



# SQM 2022

The 20th International Conference on Strangeness in Quark Matter  
13-17 June 2022 Busan, Republic of Korea

**The quark number scaling of strange quark and light quarks as well as the coupled approach to solving the  $R_{AA} \otimes v_2$  puzzle in heavy-ion collisions**

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**SQM, 06.14, 2022**

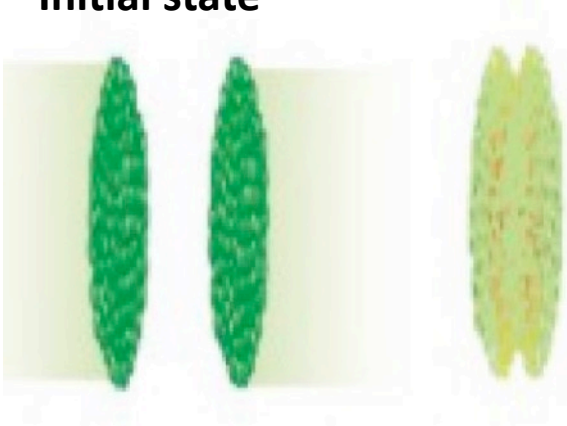


**WAYNE STATE  
UNIVERSITY**

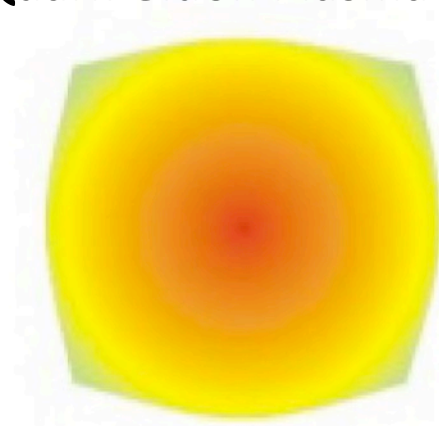


# Illustration of heavy-ion collisions

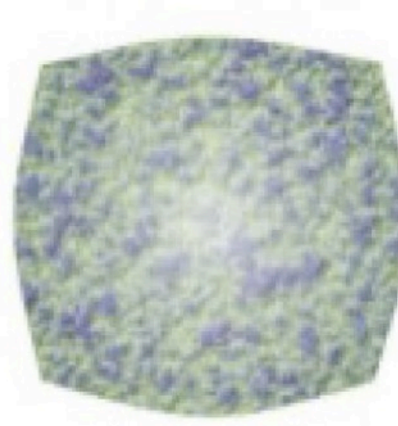
Initial state



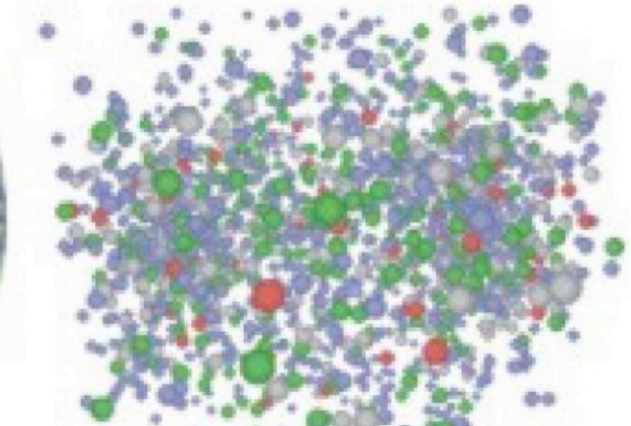
Quark Gluon Plasma



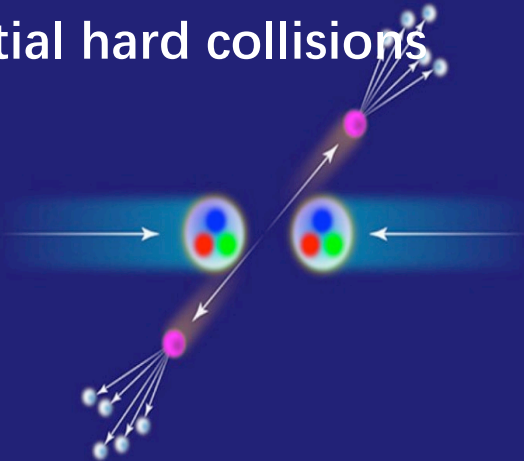
Hadronization



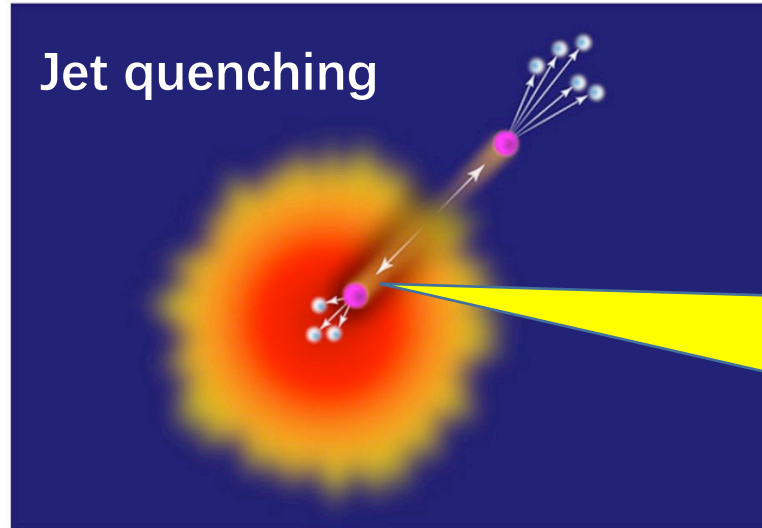
Hadron Cascade



Initial hard collisions



Jet quenching

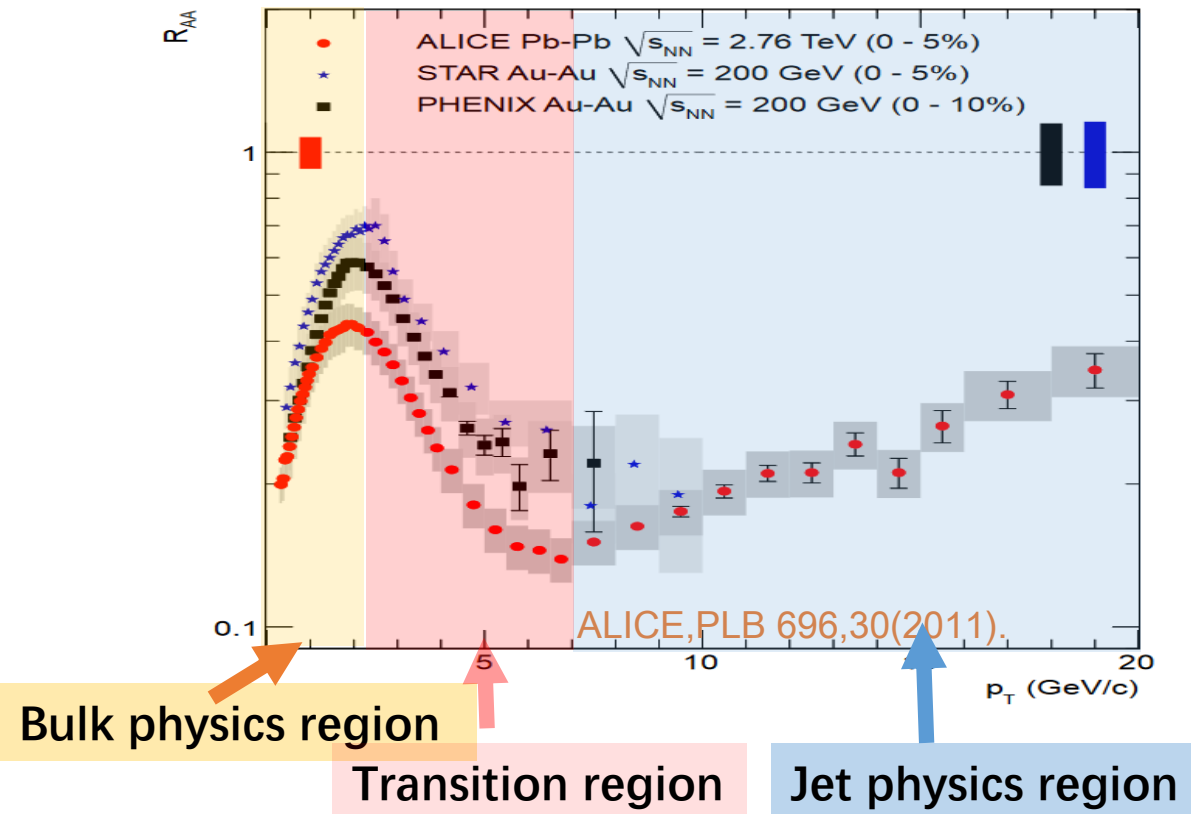


## Jet quenching

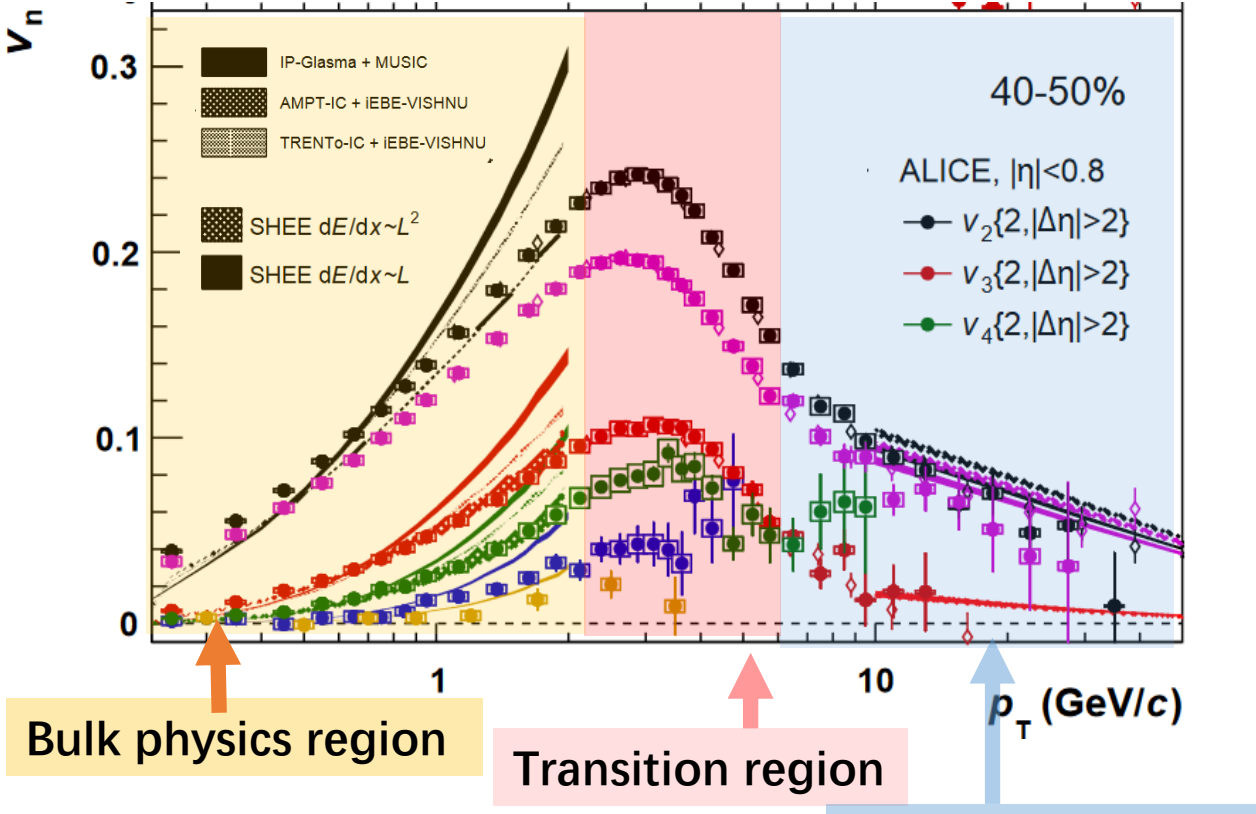
Strong interaction with the energetic partons and the medium.

# Different domains in heavy-ion collisions

ALICE JHEP 1807,103 (2018)



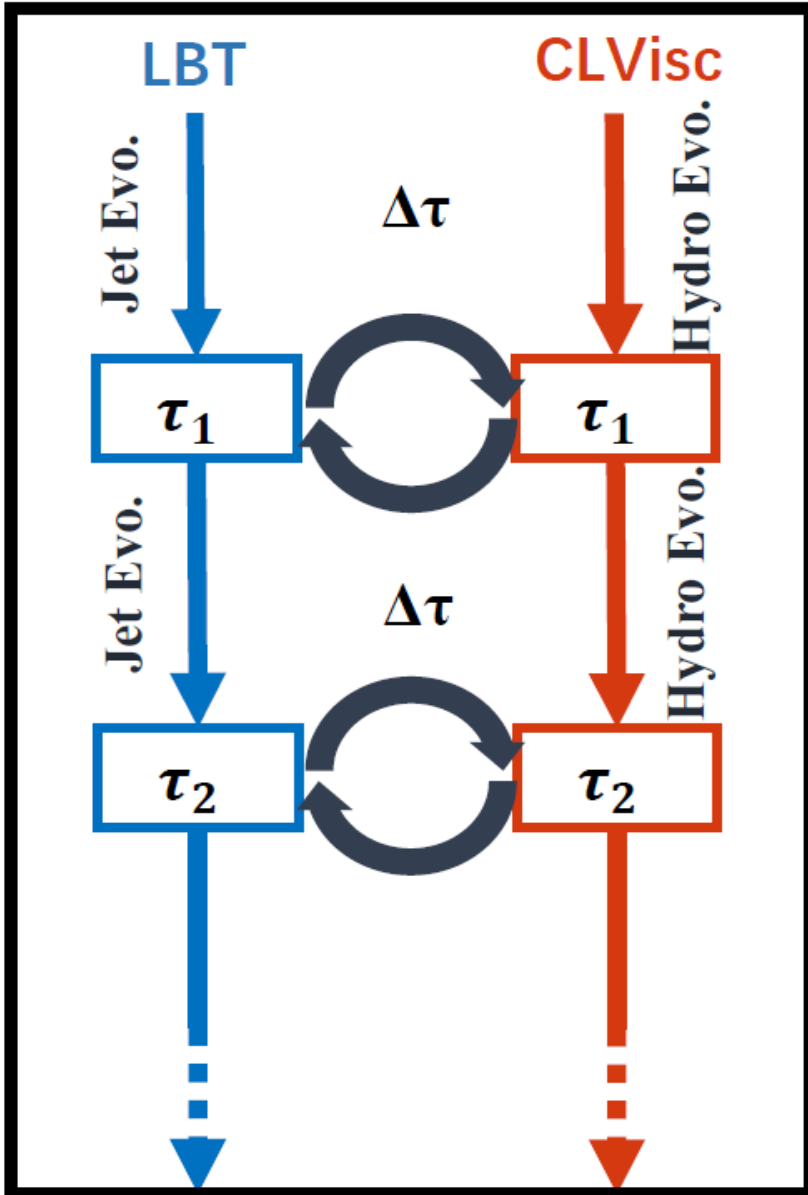
$$R_{AA}(p_T, y, \phi) \equiv \frac{1}{\langle N_{coll} \rangle} \frac{\frac{dN_{AA}}{dp_T dy d\phi}}{\frac{dN_{PP}}{dp_T dy d\phi}}$$



$$v_2(p_T, y) \equiv \frac{\int d\phi \cos(2\phi) \frac{dN}{dp_T dy d\phi}}{\int d\phi \frac{dN}{dp_T dy d\phi}}$$

- Different domains are clearly observed in data in heavy-ion collisions.
- Low  $p_T$  ( $p_T < 2-3$  GeV): bulk physics; High  $p_T$  ( $p_T > 10$  GeV): jet physics.
- Intermediate  $p_T$  ( $3 < p_T < 8-10$  GeV): transition regime; (Not well studied.)

# CoLBT-hydro model



## CoLBT-Hydro model

Linear Boltzmann Transport model + 3+1D hydrodynamic model  
(LBT) (CLVis)

Evolve the energetic partons and the bulk medium concurrently.

Hydrodynamics equations with the source terms:

