The phi meson in nuclear matter from dilepton and K^+K^- decays

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Based on work done in collaboration with
Elena Bratkovskaya (Frankfurt/GSI),
Taesoo Song (Frankfurt)
and ongoing discussions with
Su Houng Lee (Yonsei U.)
Hiroyuki Sako (JAEA)

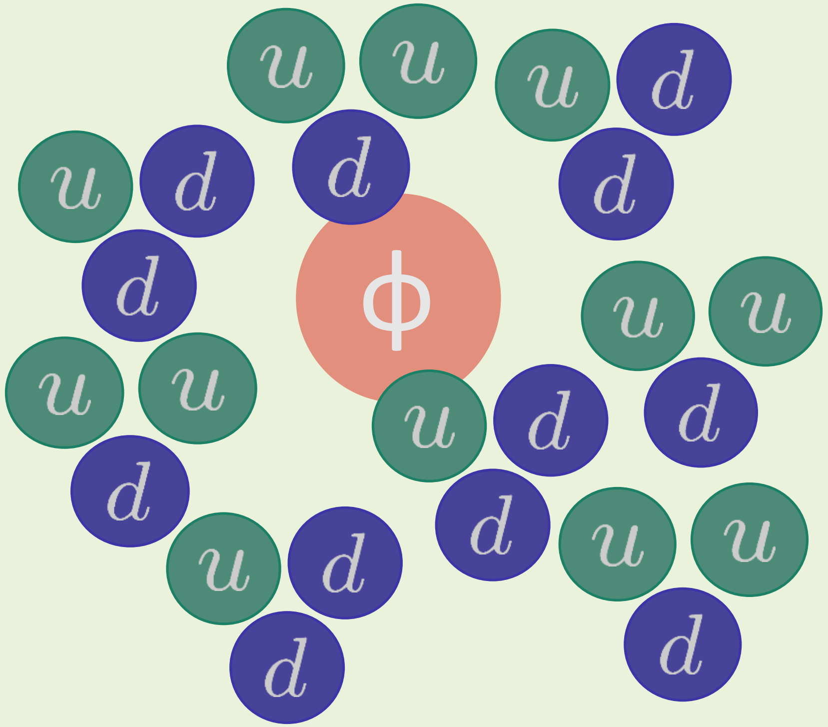
Contents

- ★ Introduction: φ meson in nuclear matter
- ★ Transport Simulations of pA reactions with density dependent vector meson spectral functions
- ★ Measuring the φ meson in nuclear matter: dilepton vs. K^+K^- channels
- ★ Considering electromagnetic and experimental rescattering effects on the dilepton spectra

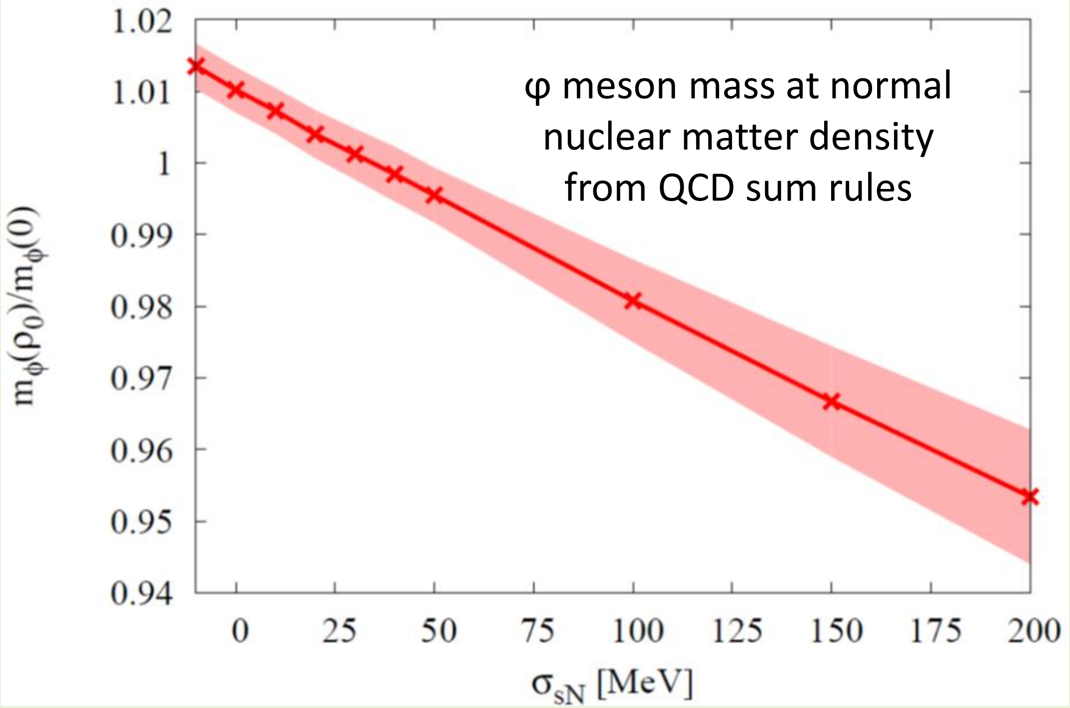
Why should we be interested?

The φ meson mass in nuclear matter probes the strange quark condensate at finite density!

$|\langle\bar{s}s\rangle_\rho|$ ↘ ⟶ m_ϕ ↘ ?



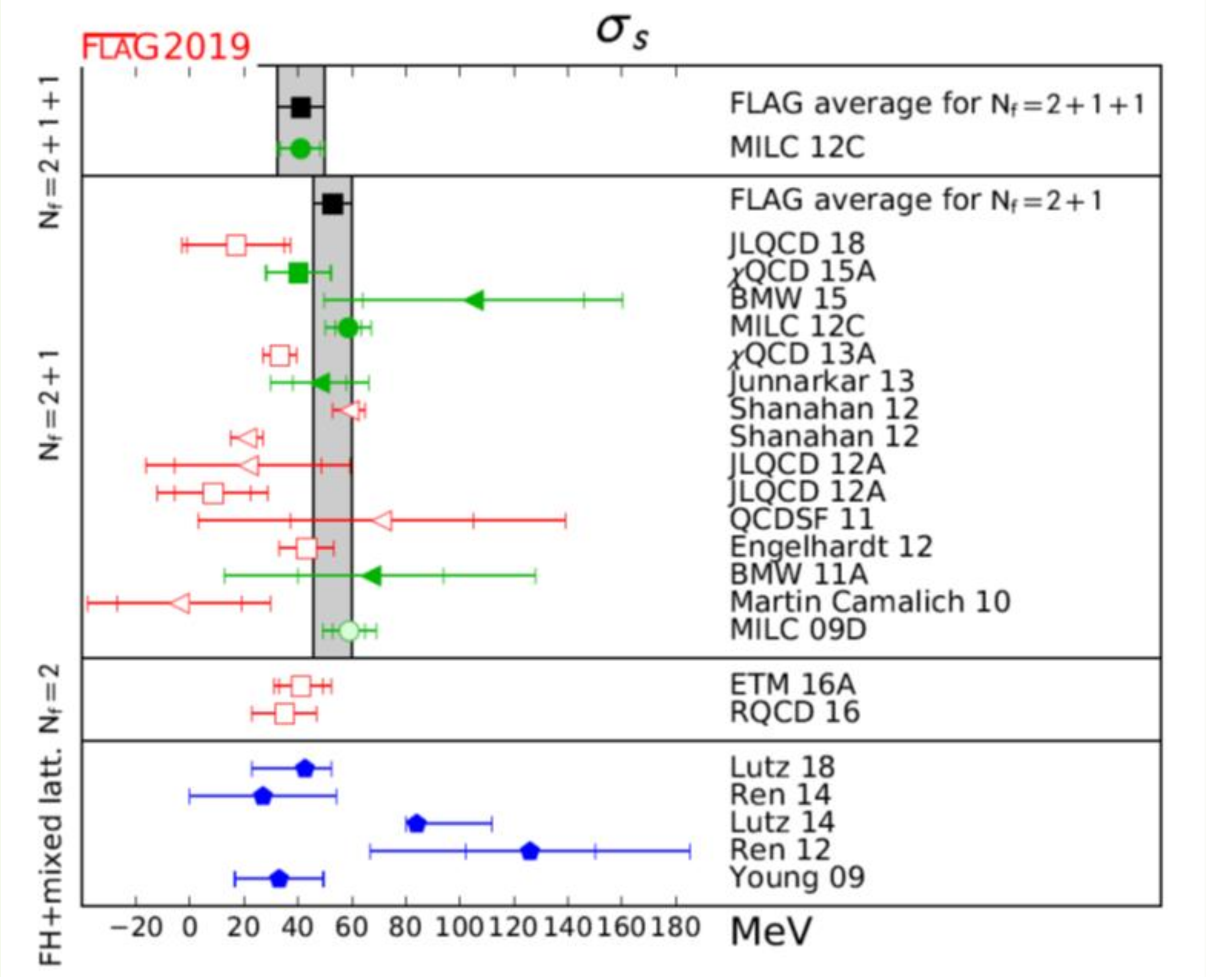
P. Gubler and K. Ohtani, Phys. Rev. D **90**, 094002 (2014).



$$|\langle\bar{s}s\rangle_\rho| = |\langle\bar{s}s\rangle_0| - \frac{\rho}{m_s}\sigma_{sN} + \dots$$

What does lattice QCD say about the strange sigma term?

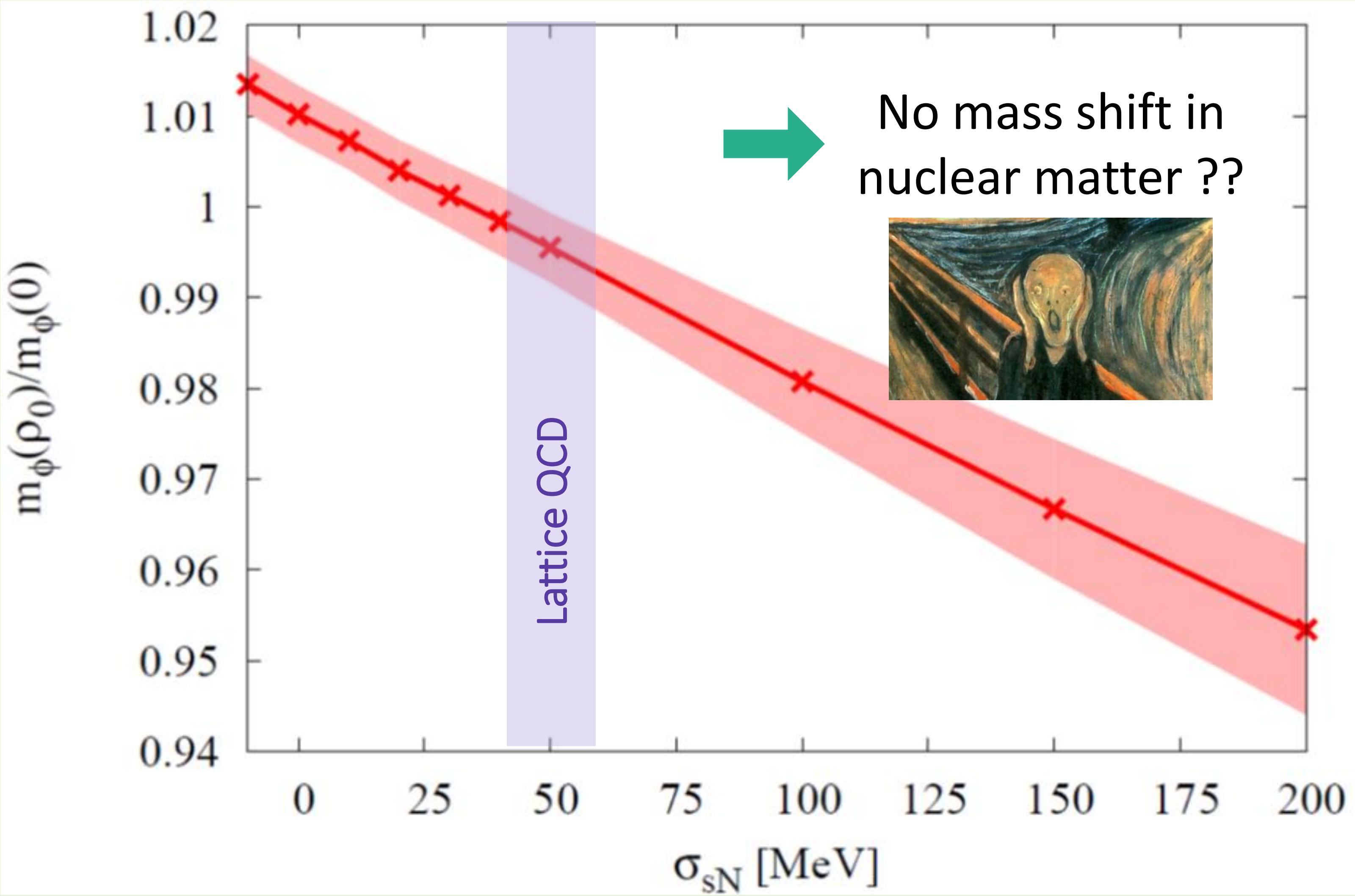
$$\sigma_{sN} = m_s \langle N|\bar{s}s|N\rangle$$



http://flag.unibe.ch/2019/

See also the most recent result of the BMW collaboration: Sz. Borsanyi et al., arXiv:2007.03319 [hep-lat].

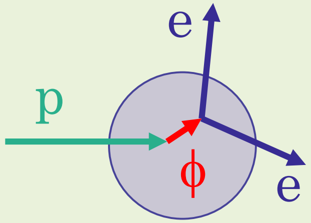
Combine QCD sum rules with lattice QCD



P. Gubler and K. Ohtani, Phys. Rev. D **90**, 094002 (2014).

Previous experimental results

KEK
E325



12 GeV
pA-reaction

Pole mass: $$\frac{m_\phi(\rho)}{m_\phi(0)} = 1 - k_1 \frac{\rho}{\rho_0}$$

k_1: 0.034 ± 0.007

Pole width: $$\frac{\Gamma_\phi(\rho)}{\Gamma_\phi(0)} = 1 + k_2 \frac{\rho}{\rho_0}$$

k_2: 2.6 ± 1.5

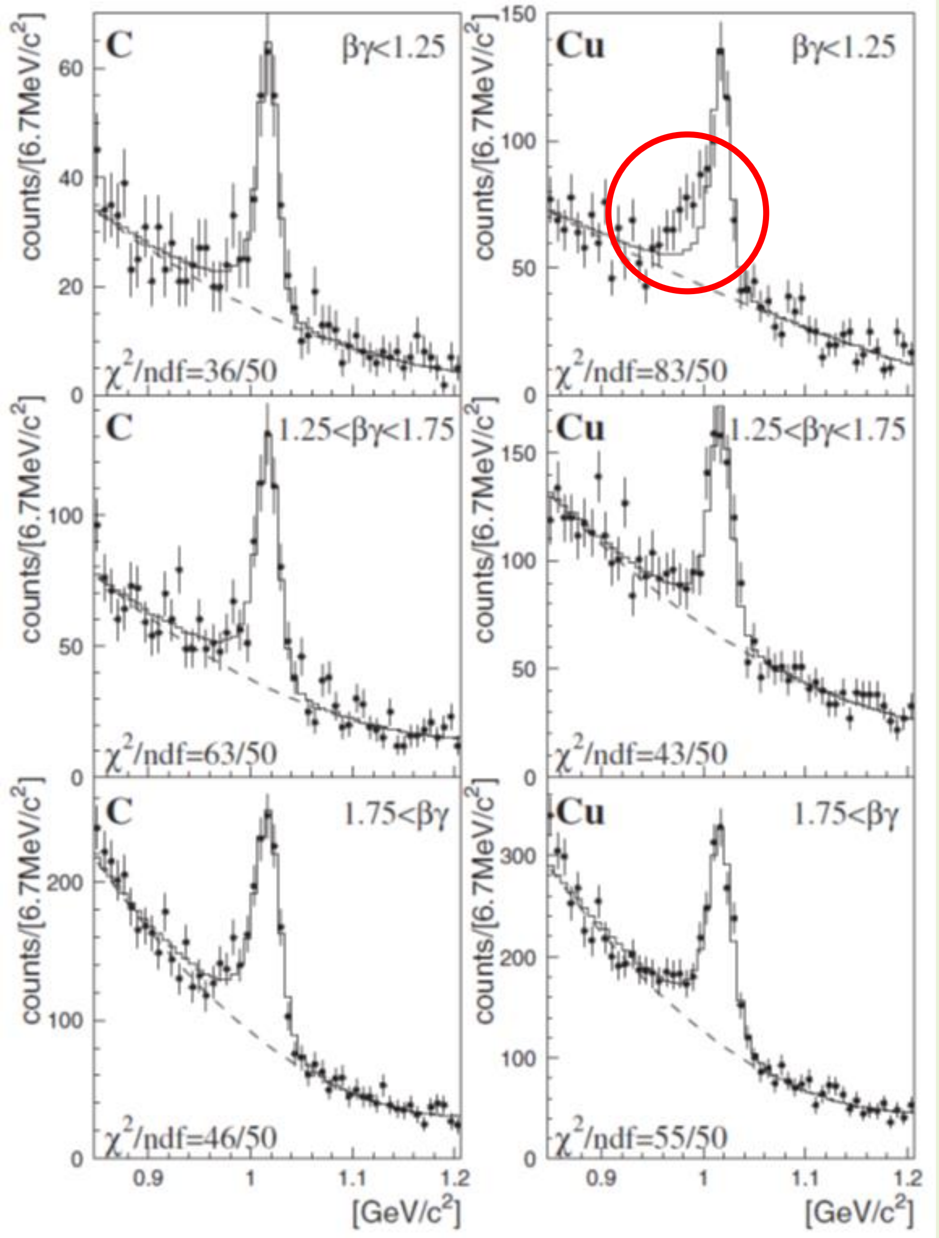
Measurement is being repeated with ~100x increased statistics at the J-PARC E16 experiment!

slow φs

intermediate φs

fast φs

$$\beta\gamma = \frac{|\vec{p}|}{m_\phi}$$

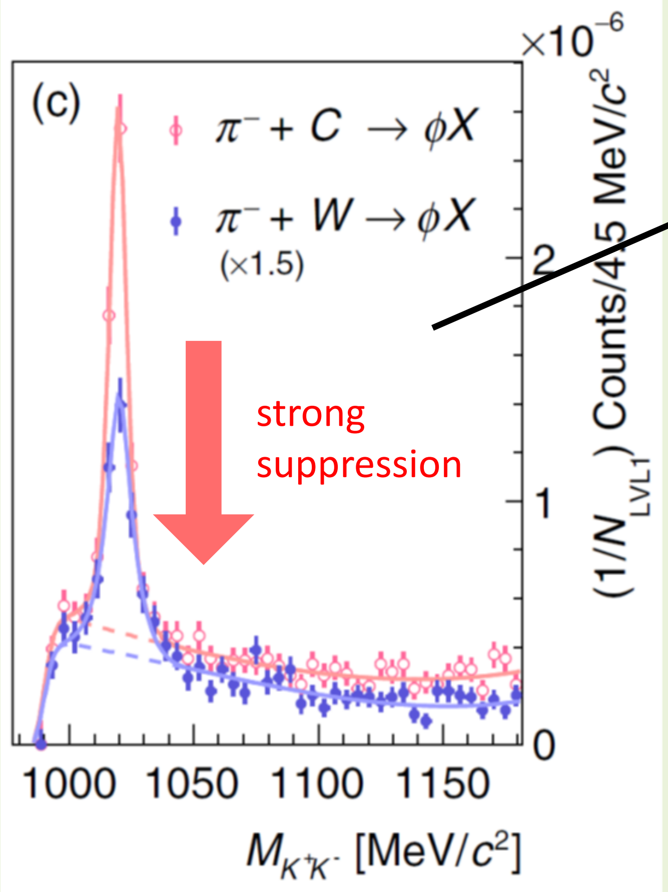


R. Muto et al. (E325 Collaboration), Phys. Rev. Lett. **98**, 042501 (2007).

More recent experiments

HADES: 1.7 GeV π^-A-reaction

K^+K^- - invariant mass spectrum



J. Adamczewski-Musch et al. (HADES Coll.),
Phys. Rev. Lett. **123**, 022002 (2019).

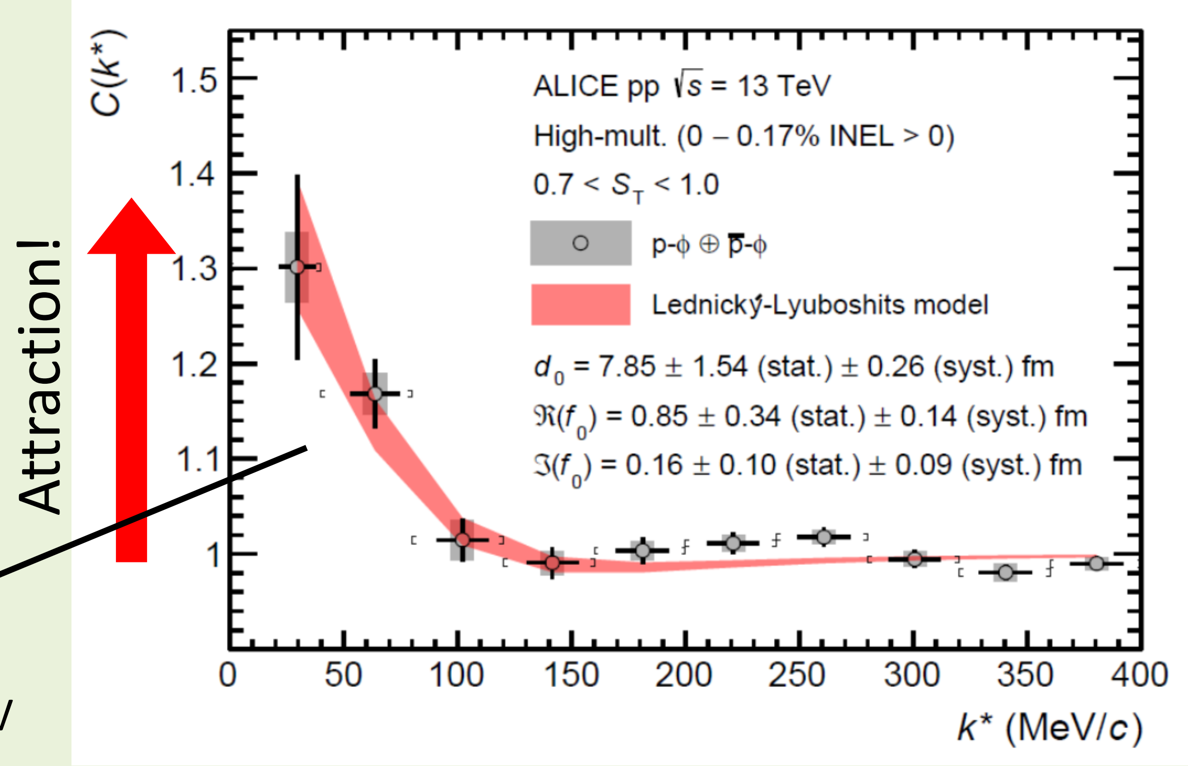
Broadening ?

