

Beam-energy dependence of the anisotropy scaling functions for identified particle species

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Takeaway

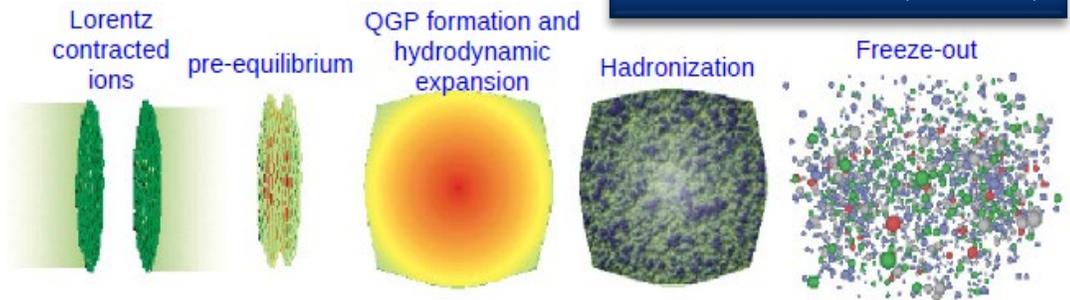
The $\sqrt{s_{NN}}$ -dependence of the anisotropy scaling functions for PID species can be used to:

- I. Delineate the respective influence of expansion dynamics and viscous attenuation*
- II. Constrain $\frac{\eta}{s}(T, \mu_B, \mu_I, \mu_S)$?*
 - ✓ Give insight on a possible critical point in the nuclear matter phase diagram*

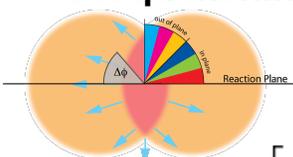
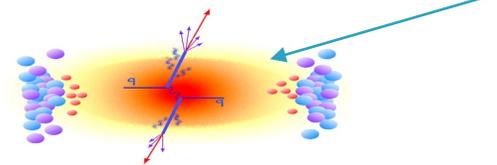


Azimuthal Anisotropy

A+A collision (S. Bass)



Drives azimuthal Anisotropy



$$R_{AA}(p_T, L) \simeq \exp \left[-\frac{2\alpha_s C_F}{\sqrt{\pi}} L \sqrt{\hat{q} \frac{\mathcal{L}}{p_T}} \right]$$

Phys. Lett. B519, 199 (2001)

$$R_{v_2}(p_T, \Delta L) = \frac{R_{AA}(90^\circ, p_T)}{R_{AA}(0^\circ, p_T)} = \frac{1 - 2v_2(p_T)}{1 + 2v_2(p_T)}$$

Specific dependencies on $\sqrt{p_T}$, ΔL and \hat{q}

Reaction Plane

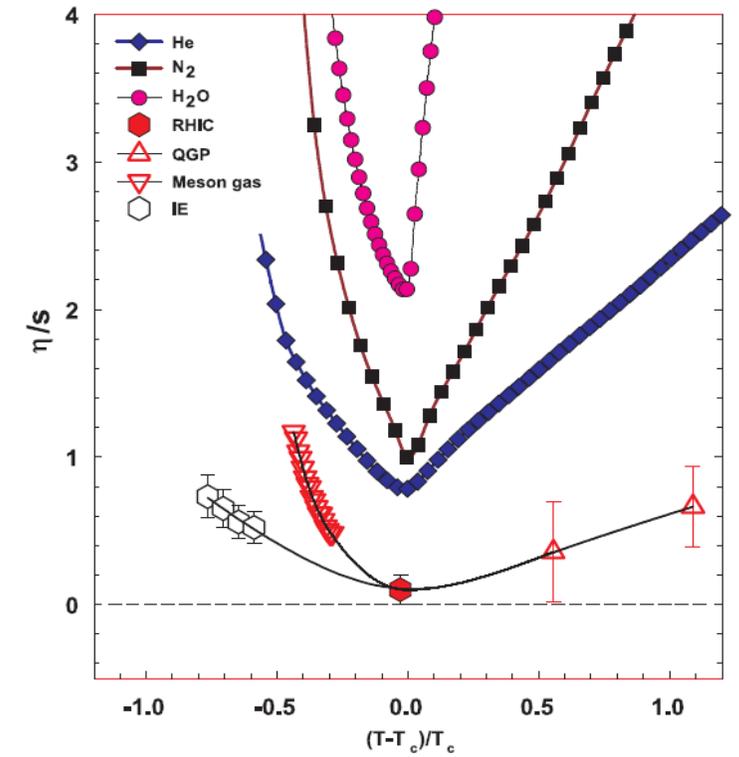
Collective flow – low p_T

$$v_n = \varepsilon_n e^{-\frac{\beta}{RT} n(n+\kappa p_T^2)}, RT \equiv \mathbb{R} \propto \langle N_{\text{chg}} \rangle^{1/3}$$