$$\partial_\mu T_{\text{fluid}}^{\mu\nu} = J^\nu$$

$T_{\text{fluid}}^{\mu\nu}$  : Energy-momentum tensor of the QGP fluid;

$J^\nu$  : Energy-momentum density deposited by energetic partons.  
with the Gaussian smearing:

$$J^\nu(\vec{x}_\perp, \eta_s) = \sum_i \frac{\theta(p_{\text{cut}}^0 - p_i \cdot u) p^\nu}{\tau (2\pi)^{3/2} \sigma_r^2 \sigma_{\eta_s} \Delta\tau} e^{-\frac{(\vec{x}_\perp - \vec{x}_{\perp i})^2}{2\sigma_r^2} - \frac{(\eta_s - \eta_{s i})^2}{2\sigma_{\eta_s}^2}}$$

$p_{\text{cut}}^0$  separates the soft and hard partons

W. Chen, S. Cao, T. Luo, L.-G. Pang, and X.-N. Wang, Phys. Lett. B 810, 135783 (2020), 2005.09678.

# Sophisticated Coalescence model

## • Coalescence model

$$\frac{dN_M}{d^3\mathbf{P}_M} = g_M \int d^3\mathbf{x}_1 d^3\mathbf{p}_1 d^3\mathbf{x}_2 d^3\mathbf{p}_2 f_q(\mathbf{x}_1, \mathbf{p}_1) f_{\bar{q}}(\mathbf{x}_2, \mathbf{p}_2) \times W_M(\mathbf{y}, \mathbf{k}) \delta^{(3)}(\mathbf{P}_M - \mathbf{p}_1 - \mathbf{p}_2),$$

$$\frac{dN_B}{d^3\mathbf{P}_B} = g_B \int d^3\mathbf{x}_1 d^3\mathbf{p}_1 d^3\mathbf{x}_2 d^3\mathbf{p}_2 d^3\mathbf{x}_3 d^3\mathbf{p}_3 f_{q_1}(\mathbf{x}_1, \mathbf{p}_1) \times f_{q_2}(\mathbf{x}_2, \mathbf{p}_2) f_{q_3}(\mathbf{x}_3, \mathbf{p}_3) W_B(\mathbf{y}_1, \mathbf{k}_1; \mathbf{y}_2, \mathbf{k}_2) \times \delta^{(3)}(\mathbf{P}_B - \mathbf{p}_1 - \mathbf{p}_2 - \mathbf{p}_3),$$

### Thermal & hard Partons:

- Thermal partons generated by hydro
- Hard partons generated by PYTHIA8, then suffered with energy loss by CoLBT

### Coalescence processes:

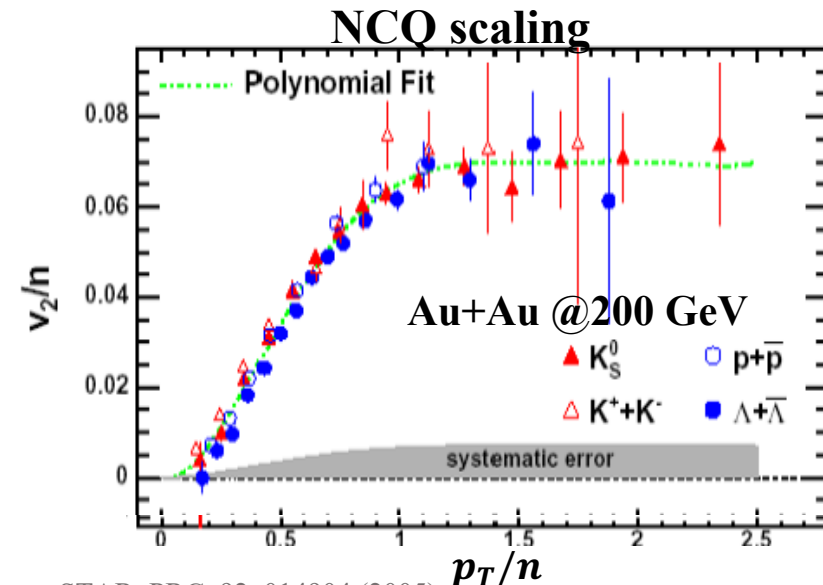
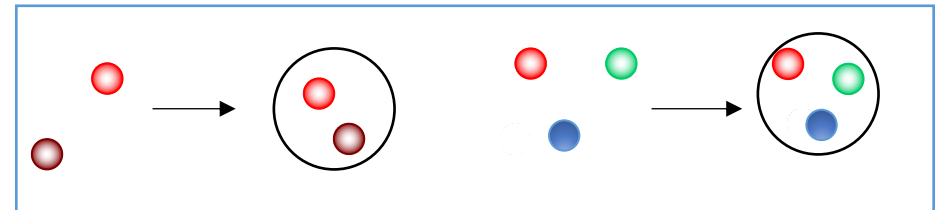
- thermal - thermal parton coalescence
- thermal - hard parton coalescence
- hard - hard parton coalescence

Han, Fries and Ko, PRC 93, 045207 (2016).