Negative
mass shift

(-40 MeV ~ -70 MeV
at ρ_0??)

ALICE: pp

Measurement of φN correlation



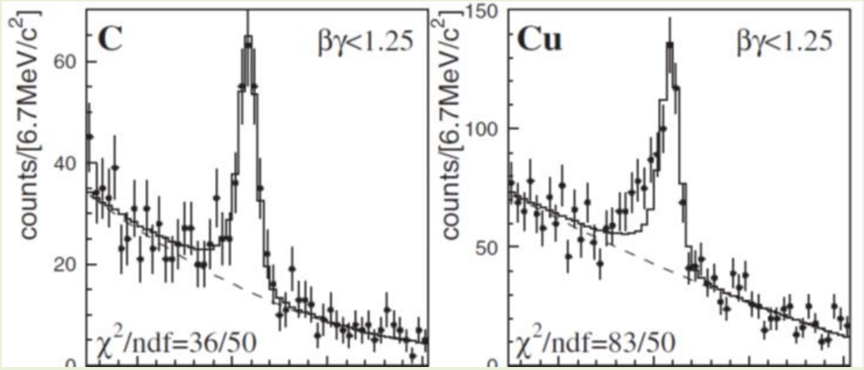
S. Acharya et al. (ALICE Coll.),
Phys. Rev. Lett. **127**, 172301 (2021).

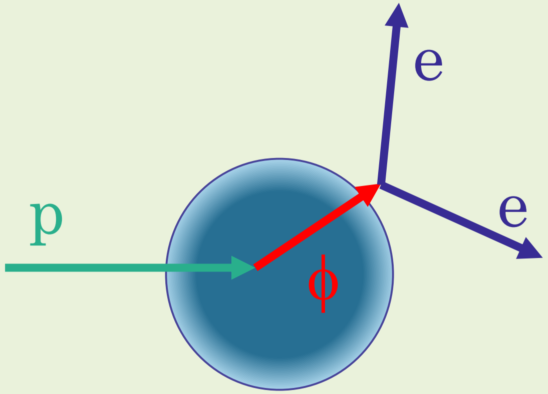
How compare theory with experiment?

Information useful for theory

- ★ Spectral function as a function of density
- ★ Mass at normal nuclear matter density
- ★ Decay width at normal nuclear matter density

Experimental data





Realistic simulation of pA reaction is needed!

Our tool: transport simulation HSD (Hadron String Dynamics)

E.L. Bratkovskaya and W. Cassing, Nucl. Phys. A **807**, 214 (2008).
W. Cassing and E.L. Bratkovskaya, Phys. Rev. C **78**, 034919 (2008).

Off-shell dynamics of vector mesons and kaons is included (dynamical modification of the mesonic spectral function during the simulated reaction)

off-shell terms

Testparticle approach:

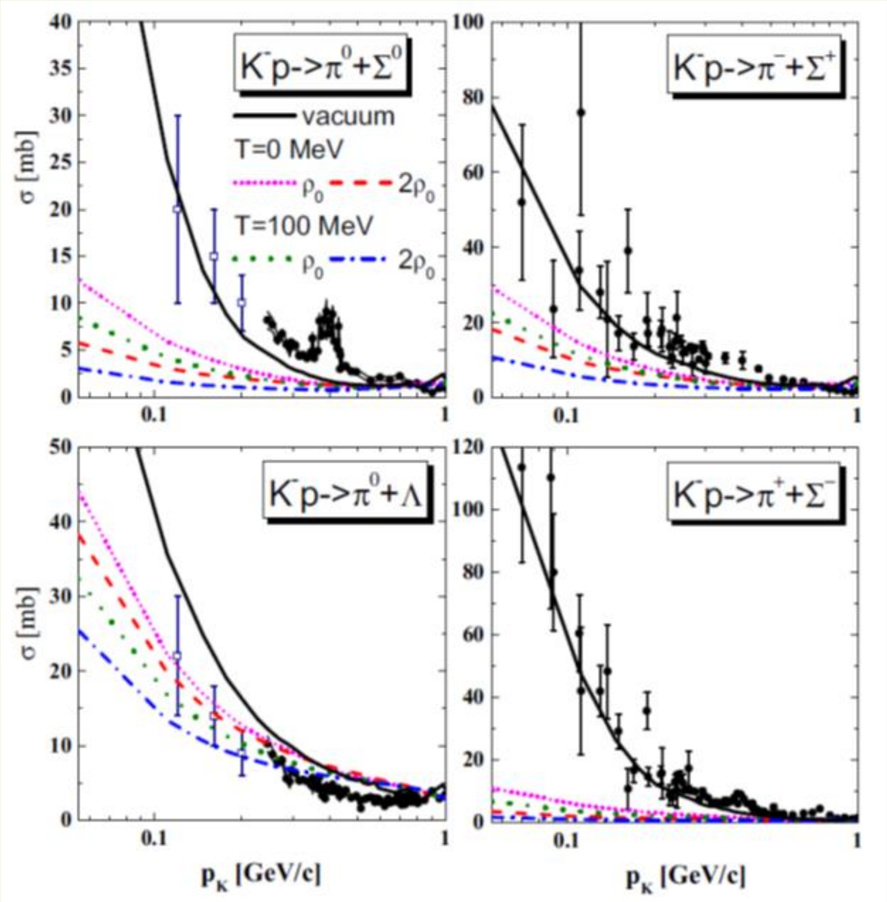
$$\frac{d\vec{X}_i}{dt} = \frac{1}{1-C_{(i)}} \frac{1}{2\varepsilon_i} \left[2\vec{P}_i + \vec{\nabla}_{P_i} \mathrm{Re}\,\Sigma^{\mathrm{ret}}_{(i)} + \frac{\varepsilon_i^2 - \vec{P}_i^2 - M_0^2 - \mathrm{Re}\,\Sigma^{\mathrm{ret}}_{(i)}}{\tilde{\Gamma}_{(i)}} \vec{\nabla}_{P_i} \tilde{\Gamma}_{(i)} \right],$$

$$\frac{d\vec{P}_i}{dt} = -\frac{1}{1-C_{(i)}} \frac{1}{2\varepsilon_i} \left[\vec{\nabla}_{X_i} \mathrm{Re}\,\Sigma^{\mathrm{ret}}_{i} + \frac{\varepsilon_i^2 - \vec{P}_i^2 - M_0^2 - \mathrm{Re}\,\Sigma^{\mathrm{ret}}_{(i)}}{\tilde{\Gamma}_{(i)}} \vec{\nabla}_{X_i} \tilde{\Gamma}_{(i)} \right],$$

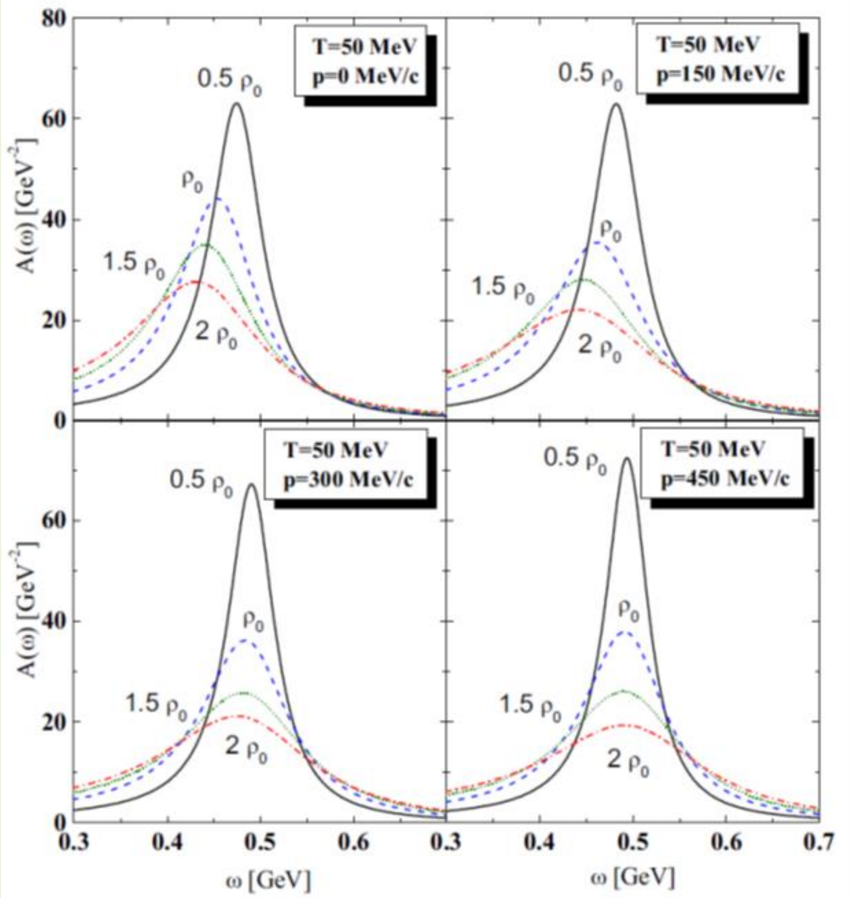
$$\frac{d\varepsilon_i}{dt} = \frac{1}{1-C_{(i)}} \frac{1}{2\varepsilon_i} \left[\frac{\partial\, \mathrm{Re}\,\Sigma^{\mathrm{ret}}_{(i)}}{\partial t} + \frac{\varepsilon_i^2 - \vec{P}_i^2 - M_0^2 - \mathrm{Re}\,\Sigma^{\mathrm{ret}}_{(i)}}{\tilde{\Gamma}_{(i)}} \frac{\partial \tilde{\Gamma}_{(i)}}{\partial t} \right],$$

Treatment of KN-interactions

Density dependent cross sections based on the chiral unitary model
(including coupled channels and s-/p-wave of $\bar{K}N$ interactions)



Vacuum and density dependent $\bar{K}N$ cross sections



Density dependent $\bar{K}$ spectral functions

T. Song et al., Phys. Rev. C **103**, 044901 (2021).

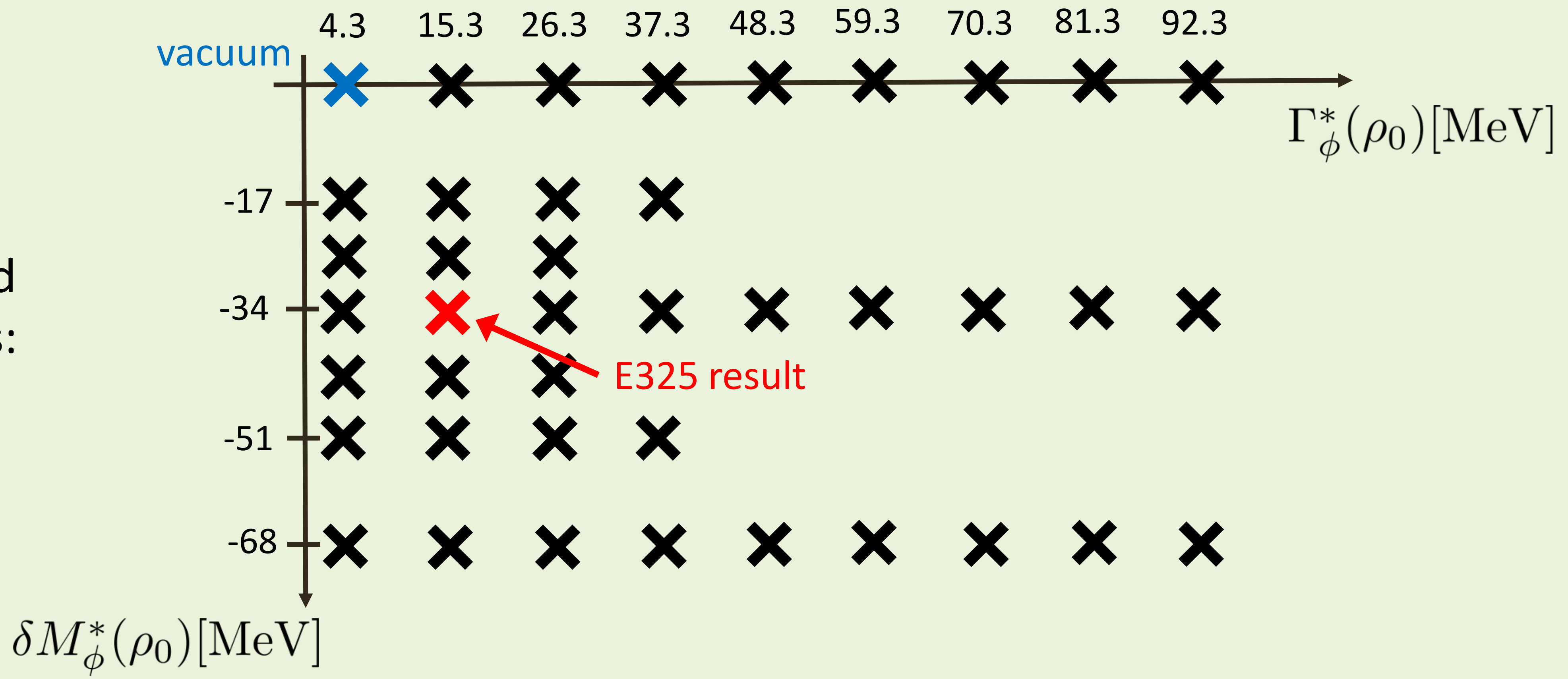
See talk by Taesoo Song

Advantage: vector meson spectra can be chosen freely

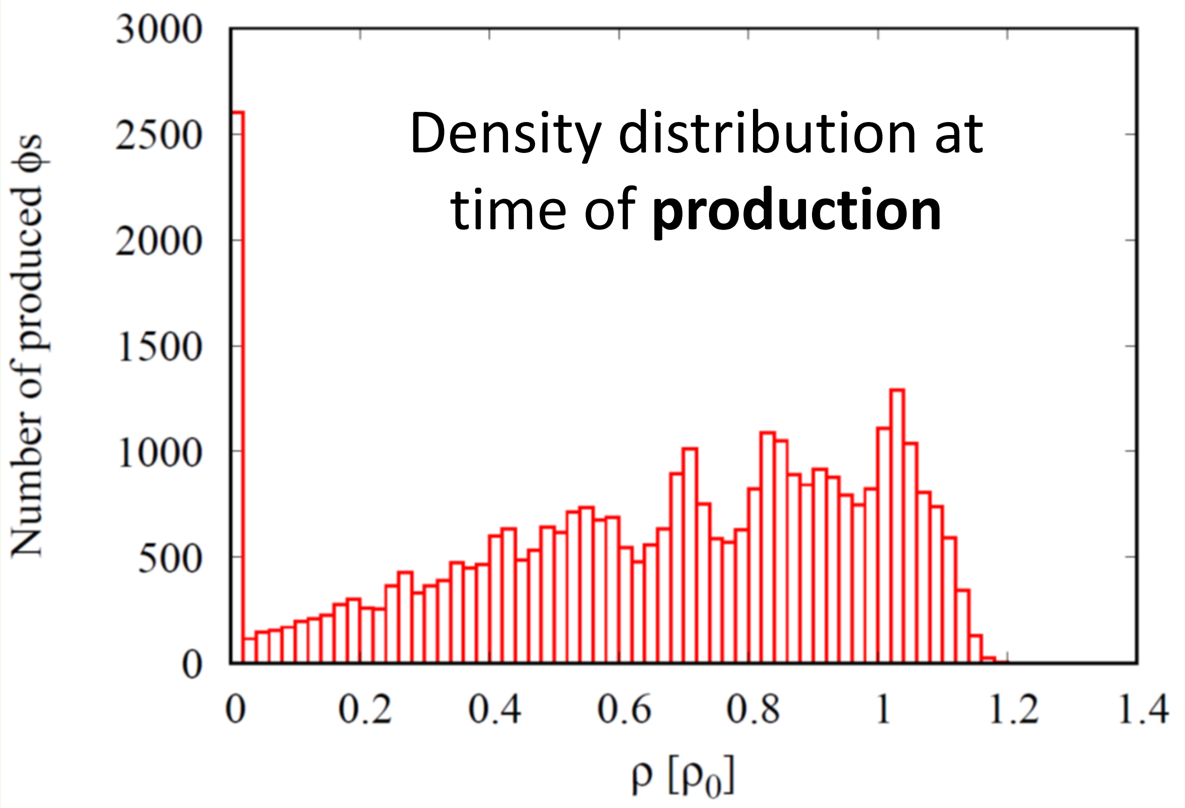
Our choice: a Breit-Wigner with density dependent mass and width

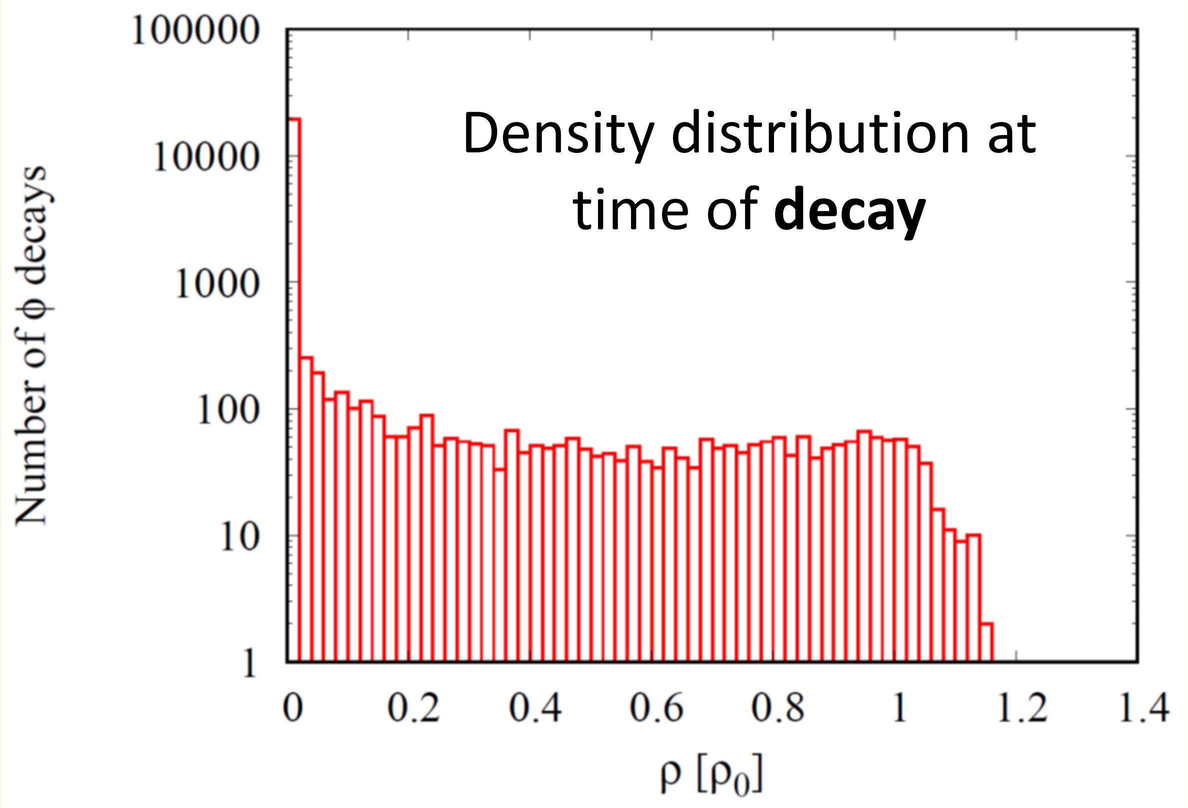
$$A_\phi(M,\rho) = C\frac{2}{\pi}\frac{M^2\Gamma^*_\phi(M,\rho)}{[M^2-M^{*2}_\phi(\rho)]^2+M^2\Gamma^{*2}_\phi(M,\rho)} \quad \text{with} \quad \begin{cases} M^*_\phi(\rho) = M^{\text{vac}}_\phi\left(1-\alpha^\phi\frac{\rho}{\rho_0}\right), \\ \Gamma^*_\phi(M,\rho) = \Gamma^{\text{vac}}_\phi + \alpha^\phi_{\text{coll}}\frac{\rho}{\rho_0} \end{cases}$$

Simulated scenarios:



What density does the φ feel in the reaction (p+Cu at 12 GeV)?