Specific dependencies on $n, \varepsilon_n, p_T, RT$ and $\frac{\eta}{s}, \frac{\xi}{s}$

Question $\frac{\eta}{s}(T, \mu_B, \mu_I, \mu_S)?$

Lacey et. al. Phys.Rev.Lett. 98 (2007) 092301



- Could give insights on:
 - ✓ the location of the critical point in the QCD phase diagram
 - ✓ Viscosity of particles vs. antiparticles? (influence of charged currents)

➤ Anisotropy Scaling Functions (ASF) for unidentified and identified particle species are used as constraints

Anisotropy Scaling Functions

$$v_n(p_T, cent) = \varepsilon_n e^{-\frac{\beta}{RT} [n(n + \kappa p_T^2)]}, RT \equiv \mathbb{R} \propto \langle N_{\text{chg}} \rangle^{1/3}$$

Same harmonic with variable centrality

$$\frac{v_n(p_T, cent)}{\varepsilon_n} = \left(\frac{v'_n(p_T, cent)}{\varepsilon'_n} \right) e^{-\frac{\beta}{\mathbb{R}} [n(n + \kappa p_T^2)] \left(1 - \frac{\mathbb{R}}{\mathbb{R}'} \right)},$$

For two harmonics at a fixed centrality

$$\frac{v_n(p_T)}{\varepsilon_n} = \left(\frac{v_m(p_T)}{\varepsilon_m} \right)^{\frac{n}{m}} e^{\frac{\beta}{\mathbb{R}}(m-n)},$$

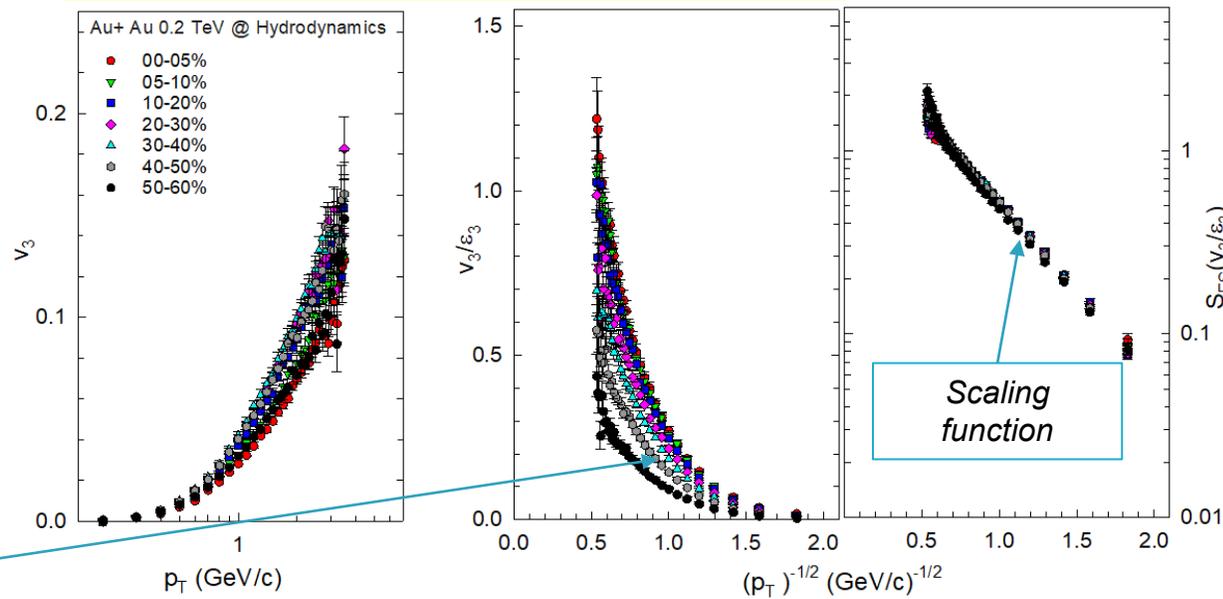
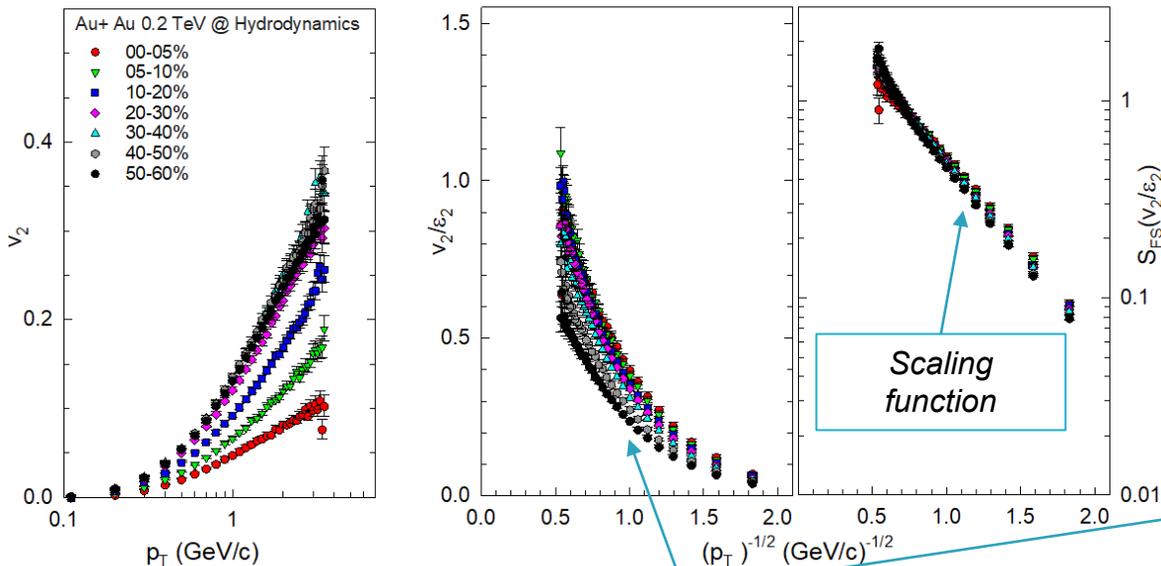
Anisotropy Scaling Functions leverage the specific dependencies of ε_n , viscous attenuation and the expansion dynamics!