Zhao, Ko, Liu, Qin and Song, Phys.Rev.Lett. 125, 072301 (2020).

$g_{B(M)}$  is statistic factor,  $f_{q/\bar{q}}$  is the phase-space distribution of (anti)quarks,  $W_{M/B}$  is Wigner function of meson(baryon).

Here, we use the harmonic oscillator for wave functions of hadrons, then do the Wigner transformation to get the  $W_{M/B}$ .



STAR, PRC, 92, 014904 (2005).

# String fragmentation

- **Colorless string fragmentation**
- Shower partons that do not coalesce will hadronize through string fragmentation using PYTHIA8.
- Shower partons have lost their original color configurations. We connect strings to minimize the distance

$$\Delta R = \sqrt{(\Delta\eta)^2 + (\Delta\phi)^2}$$

JETSCAPE framework: Phys. Rev. C 102, 054906 (2020).

W. Zhao, W. Ke, W. Chen, T. Luo and X. N.Wang, PRL. 128, 022302 (2022).

# Framework of calculations

## Hydro-Coal-Frag hadronization

### Thermal hadrons, low $p_T$ (CLVis):

- generated by hydro. with Cooper-Frye.  
Meson:  $p_T < 2p_{T1}$ ; baryon:  $p_T < 3p_{T1}$ .

- initial shower partons from pythia8 with  $p_T > p_{T2}$

### Coalescence hadrons (Coal Model):

- generated by coalescence model including thermal-thermal, thermal-hard & hard-hard coalescence.

### Fragmentation hadrons :

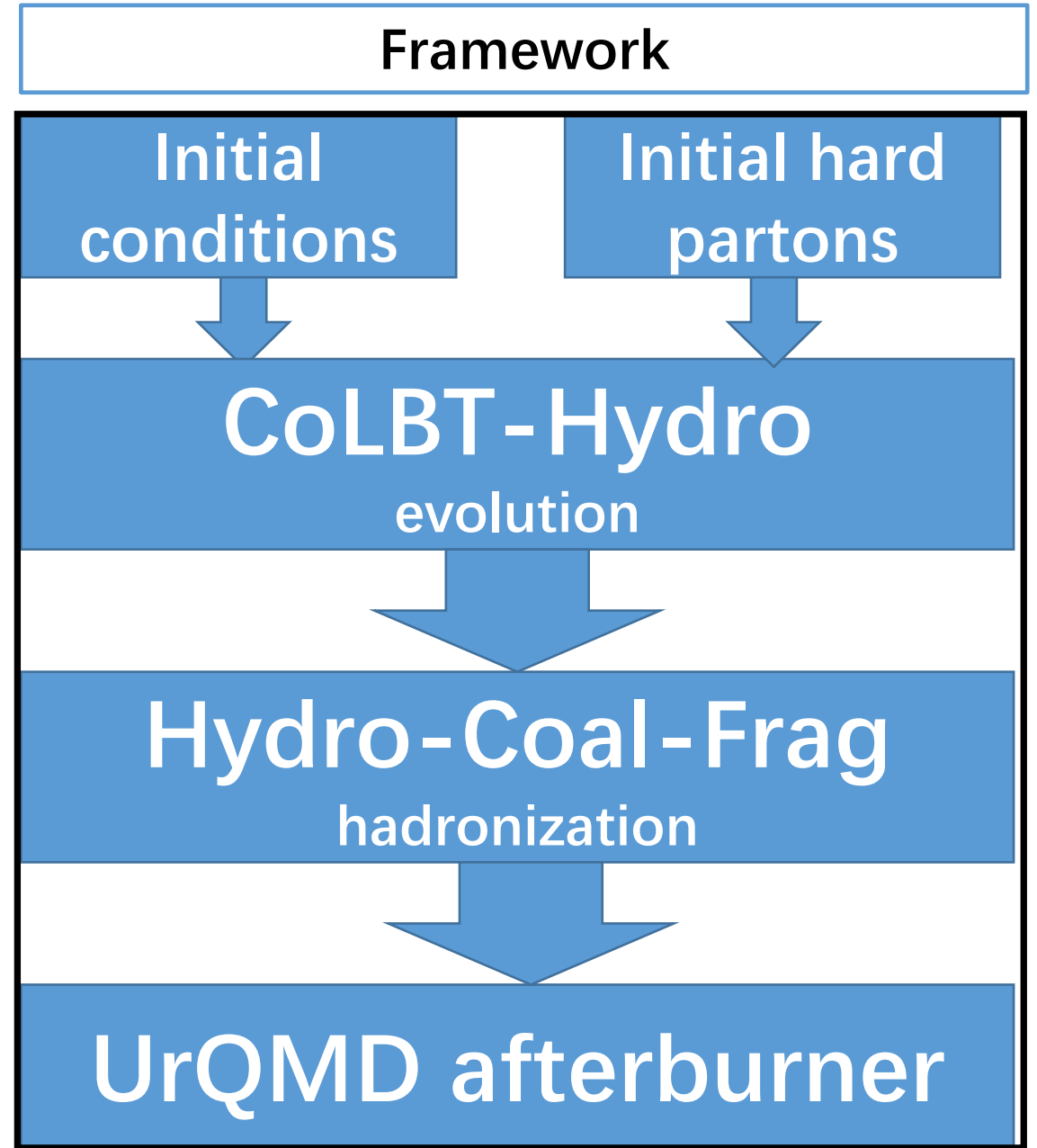
- the remnant hard quarks feed to fragmentation.

### UrQMD afterburner:

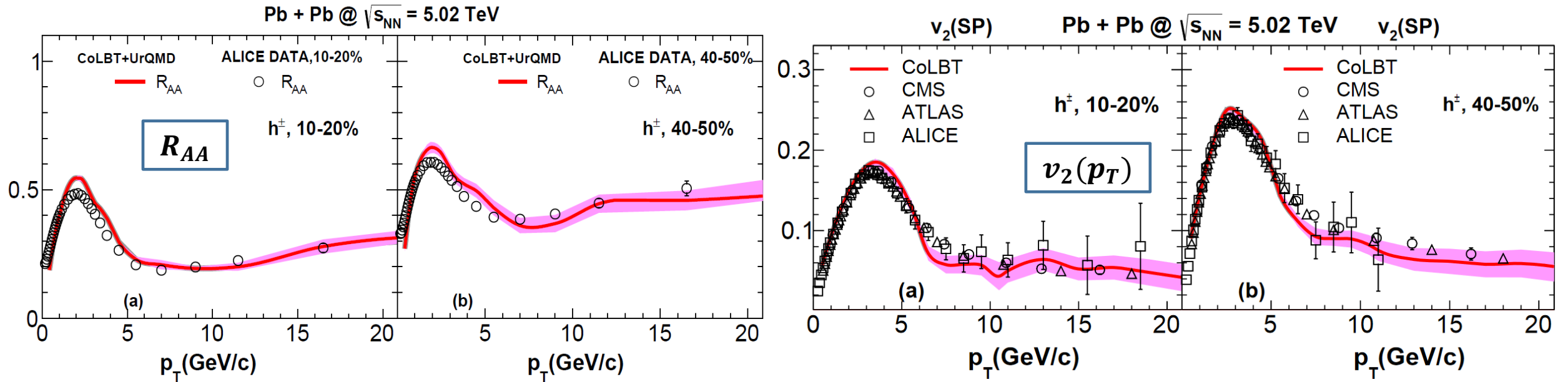
- All hadrons are feed into UrQMD for hadronic evolution, scatterings and decays

Hydro.    Coalescence, fragmentation    fragmentation

0                      3GeV                      7GeV                       $P_T$



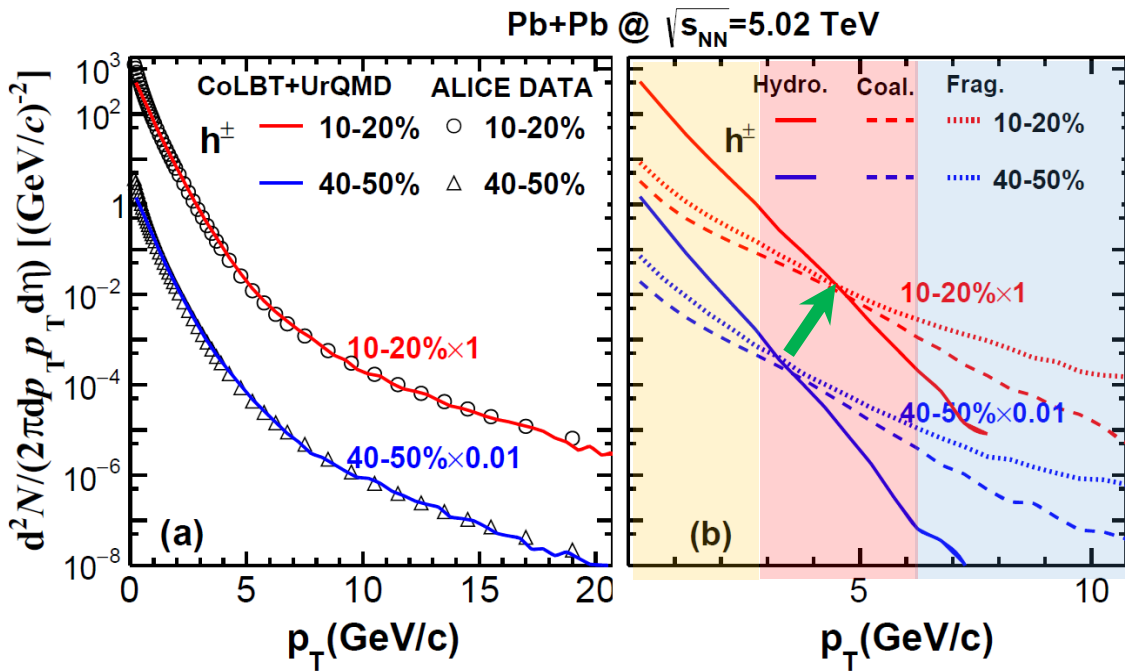
# $R_{AA}$ v.s. $v_2(p_T)$ from low $p_T$ to high $p_T$



