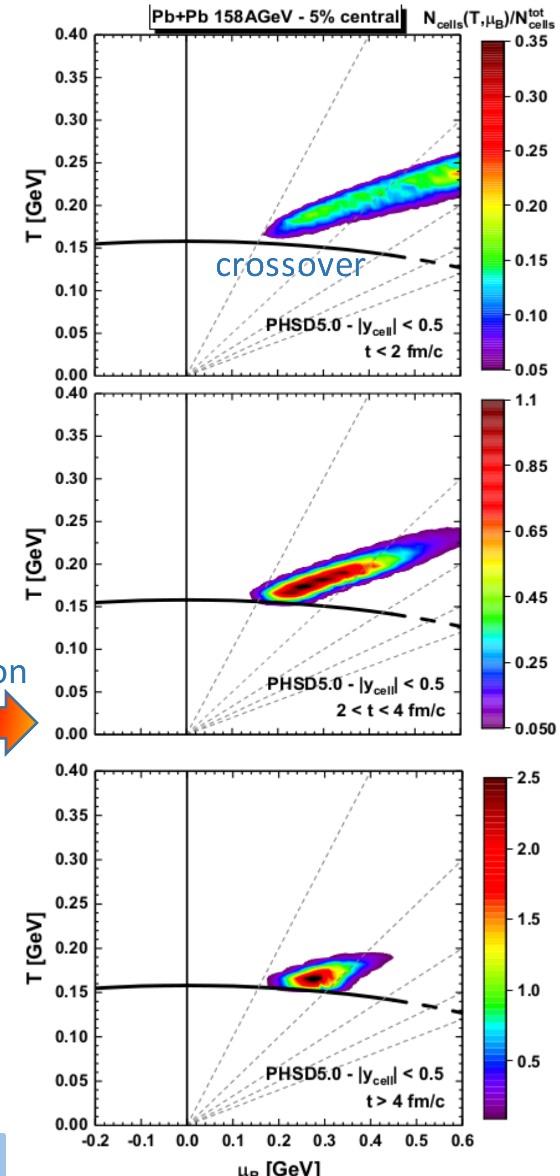
T, μ_B

The T -profile in (x;y) & μ_B profile in (x;y)

at midrapidity ($|y_{cell}|$< 1) at fixed times(1 and 4 fm/c)



time evolution



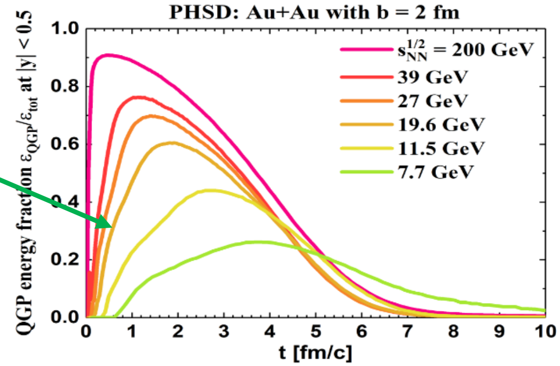
Path through the phase diagram is not trivial and localized

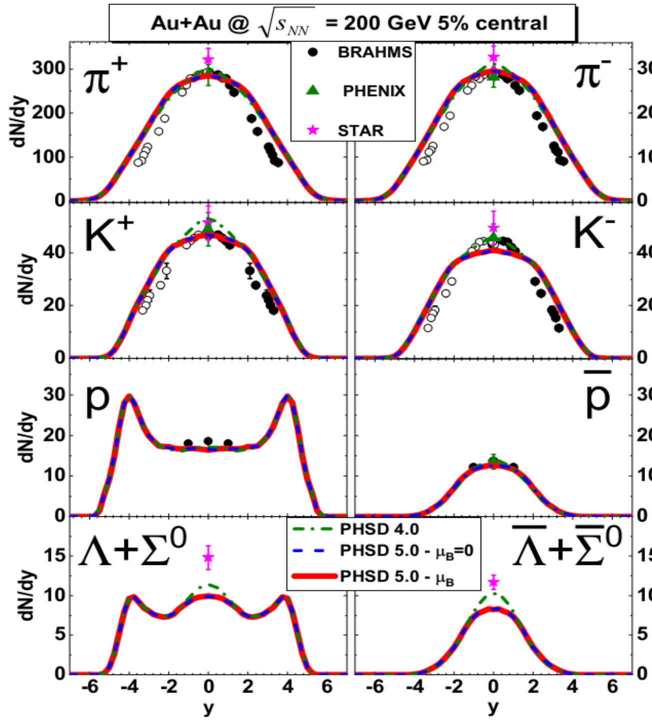
Results for ($\sqrt{s_{NN}}$ = 200 GeV - 7 GeV)

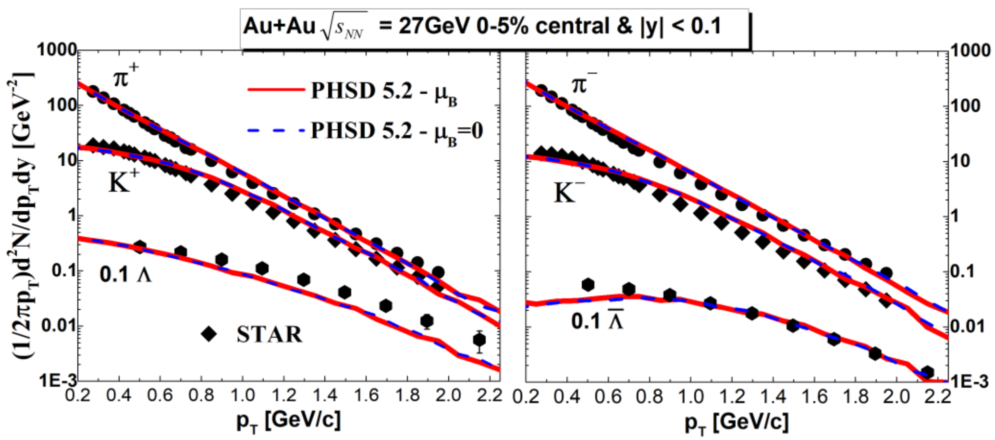
- No visible effects on p_T-spectra, dN/dy of μ_B-dependence
- Small effect of the angular dependence of **dσ/dcosθ**