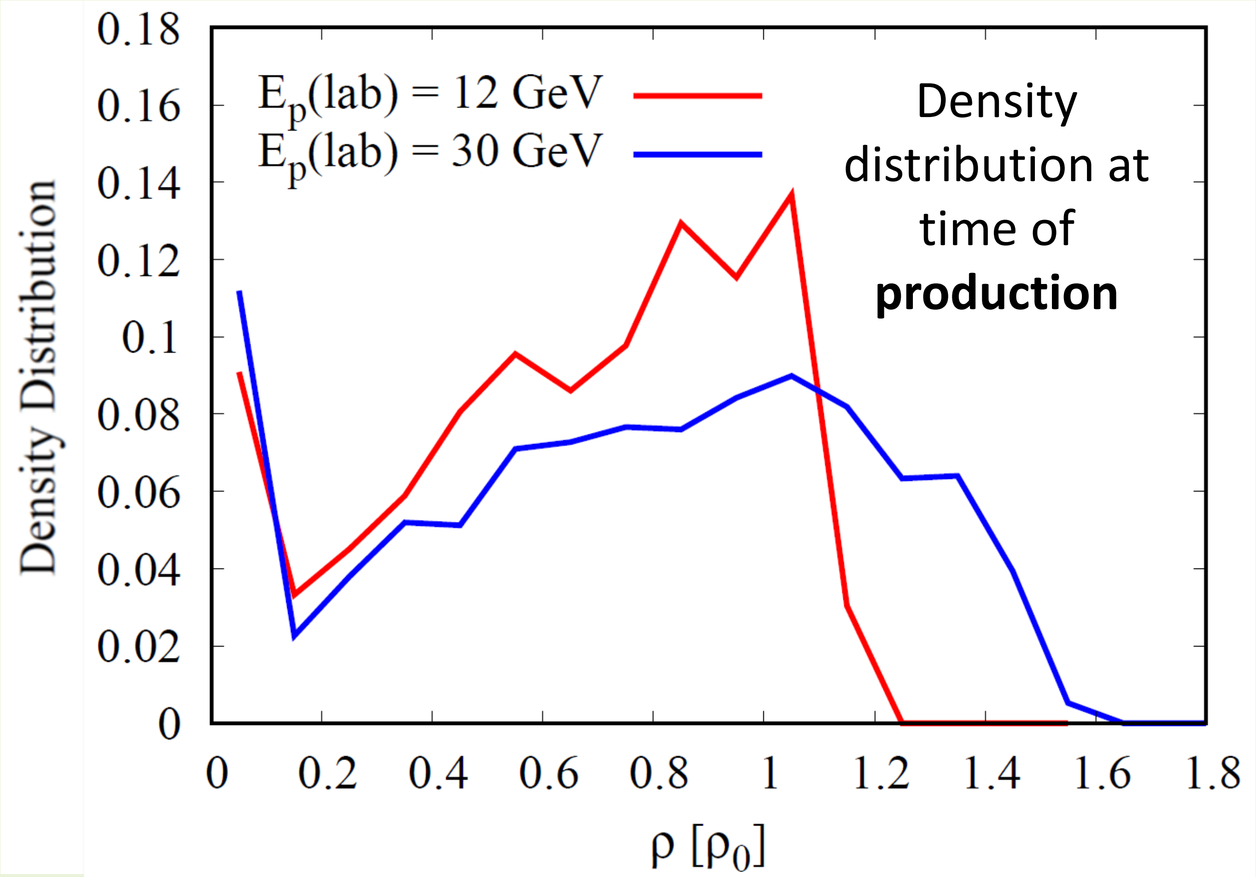


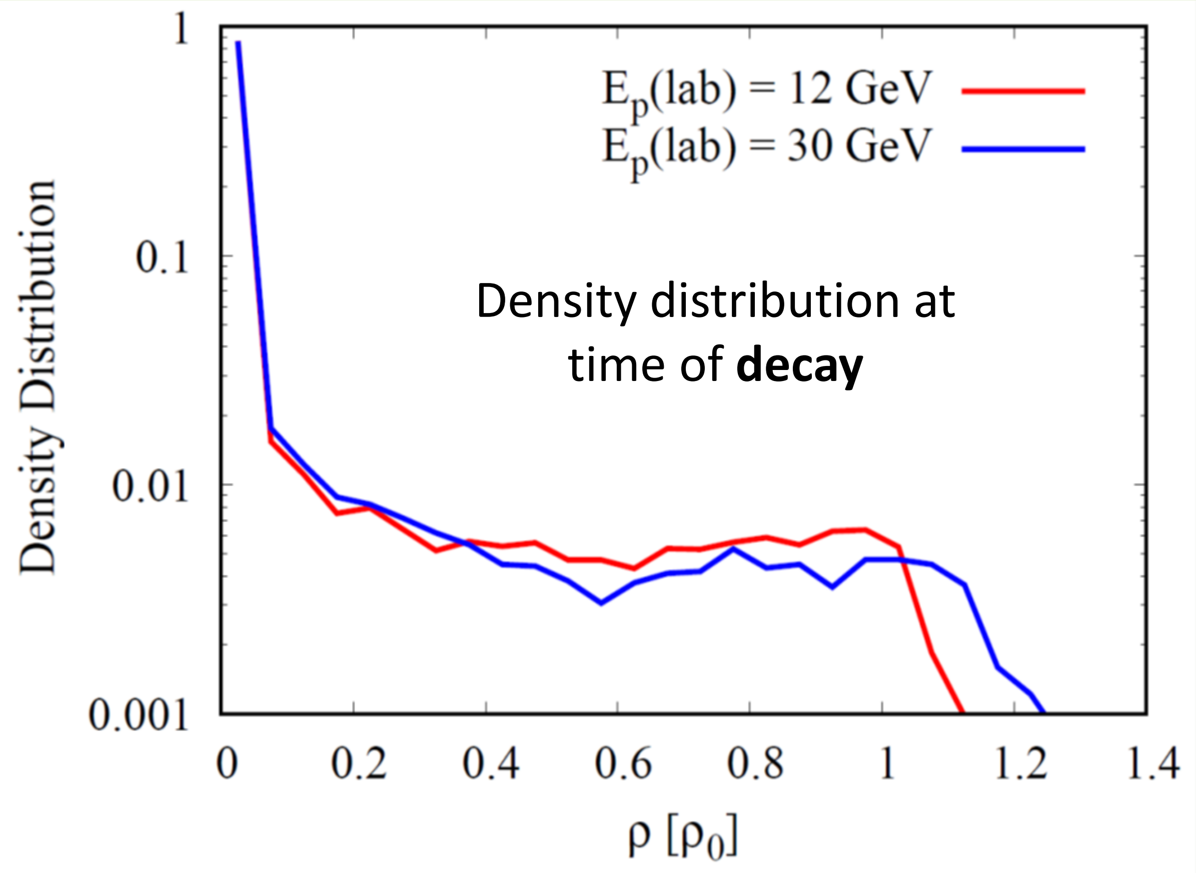


Majority of φ mesons are produced at densities around ρ_0

Majority of φ mesons decay in free space (note the log-scale!)

What density does the φ feel in different pA (p+Cu) reactions?

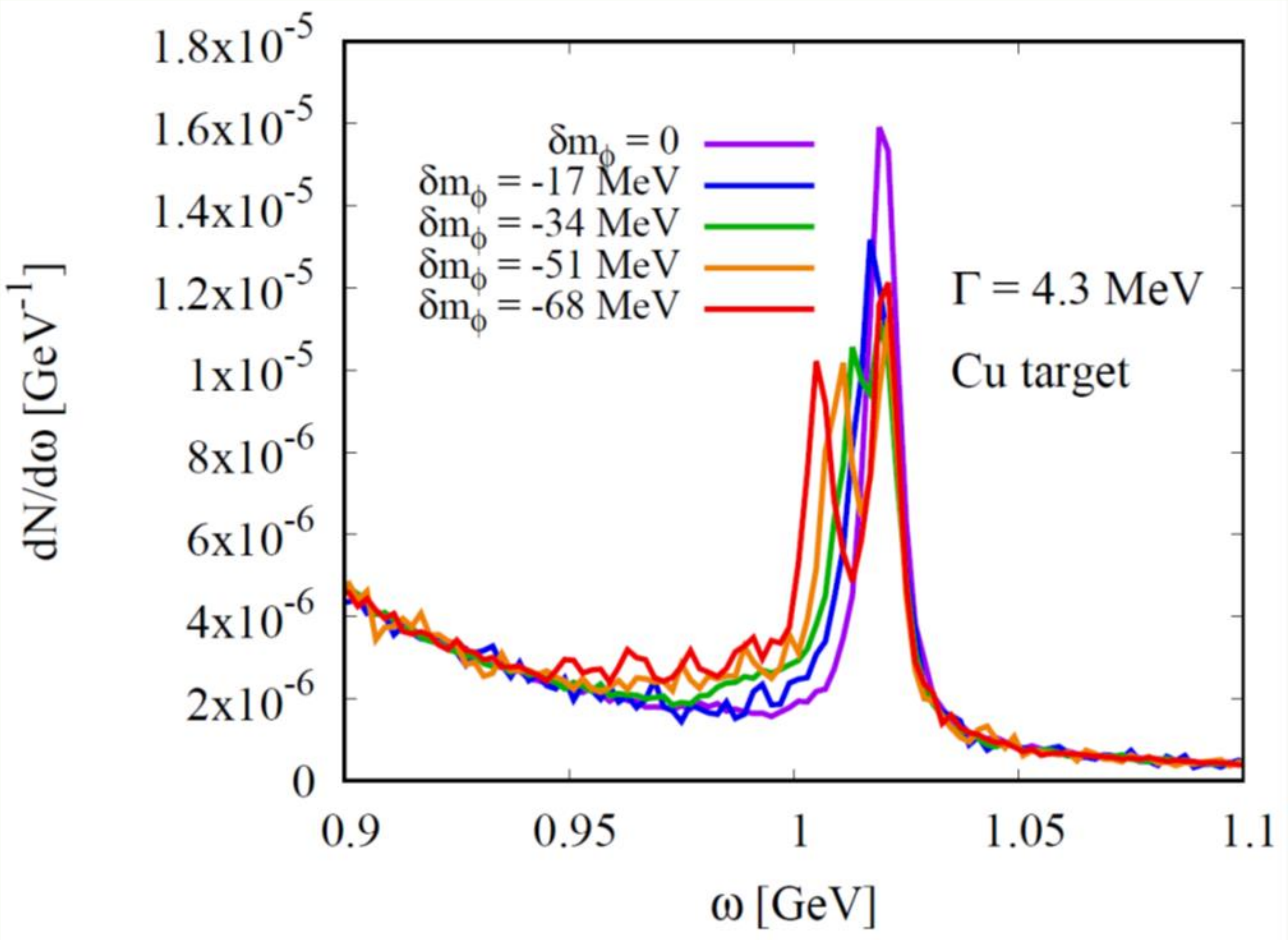




Larger densities are reached for larger incoming proton energy

Majority of φ mesons decay in free space (note the log-scale!)

The dilepton spectrum in the φ meson region



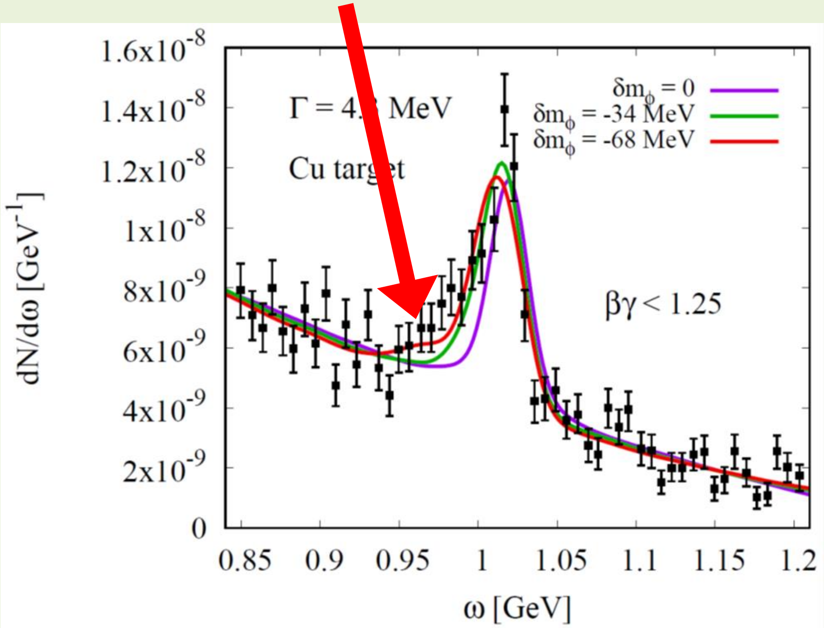
p + Cu at 12 GeV

No acceptance corrections!

No finite resolution effects!

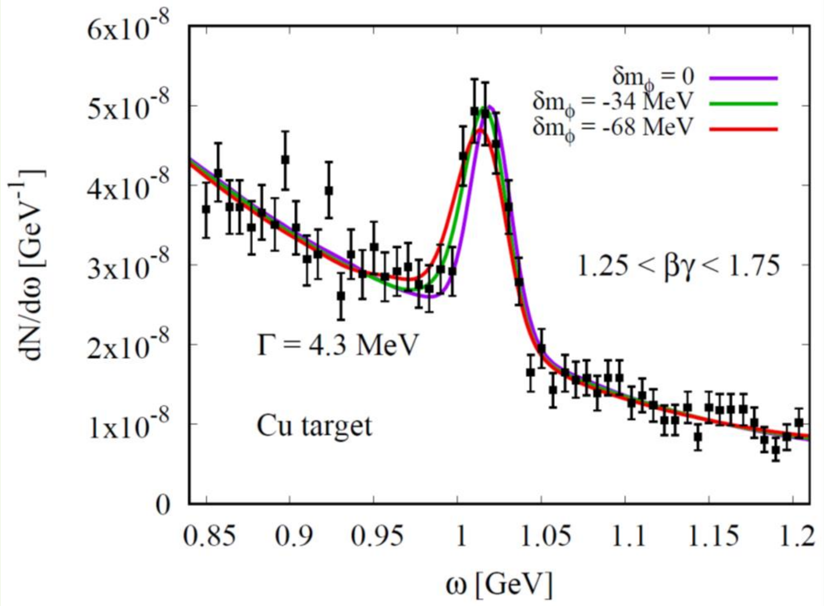
Fits to experimental Copper target data (E325)

A significant negative mass shift is needed to reproduce the slow φ data



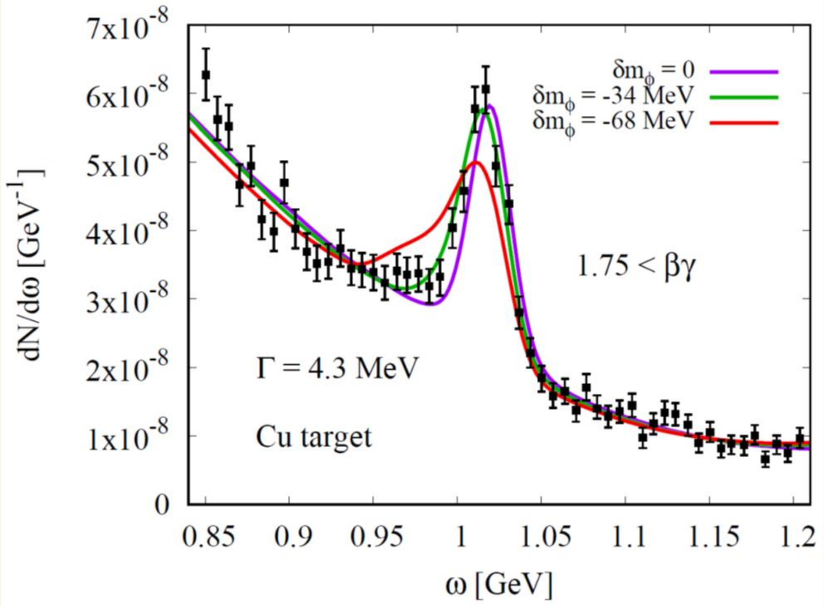
slow φs

Favors relatively large negative mass shift



intermediate φs

No strong constraints for any modification scenario

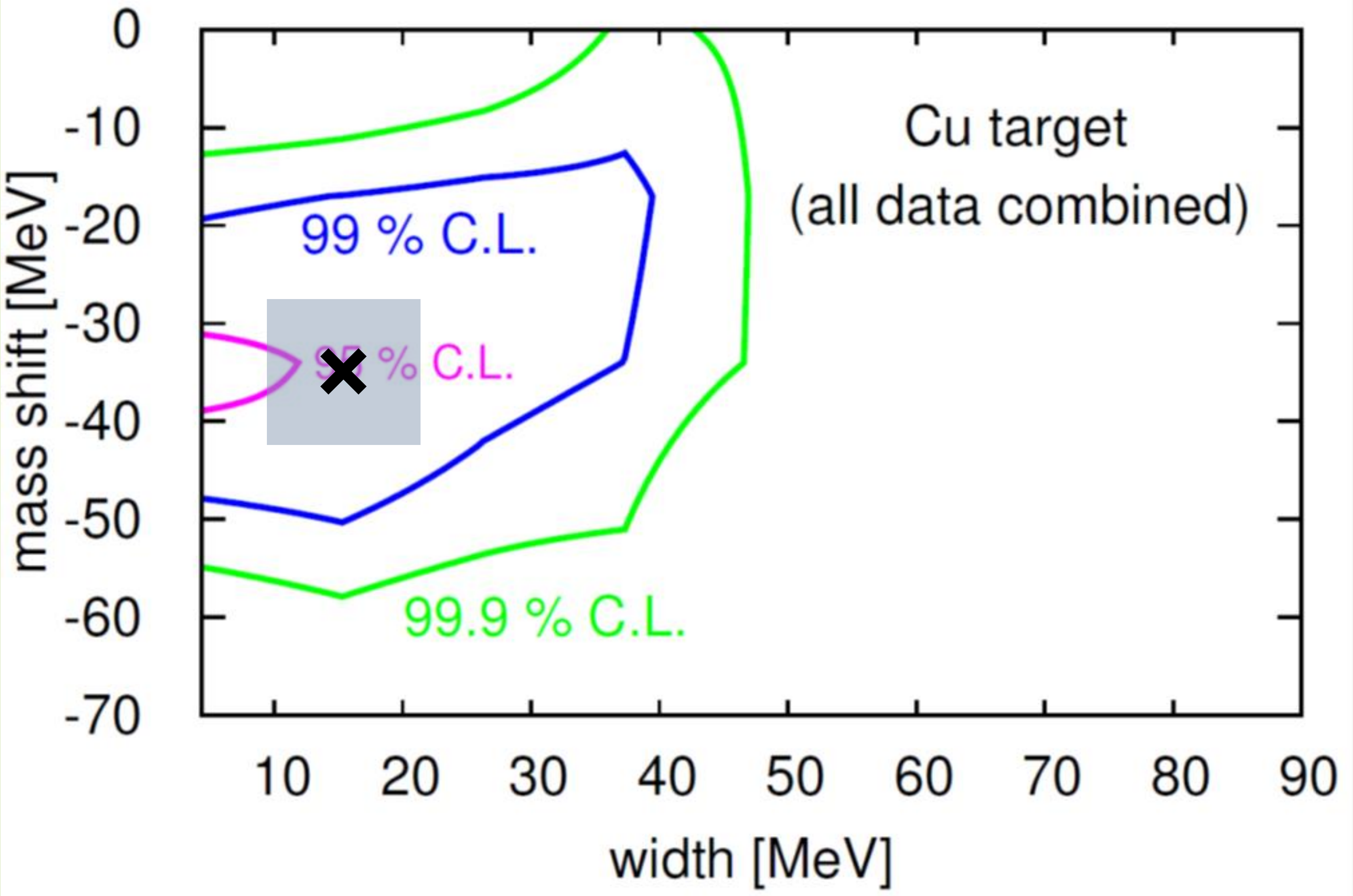


fast φs

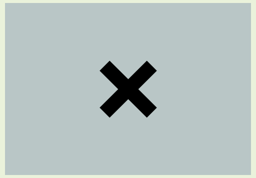
Favors small mass shift

Fits to experimental Copper target data (KEK E325)