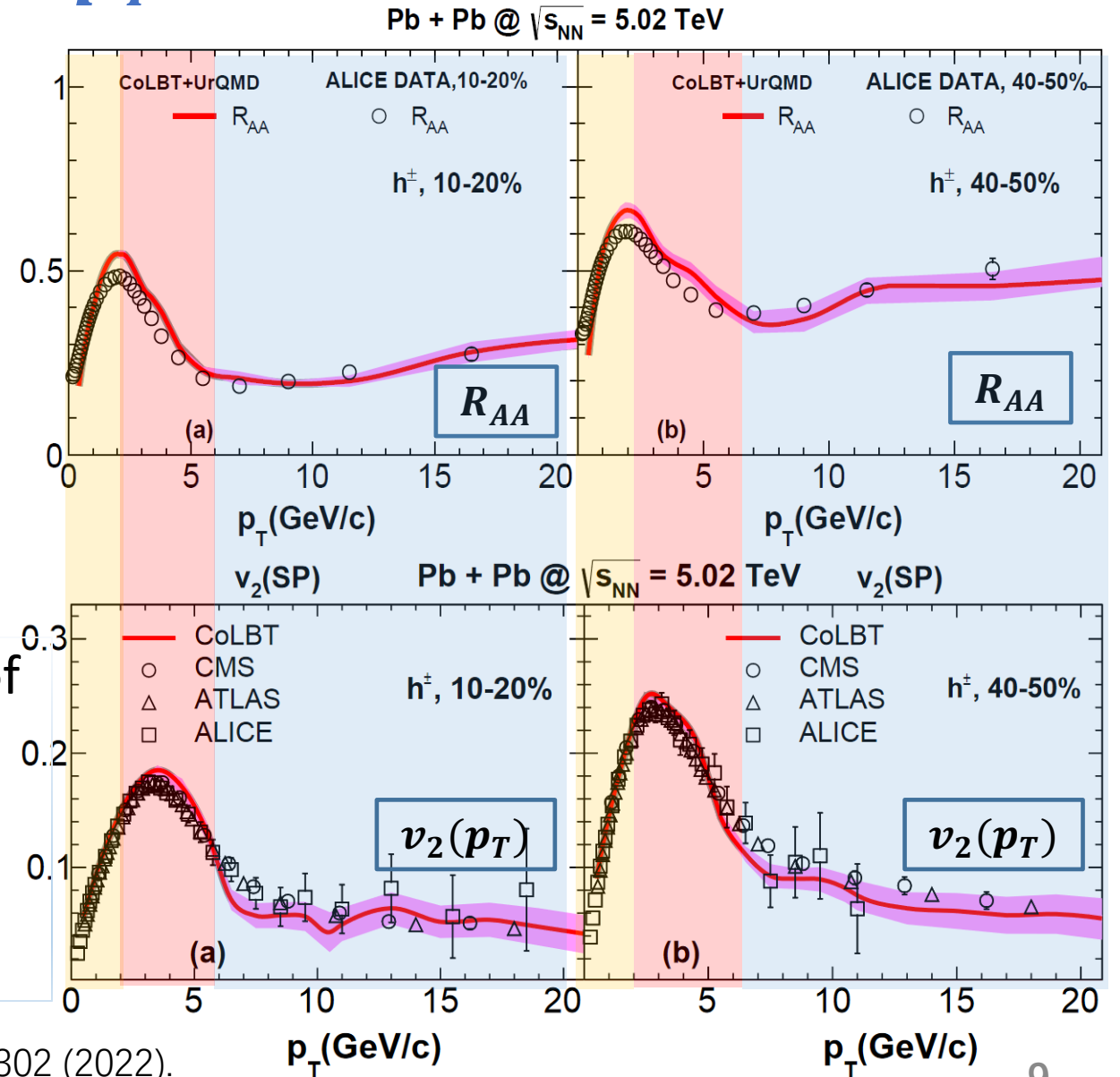
- CoLBT-hydro with Hydro-Coal-Frag hadronizations can simultaneously describe the  $R_{AA}$  and collective flow from low  $p_T$  to high  $p_T$  regions in Pb+Pb collisions.



# Transition from low $p_T$ to high $p_T$

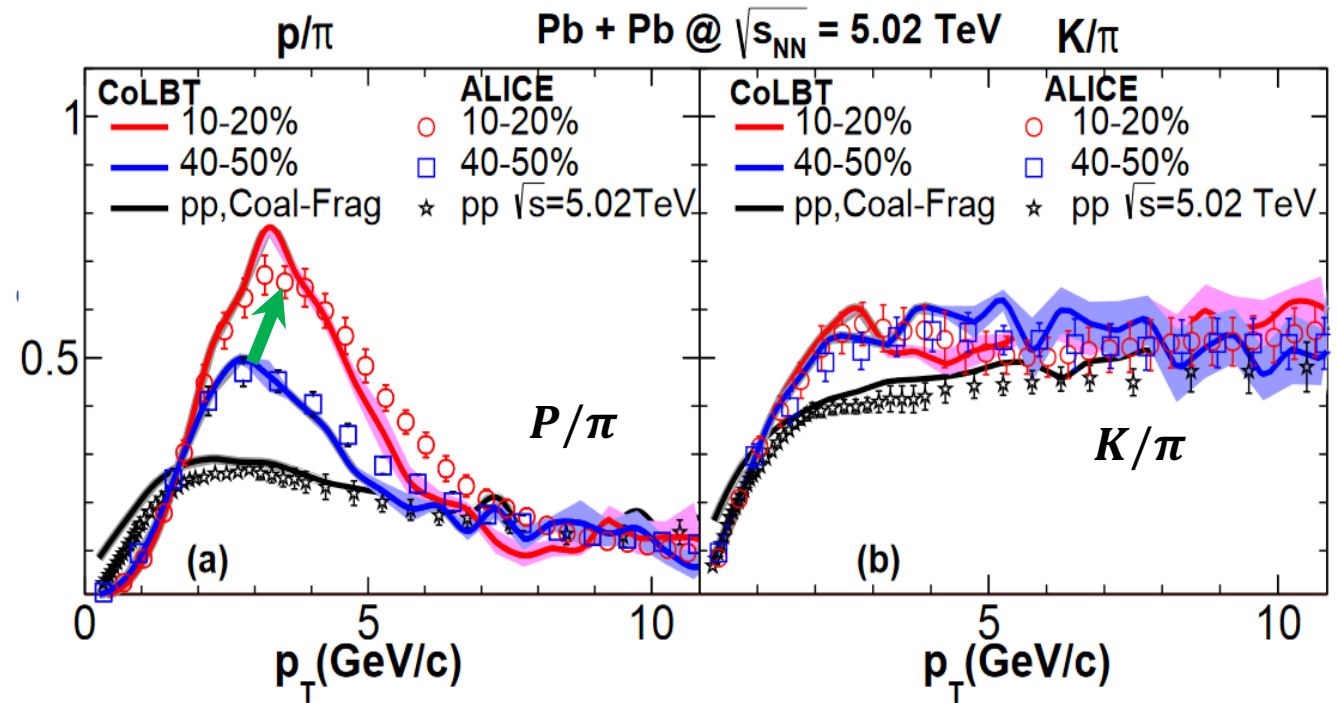
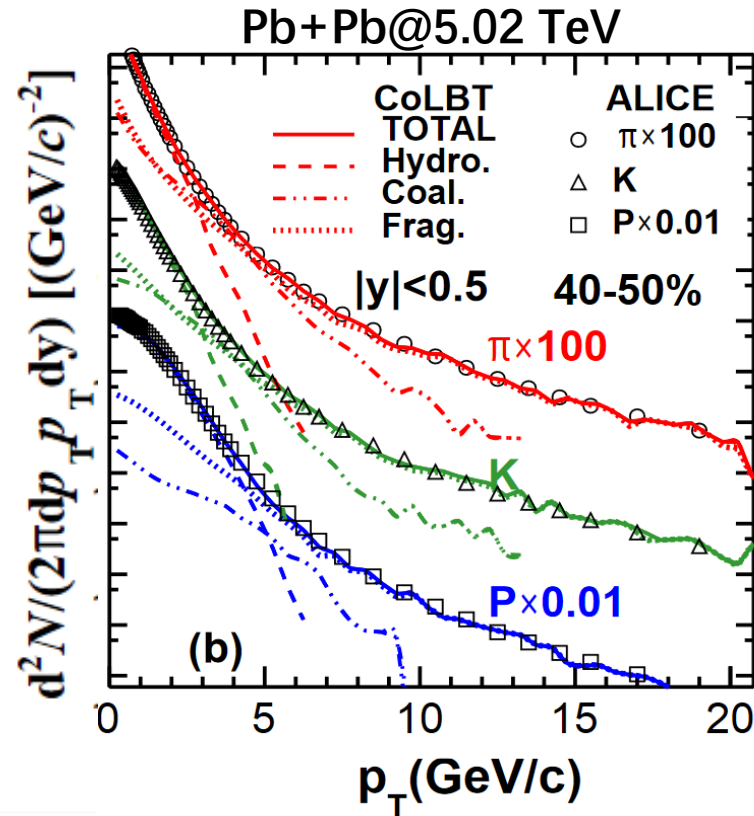


- CoLBT-hydro nicely describes the spectra of charged from 0 to 20 GeV.
- Low  $p_T$ : hydro; Intermediate  $p_T$ : transition regime; High  $p_T$ : jet physics.
- Transition  $p_T$  is higher in central collisions.