at high $\sqrt{s_{NN}}$ - low μ_B

! QGP fraction is small at low $\sqrt{s_{NN}}$







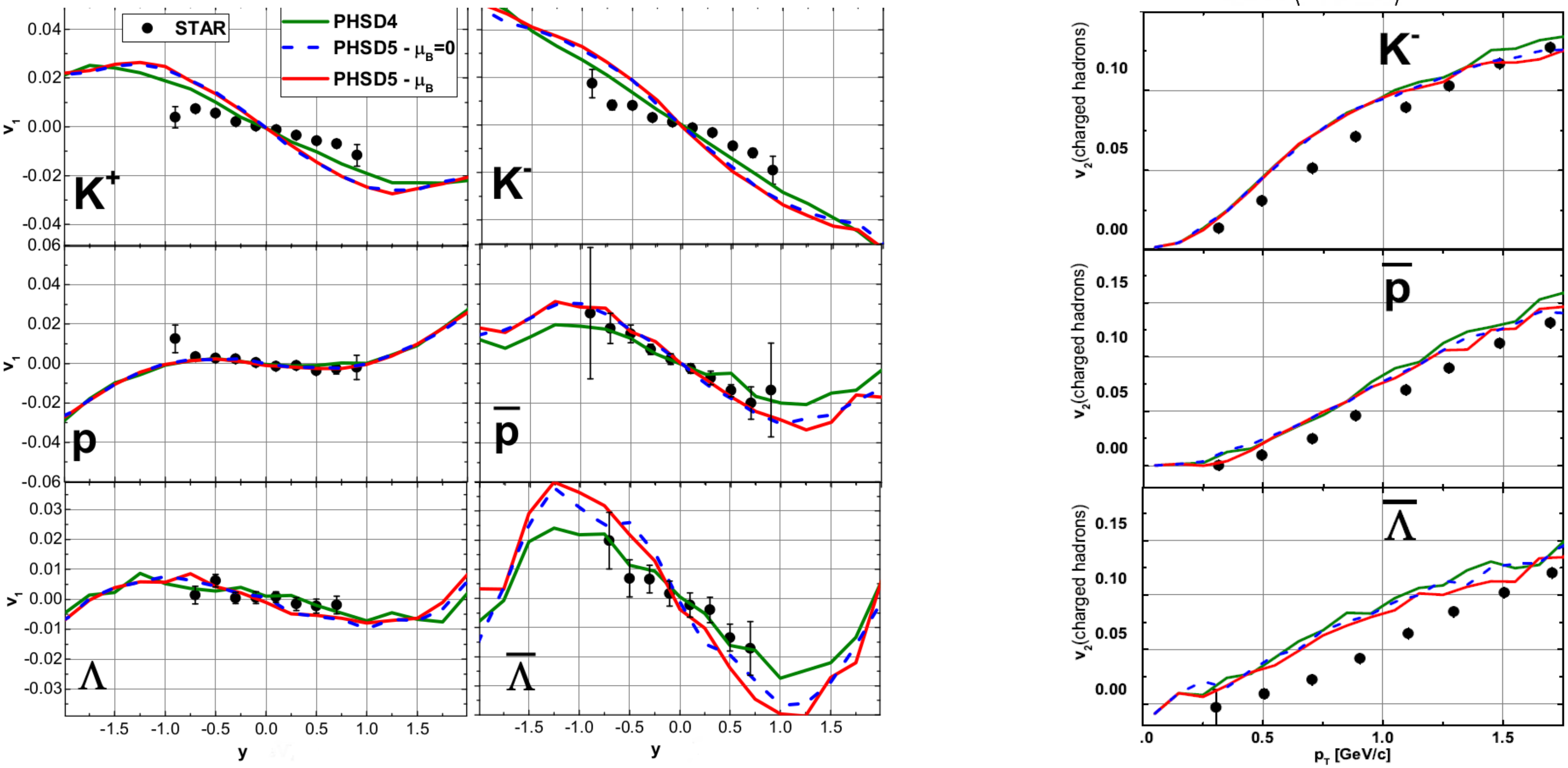
P. Moreau, O. Soloveva , L. Oliva, T. Song, W. Cassing, E. Bratkovskaya, PRC 100 (2019) , 014911;
O. Soloveva, P. Moreau, L. Oliva, V. Voronyuk, V. Kireyeu, T. Song, E. Bratkovskaya, Particles 3 (2020), 178-192

Elliptic flow ($\sqrt{s_{NN}} = 200\ GeV\ - 27GeV$)

- Weak μ_B –dependence – small fraction of QGP or low μ_B
- Small effect of the angular dependence of **dσ/dcosθ**
- Strong flavor dependence

27GeV, 10-40%

$$v_1 = \left\langle \frac{p_x}{p_T} \right\rangle \qquad v_2 = \left\langle \frac{p_x^2 - p_y^2}{p_T^2} \right\rangle$$



O. Soloveva, P. Moreau, L. Oliva, V. Voronyuk, V. Kireyeu, T. Song, E. Bratkovskaya, Particles 3 (2020), 178-192

Summary

Transport properties of the strongly-interacting QGP matter at finite **T** and $\boldsymbol{\mu_B}$ have been investigated.
Influence of an order of a phase transition on thermodynamic and transport properties has been studied.

- Transport coefficients can differ among the models, which have similar phase structures and EoS

Evolution of the QGP matter created in HICs and the sensitivity of the bulk and flow observables on the QGP interactions and transport properties have been explored by the simulations within the PHSD transport approach

- High-$\boldsymbol{\mu_B}$ regions are probed at low $\sqrt{\boldsymbol{s_{NN}}}$ or high rapidity regions
  Moreover, QGP fraction is small at low $\sqrt{\boldsymbol{s_{NN}}}$: small effect seen in observables
- $\boldsymbol{\mu_B}$-dependence of QGP interactions is more pronounced in observables for strange hadrons and antiprotons

Summary

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Thank you for your attention!

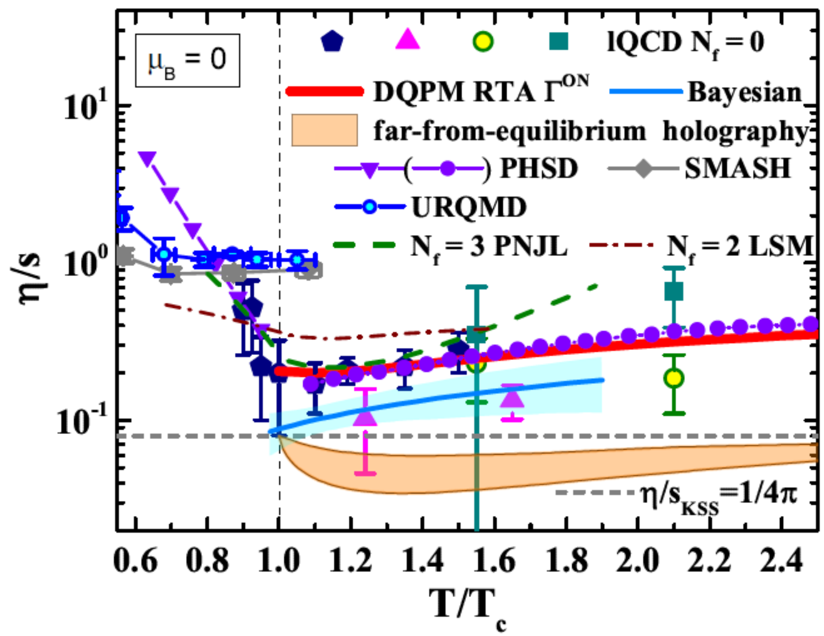
SQM2022

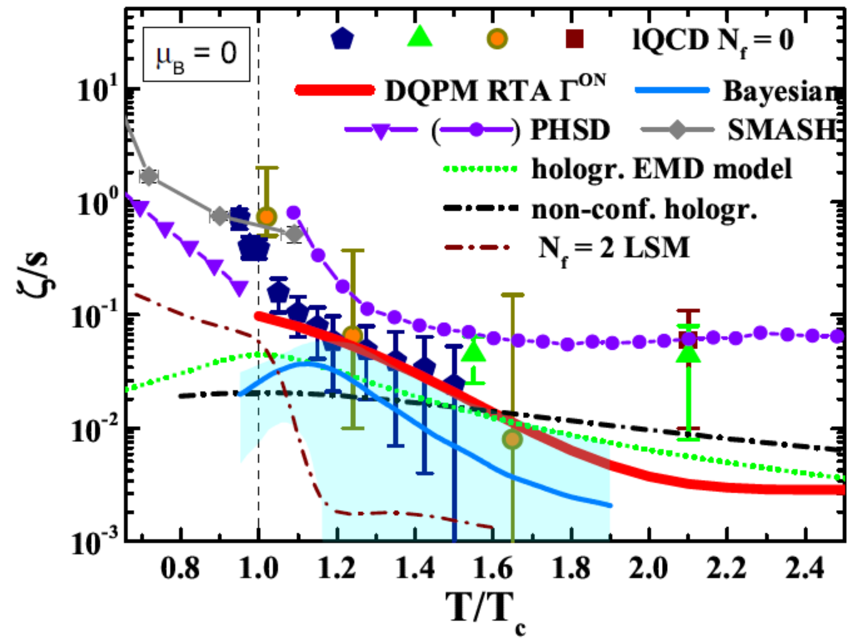
The 20th International Conference on Strangeness in Quark Matter
13-17 June 2022 Busan, Republic of Korea

Back up slides

Specific shear viscosity compilation

Model predictions:





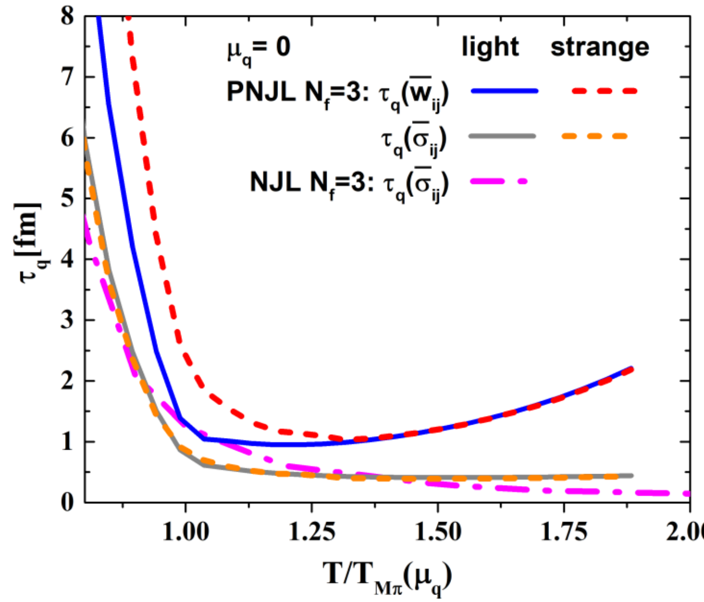
! Different models using the same EoS can have completely different transport coefficients!

