

- $T_{\mu\nu}$: Noether current for translational symmetry

Energy density

Momentum density

$$T_{\mu\nu} = \begin{pmatrix} T_{00} & T_{01} & T_{02} & T_{03} \\ T_{01} & T_{11} & T_{12} & T_{13} \\ T_{02} & T_{21} & T_{22} & T_{23} \\ T_{03} & T_{23} & T_{32} & T_{33} \end{pmatrix}$$

stress

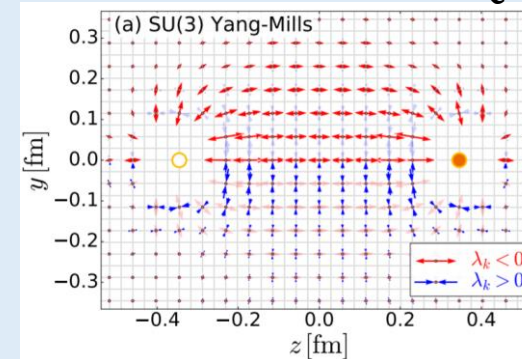
Advantage

- ✓ Gauge invariance
- ✓ Conserved quantity
- ✓ Local interaction

- **Distribution of EMT** in $Q\bar{Q}$

Lattice QCD

Stress distribution in $Q\bar{Q}$



Yanagihara, et al. 2019

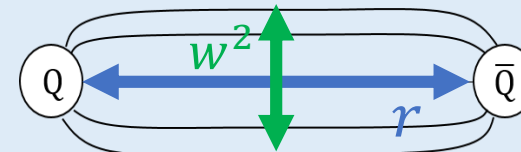
Forming fluxtube

Pulling force
(parallel to field)

Pulling force
(vertical to field)

- **Quantum effect** in $Q\bar{Q}$

Theoretical model



Fattening due to
string vibration

$$w^2 \sim \log r$$

Lüscher, Münster, and Weisz 1981



Motivate **theoretical analysis of quantum correction to EMT distribution in $Q\bar{Q}$ system**

ϕ^4 Model in 1+1 d

Purpose: Analysis of **quantum correction to EMT distribution** around a soliton in **1+1d real scalar ϕ^4 model**

$$S = \int dx^2 \left\{ \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - \frac{\lambda}{4} \left(\phi^2 - \frac{m^2}{\lambda} \right)^2 \right\}$$

$U = \left(\phi^2 - \frac{m^2}{\lambda} \right)^2$

vacuum $\frac{m}{\sqrt{\lambda}}$ $\frac{m}{\sqrt{\lambda}}$ ϕ

Classical

- Kink solution of EOM

$$\phi_{cl} = + \frac{m}{\sqrt{\lambda}} \tanh \left(\frac{m(x-X)}{\sqrt{2}} \right)$$

X

$\frac{m}{\sqrt{\lambda}}$

$-\frac{m}{\sqrt{\lambda}}$

x

X : free parameter

- T_{00} under kink solution

$$T_{00} = \frac{m^4}{\lambda} \cosh^{-4} \left(\frac{m(x-X)}{\sqrt{2}} \right)$$

T_{00}

$\sim m^{-1}$

X

x

QM

$$E_{LO} \sim O(\lambda^0) \quad T_{00}^{LO} \sim O(\lambda^0)$$

Dashen, et al 1974

Goldhaber, et al. 2003

We calculated $T_{00}^{LO}(x)$ and $T_{11}^{LO}(x)$!