Confidence levels of combined Copper data



Conclusion of the E325 Collaboration

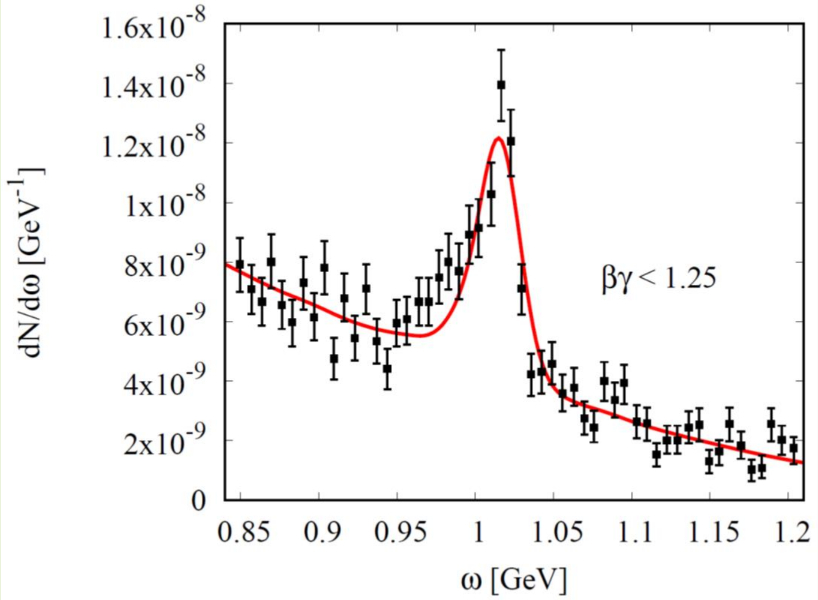


Best fit to E325 data

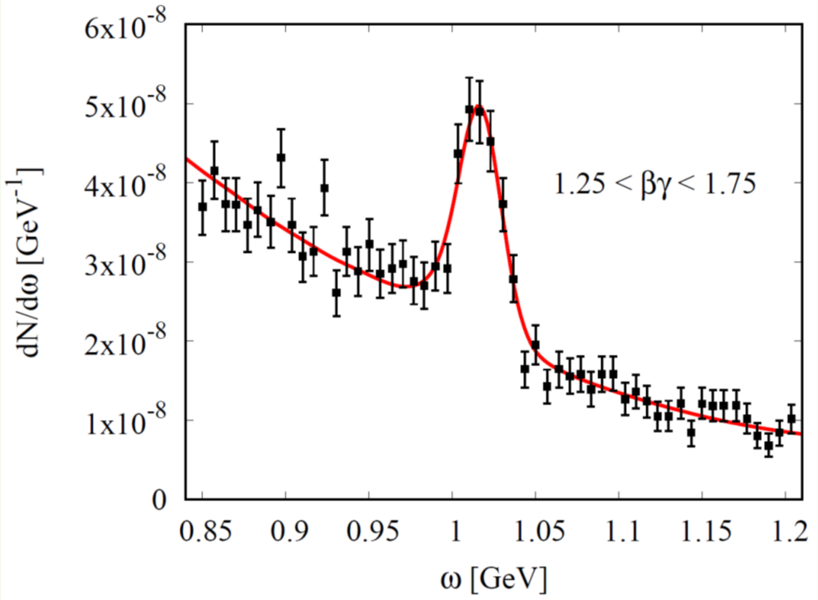
(p + Cu at 12 GeV)

vacuum value

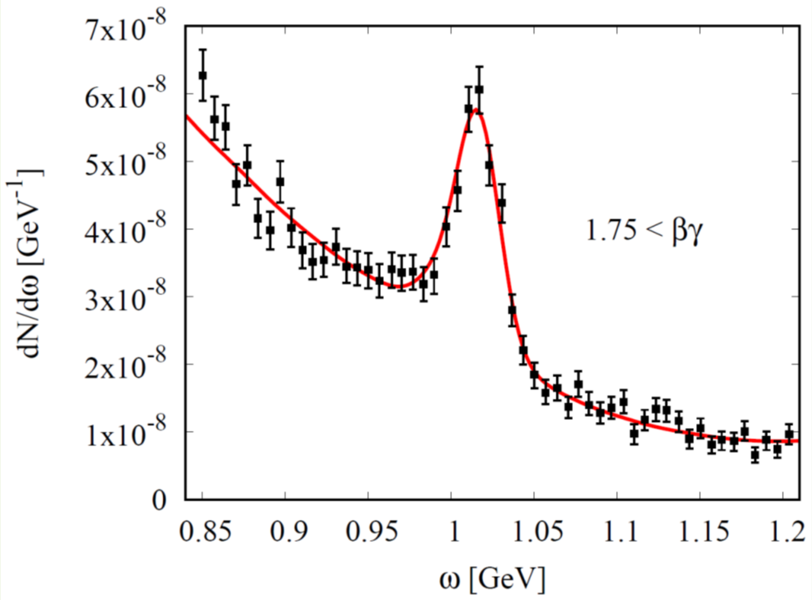
$$\delta m_\phi(\rho_0) = -34\,\mathrm{MeV} \qquad \Gamma(\rho_0) = 4.3\,\mathrm{MeV}$$



slow φs



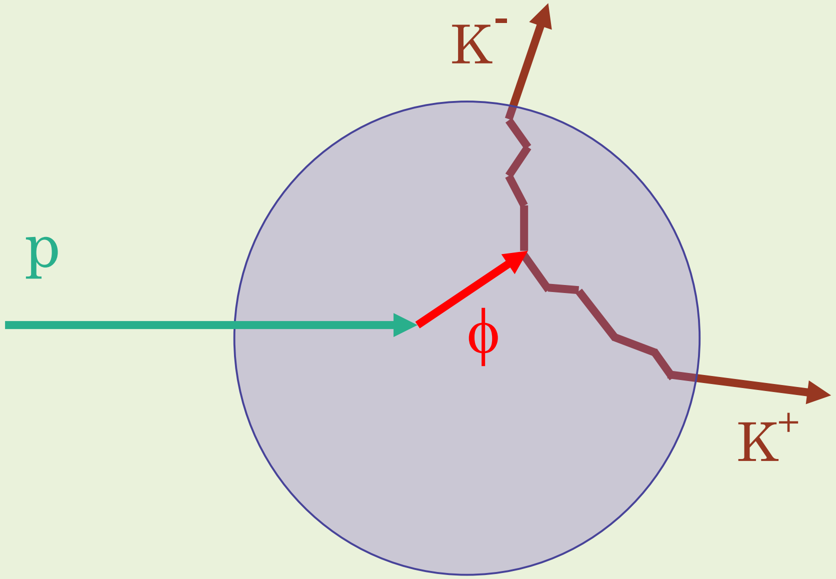
intermediate φs

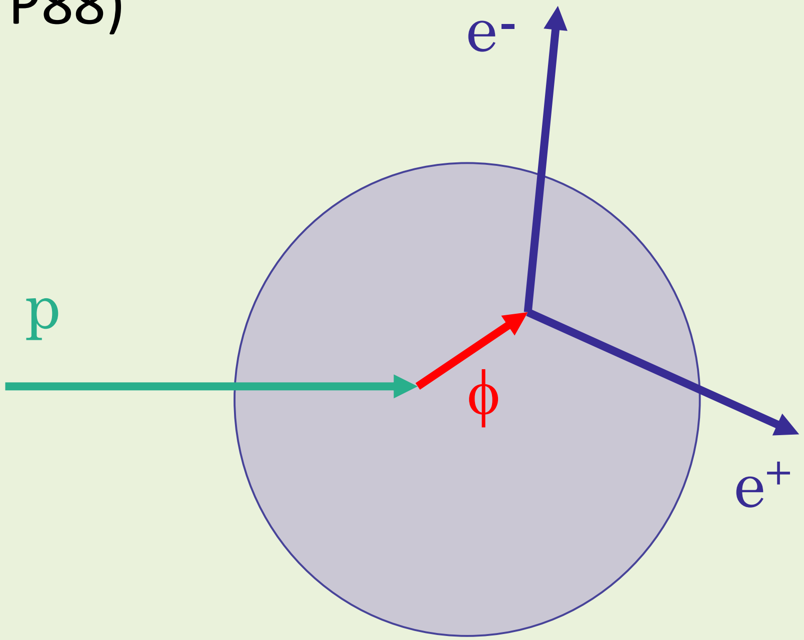


fast φs

What about the K⁺K⁻ decay channel?

(new J-PARC proposal P88)





Kaons feel the strong interaction → Distorted in-medium φ meson signal

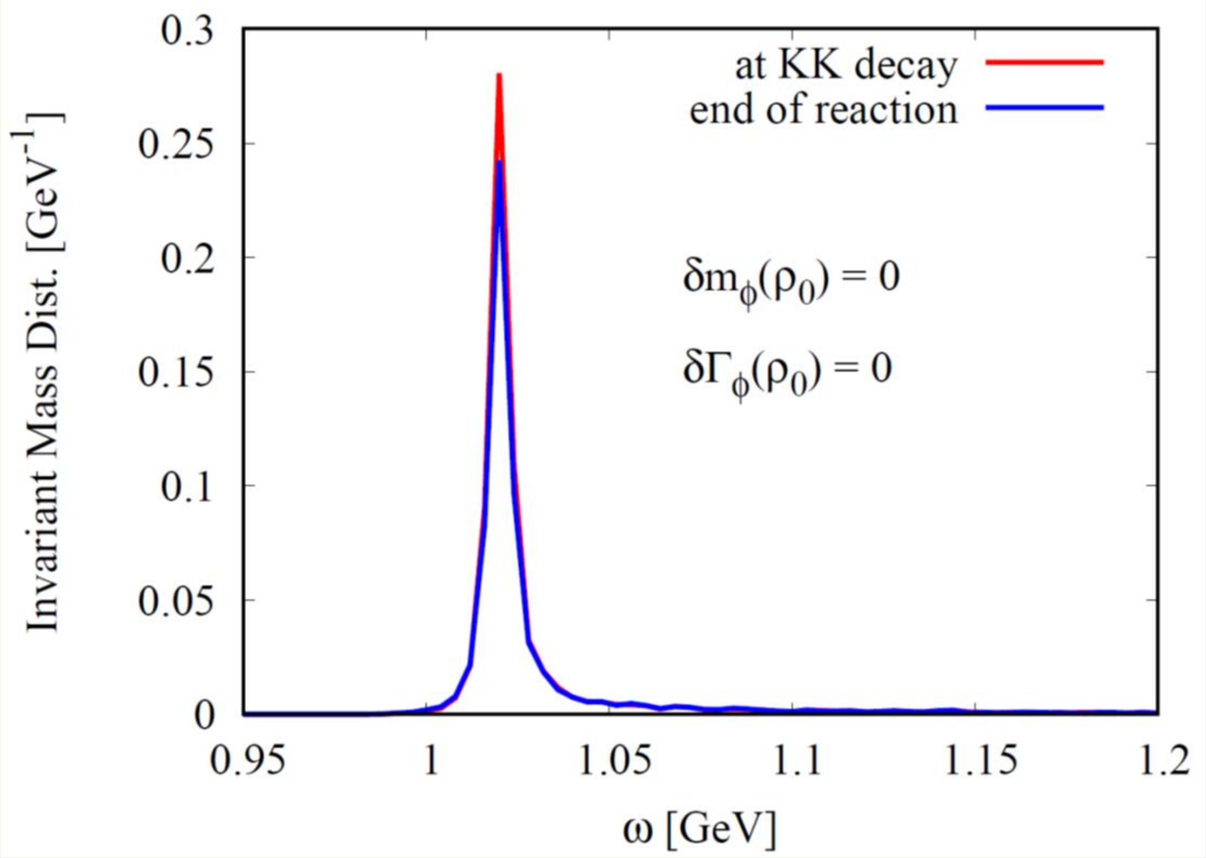
Large branching ratio → Good statistics

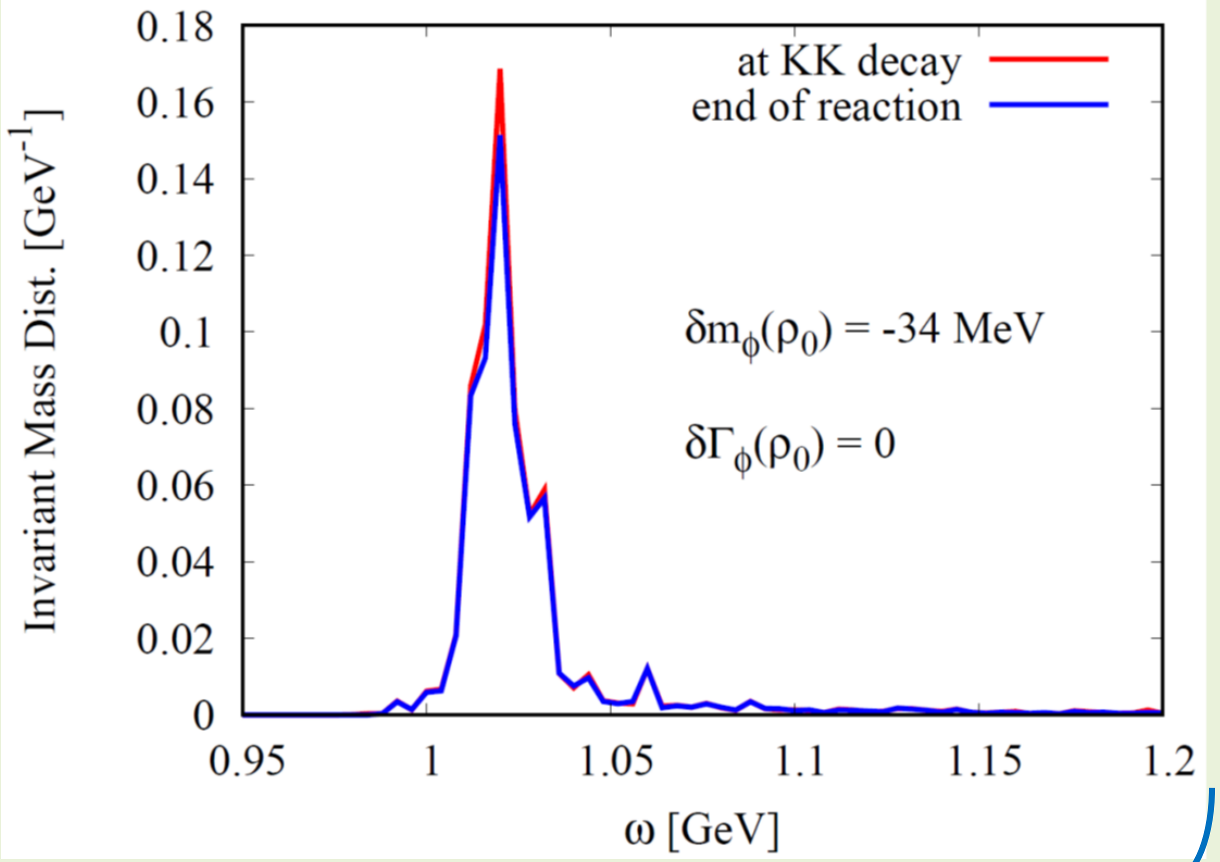
Kaons do not feel the strong interaction → Clear in-medium φ meson signal

Small branching ratio → Bad statistics

Distortion of the in-medium φ meson signal in the K^+K^- channel

(p + Cu at 30 GeV)

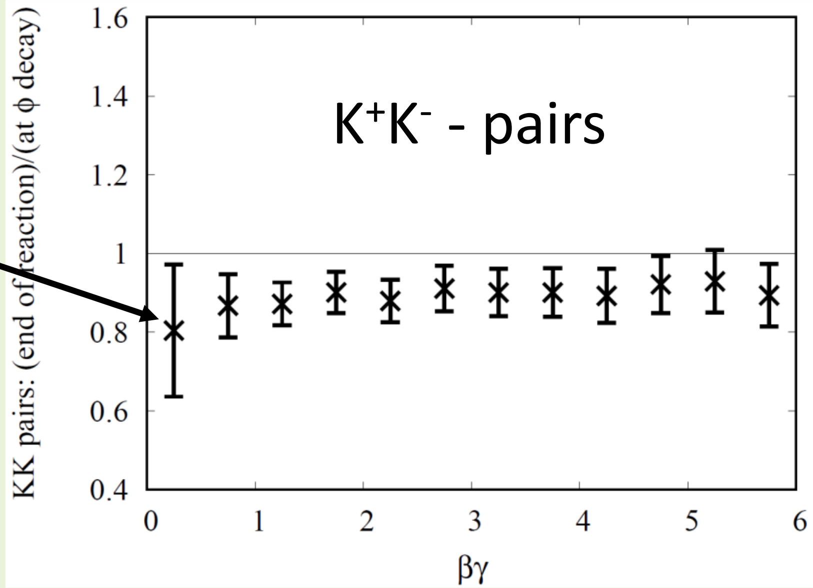




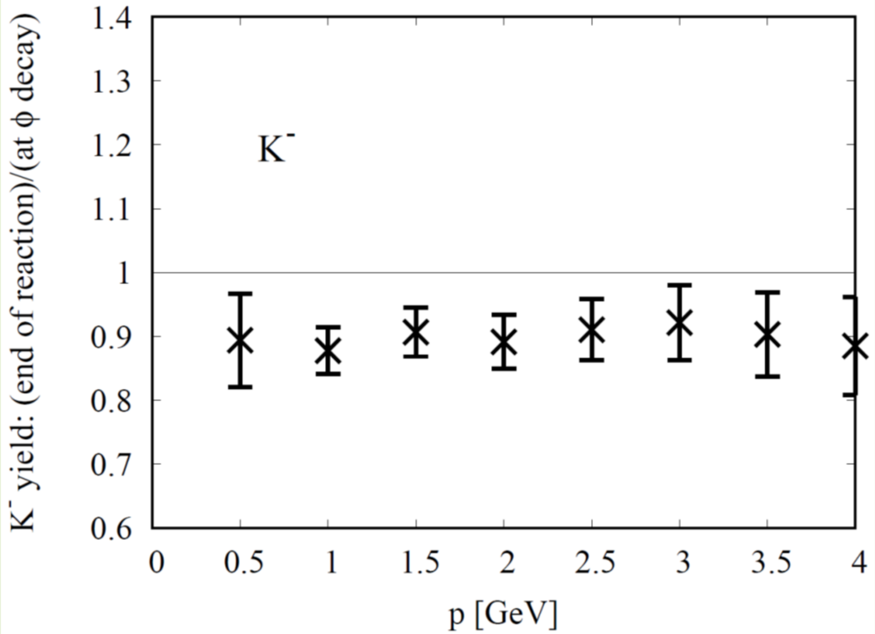
Small distortion effect from the strong KN interaction !?