PNJL relaxation times

$$\boxed{\tau_i(\mathbf{p},T,\mu_B) = \frac{1}{\Gamma_i(\mathbf{p},T,\mu_B)}}$$

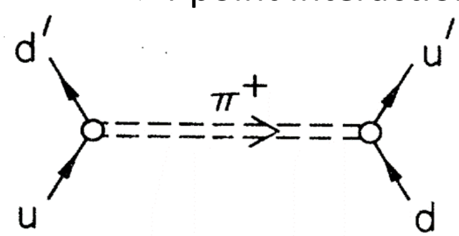
$$\Gamma_i^{\rm on}(\mathbf{p}_i,T,\mu_q) = \frac{1}{2E_i}\sum_{j=q,\bar{q},g}\int\frac{d^3p_j}{(2\pi)^3 2E_j}\, d_j\; f_j(E_j,T,\mu_q)$$
$$\int\frac{d^3p_3}{(2\pi)^3 2E_3}\int\frac{d^3p_4}{(2\pi)^3 2E_4}(1\pm f_3)(1\pm f_4)$$
$$\boxed{|\bar{\mathcal{M}}|^2(p_i,p_j,p_3,p_4)}\,(2\pi)^4\delta^{(4)}\,(p_i+p_j-p_3-p_4)$$

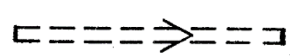
Relaxation times(PNJL vs NJL)



qq - interactions:

4 point interaction -> meson exchange(π,Ϭ,η,η',K,.. for s,t,u channels)

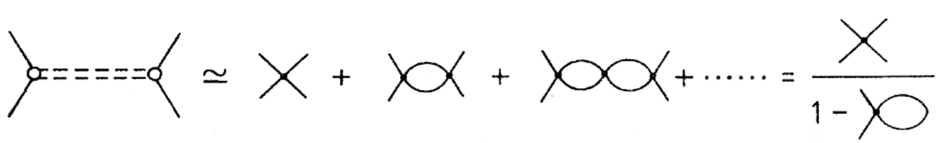




$$=(i\gamma_5)\tau^{(-)}\frac{-ig^2_{\pi qq}}{k^2-m^2_\pi}(i\gamma_5)\tau^{(+)}$$

meson propagator $\mathscr{D} = \dfrac{2ig_m}{1-2g_m\Pi^{\pm}_{ff'}(k_0,\vec{k})}$

Effective interaction in RPA



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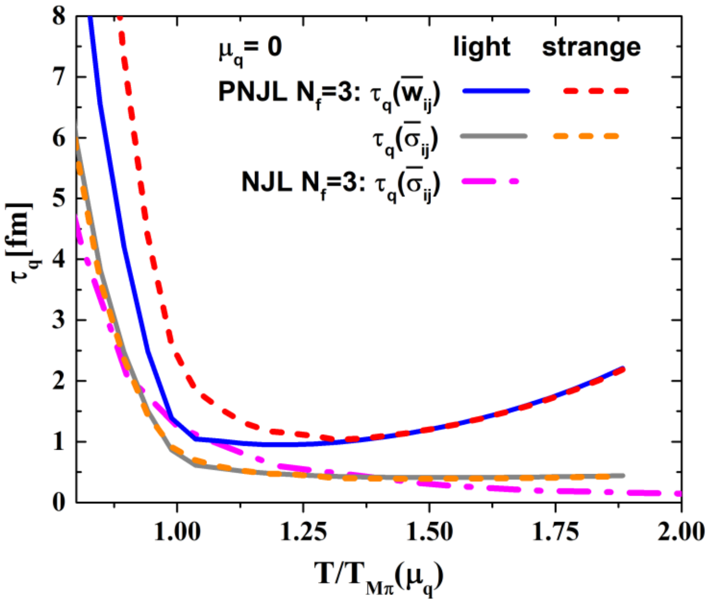
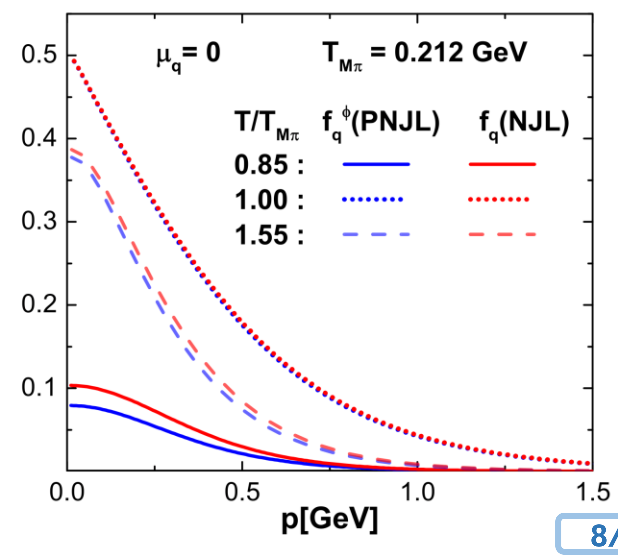


Modified distribution functions:
Polyakov loop contributions

$$f_q \to f_q^{\Phi}(\mathbf{p},T,\mu) = \frac{(\bar{\Phi}+2\Phi e^{-(E_\mathbf{p}-\mu)/T})e^{-(E_\mathbf{p}-\mu)/T}+e^{-3(E_\mathbf{p}-\mu)/T}}{1+3(\bar{\Phi}+\Phi e^{-(E_\mathbf{p}-\mu)/T})e^{-(E_\mathbf{p}-\mu)/T}+e^{-3(E_\mathbf{p}-\mu)/T}},$$



QGP in the Polyakov extended NJL model

- PNJL model based on effective Lagrangian with the same symmetries for the quark dof as QCD

$$
\begin{aligned}
\mathscr{L}_{PNJL} = & \sum_i \bar{\psi}_i (i\not{D} - m_{0i} + \mu_i \gamma_0)\psi_i \\
& + G \sum_a \sum_{ijkl} \left[(\bar{\psi}_i \; i\gamma_5 \tau^a_{ij} \psi_j)\; (\bar{\psi}_k \; i\gamma_5 \tau^a_{kl} \psi_l) + (\bar{\psi}_i \tau^a_{ij} \psi_j)\; (\bar{\psi}_k \tau^a_{kl} \psi_l) \right] \\
& - K \det_{ij} \left[\bar{\psi}_i \; (-\gamma_5)\psi_j \right] - K \det_{ij} \left[\bar{\psi}_i \; (+\gamma_5)\psi_j \right] \\
& - \mathcal{U}(T; \Phi, \bar{\Phi})
\end{aligned}
$$

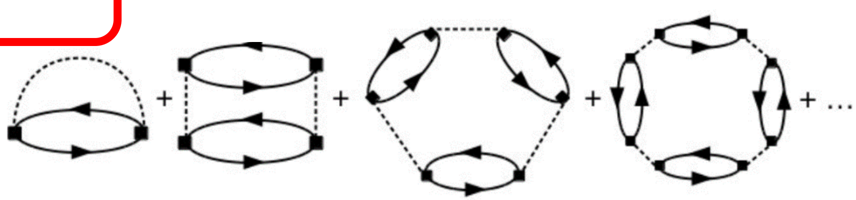
Polyakov-loop effective potential fitted to the YM

5 parameters fixed by vacuum values K,π masses, η-η'mass splitting , π decay constant, Chiral condensate

Improvements:

- Next to leading order in Nc($O(1/Nc)^{0}$) of the grand-canonical potential : presence of the mesons below Tc

