

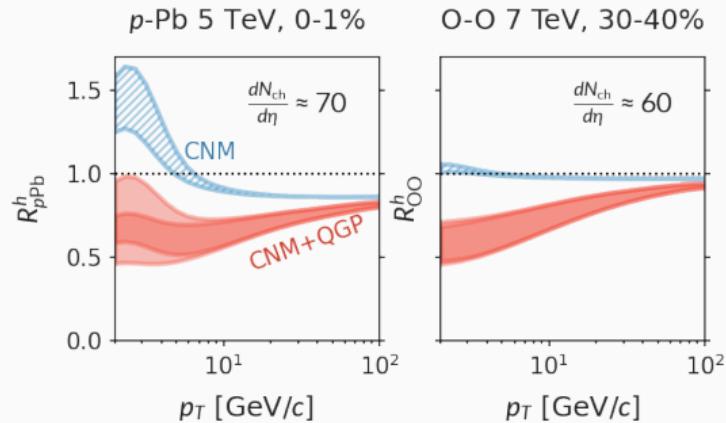


## Impacts of heavy-flavor probes in small colliding systems

Weiyao Ke & Ivan Vitev, Los Alamos National Laboratory, arXiv:2204.00634

### Bottom line up front

- Dynamical CNM effects for no-QGP baselines.
- QGP-modified splitting functions from SCET<sub>G</sub> + modified DGLAP evolution.
- Collisional energy loss more important in small systems.

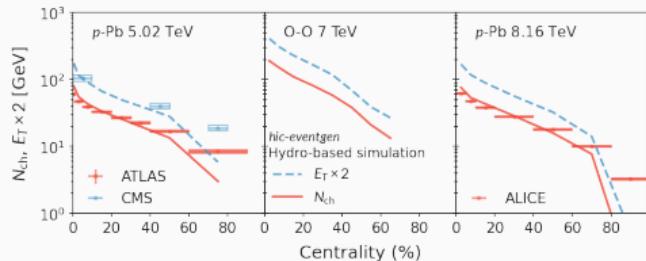


# Challenges of identifying QGP signals in small systems with high- $p_T$ hadrons

- Potential cold nuclear matter (CNM) effects distort the baseline  
**This work** → dynamical calculations of CNM modifications [1].
- Nuclear geometry / centrality definition:  
**This work** → study both  $p$ -Pb,  $d$ -Au and light-ion systems O-O.
- Is there a QGP? How it evolves in small systems?  
**This work** → test two extremes:
  - No QGP formation, color density=0 (medium scattering centers=0).
  - QGP formation with maximum color density  $\propto T^3$ , given by hydrodynamic simulation [2] ▽.

Max.  $T$  of QGP using lattice EoS [3]

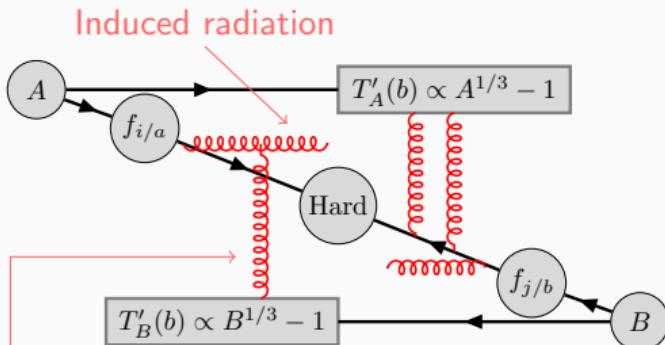
Systems	$p$ -Pb 5 TeV		O-O 7 TeV	
	0-1%	60-90%	0-10%	30-50%
$T_{\max}$ [GeV]	0.315	0.174	0.325	0.263



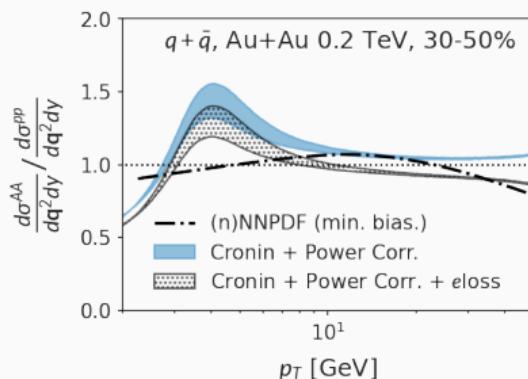
# Dynamical cold nuclear matter (CNM) effects