- **Data collapse to a single curve for fully constrained scaling coefficients**
 - ✓ **Scaling coefficients give access to transport coefficients**
 $\frac{\eta}{s}(T, \mu_B)$, *etc.*

Anisotropy Scaling Function – proof of principle

$$\frac{v_n(p_T, cent)}{\varepsilon_n} = \left(\frac{v'_n(p_T, cent)}{\varepsilon'_n} \right) e^{-\frac{\beta}{R} \left[n(n+\kappa p_T^2) \right] \left(1 - \frac{R}{R'} \right)},$$

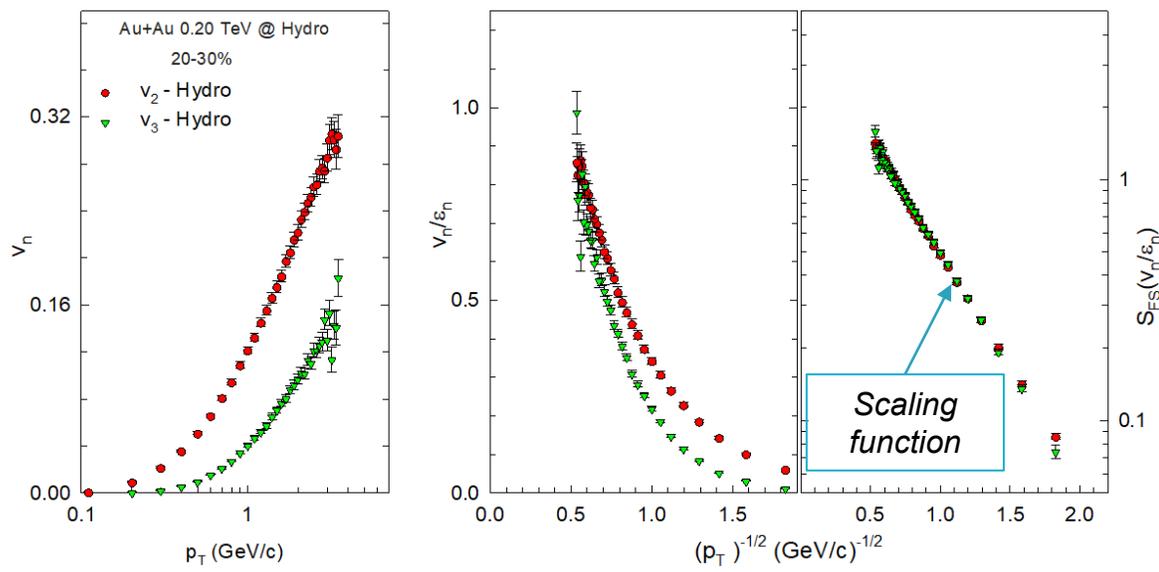
Simulated data [for charged hadrons] from Bjoern Schenke et al.