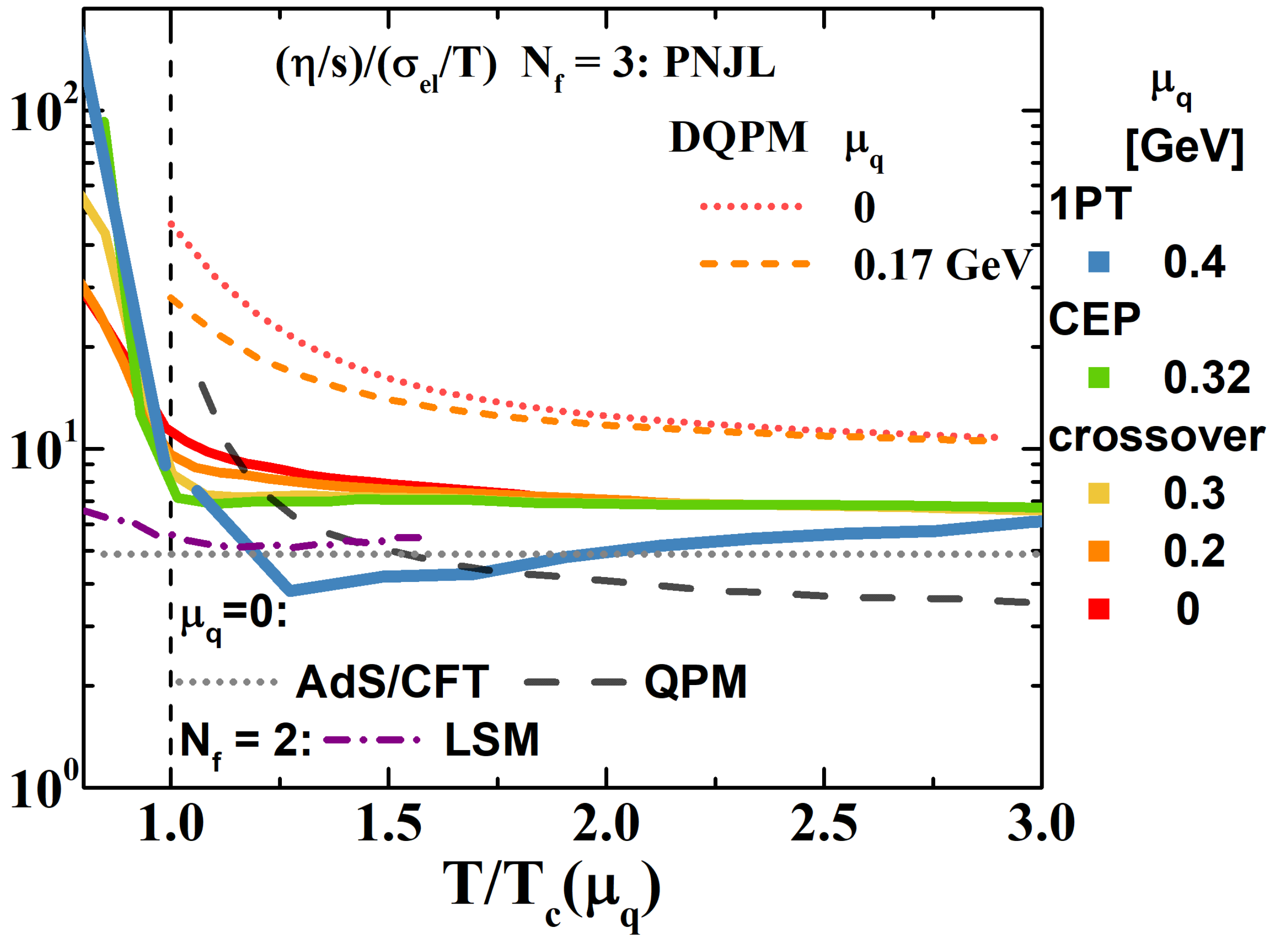
$$
\Omega_{\rm PNJL}(T,\mu_i) = \Omega_q^{(-1)}(T,\mu_i) + \boxed{\sum_{M \in J^\pi = \{0^+, 0^-\}} \Omega_M^{(0)}(T, \mu_M(\mu_i))} + \mathcal{U}_{glue}(T) \;,
$$



J. M. Torres-Rincon, J. Aichelin PRC 96 (2017) 4 045205
D. Fuseau, T. Steinernert, J. Aichelin PRC 101 (2020) 6 065203

- Modification of the gluon potential due to the presence of the quark

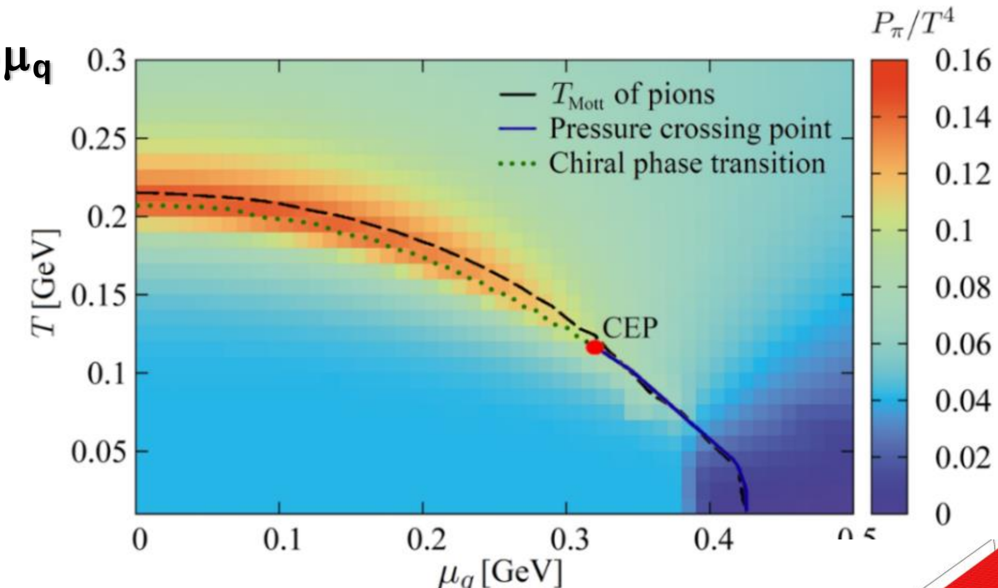
Specific shear viscosity to conductivity



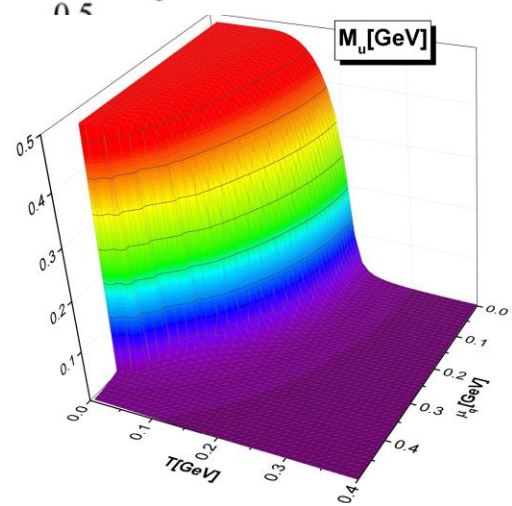
QGP in the Polyakov extended NJL model

- PNJL allows for prediction of macroscopic properties of QGP at finite **T** and μ_B
- & QGP transport coefficients for $0 \leq \boldsymbol{\mu_B} \leq 1.2$ GeV

$\mu_B = 3\mu_q$



D. Fuseau, T. Steinert, J. Aichelin PRC 101 (2020) 6 065203



- CEP: (T , $\boldsymbol{\mu_B}$) = (110,960) MeV , $\boldsymbol{\mu_B}$/T = 8.73
- 1st order PT at high $\boldsymbol{\mu_B}$
- same symmetries for the quarks as QCD

Chiral masses ($\boldsymbol{M_l}, \boldsymbol{M_s}$)

$$m_i = m_{0i} - 4G\langle\langle\bar{\psi}_i\psi_i\rangle\rangle + 2K\langle\langle\bar{\psi}_j\psi_j\rangle\rangle\langle\langle\bar{\psi}_k\psi_k\rangle\rangle$$