W. Zhao, W. Ke, W. Chen, T. Luo and X. N.Wang, PRL. 128, 022302 (2022).

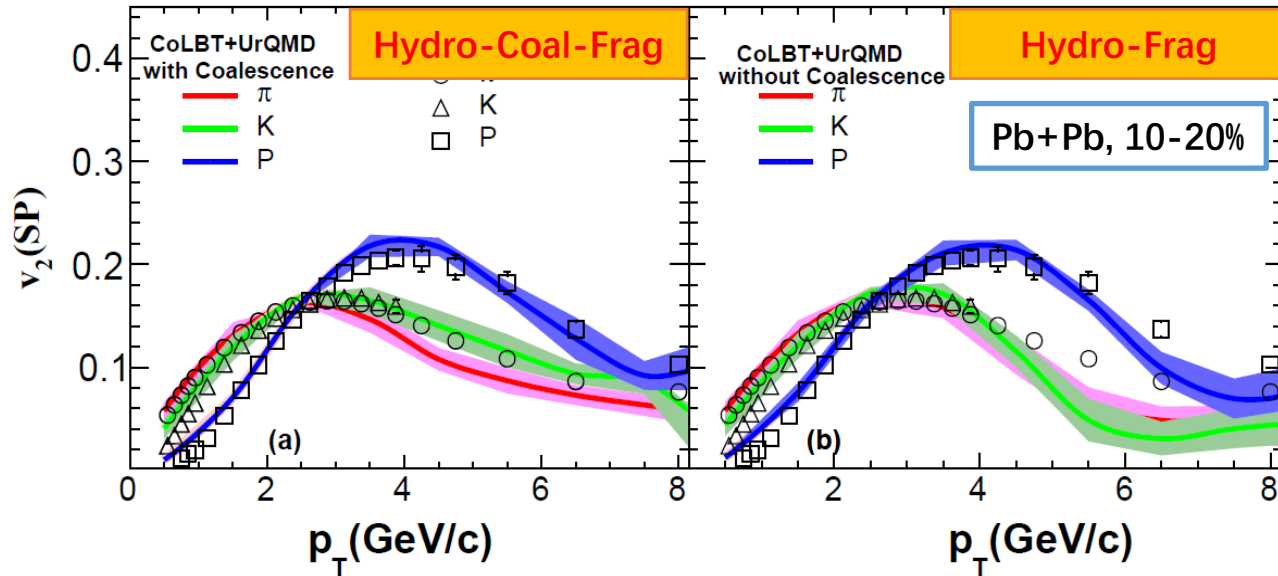
# Transverse momentum spectra of identified hadrons



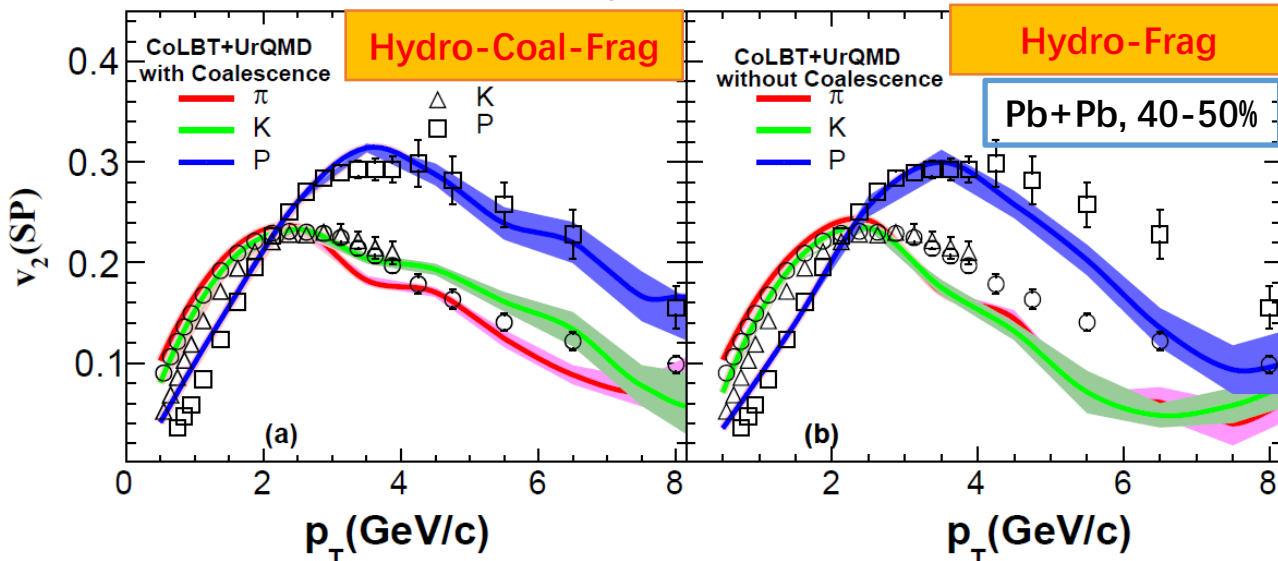
- CoLBT-hydro nicely describes the spectra of identified hadrons,  $P/\pi$  and  $K/\pi$  from 0 to 20 GeV.
- $P/\pi$  in Pb+Pb is higher than pp;  $P/\pi$  peak moves to higher  $p_T$  in central collision.
- $P/\pi$  and  $K/\pi$  approach to the p-p value at high  $p_T$ .

# Collective flow of identified hadrons

Pb + Pb @  $\sqrt{s_{NN}} = 5.02$  TeV, 10-20%



Pb + Pb @  $\sqrt{s_{NN}} = 5.02$  TeV, 40-50%

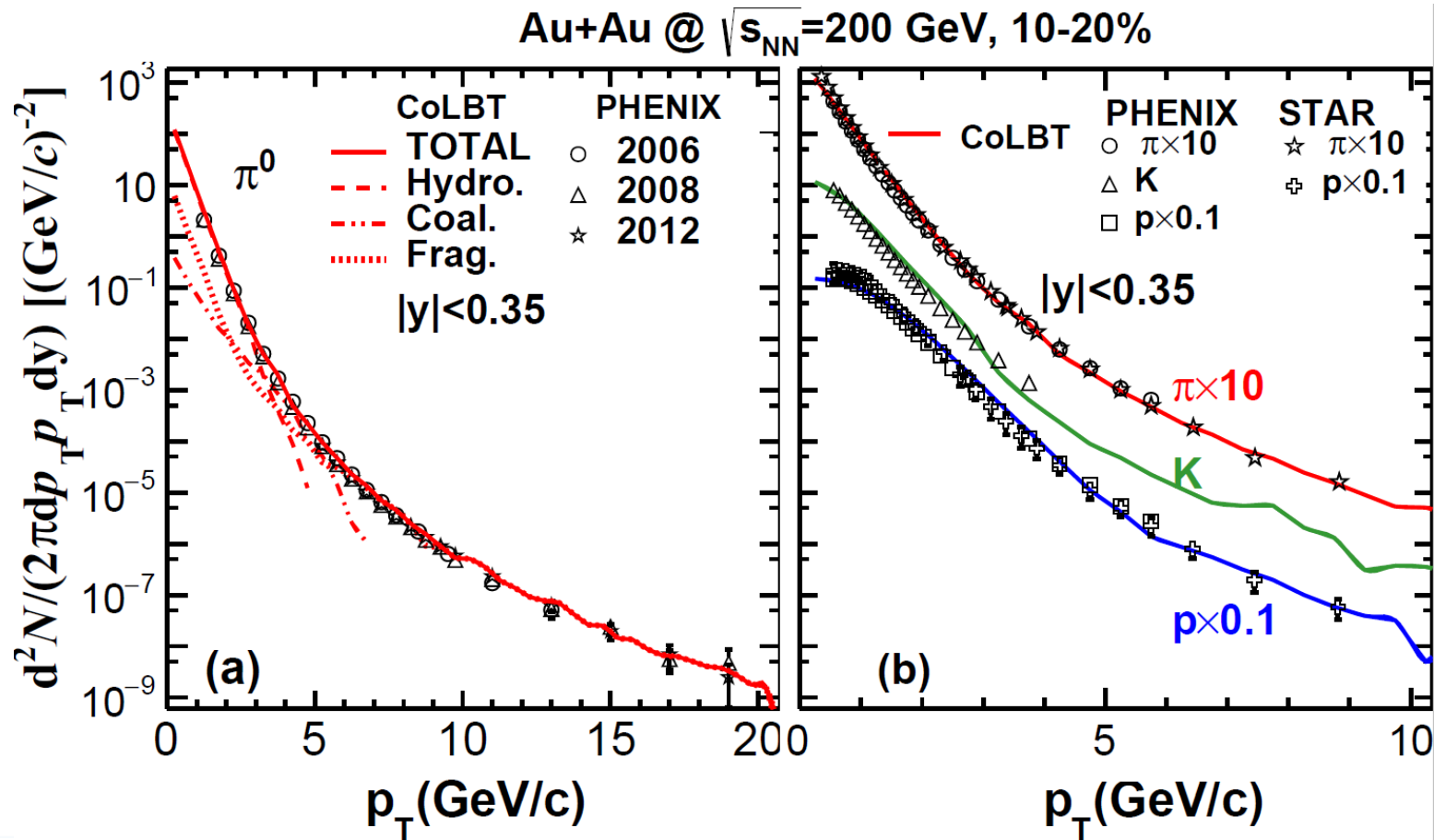


- CoLBT-hydro with Hydro-Coal-Frag works well for PID flow from 0 to 8 GeV.
- $v_2(p_T)$  of P larger than  $\pi$  and K at 3 GeV, caused by interplay between hydro. Coal. and frag.
- Quark coalescence is important for Pb+Pb collisions at intermediate  $p_T$  range.