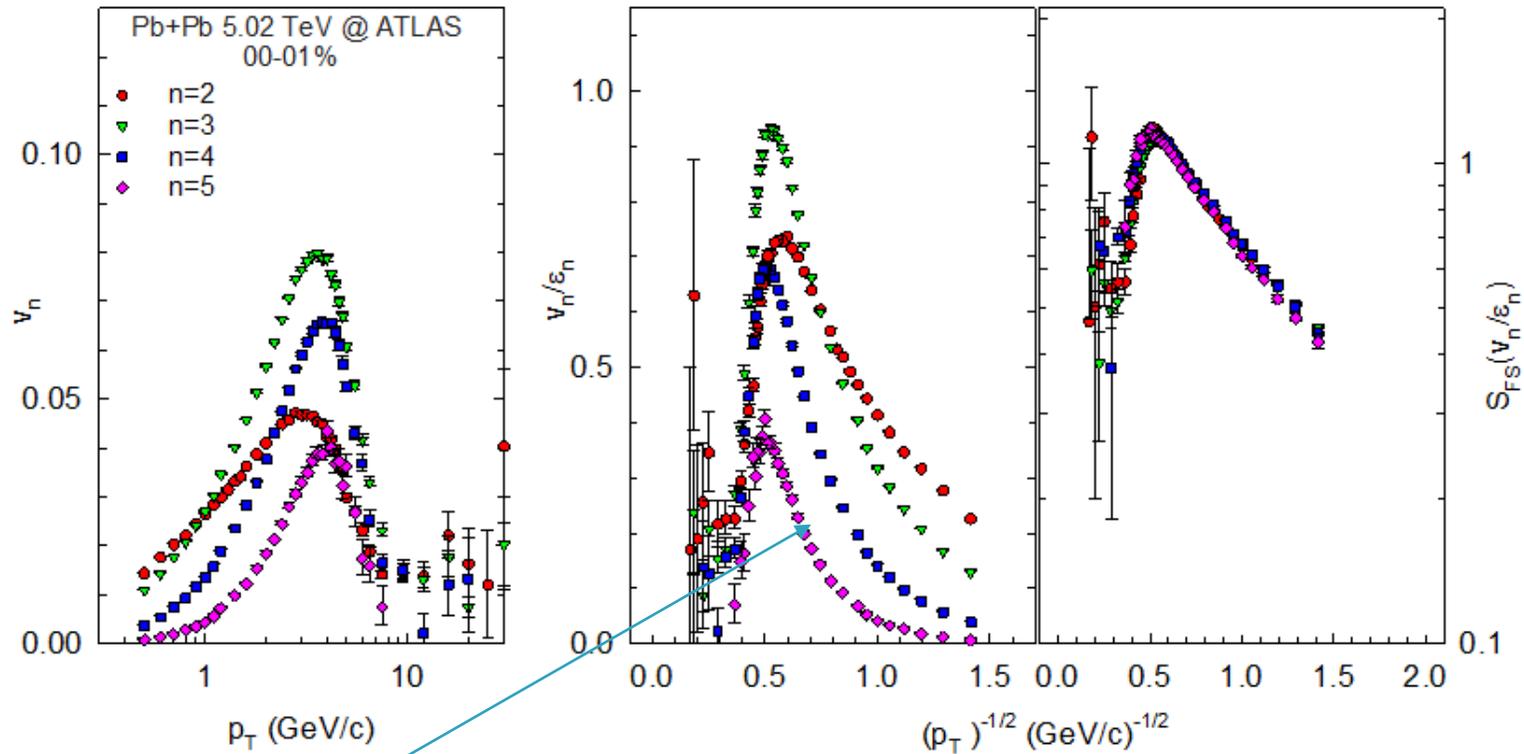
- ✓ **Eccentricity scaling (alone) is insufficient**
- Final-state interactions are crucial**
- ✓ **Same $\frac{\eta}{s}$ for v_2 & v_3**