Specific shear viscosity compilation

- **Kubo formalism: transport coefficients are expressed through correlation functions of stress-energy tensor**

used in lattice QCD, transport approaches(hadrons), effective models

$$\eta = \frac{1}{20}\lim_{\omega\to 0}\frac{1}{\omega}\int d^4x\, e^{i\omega t}\left\langle\left[\mathcal{S}^{ij}(t,\mathbf{x}),\, \mathcal{S}^{ij}(0,\mathbf{0})\right]\right\rangle\theta(t) \qquad \mathcal{S}^{ij} = T^{ij} - \delta^{ij}\mathcal{P}$$

$$\zeta = \frac{1}{2}\lim_{\omega\to 0}\frac{1}{\omega}\int d^4x\, e^{i\omega t}\langle[\mathcal{P}(t,\mathbf{x}),\, \mathcal{P}(0,\mathbf{0})]\rangle\theta(t) \qquad \mathcal{P} = -\tfrac{1}{3}T^{i}{}_{i}$$

R. Lang and W. Weise, EPJ. A 50, 63 (2014) (NJL model)
A. Harutyunyan et al,PRD 95, 114021, (2017)

Kinetic theory:

- **Relaxation time approximation(RTA) : consider relaxation time** $\dfrac{df_a^{\mathrm{eq}}}{dt} = \mathcal{C}_a = -\dfrac{f_a^{\mathrm{eq}}\phi_a}{\tau_a}$

P. Chakraborty and J. I. Kapusta, PRC 83,014906 (2011)

- **Chapman-Enskog : expand the distribution in terms of the Knudsen number**

J. A. Fotakis et al, PRD 101 (2020) 7, 076007 (HRG)

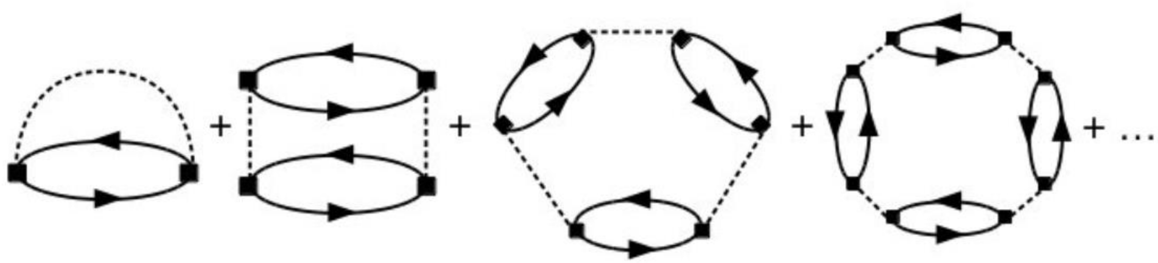
And more!

Holographic models: AdS/CFT correspondence

D. T. Son and A. O. Starinets, JHEP 0603, 052 (2006)
M. Attems et al , JHEP 10 (2016), 155.

PNJL improvements

- **Next to leading order in Nc ($O(1/Nc)^0$) of the grand-canonical potential :** **presence of the mesons below Tc**



J. M. Torres-Rincon, J. Aichelin PRC 96 (2017) 4 045205

- **Modification of the gluon potential due to the presence of the quark**

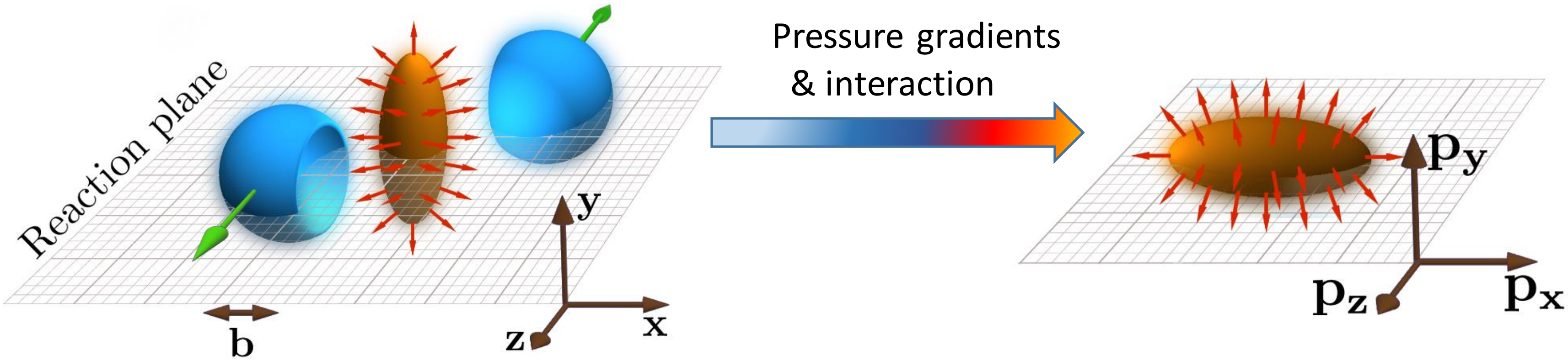
$$\frac{U(\phi,\bar\phi,T)}{T^4} = -\frac{b_2(T)}{2}\bar\phi\phi - \frac{b_3}{6}(\bar\phi^3+\phi^3) + \frac{b_4}{4}(\bar\phi\phi)^2$$

$$b_2(T) = a_0 + \frac{a_1}{1+\tau} + \frac{a_2}{(1+\tau)^2} + \frac{a_3}{(1+\tau)^3}$$ where $$\tau_{\text{phen}} = 0.57\frac{T - T^{\text{cr}}_{\text{phen}}(T)}{T^{\text{cr}}_{\text{phen}}(T)}$$

$$T_{phen}(T) = a + bT + cT^2 + dT^3 + \boxed{e\frac{1}{T}}$$

D. Fuseau, T. Steinernert, J. Aichelin PRC 101 (2020) 6 065203

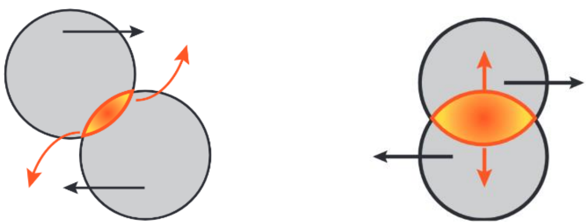
Anisotropic flow coefficients



Quantify the anisotropic flow using Fourier expansion

$$\frac{dN}{d\varphi} \propto \left(1 + 2\sum_{n=1}^{+\infty} v_n \cos[n(\varphi - \psi_n)]\right)$$

$$v_{\mathrm{n}} = \left\langle \cos n\left(\varphi - \psi_n\right)\right\rangle, \quad n = 1,2,3..,$$



$$v_1 = \left\langle \frac{p_x}{p_T} \right\rangle \qquad v_2 = \left\langle \frac{p_x^2 - p_y^2}{p_T^2} \right\rangle$$

Anisotropic flow

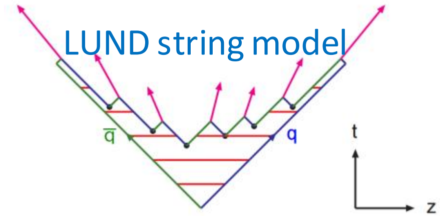
- Assess the transport properties of the QGP
- Sensitive to the QGP EoS and initial state
- Validate models of bulk evolution that are used in the computation of other observables

Stages of collisions in PHSD

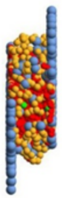
Initial A+A collision

- **String formation** in primary NN collisions

  → **decays** to pre-hadrons (baryons and mesons)



Partonic phase



- Formation of a **QGP** state if $\varepsilon > \varepsilon_{critical}$:
  Dissolution of pre-hadrons → DQPM

  → massive quarks/gluons and mean-field energy

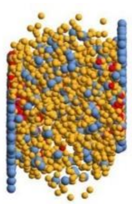
  (quasi-)elastic collisions:

  $$q+q \to q+q \quad g+q \to g+q$$
  $$q+\bar{q} \to q+\bar{q} \quad g+\bar{q} \to g+\bar{q}$$
  $$\bar{q}+\bar{q} \to \bar{q}+\bar{q} \quad g+g \to g+g$$

  inelastic collisions:

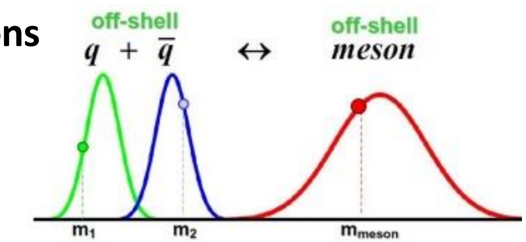
  $$q+\bar{q} \to g$$
  $$g \to q+\bar{q}$$