Absorption of kaons in nuclear matter

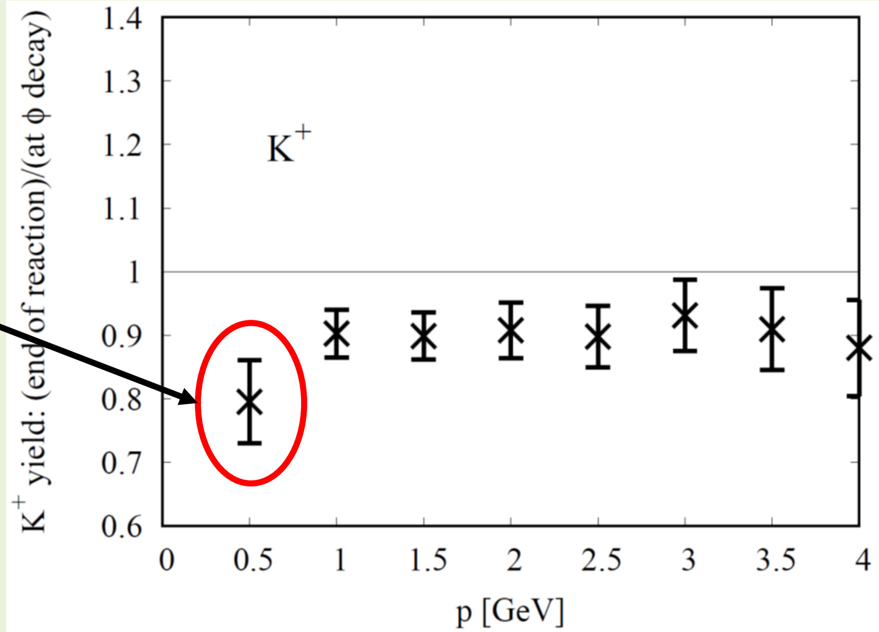
Stronger absorption for lower momenta



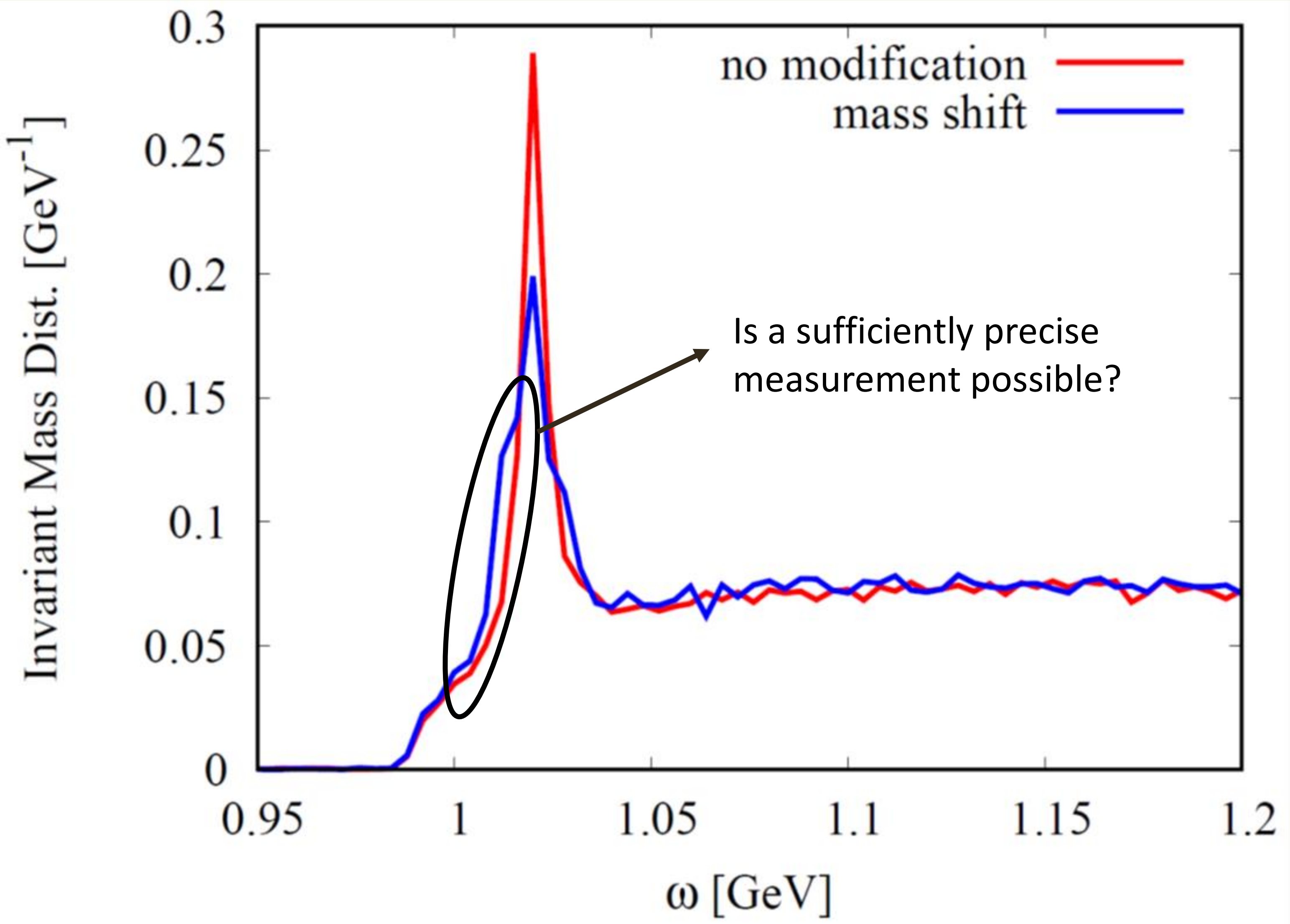
(p + Cu at 30 GeV)



Suppression due to repulsive K^+N interaction??



Expected K^+K^- invariant mass spectrum (incl. background)



p + Cu at 30 GeV

No acceptance corrections!

No finite resolution effects!

What about other effects?

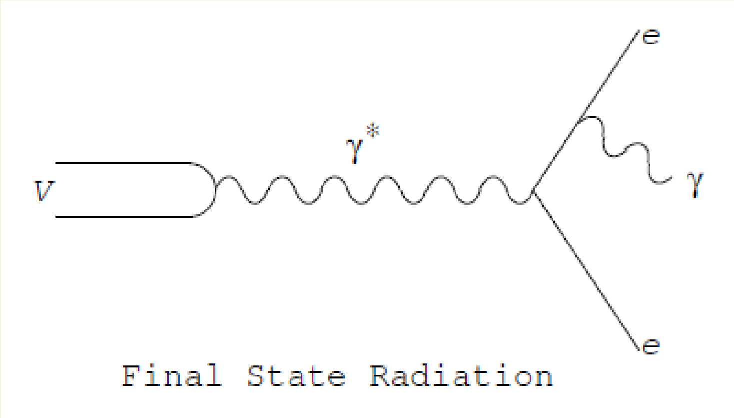
- ★ Electromagnetic corrections to the dilepton spectrum?
- ★ Rescattering effects of dileptons on experimental environment?

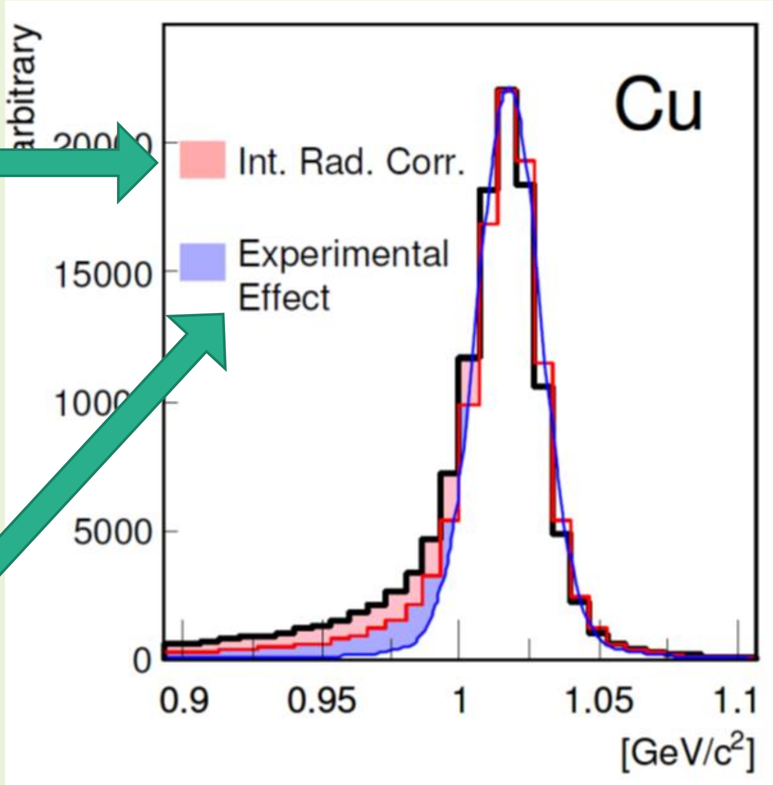
Detailed information about the experiment is needed

Evidence for in-medium modification of the ϕ meson at normal nuclear density

Ryotaro Muto

Old PhD Thesis from the early 2000s





Rescattering effect

Both effects contribute roughly equally

Evidence for in-medium modification of the ϕ meson at normal nuclear density

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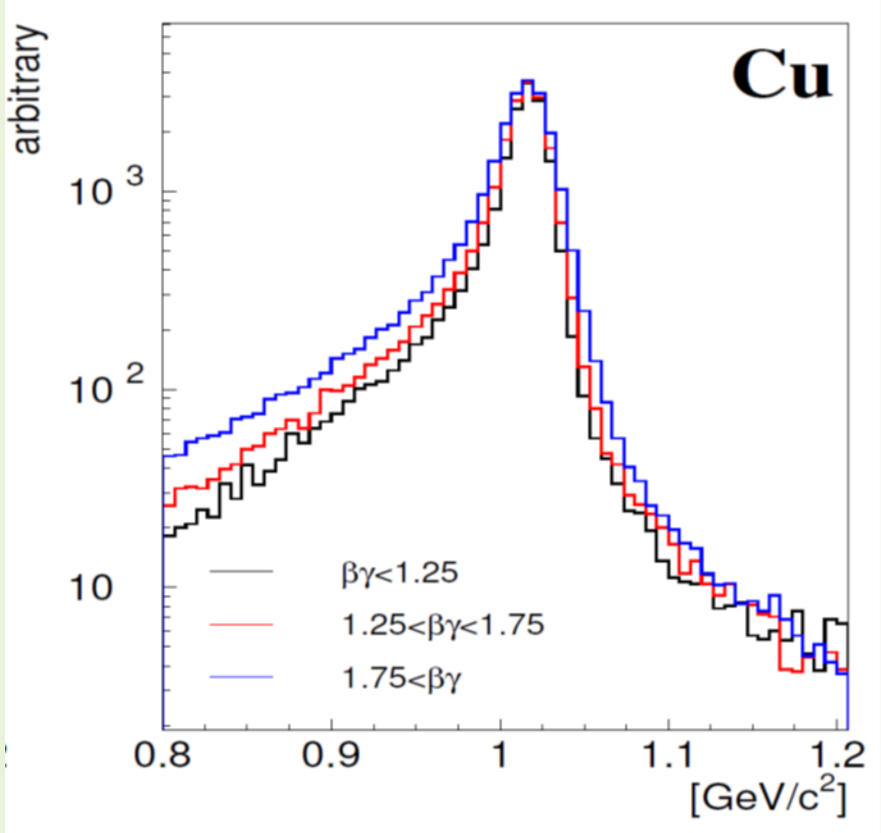
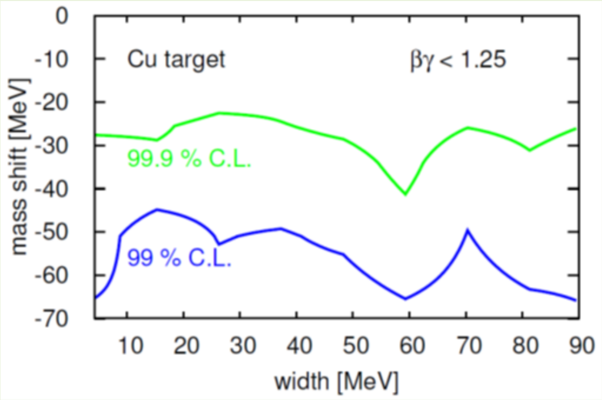


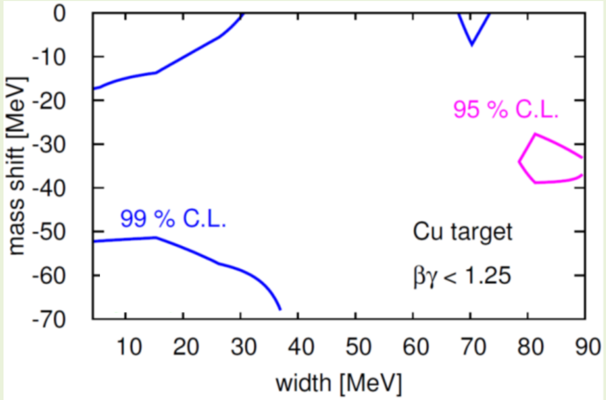
βγ-dependence of electromagnetic + rescattering effects

How do the electro + rescattering effects change the fits (Cu target)?

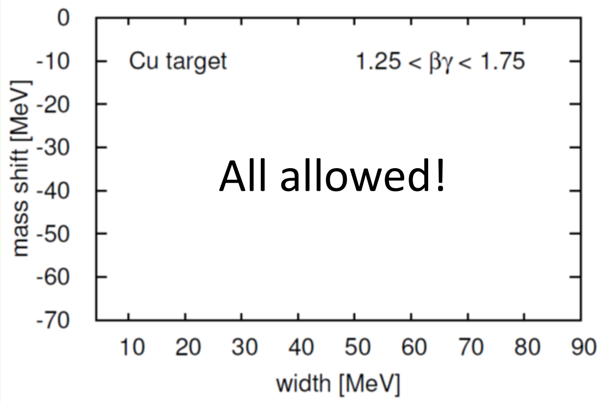
Before → **After**

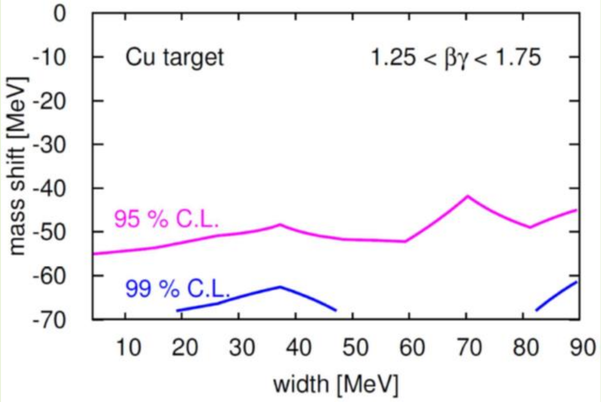
slow φs



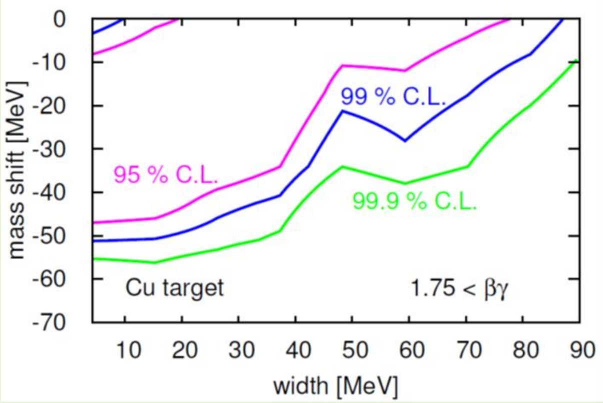


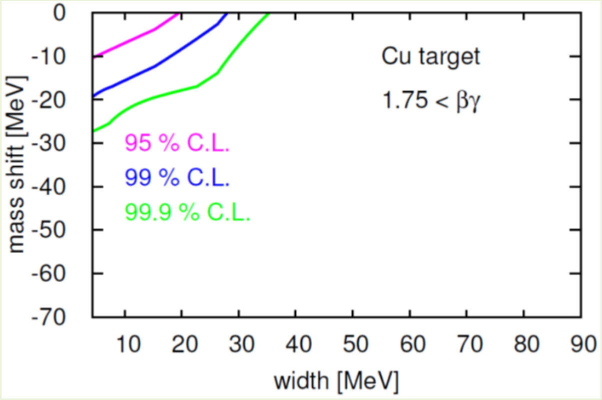
intermediate φs





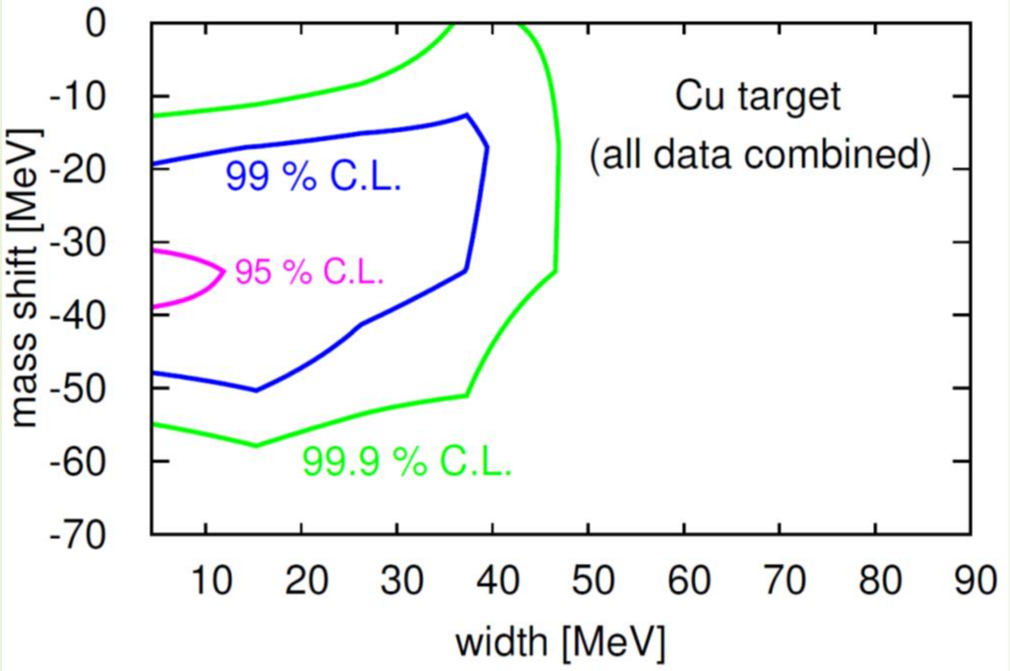
fast φs



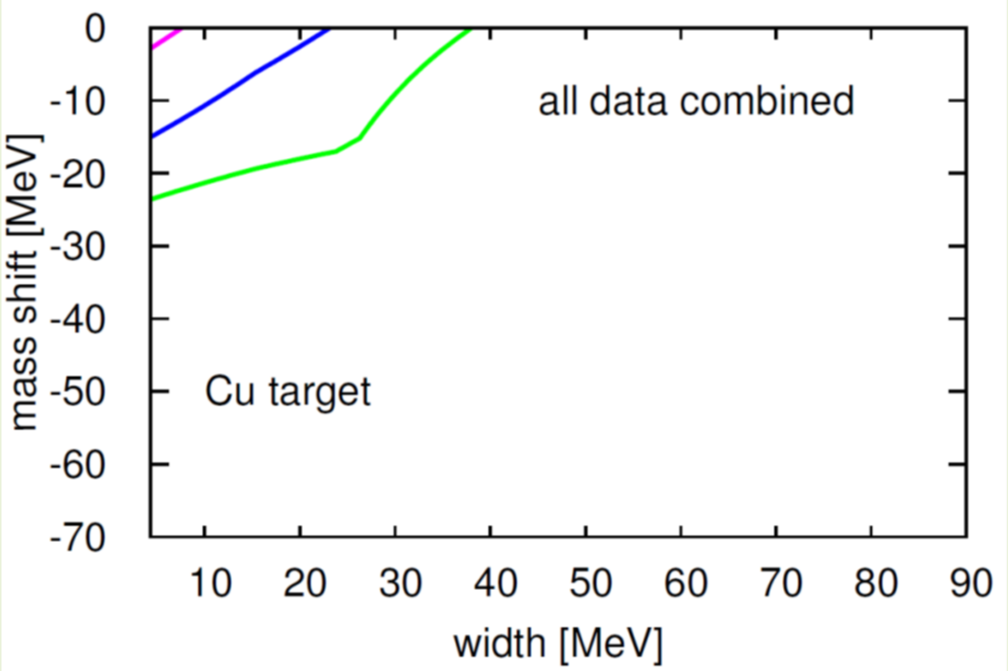


All βγ-regions combined (Cu target)