Eigenmode of quantum fluctuation

Expanding ϕ^4 action $S[\phi]$ around a soliton

Substituting $\phi(x) = \phi_{cl} + \eta(x)$

$$S[\eta] = S_{cl} + \int dx^2 \left[\frac{1}{2} (\partial_0 \eta)^2 - \frac{1}{2} \eta \left(-\frac{\partial^2}{\partial x^2} - m^2 + \frac{3\phi_{cl}^2}{\lambda} \right) \eta + O(\lambda^{\frac{1}{2}}) \right]$$

ω

$$\left(-\frac{\partial^2}{\partial x^2} - m^2 + \frac{3\phi_{cl}^2}{\lambda} \right) \eta_n = \omega_n^2 \eta_n$$

Eigenvalues

Eigenfunctions

$$\omega_q^2 = q^2 + 2m^2 \quad \eta_q(x) \xrightarrow{x \rightarrow \pm \infty} e^{i \left(qx \pm \frac{1}{2} \delta(q) \right)}$$

Phase shift

$$\omega_1^2 = \frac{3}{2} m^2$$

$$\eta_1(x)$$

$$\omega_0^2 = 0$$

$$\eta_0(x) = \partial_x \phi_{cl}$$

Translational mode \rightarrow IR divergence

Collective coordinate method

Gervais, Jevicki, Sakita 1975 Tomboulis 1975

Rewriting Lagrangian

Remove translational mode($\eta_0(x)$)
and IR divergence

$$\phi(x, t) = \phi_{cl}(x - X) + \eta(x - X)$$

X is promoted to a dynamical variable

$X \rightarrow X(t)$ Dynamical variable

$P \leftrightarrow X$ Canonical conjugate momentum

$$L[\pi, \phi] \rightarrow L[\tilde{\pi}, \tilde{\eta}, X, P]$$

$\tilde{\pi}, \tilde{\eta}$: without translational mode

Convert \rightarrow Degree of freedom of $\eta_0(x)$

\rightarrow Degree of freedom of
Soliton center of motion

Rewriting EMT

- Translational invariance Gervais, Jevicki, Sakita 1975
- Lorentz symmetry Goldstone and Jackiw 1975
Tomboulis 1975

ϕ^4 Model in 1+1 d

$$T^\mu_\nu[\phi] = \frac{\partial \mathcal{L}}{\partial(\partial_\mu \phi)} \partial_\nu \phi - \mathcal{L} \delta^\mu_\nu$$



$$T^\mu_\nu[\tilde{\pi}, \tilde{\eta}, X, P] = T^\mu_\nu[X, P] + T^\mu_\nu[\tilde{\pi}, \tilde{\eta}]$$

$O(\lambda^1)$

Collective
coordinate
method

We calculate $\langle T^\mu_\nu \rangle$

$$\langle T^\mu_\nu \rangle \equiv \frac{1}{Z} \int \mathcal{D}\tilde{\pi} \mathcal{D}\tilde{\eta} T^\mu_\nu[\tilde{\pi}, \tilde{\eta}] e^{-i \int dt L[\tilde{\pi}, \tilde{\eta}]}$$

EMT regularization

Rebhan and Nieuwenhuizen 1997
Dashen, et al 1974

Counter terms

$$\langle T_{\mu\nu} \rangle_{\text{soliton}} - \langle T_{\mu\nu} \rangle_{\text{vac}} + (\langle \delta T_{\mu\nu} \rangle_{\text{soliton}} - \langle \delta T_{\mu\nu} \rangle_{\text{vac}}) = \text{finite}$$

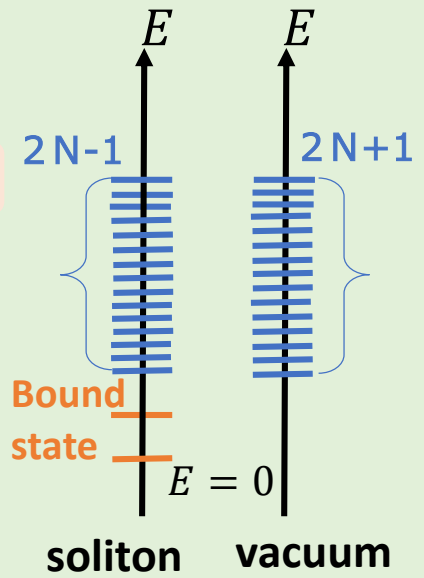
Mode number cutoff (MNC)

- ✓ Finite system
- ✓ Phase shift
- ✓ "Mode number" is the same for soliton and vacuum