Hadronization



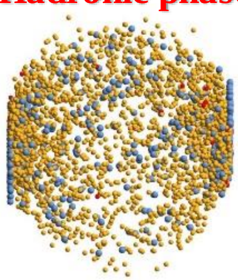
- **Hadronization** to colorless off-shell mesons and baryons

  $$g \to q+\bar{q}, \qquad q+\bar{q} \leftrightarrow meson\ ('\ string\ ')$$
  $$q+q+q \leftrightarrow baryon\ ('\ string\ ')$$



Strict 4-momentum and quantum number conservation

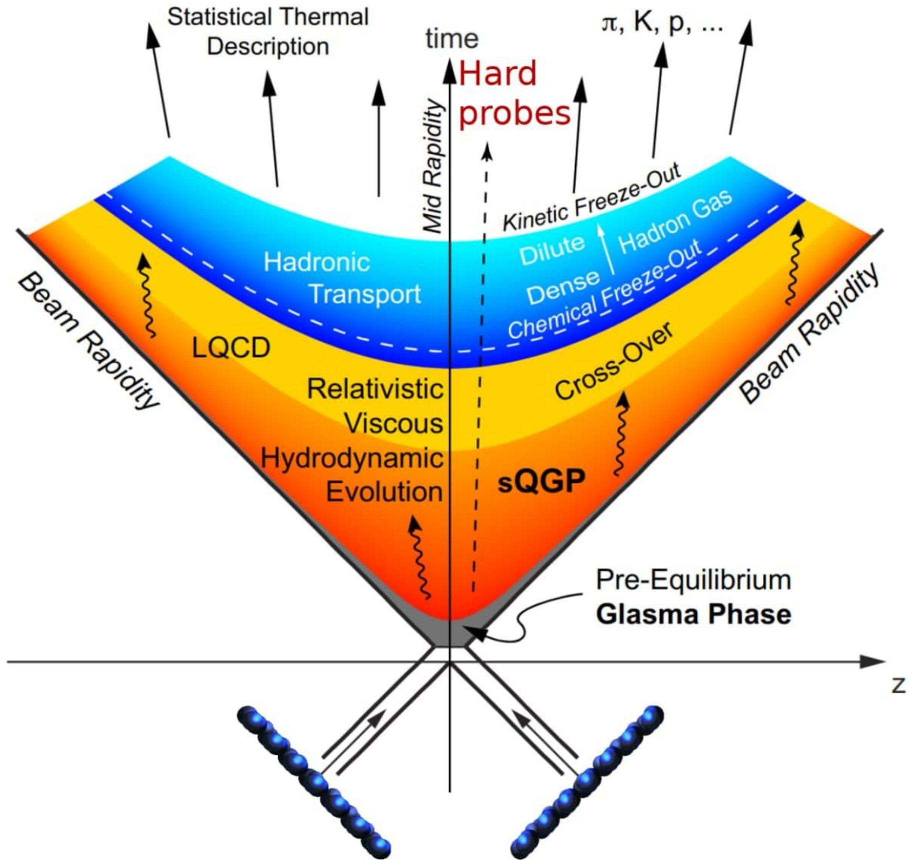
Hadronic phase



- **Hadron-string interactions** – off-shell HSD

W. Cassing, E. Bratkovskaya, PRC 78 (2008) 034919; NPA831 (2009) 215;
W. Cassing, EPJ ST 168 (2009) 3

Stages of HIC



t > 10 fm - hadronisation and free streem to detectors

10 fm > t > 1 fm - QGP expansion

t ≈ 1 fm - Equilibration

t << 1 fm - Initial state

QGP out-of equilibrium ←→ HIC

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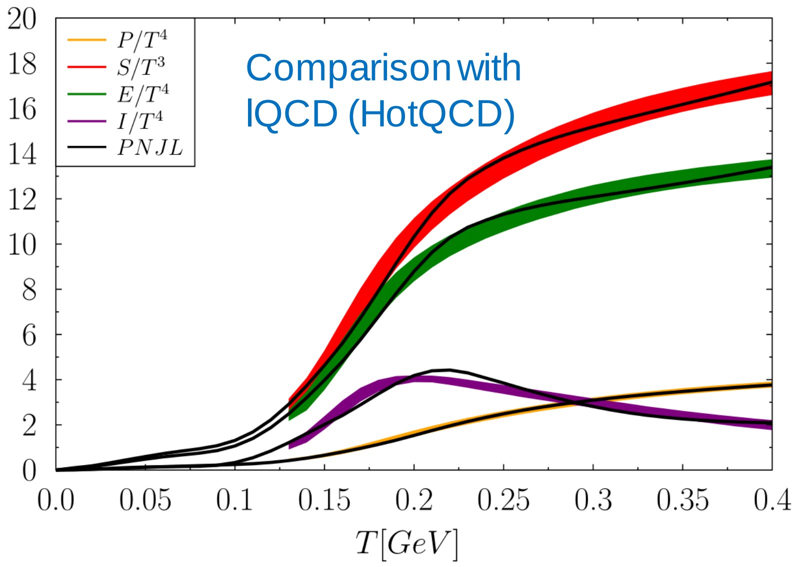
- PNJL allows for prediction of macroscopic properties of QGP at finite **T** and μ_B
- & QGP transport coefficients for $0 \leq \boldsymbol{\mu_B} \leq 1.2$ GeV

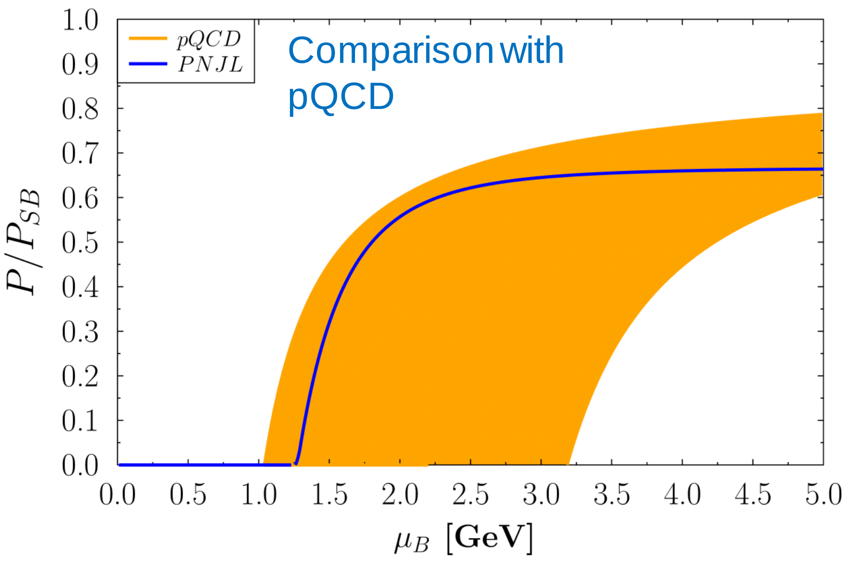
➢ Parameters fixed, EoS at $\mu_B = 0$:

HotQCD PRD 90 (2014) 094503



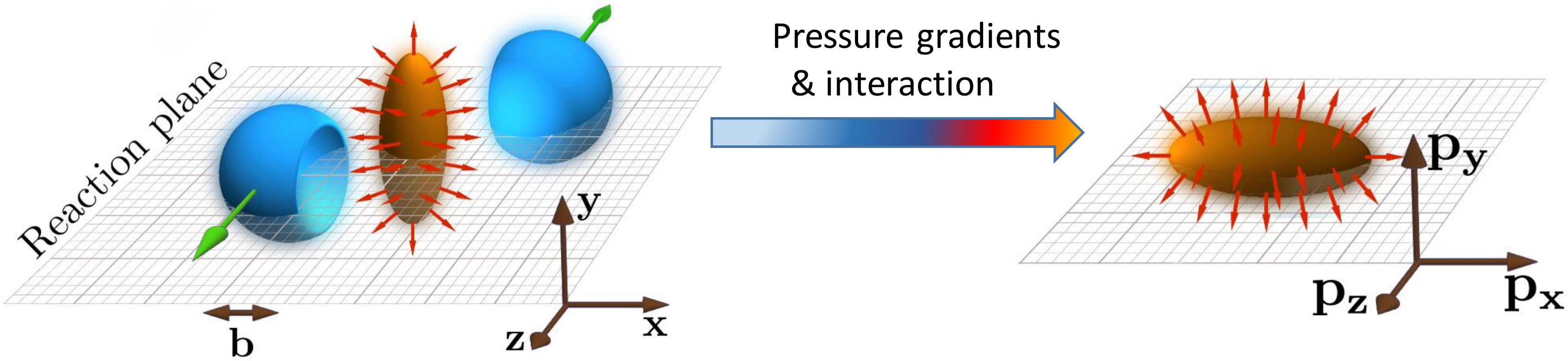
➢ EoS at high μ_B :