Before



After

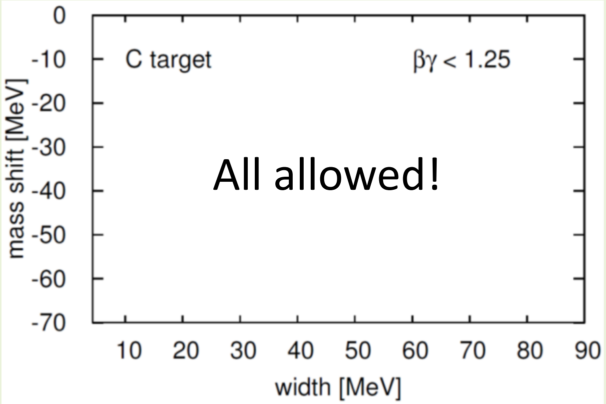


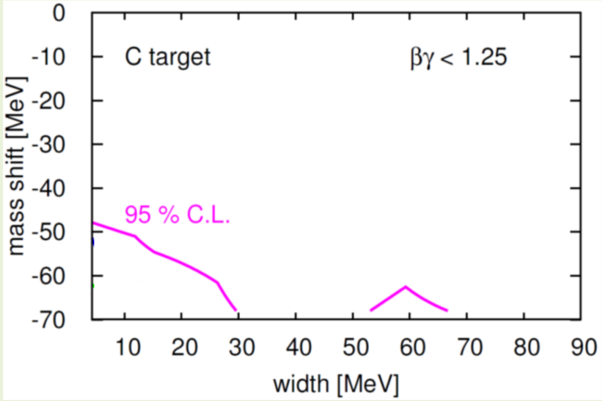
No modification
scenario favored??

How do the electro + rescattering effects change the fits (C target)?

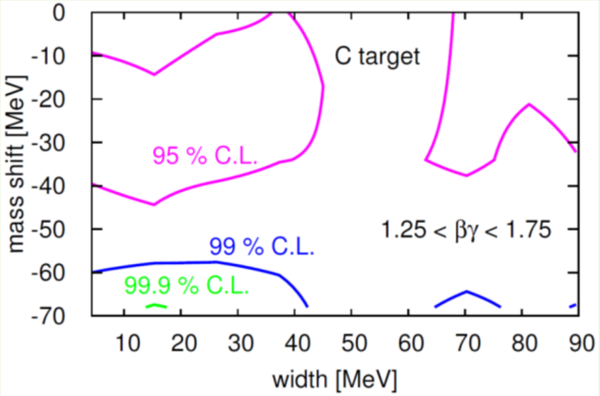
Before | **After**

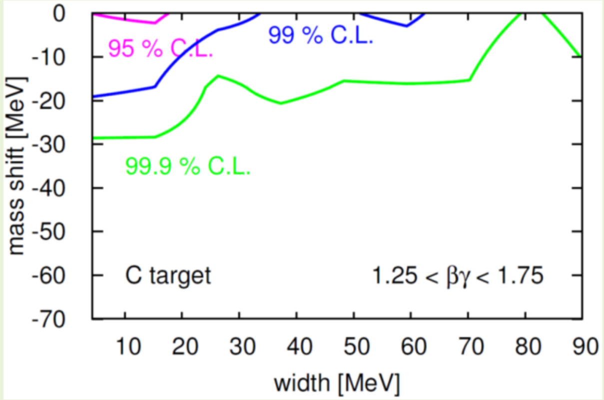
slow φs



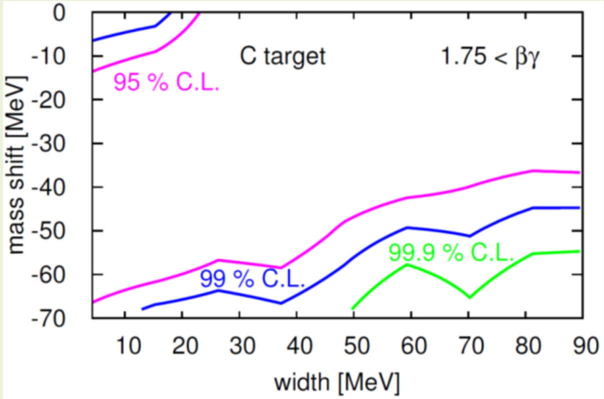


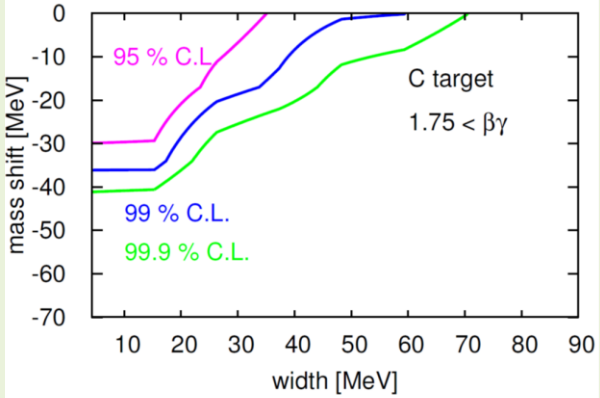
intermediate φs





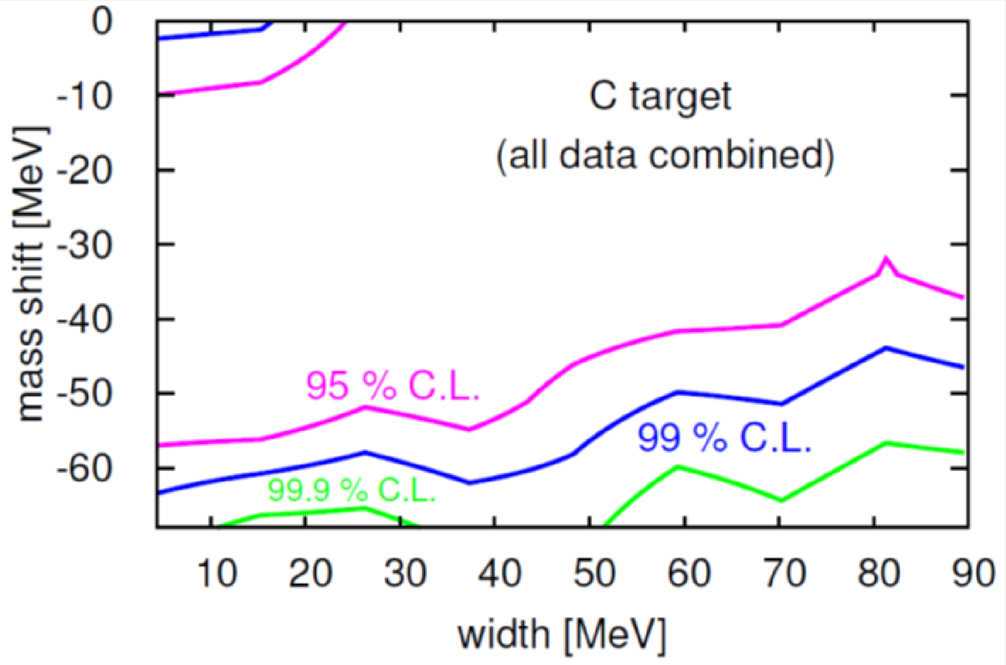
fast φs



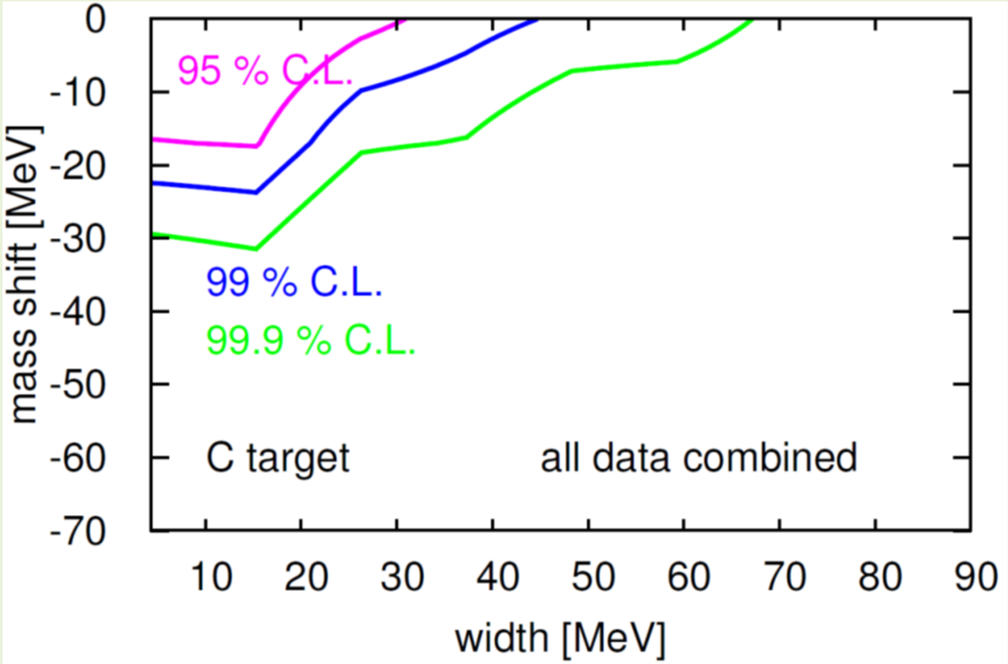


All βγ-regions combined (C target)

Before

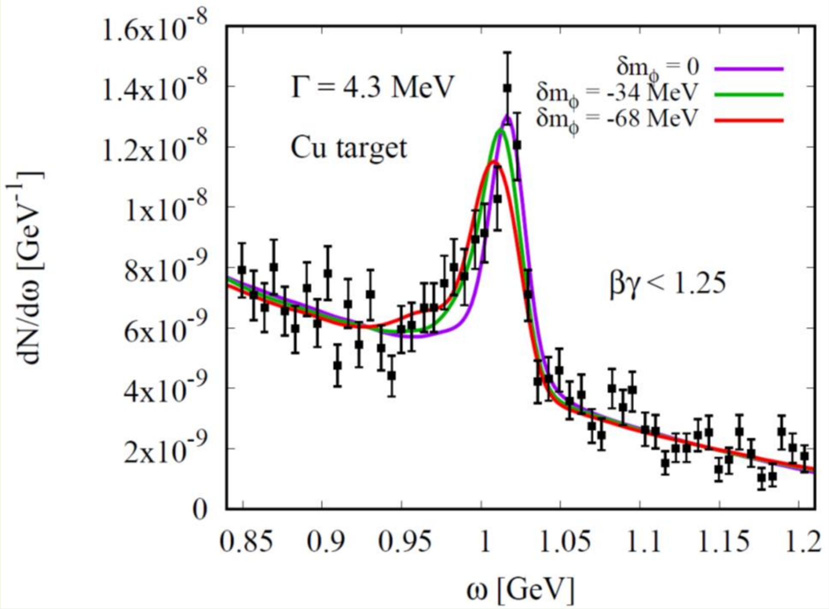


After

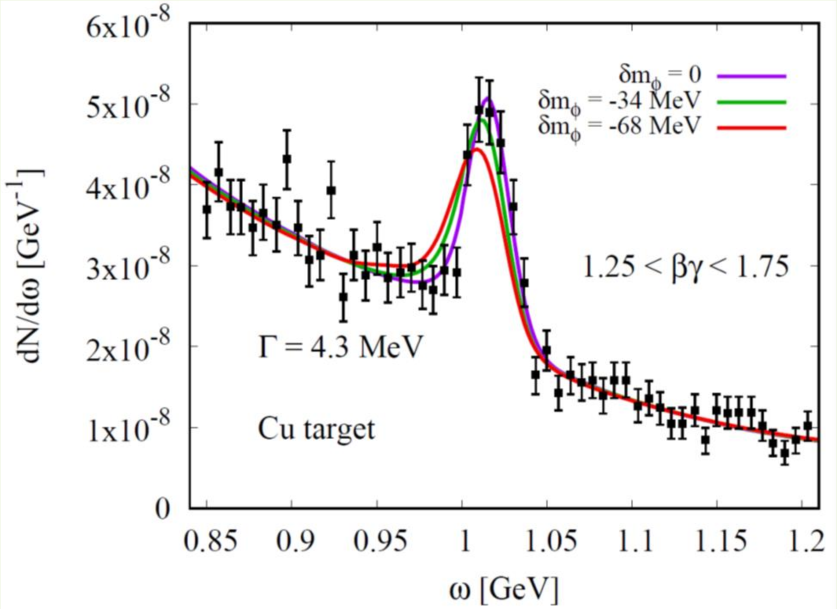


Small modification scenario favored?

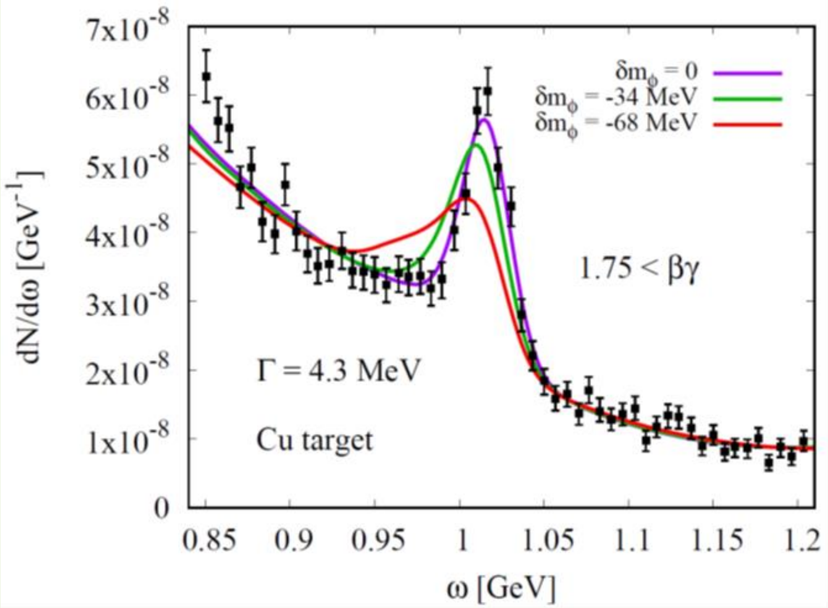
New fits to experimental Copper target data (E325)



slow φs



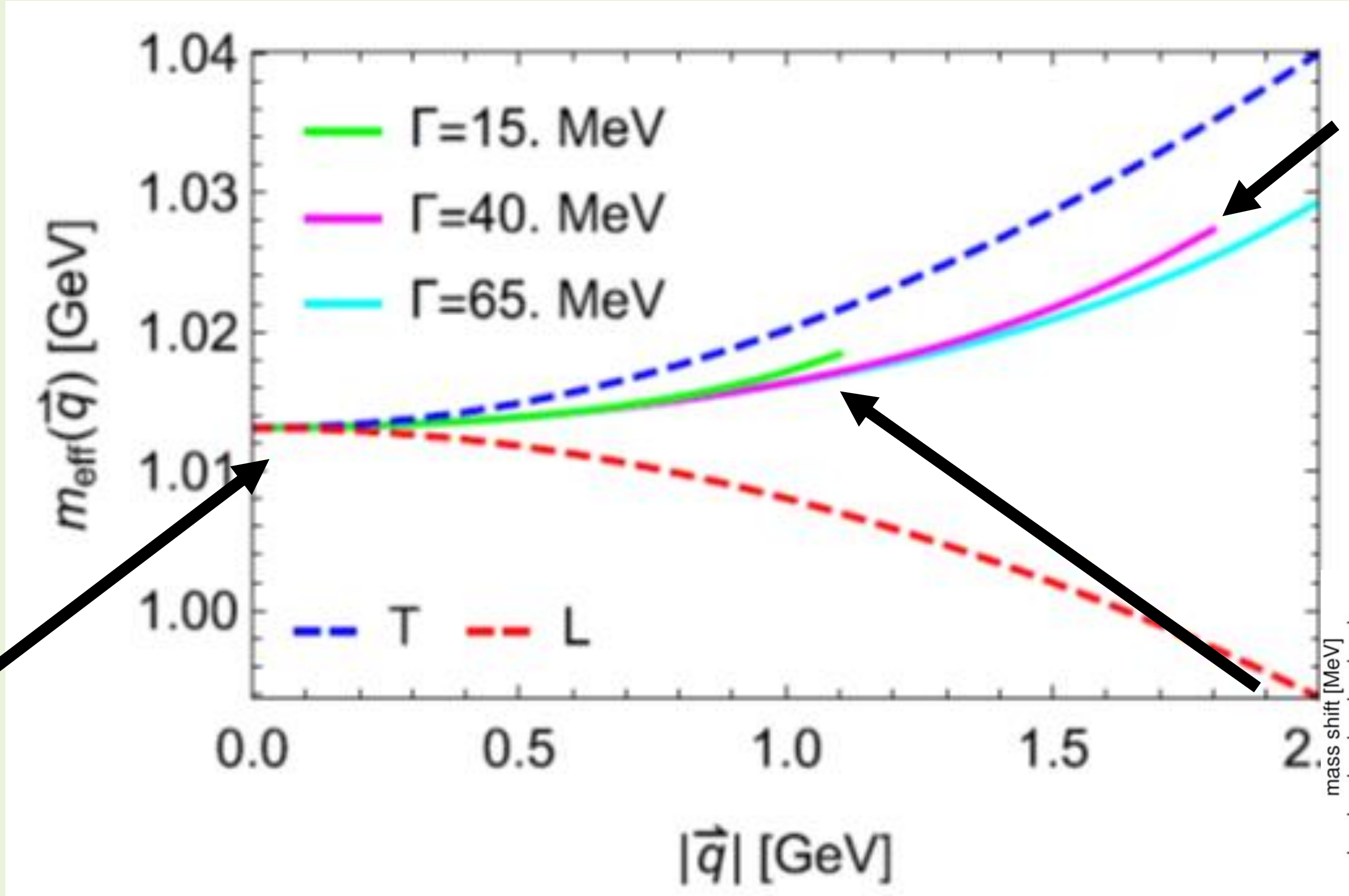
intermediate φs



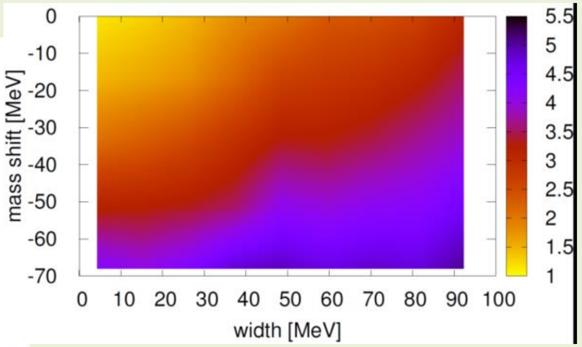
fast φs

Possible solution?