Soliton continuous mode vacuum

$$\sum_{n=-\infty}^{n=\infty} \sqrt{q_n^2 + 2m^2} - \sum_{n=-\infty}^{n=\infty} \sqrt{k_n^2 + 2m^2}$$

$$= \sum_{n=-(N-1)}^{N-1} \sqrt{q_n^2 + 2m^2} - \sum_{n=-N}^N \sqrt{k_n^2 + 2m^2}$$



subtraction $R \rightarrow \infty$ Infinite system

Mass renormalization

- ✓ Only 1-loop mass renormalization
- ✓ Mass renormalization in **vacuum** sector
 ➔ Counter terms appear in **soliton** sector

Vacuum sector $m^2 \rightarrow m^2 + \delta m^2$

$$\text{Diagram with } \delta m^2 \text{ and } \lambda \text{ counterterms} = 0$$

Soliton sector

$$\text{Diagram with } \lambda \phi_{cl} \text{ and } \delta m^2 \phi_{cl} \text{ counterterms} = \text{finite}$$

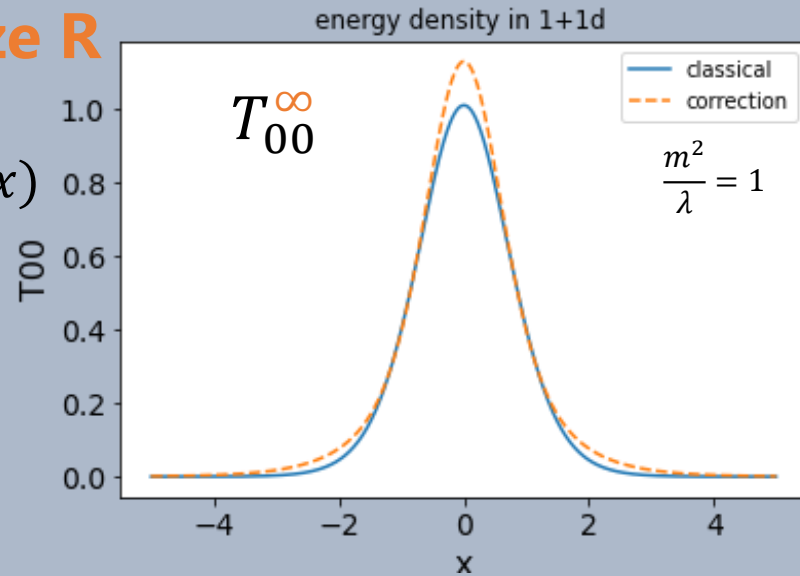
$$\text{Diagram with } \delta m^2 \phi_{cl}^2 \text{ counterterm} = \text{finite}$$

Result, summary, and future work

In a space of finite size R

$$T_{00}^R(x) = -\frac{3\sqrt{2m}}{2\pi} \frac{1}{R} + T_{00}^\infty(x)$$

$$T_{11}^R = -\frac{3\sqrt{2m}}{2\pi} \frac{1}{R}$$



■ T_{00} distribution

$$\lim_{R \rightarrow \infty} \int_{-R/2}^{R/2} dx T_{00}^R = E \quad \times \quad \int_{-\infty}^{\infty} dx \lim_{R \rightarrow \infty} T_{00}^R$$

reproduce the E_{MNC} calculated by Dashen et al.

■ T_{11} distribution

$$T_{11}^R \text{ is constant} \quad \rightarrow \quad \partial_0 T^{01} + \partial_1 T^{11} = 0$$

EMT conservation law is satisfied

Summary

1-loop calculation of EMT distribution around a soliton in 1+1d real scalar ϕ^4 model

- Removing IR divergences by collective coordinate method

- Removing UV divergences by vacuum subtraction (Mode number cutoff) and mass renormalization

- The spatial integrals of obtained T_{00} reproduce the known total energies. T_{11} are consistent with the EMT conservation laws.

Future work

- 2+1d ϕ^4 model
- finite temperature
- sine-Gordon model