pQCD: A.Kurkela, A.Vuorinen, PRL 117 (2016)4 042501



D. Fuseau, T. Steinert, J. Aichelin PRC 101 (2020) 6 065203

Anisotropic flow coefficients



$\boxed{v_n = \langle \cos(n(\phi - \Psi_r)) \rangle}$

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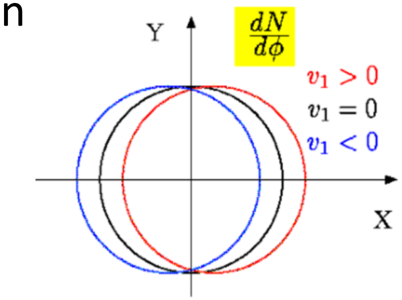
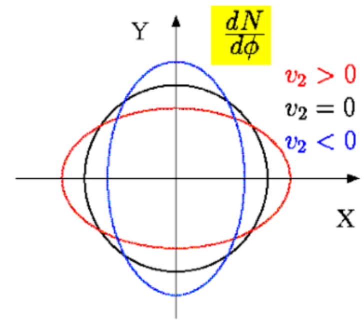


$$v_1 = \left\langle \frac{p_x}{p_T} \right\rangle$$



$$v_2 = \left\langle \frac{p_x^2 - p_y^2}{p_T^2} \right\rangle$$

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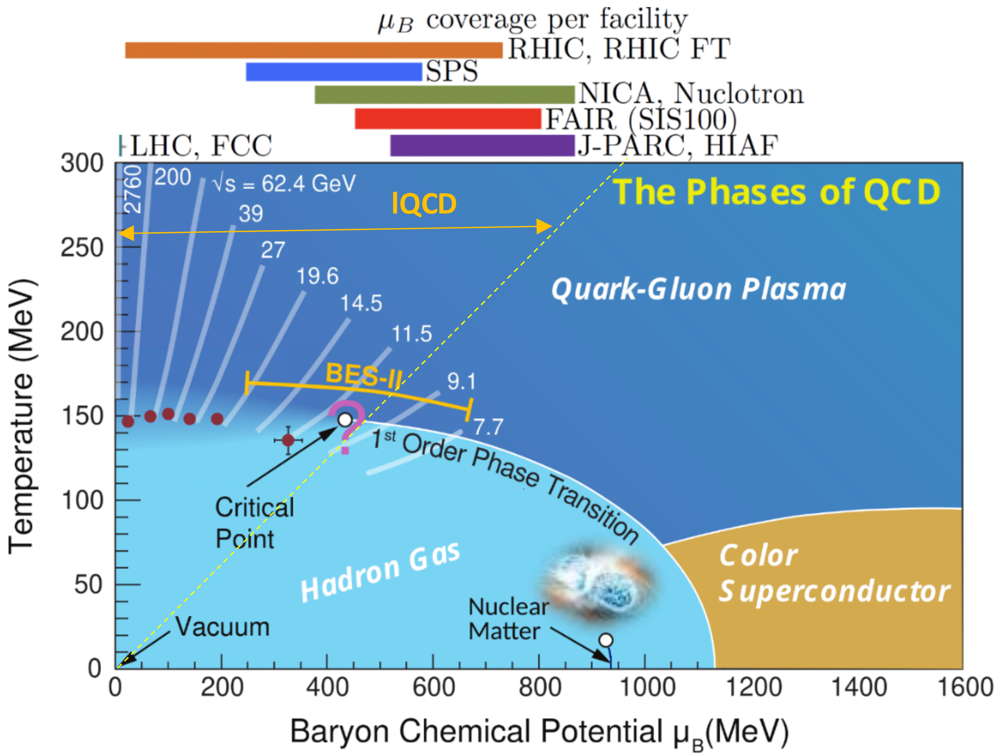
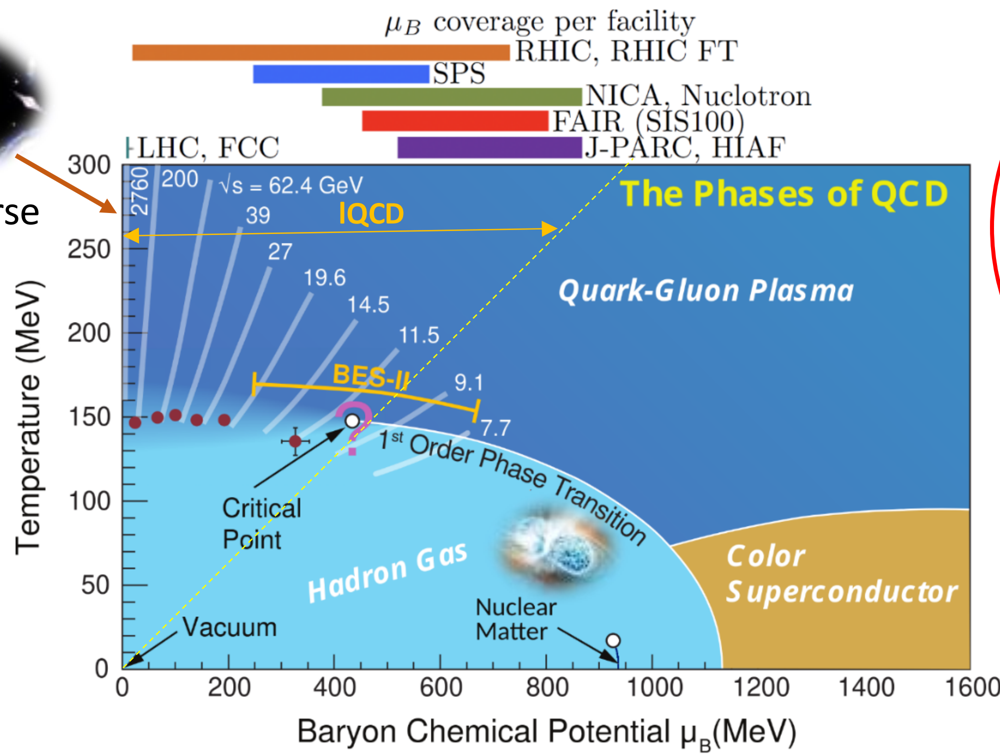
Motivation: the QCD phase diagram

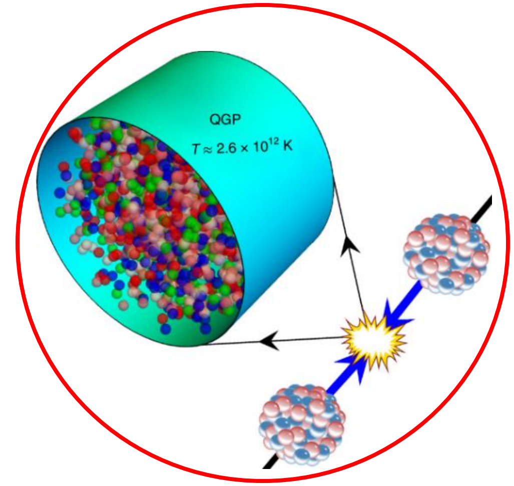


- Explore the QCD phase diagram at finite T and μ_B through heavy-ion collisions
- Search for a possible Critical End Point (CEP) and 1st order **phase transition**
- Quantify macroscopic properties of the QCD matter at finite T and μ_B and relate them to its microscopic structures