W. Zhao, W. Ke, W. Chen, T. Luo and X. N.Wang, PRL. 128, 022302 (2022).

# Predictions for Au-Au at RHIC

# Spectra at Au-Au at RHIC

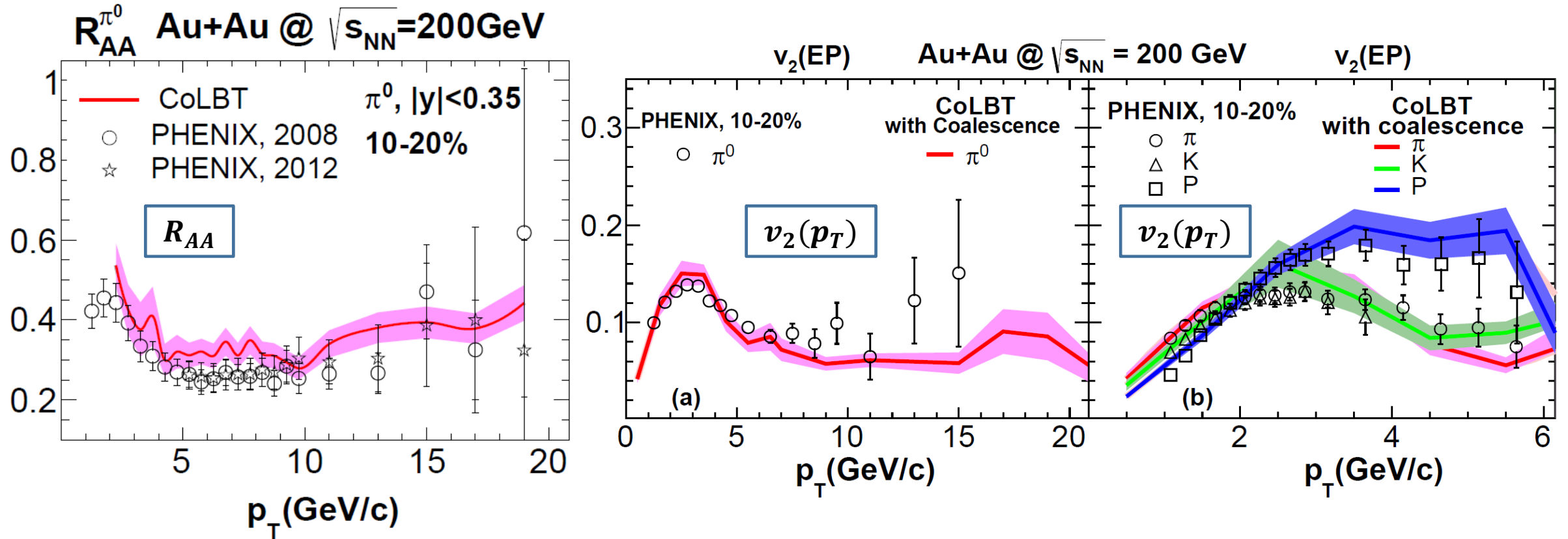


- With parameters fixed at LHC, CoLBT-hydro nicely predicts the spectra of  $\pi^0$  and of  $\pi^\pm$ , K and P from low  $p_T$  to high  $p_T$  in Au-Au at 200 GeV.
- Low  $p_T$ : hydro; Intermediate  $p_T$ : transition region; High  $p_T$ : fragmentation.

W. Zhao, W. Ke, W. Chen, T. Luo and X. N.Wang, PRL. 128, 022302 (2022).

W. Zhao, W. Ke, W. Chen, T. Luo and X. N.Wang, in preparation.

# $R_{AA}$ and $v_2(p_T)$ at Au-Au at RHIC

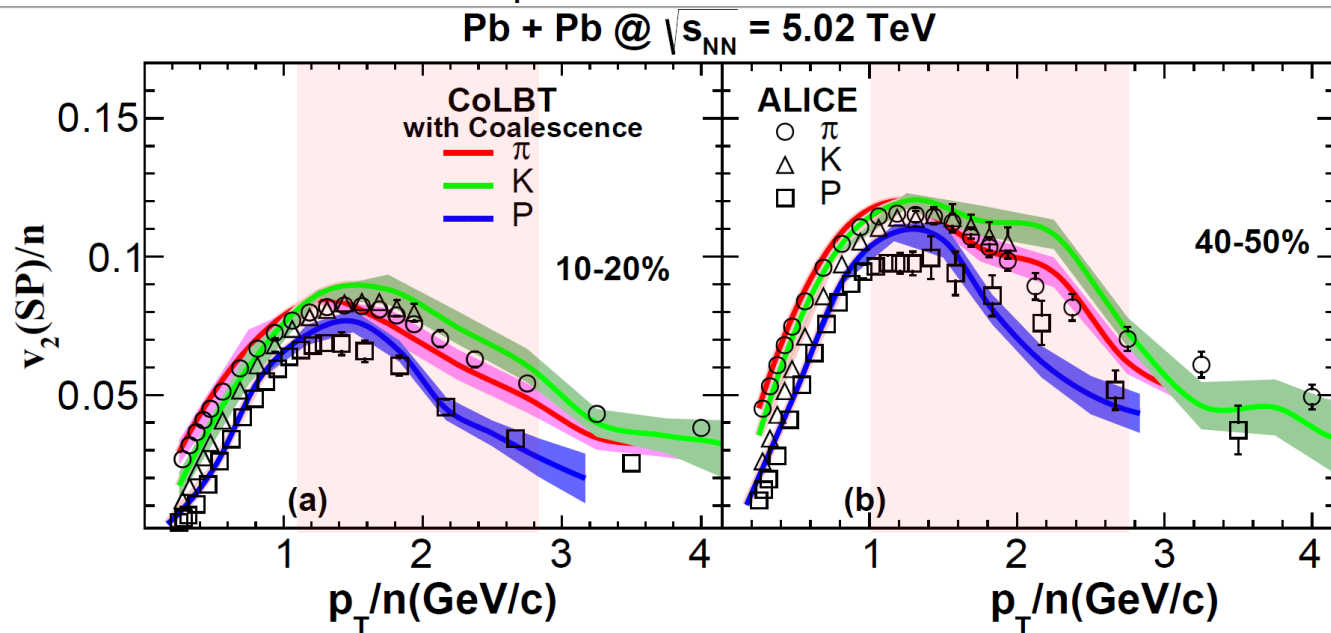
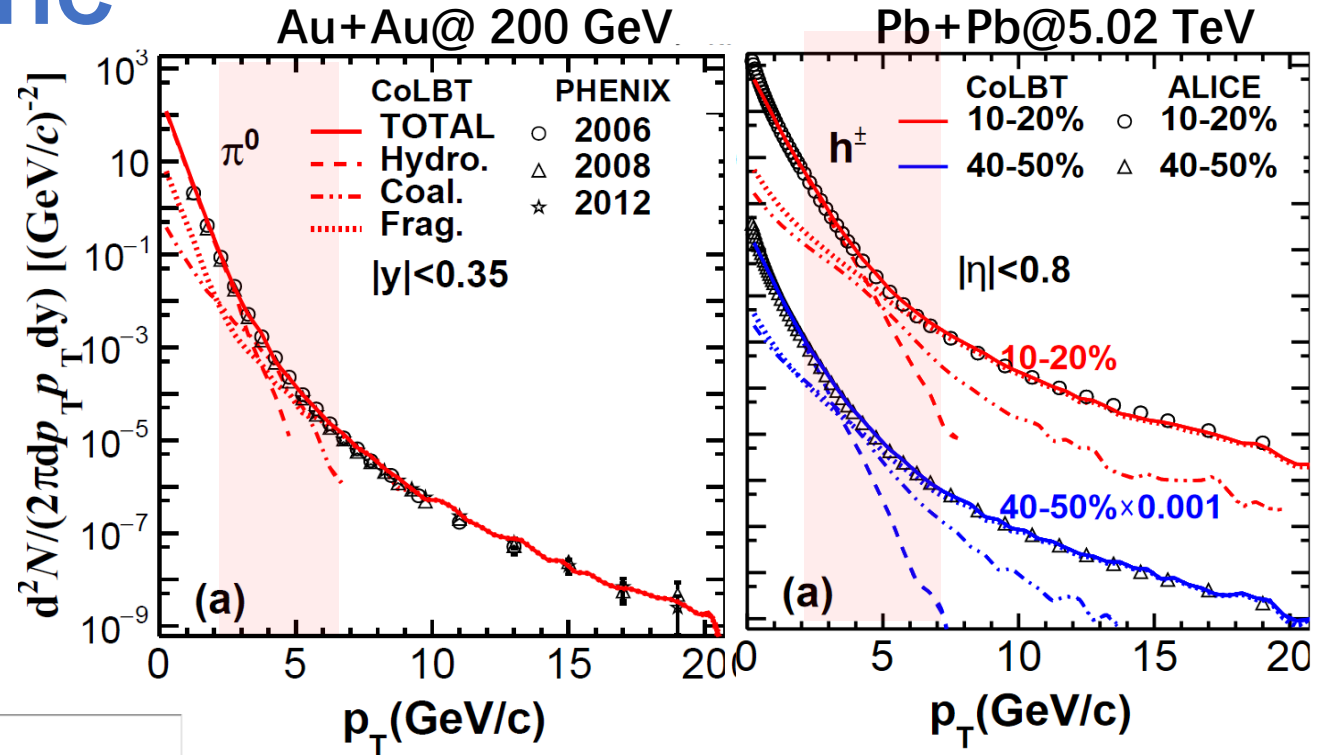
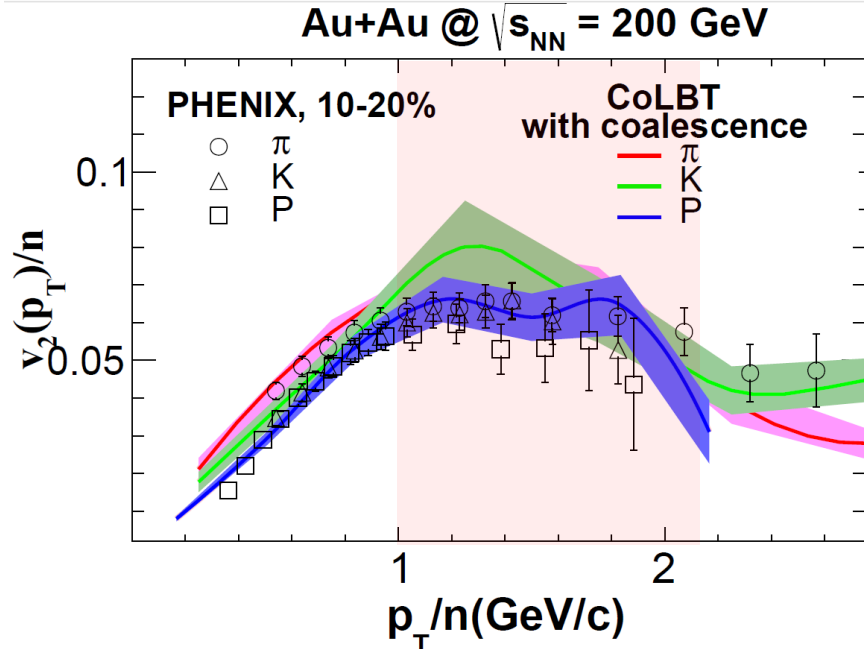


- With parameters fixed at LHC, CoLBT-hydro nicely predicts the  $R_{AA}$  and  $v_2(p_T)$  from 0 to 20 GeV in Au-Au at 200 GeV.
- CoLBT-hydro nicely predicts the  $v_2(p_T)$  of  $\pi$ , K and P from 0 to 6 GeV in RHIC.