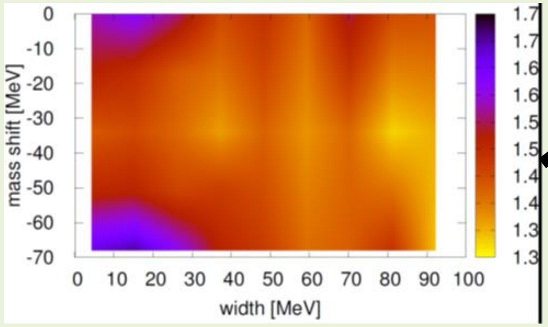
Momentum dependent mass shift



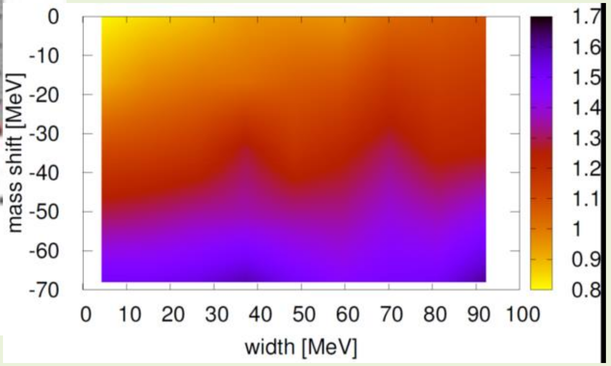
fast φs



slow φs



intermediate φs



H.J. Kim and P. Gubler, Phys. Lett. B **805**, 135412 (2020).

Summary and Conclusions

★ Relating modification of QCD condensates with hadron properties in nuclear matter is a non-trivial multi-step process

QCD condensates ⟷ Hadronic spectrum ⟷ Experimental data

★ For studying the modification of the φ meson spectral function experimentally at finite density, a good understanding of the underlying reactions is needed

★ We conducted numerical simulations of the pA reactions measured at the E325 experiment at KEK, using the HSD transport code

⇓

Estimation of electromagnetic and rescattering effect is ongoing

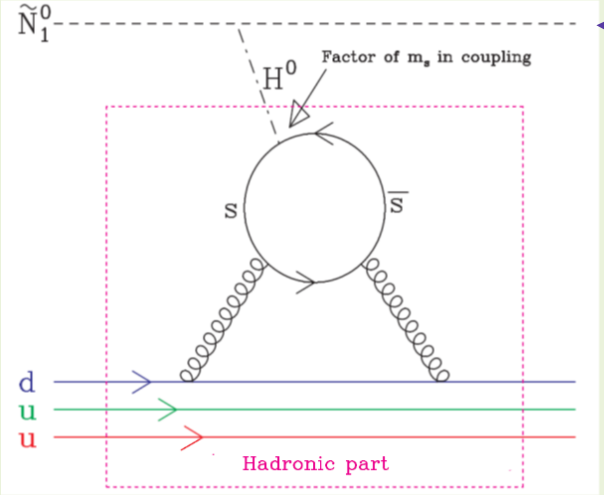
★ New J-PARC proposal P88 to measure the φ meson K^+K^- decay channel

⇒ Distortion effects due to the strong KN interaction appears to be small

Backup slides

The strangeness content of the nucleon: $\sigma_{sN} = m_s \langle N|\bar{s}s|N\rangle$

Important parameter for dark-matter searches!



Neutralino:
Linear superposition of the Super-partners of the Higgs, the photon and the Z-boson

Adapted from:
W. Freeman and D. Toussaint (MILC Collaboration), Phys. Rev. D **88**, 054503 (2013).

$$\sigma^{(\text{nucleon})}_{\text{scalar}} = \frac{8G_F^2}{\pi} M_Z^2 m_{\text{red}}^2 \left[\frac{F_h I_h}{m_h^2} + \frac{F_H I_H}{m_H^2} \frac{M_Z}{2} \sum_q \langle N|\bar{q}q|N\rangle \sum_i P_{\tilde{q}_i} (A^2_{\tilde{q}_i} - B^2_{\tilde{q}_i}) \right]^2$$

most important contribution

$$I_{h,H} = k^{h,H}_{u-\text{type}} g_u + k^{h,H}_{d-\text{type}} g_d$$

dominates

$$g_d = \frac{2}{27}\left(m_N + \frac{23}{4}\sigma_{\pi N} + \frac{25}{2}\sigma_{sN}\right)$$

A. Bottino, F. Donato, N. Fornengo and S. Scopel, Asropart. Phys. **18**, 205 (2002).

Structure of QCD sum rules for the φ meson channel

(after application of the Borel transform)

$$\chi(x) = \overline{s}(x)\gamma_\mu s(x)$$

$$\frac{1}{M^2}\int_0^\infty ds e^{-\frac{s}{M^2}}\rho(s) = c_0(\rho) + \frac{c_2(\rho)}{M^2} + \frac{c_4(\rho)}{M^4} + \frac{c_6(\rho)}{M^6} + \dots$$

In Vacuum

Dim. 0: $c_0(0) = 1 + \frac{\alpha_s}{\pi}$

Dim. 2: $c_2(0) = -6m_s^2$

Dim. 4: $c_4(0) = \frac{\pi^2}{3}\langle 0|\frac{\alpha_s}{\pi}G^2|0\rangle + 8\pi^2 m_s\langle 0|\overline{s}s|0\rangle$

Dim. 6: $c_6(0) = -\frac{448}{81}\kappa\pi^3\alpha_s\langle 0|\overline{s}s|0\rangle^2$

Structure of QCD sum rules for the φ meson

$$\frac{1}{M^2}\int_0^\infty ds e^{-\frac{s}{M^2}}\rho(s) = c_0(\rho) + \frac{c_2(\rho)}{M^2} + \frac{c_4(\rho)}{M^4} + \frac{c_6(\rho)}{M^6} + \dots$$

At finite density

(within the linear density approximation)

$$\langle \overline{s}s\rangle_\rho = \langle 0|\overline{s}s|0\rangle + \langle N|\overline{s}s|N\rangle\rho + \dots$$

Dim. 0: $$c_0(\rho) = c_0(0)$$

Dim. 2: $$c_2(\rho) = c_2(0)$$

Dim. 4: $$c_4(\rho) = c_4(0) + \rho[-\frac{2}{27}M_N + \frac{56}{27}m_s\langle N|\overline{s}s|N\rangle + \frac{4}{27}m_q\langle N|\overline{q}q|N\rangle + A_2^s M_N - \frac{7}{12}\frac{\alpha_s}{\pi}A_2^g M_N]$$

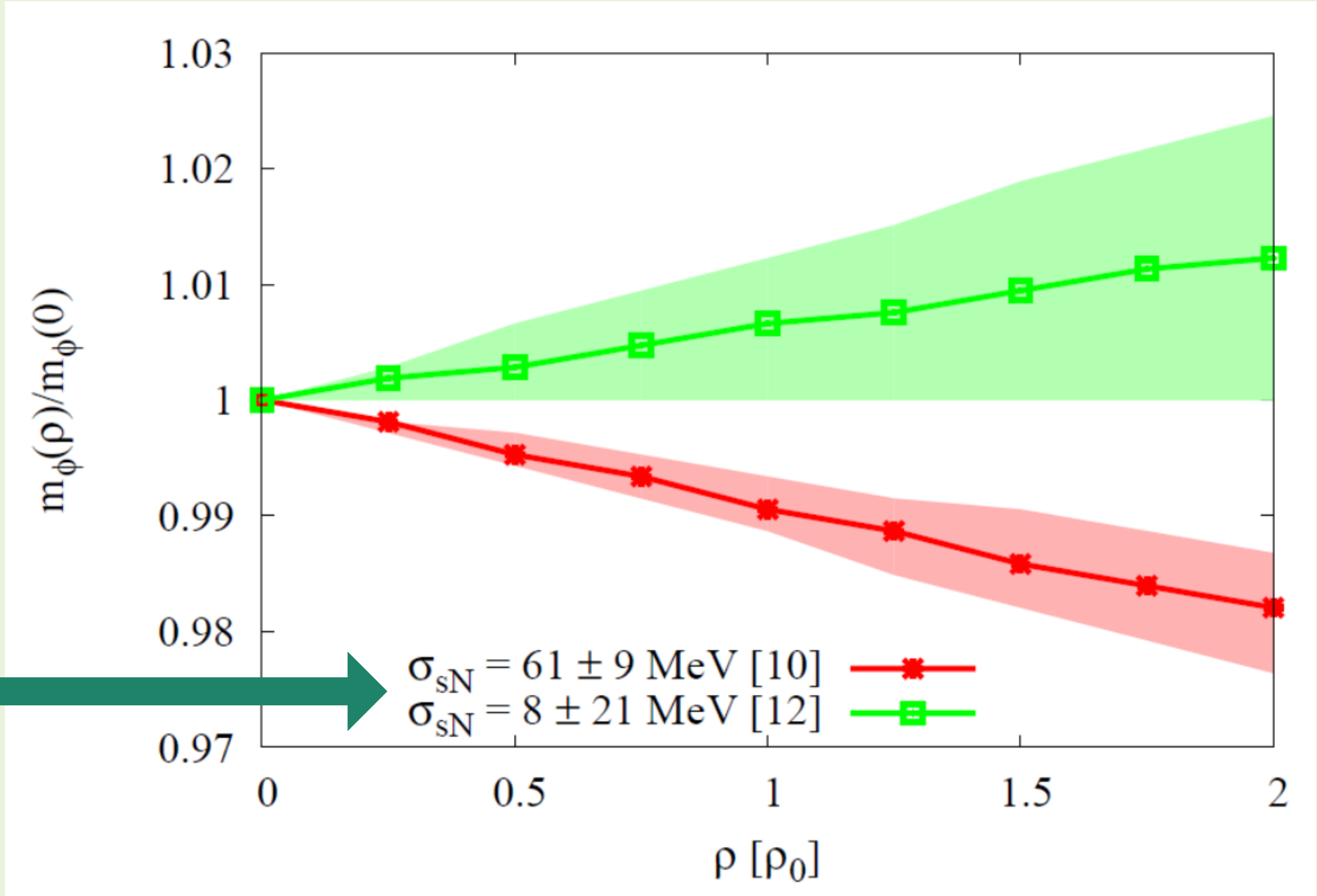
Dim. 6: $$c_6(\rho) = c_6(0) + \rho[-\frac{896}{81}\kappa_N \pi^3 \alpha_s \langle \overline{s}s\rangle\langle N|\overline{s}s|N\rangle - \frac{5}{6}A_4^s M_N^3]$$

Results for the φ meson mass at rest

Most important parameter, that determines the behavior of the φ meson mass at finite density:

Strangeness content of the nucleon

$$\sigma_{sN} = m_s \langle N|\bar{s}s|N\rangle$$



P. Gubler and K. Ohtani, Phys. Rev. D **90**, 094002 (2014).

Recent experimental results

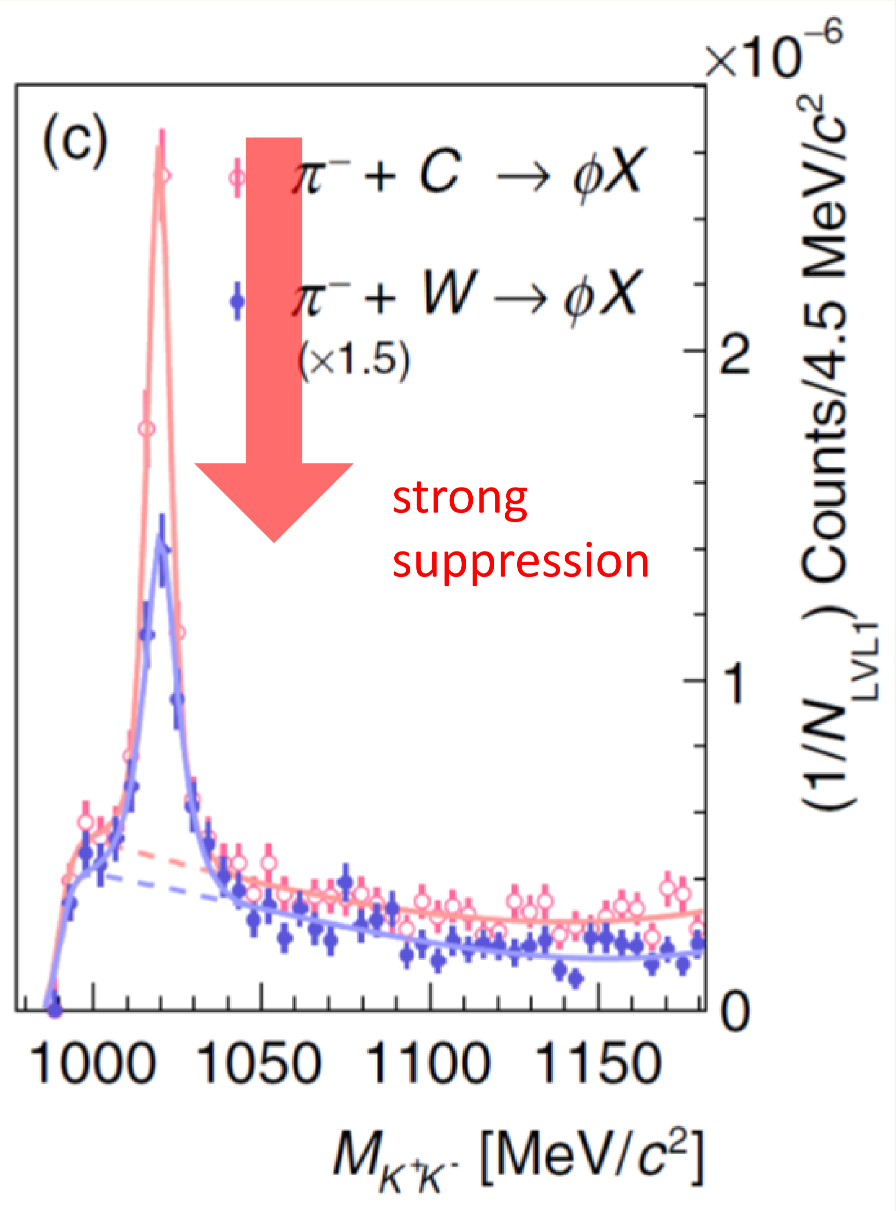
HADES: 1.7 GeV π^-A-reaction

★ Larger suppression of K^- in the Tungsten target compared to the Carbon target

★ K^-/φ ratio is similar for both Tungsten and Carbon targets

★ Observation of large suppression (broadening?) of the φ meson in large nuclei

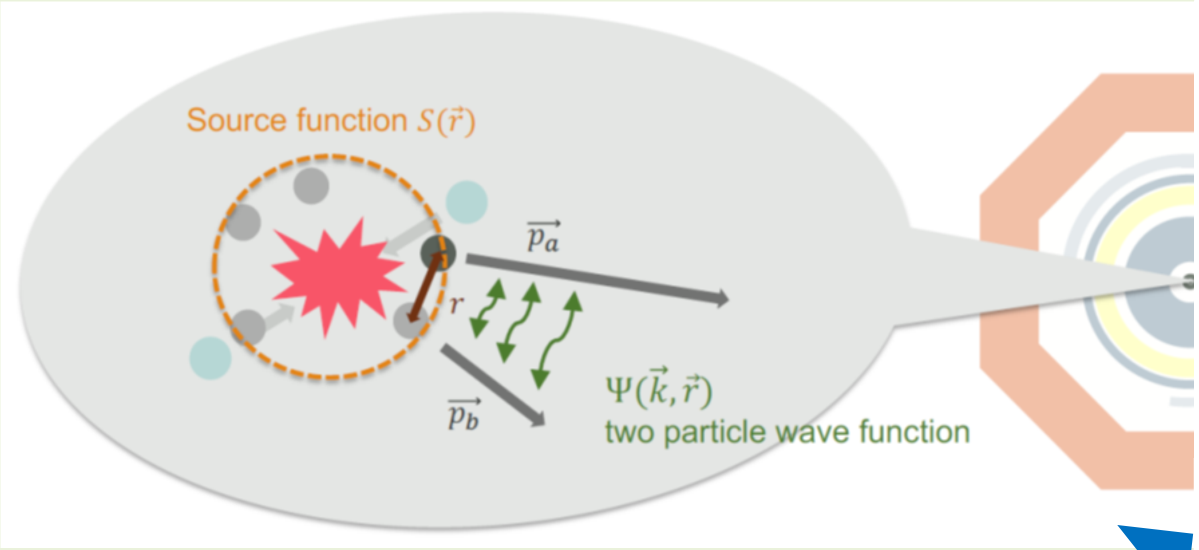
K^+K^- - invariant mass spectrum



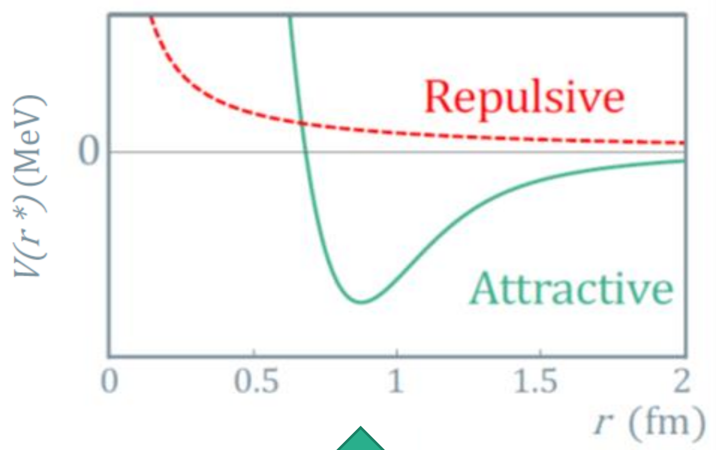
J. Adamczewski-Musch et al. (HADES Collaboration), Phys. Rev. Lett. **123**, 022002 (2019).