- Cronin (broadening) effect:  $\Delta k_\perp^2 \propto \mu^2 A^{1/3}$ ; dynamical shadowing:  $\delta x_a/x_a \propto \frac{\mu^2 B^{1/3}}{-u}$  [4]
- Cold nuclear matter energy loss from collision-induced soft gluon emission [5]

$$\frac{\Delta x_a}{x_a} \approx \frac{\alpha_s C_R}{\pi^2} \frac{L_B}{\lambda_g} \int_0^{\frac{\mu p^+}{4}} d^2 \mathbf{q}_\perp \frac{\mu^2}{\pi(\mathbf{q}_\perp^2 + \mu^2)^2} \left[ \frac{\mathbf{q}_\perp^2}{\mathbf{k}_\perp^2 (\mathbf{k}_\perp - \mathbf{q}_\perp)^2} - \frac{2(\mathbf{q}_\perp^2 - \mathbf{q}_\perp \cdot \mathbf{k}_\perp)}{\mathbf{k}_\perp^2 (\mathbf{k}_\perp - \mathbf{q}_\perp)^2} \frac{\sin \frac{\mathbf{k}_\perp^2 L_B}{xp^+}}{\frac{\mathbf{k}_\perp^2 L_B}{xp^+}} \right]$$

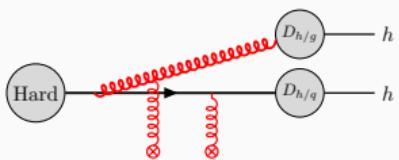


Multiple collision broadening  
and power corrections



△ Dynamical approach compared to nNNPDF [6]

# QGP effects in large and small systems



- HTL collisional energy loss of hard parton [7]:

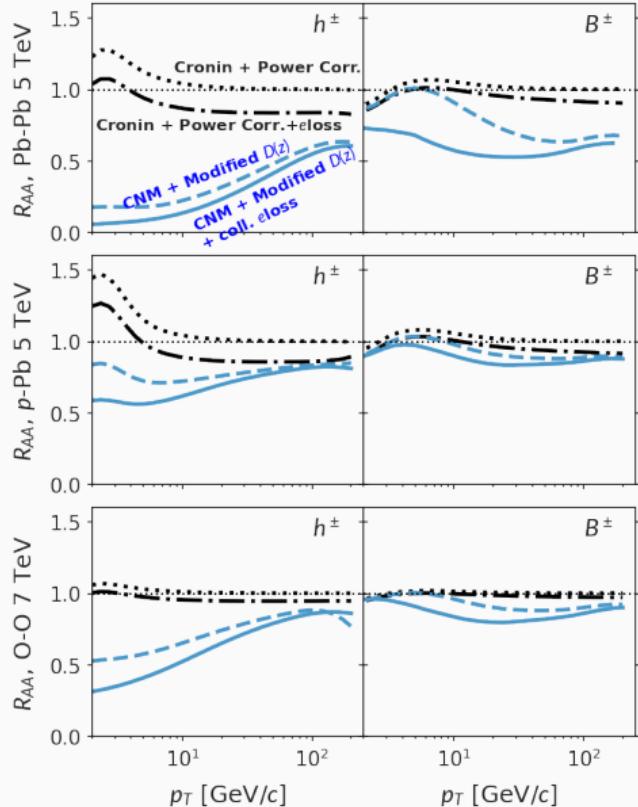
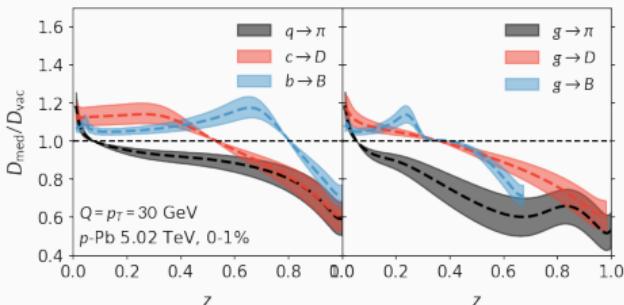
$$\frac{dE_{el}}{d\Delta z} = \frac{C_R}{4} \alpha_s(ET) m_D^2 \ln \left( \frac{ET}{m_D^2} \right) \left( \frac{1}{v} - \frac{1-v^2}{2v^2} \ln \frac{1+v}{1-v} \right)$$