W. Zhao, W. Ke, W. Chen, T. Luo and X. N.Wang, PRL. 128, 022302 (2022).

W. Zhao, W. Ke, W. Chen, T. Luo and X. N.Wang, in preparation.

# NCQ scaling at RHIC and LHC



- NCQ scaling at intermediate  $p_T$  are caused by interplay of hydro, coal. and frag.

W. Zhao, W. Ke, W. Chen, T. Luo and X. N.Wang, PRL. 128, 022302 (2022).  
 W. Zhao, W. Ke, W. Chen, T. Luo and X. N.Wang, in preparation.

# Summary

- CoLBT-hydro with Hydro-Coal-Frag hadronization simultaneously describe the  $R_{AA}$  and collective flow from low  $p_T$  to high  $p_T$  in Pb+Pb collisions.
- CoLBT-hydro also nicely describes the collective flow of identified hadrons with  $p_T$  from 0 to 8 GeV.
- Quark coalescence is important in heavy-ion collisions.
- With parameters fixed at LHC, CoLBT-hydro excellently predicts the  $R_{AA}$  and collective flow from low  $p_T$  to high  $p_T$  in Au+Au collisions at RHIC.

**Thanks for Your Attention!**

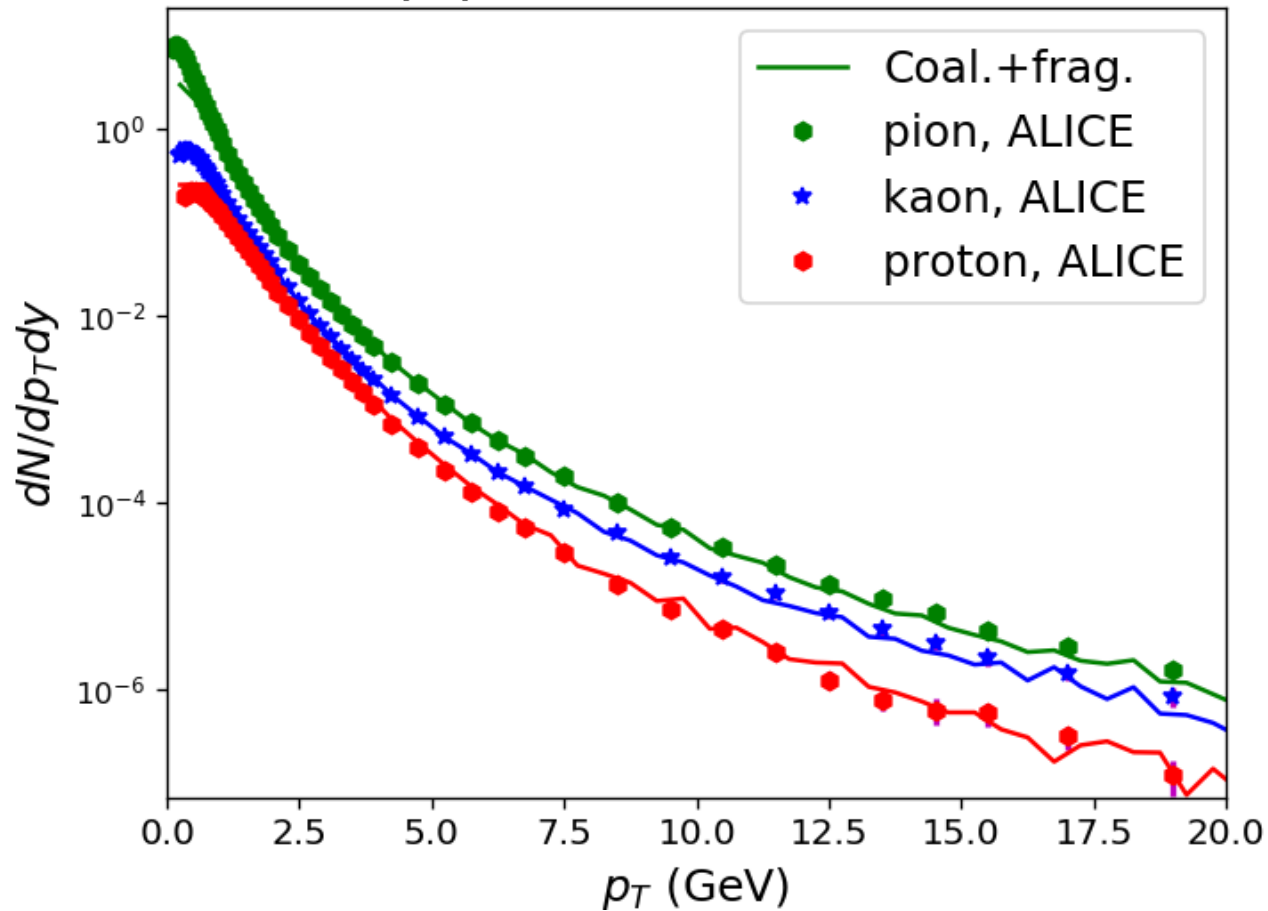


# Back Up

Back Up

# Verification of the hadronization code in p-p

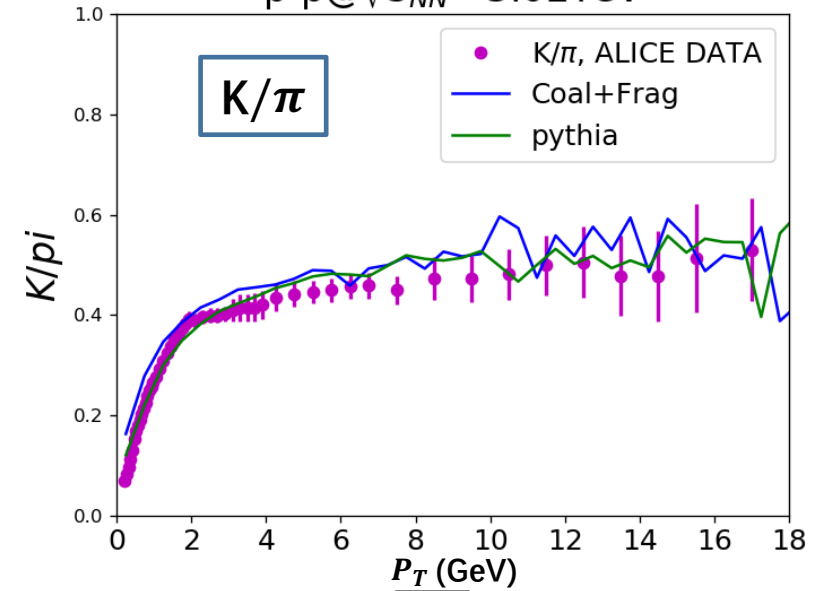
p-p@ $\sqrt{s_{NN}}=5.02\text{TeV}$



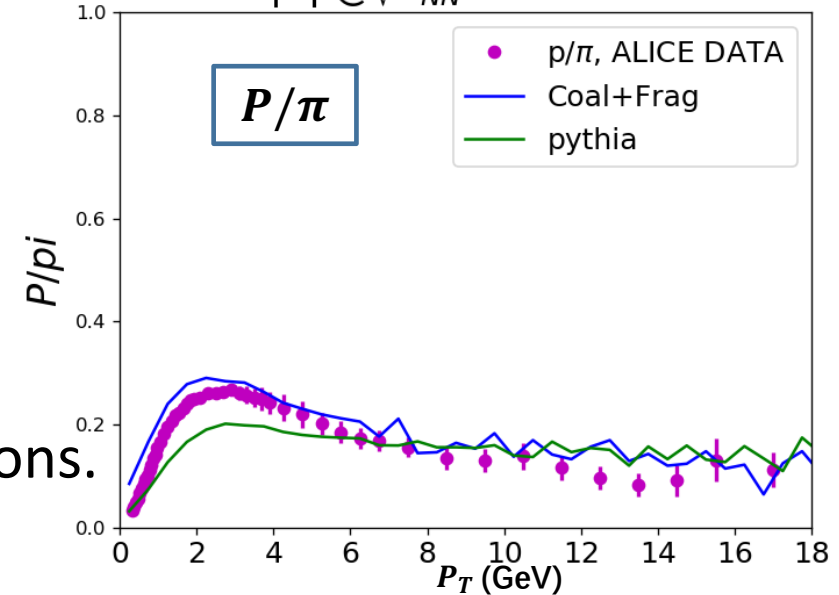
- Coal-Frag model can well reproduce the  $p_T$ -spectra of  $\pi$ , K and P as well as the  $K/\pi$  and  $P/\pi$  in p-p collisions.