New experimental results

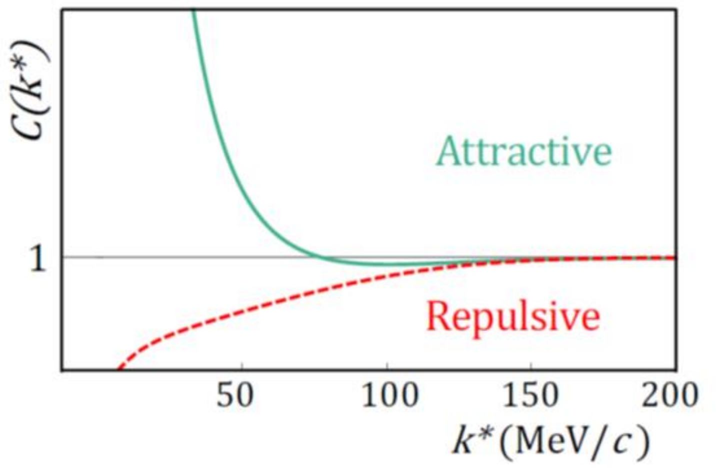
ALICE (Femtoscopy)



Interaction potential



Correlation function



The observable to be measured: the correlation function:

$$C(k) = \mathcal{N}\frac{N_{\text{Same}}}{N_{\text{Mixed}}} = \int S(\vec{r})|\Psi(\vec{k},\vec{r})|^2 d^3\vec{r}$$

Emission source (Gaussian)

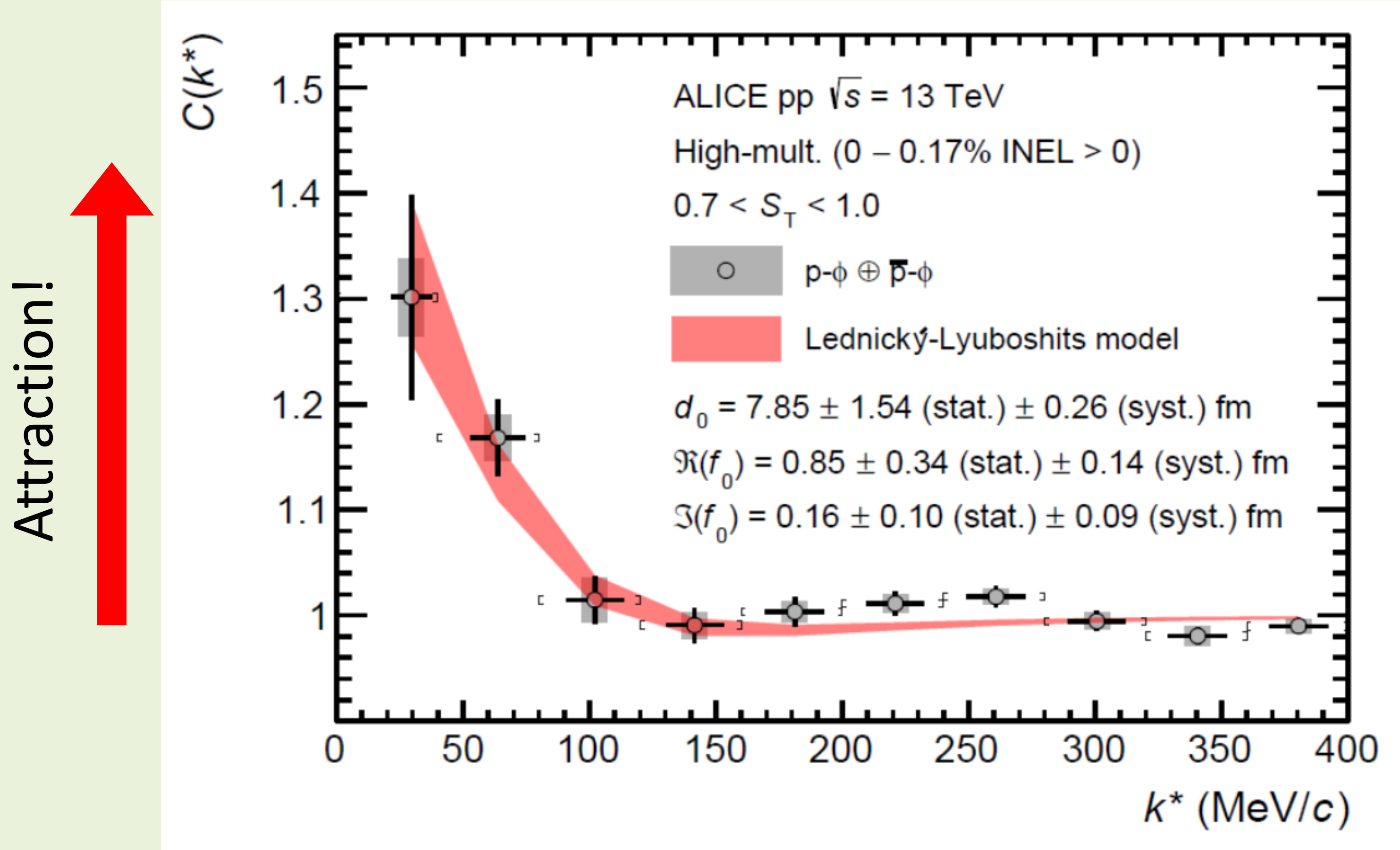
Relative momentum of the particle pair

S. Acharya et al. (ALICE Collaboration), Phys. Rev. Lett. **127**, 172301 (2021).

New experimental results

ALICE

Measurement of φN correlation



Extracted φN scattering length

Real part:

$$\mathrm{Re}(f_0) = 0.85 \pm 0.34(\mathrm{stat.}) \pm 0.14(\mathrm{syst.})\ \mathrm{fm}$$

→ Attractive

Imaginary part:

$$\mathrm{Im}(f_0) = 0.16 \pm 0.10(\mathrm{stat.}) \pm 0.09(\mathrm{syst.})\ \mathrm{fm}$$

→ Small absorption/broadening ?

S. Acharya et al. (ALICE Collaboration), Phys. Rev. Lett. **127**, 172301 (2021).

New experimental results

ALICE

Fit of the correlation function data to two simple phenomenological potentials

$$E_{\rm int} = \int d^3\vec{r} \int d^3\vec{r}\,' \rho_N(\vec{r}) V(\vec{r}-\vec{r}\,')\rho_\phi(\vec{r}\,')$$

$\rho_N(\vec{r}) \rightarrow \rho_0$, $\rho_\phi(\vec{r}\,') \rightarrow \delta^{(3)}(\vec{r}\,')$

$$V_{\rm Yukawa}(r) = -\frac{A}{r}e^{-\alpha r}$$

$$A = 0.021 \pm 0.009\,({\rm stat.}) \pm 0.006\,({\rm syst.})$$

$$\alpha = 65.9 \pm 38.0\,({\rm stat.}) \pm 17.5\,({\rm syst.})\,{\rm MeV}$$

$\rightarrow$

$$E_{\rm int} = -\frac{4\pi A \rho_0}{\alpha^2} = -79.3 \pm 108.8\,{\rm MeV}$$

$$V_{\rm Gaussian}(r) = -V_{\rm eff}\, e^{-\mu r^2}$$

$$V_{\rm eff.} = 2.5 \pm 0.9\,({\rm stat.}) \pm 1.4\,({\rm syst.})\,{\rm MeV}$$

$$\mu = 0.14 \pm 0.06\,({\rm stat.}) \pm 0.09\,({\rm syst.})\,{\rm fm}^{-2}$$

$\rightarrow$

$$E_{\rm int} = -\frac{\pi^{3/2} V_{\rm eff}\,\rho_0}{\mu^{3/2}} = -45.2 \pm 61.5\,{\rm MeV}$$

S. Acharya et al. (ALICE Collaboration), arXiv:2105.05578 [nucl-ex].

$\rightarrow$ Larger attraction than what was observed at KEK 325, but large statistical and systematic uncertainties

Our tool: a transport approach

Basic Ingredient 1: Solve a Boltzmann-Uehling-Uhlenbeck (BUU) type equation for each particle type

$$\left(\frac{\partial}{\partial t} + \vec{\nabla}_p \epsilon \cdot \vec{\nabla}_r - \vec{\nabla}_r \epsilon \cdot \vec{\nabla}_p\right) f_a(\vec{r}, \vec{p}; t) = I_{\text{coll}}[f_a(\vec{r}, \vec{p}; t)]$$

Includes mean field (tuned to reproduce nuclear matter properties)

particle distribution function

Basic Ingredient 2: „Testparticle“ approach

$$f_h(\boldsymbol{r}, \boldsymbol{p}; t) = \frac{1}{N_{\text{test}}} \sum_{i}^{N_h(t) \times N_{\text{test}}} \delta(\boldsymbol{r} - \boldsymbol{r}_i(t))\ \delta(\boldsymbol{p} - \boldsymbol{p}_i(t))$$

Example of a transport calculation

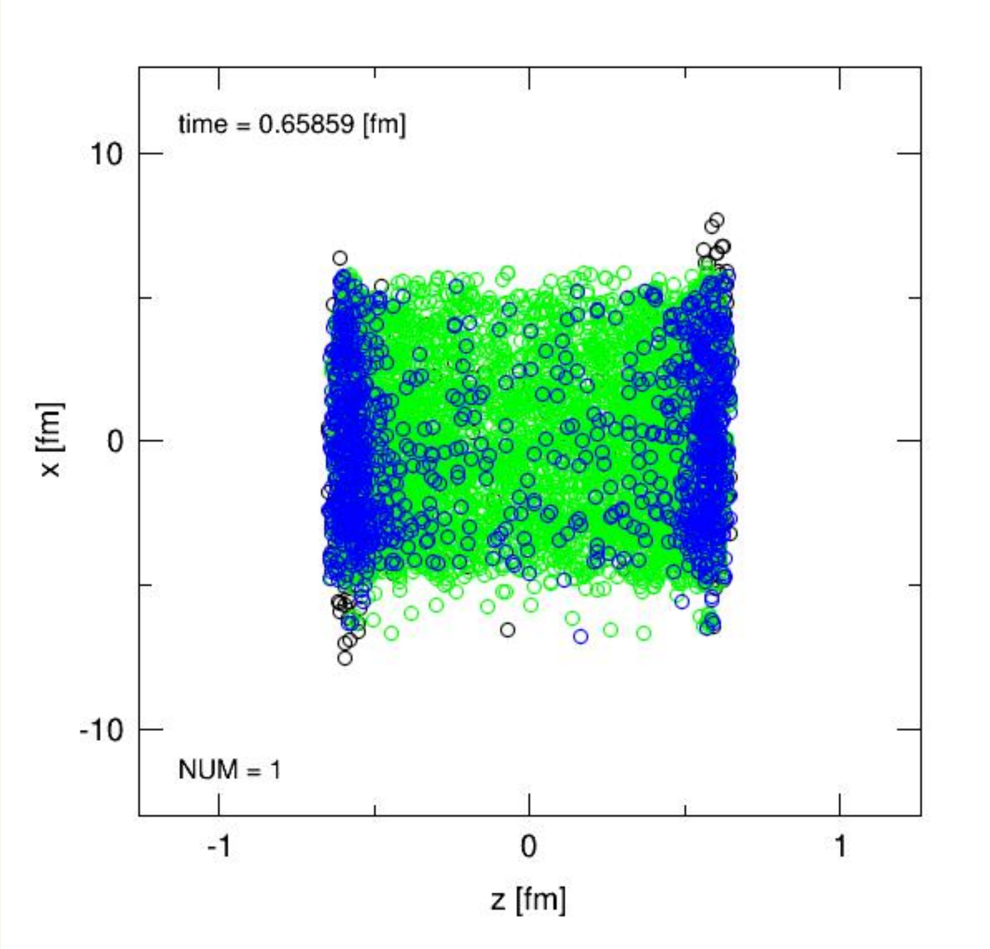
Au+Au collision at $s^{1/2}$ = 200 GeV, b = 2 fm

nucleons

quarks

gluons

will not be included in the simulations shown in this talk



Final step: comparison to experimental data

Potential issues:

- ★ Experimental background is not included in the simulation
- ★ Normalization of the experimental dilepton spectrum is not given

Fit to experimental data is necessary!

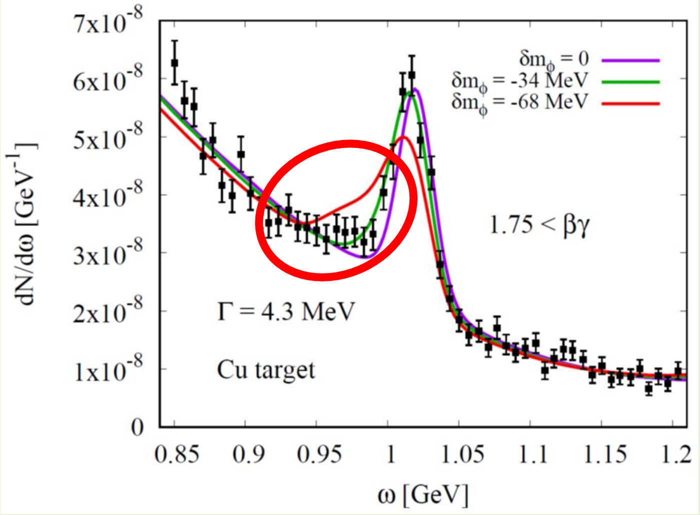
Background — ϕ meson signal

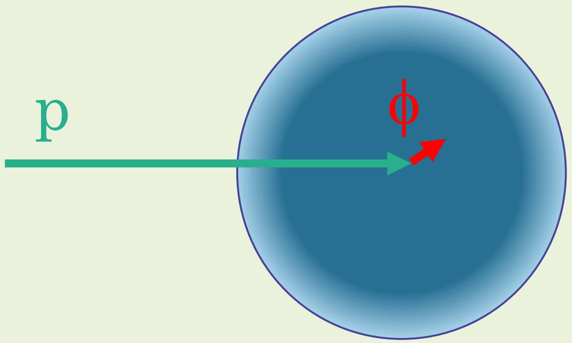
Dilepton spectrum: $$\rho(\omega) = \underbrace{a\omega^2 + b\omega + c}_{\text{Background}} + \underbrace{A\rho_{\phi,\mathrm{HSD}}(\omega)}_{\phi \text{ meson signal}}$$

a, b, c, A: Fitted to the experimental dilepton spectrum independently for each βγ-region

Reason for large modification for fast φ mesons

→ Initial stage of φ meson production

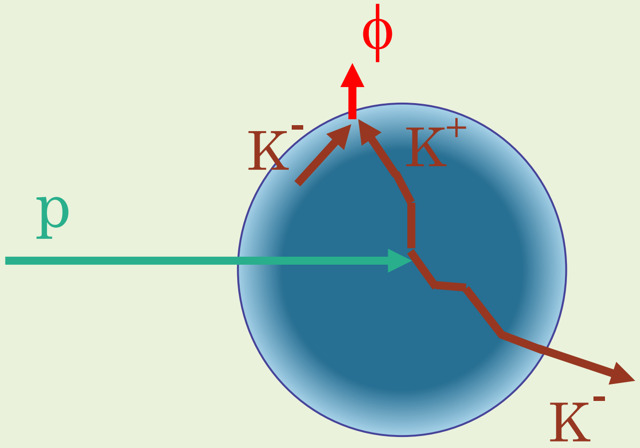




φ mesons are generated from high energy collisions (via strings)

→ **large momentum**

→ **high density**



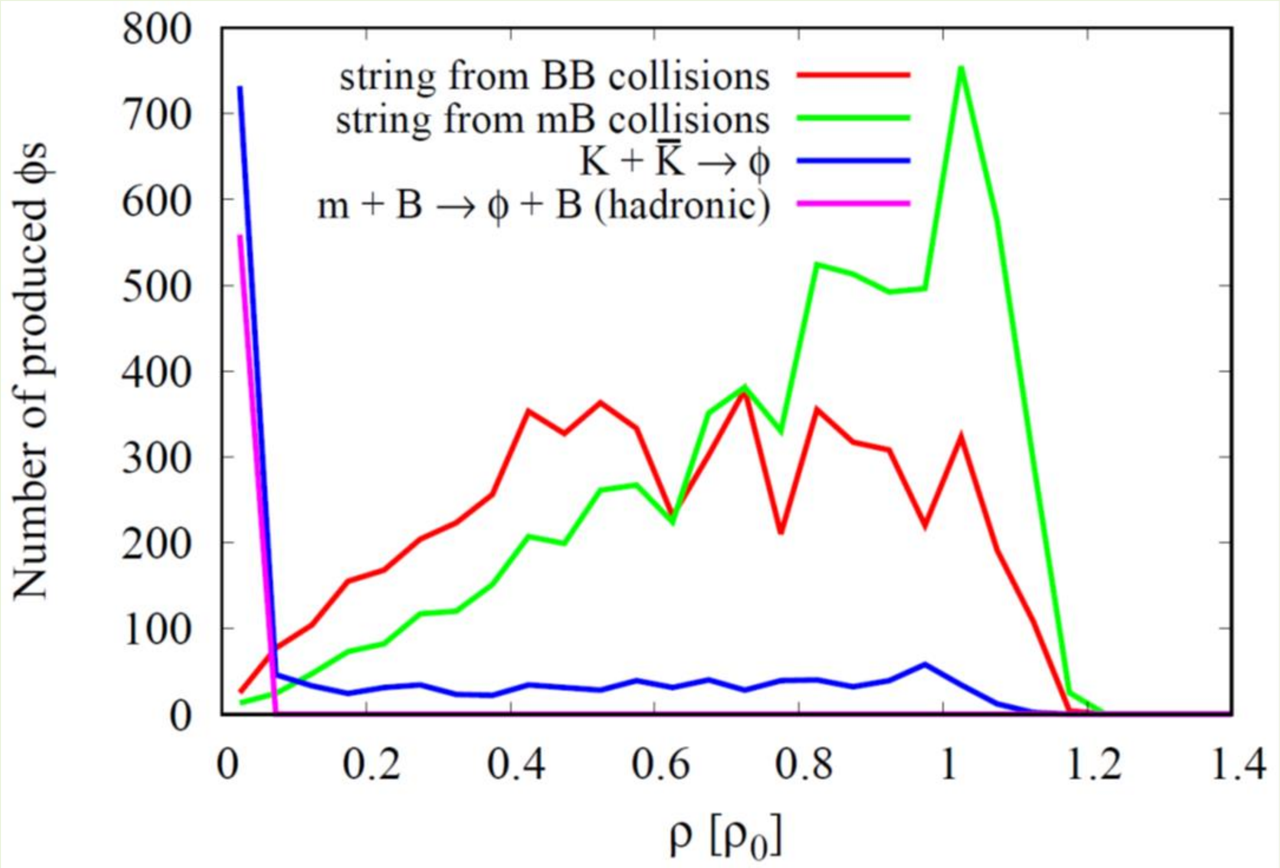
φ mesons are generated from low energy hadronic collisions

→ **small momentum**

→ **low density**

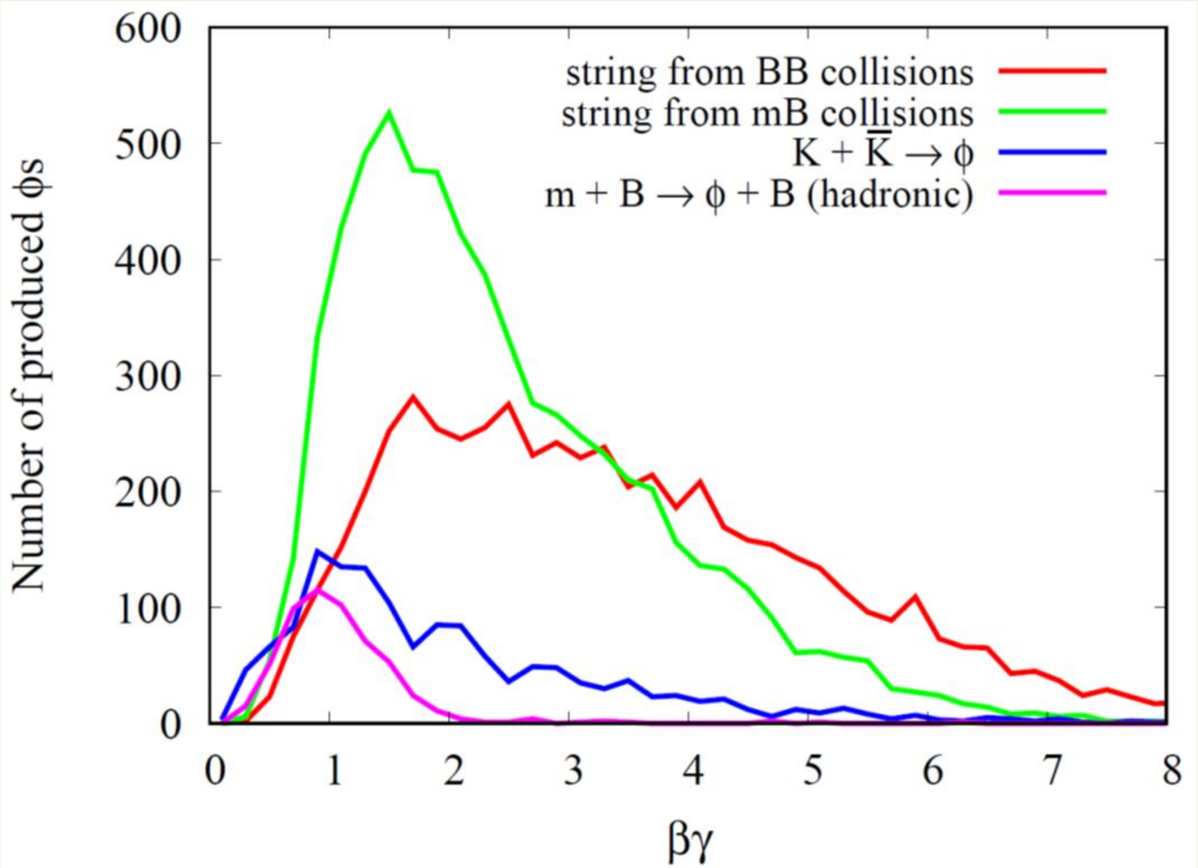
Density and βγ distributions for the different production mechanisms

Density distribution at production



Low energy hadronic production occurs dominantly at the nuclear surface

βγ distribution at production



For βγ>1.5, high energy φ meson production via strings dominates