Motivation: the QCD phase diagram

Early Universe





heavy-ion collisions

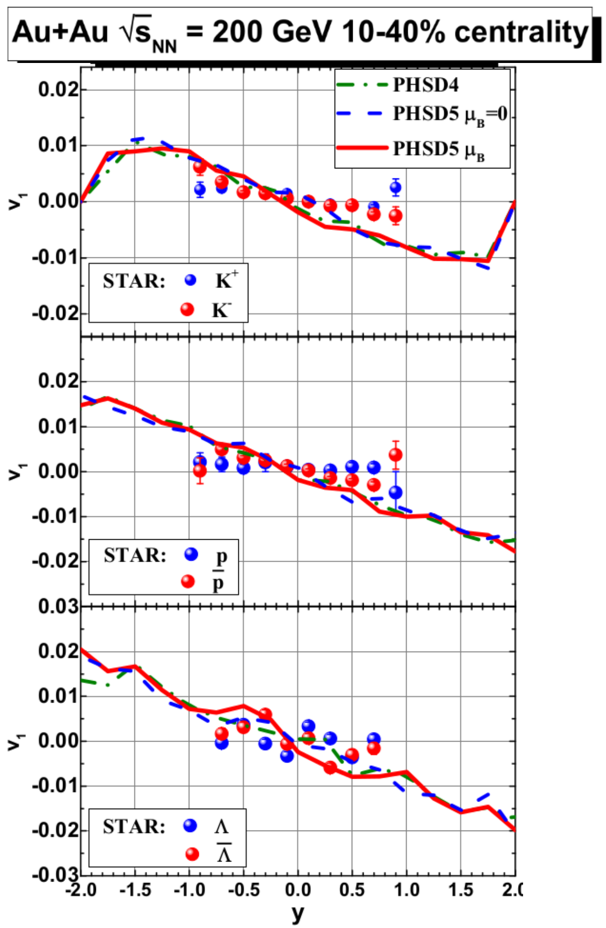
- Explore the QCD phase diagram at finite T and μ_B through heavy-ion collisions
- Search for a possible Critical End Point (CEP) and 1st order **phase transition**
- Quantify macroscopic properties of the QCD matter at finite T and μ_B and relate them to its microscopic structures

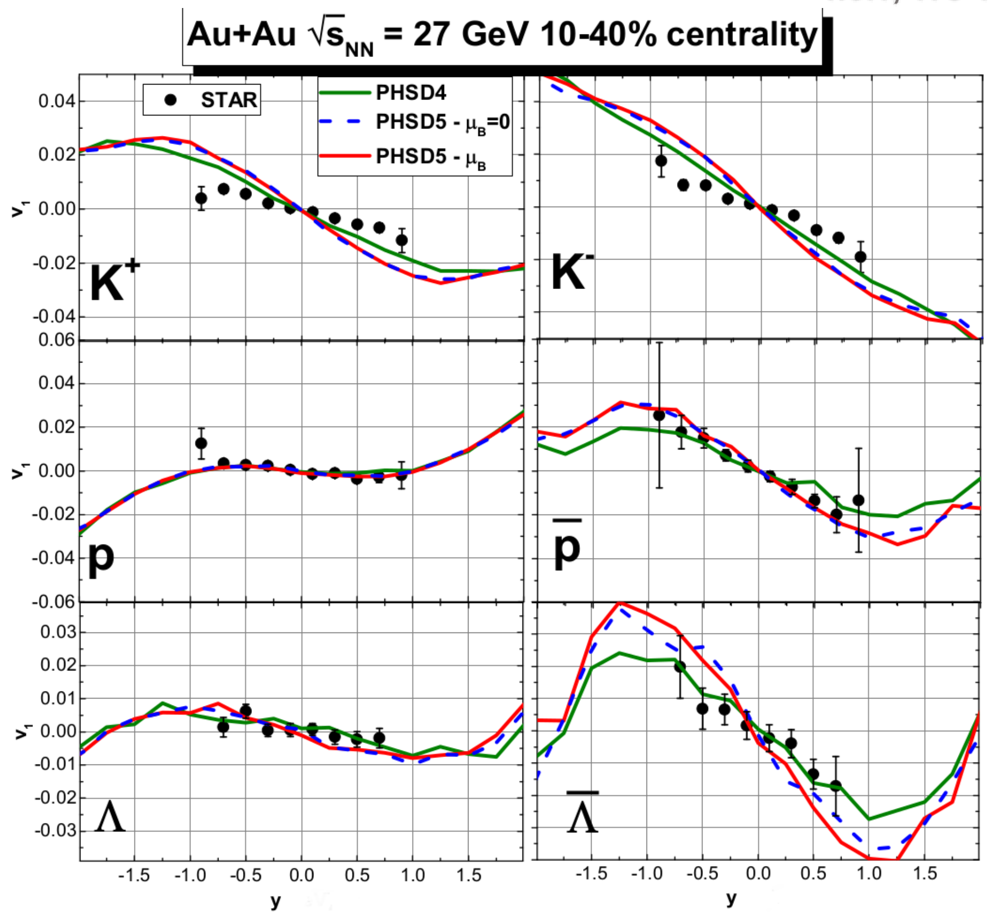
Directed flow ($\sqrt{s_{NN}} = 200\ GeV\ vs\ 27\ GeV$)

$$v_1 = \left\langle \frac{p_x}{p_T} \right\rangle$$

- Influence of the QGP dynamics on final particles observables
- Weak μ_B –dependence – small fraction of QGP or low μ_B

Particles 3 (2020) no.1, 178-192



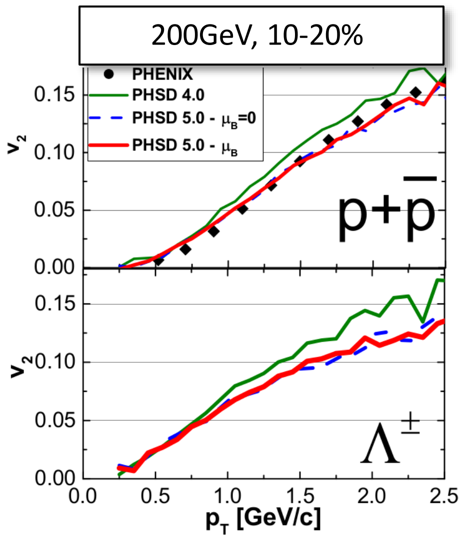


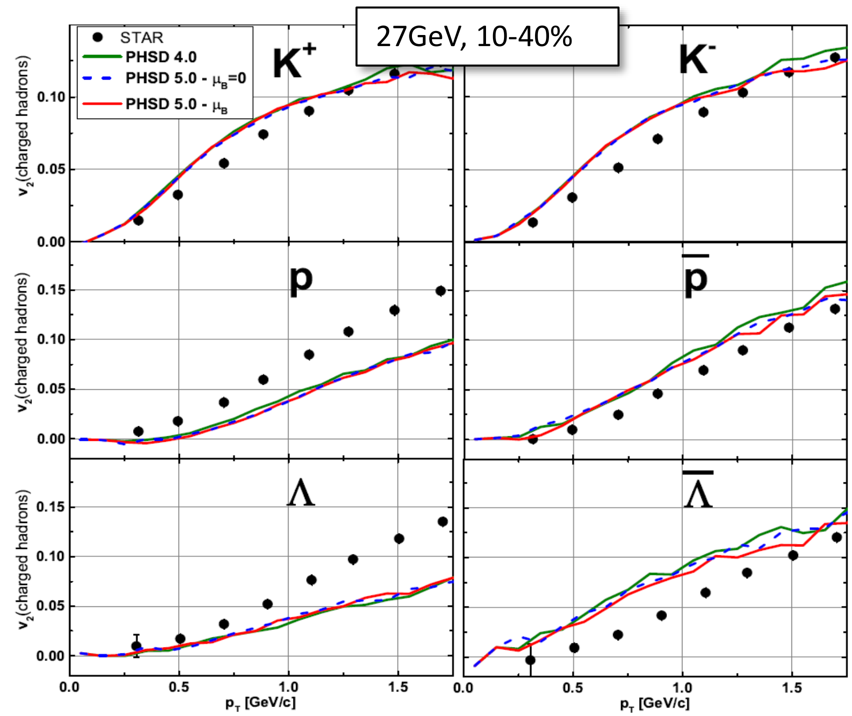
Elliptic flow ($\sqrt{s_{NN}} = 200\ GeV\ vs\ 27GeV$)

$$v_2 = \left\langle \frac{p_x^2 - p_y^2}{p_T^2} \right\rangle$$

- Weak μ_B-dependence
- Small effect of the angular dependence of **dσ/dcosθ**
- Strong flavor dependence

Particles 3 (2020) no.1, 178-192





Channel decomposition :