W. Zhao, W. Ke, W. Chen, T. Luo and X. N.Wang, PRL. 128, 022302 (2022).

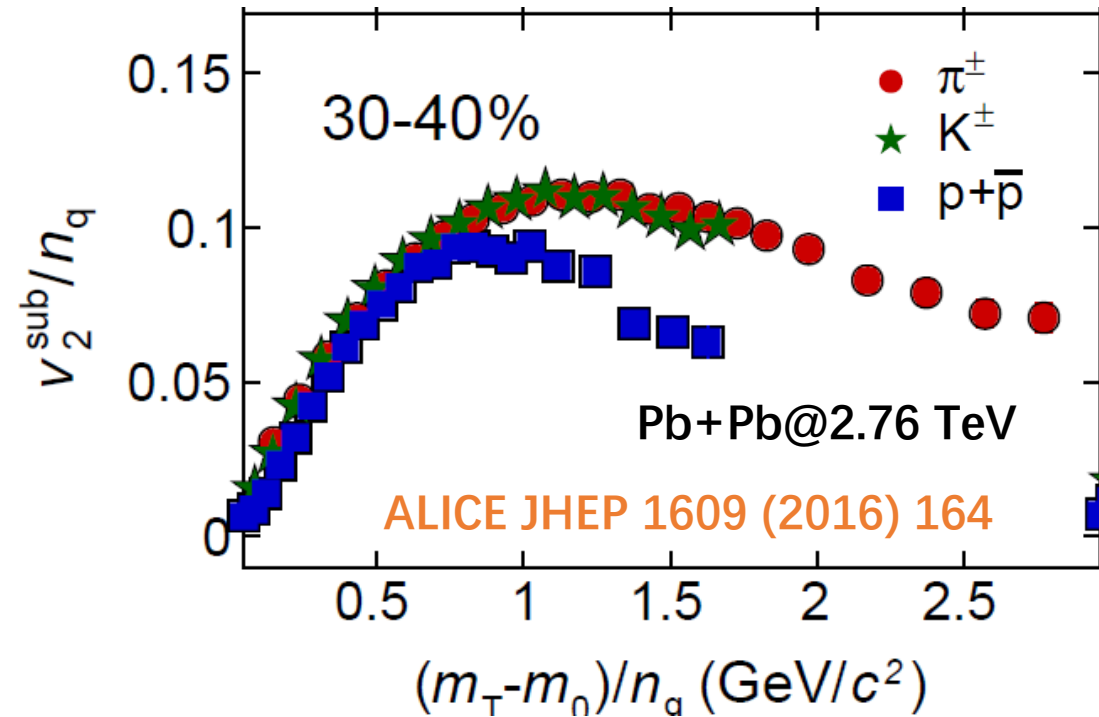
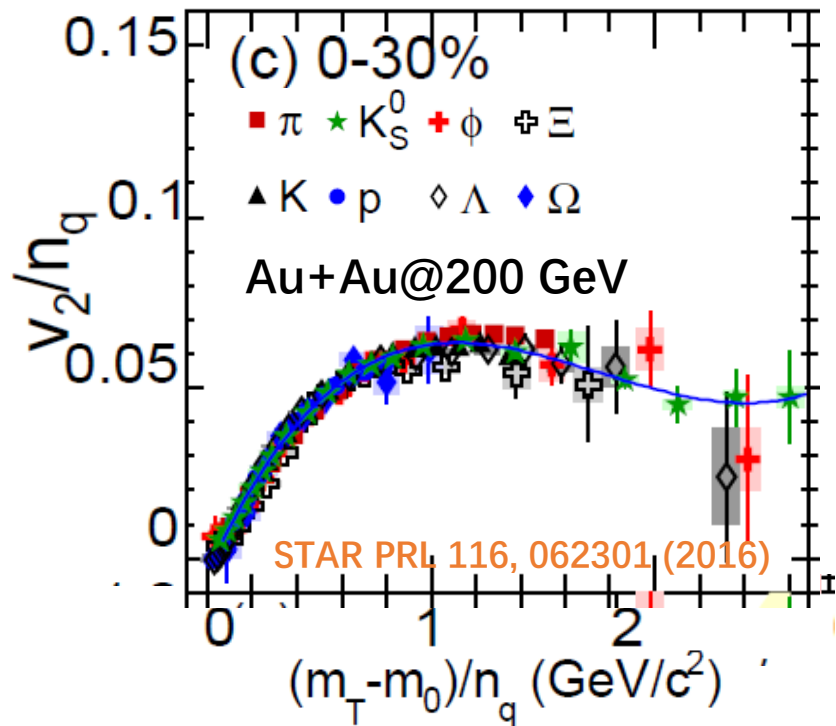
p-p@ $\sqrt{s_{NN}}=5.02\text{TeV}$



p-p@ $\sqrt{s_{NN}}=5.02\text{TeV}$



# “NCQ scaling puzzle” between RHIC and LHC



- At intermediate  $p_T$  range, NCQ scaling behavior is perfect at Au-Au@200 GeV collisions, how it's greatly violated at Pb-Pb@2.76 TeV.
- What causes such discrepancy between RHIC and LHC?

# Initialization of hard partons

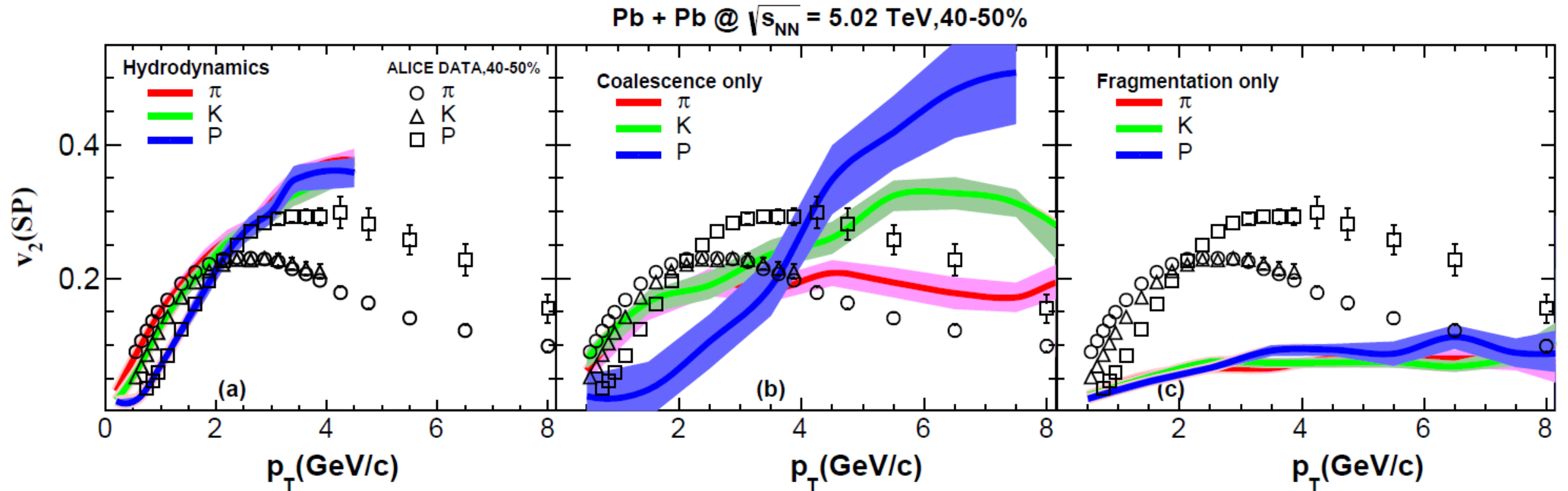
- Transverse locations of hard collisions  $\mathbf{r}_\perp$  are sampled from the binary collision density

$$\frac{dN_{\text{coll}}}{d\mathbf{r}_\perp^2}(\mathbf{r}_\perp; b) = T_{\text{Pb}}(\mathbf{r}_\perp + \mathbf{b}/2)T_{\text{Pb}}(\mathbf{r}_\perp - \mathbf{b}/2)$$

- Initial partons in the initial vacuum showers free-streams during the formation time of vacuum splittings

$$\tau_f = 2x(1 - x)E/k_\perp^2$$

# $v_2(p_T)$ of hydro. Coal. and Frag. parts



W. Zhao, W. Ke, W. Chen, T. Luo and X. N. Wang, arXiv:2103.14657.

- Hydro. works at low  $p_T$  range ( $p_T < 2-3$  GeV).
- Quark coalescence generates large  $v_2$  at intermediate  $p_T$  ( $3 < p_T < 8$  GeV)
- Fragmentation can't generate enough  $v_2$  below 8 GeV.

# Wigner functions of hadrons

To guarantee positive value of Wigner function for stable Monte Carlo sampling, the Wigner function replaced by the overlap of hadron Wigner function  $W_M$  with parton's Wigner function,  $W_{q,\bar{q}}$ :

$$\begin{aligned} \overline{W}_M(\mathbf{y}, \mathbf{k}) &= \int d^3\mathbf{x}'_1 d^3\mathbf{k}'_1 d^3\mathbf{x}'_2 d^3\mathbf{k}'_2 \\ &\times W_q(\mathbf{x}'_1, \mathbf{k}'_1) W_{\bar{q}}(\mathbf{x}'_2, \mathbf{k}'_2) W_M(\mathbf{y}', \mathbf{k}'). \end{aligned} \quad (3)$$

Using harmonic oscillator for wave functions of excited states of hadrons,

$$\phi_n(x) = \left(\frac{m\omega}{\pi\hbar}\right)^{1/4} \frac{1}{\sqrt{2^n n!}} H_n(\xi) e^{-\xi^2/2}, \quad (4)$$

$\xi = \sqrt{\frac{m\omega}{\hbar}}x$ ,  $H_n(\xi)$  are Hermite polynomials,  $\omega$  is the oscillator frequency.

K. C. Han, R. J. Fries and C. M. Ko, Phys. Rev. C 93, no. 4, 045207 (2016).

# Wigner functions of hadrons

The quark wave function to be Gaussian wave packet, the wigner function of a meson in  $n$ -th excited state is

$$\overline{W}_{M,n}(\mathbf{y}, \mathbf{k}) = \frac{v^n}{n!} e^{-v}. \quad (5)$$

with

$$v = \frac{1}{2} \left( \frac{\mathbf{y}^2}{\sigma_M^2} + \mathbf{k}^2 \sigma_M^2 \right). \quad (6)$$

Similarly, the Gaussian smeared Wigner function for baryon is:

$$\overline{W}_{B,n_1,n_2}(\mathbf{y}_1, \mathbf{k}_1; \mathbf{y}_2, \mathbf{k}_2) = \frac{v_1^{n_1}}{n_1!} e^{-v_1} \cdot \frac{v_2^{n_2}}{n_2!} e^{-v_2}, \quad (7)$$

with

$$v_i = \frac{1}{2} \left( \frac{\mathbf{y}_i^2}{\sigma_{Bi}^2} + \mathbf{k}_i^2 \sigma_{Bi}^2 \right), \quad i = 1, 2. \quad (8)$$

# Positions of fragmented hadron

the center of mass of a string  $\mathbf{r}_s$  and the time of the latest born remnant parton  $t_s$ .  
The formation time of a hadron with  $E$  and  $m_T$  is

$$t_f = E/m_T^2,$$

The position and time of each hadron produced from the string fragmentation are written by

$$t_H = t_s + t_f$$

$$\mathbf{r}_H = \mathbf{r}_s + \mathbf{v}t_f$$

Where  $\mathbf{v}$  is the speed of the center-of-mass of the partons in the lab frame

Han, Fries and Ko, PRC 93, 045207 (2016).

Zhao, Ko, Liu, Qin and Song, Phys.Rev.Lett. 125, 072301 (2020).