More important in small systems

- QGP-modified fragmentation functions:

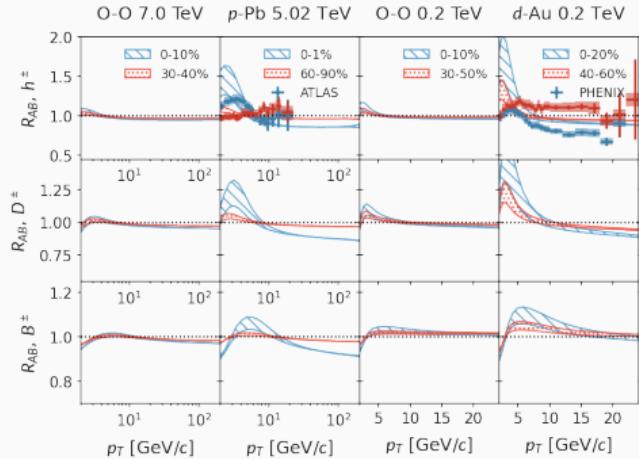
$$\partial_{\ln Q^2} D_{h/i} = [P_{ji}^{\text{vac}} + \Delta P_{ji}^{\text{med}}] \otimes D_{h/j}$$

↑ obtained in SCET<sub>G</sub> [8]

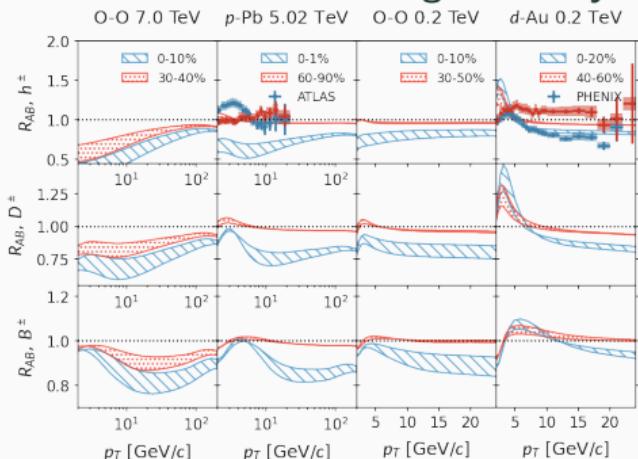


# Results

## Scenario I: no QGP formation, only cold nuclear matter effects



## Scenario II: QGP formed with a local-thermal color charged density



- $d$ -Au data: PHENIX [9];  $p$ -Pb data: ATLAS [10] (with  $\langle T_{pA} \rangle$  from the Glauber-Gribov model)
- Compare to  $d$ -Au &  $p$ -Pb, CNM effects in O-O are small → ideal to searching for QGP.
- QGP effects in  $d$ -Au at RHIC energy is small.
- In  $p$ -Pb at LHC energy, QGP effects with color density  $\propto T^3$  is inconsistent with data.

# Conclusion & references

## CNM effects

- are strong in small asymmetry systems such as  $p$ -Pb and  $d$ -Au;
- are much weaker in light-ion system O-O;
- can introduce centrality-dependent suppression at large  $p_T$  due to CNM energy loss.

## QGP effects?

- QGP effects with color-source density give by hydrodynamics ( $\propto T^3$ ) is inconsistent with  $R_{AA}^h$  in  $p$ -Pb.
- $\Delta E^{el}/\Delta E^{rad}$  increases in small systems: a change in the flavor separation of  $R_{AA}$ .
- Future scenario III: jet quenching with non-equilibrium color charged density.

## References:

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- [1] Z.-B. Kang, I. Vitev, H. Xing, Phys. Lett. B 718, 482–487 (2012).
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