➤ **Anisotropy data as a function of control variables collapse on to a single curve for fully constrained scaling coefficients**

✓ **Scaling coefficients are proportional to the respective transport coefficients**



n & p_T dependence



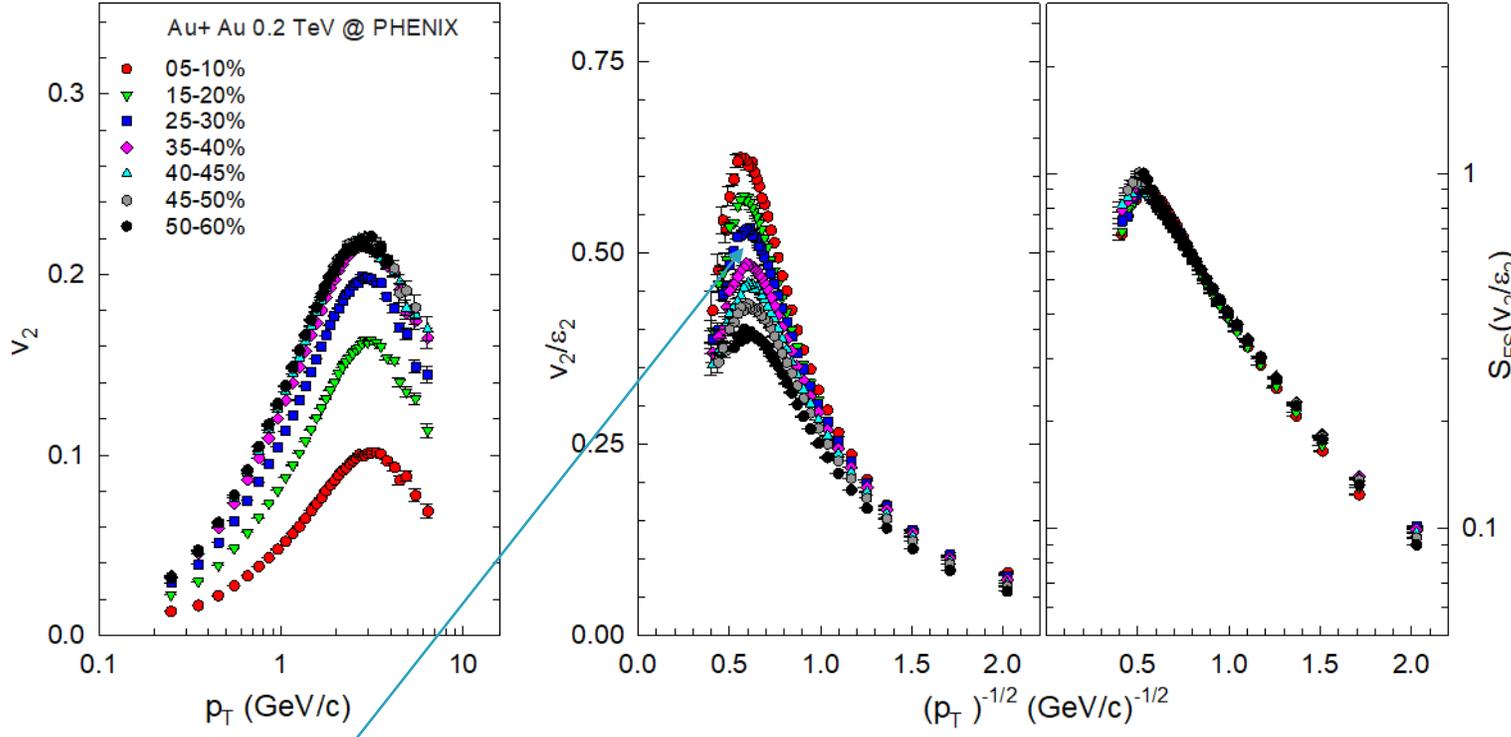
n^2 -dependent viscous attenuation apparent

- **Characteristic patterns of viscous attenuation validated**
 - ✓ **n^2 dependence of the viscous coefficient confirmed**
 - ✓ **Small influence from expansion dynamics**
 - ✓ **Scaling breaks when mode-coupled harmonics contribute**
- **Scaling coefficient provides constraint for $\frac{\eta}{s}$ (T)**

Centrality & p_T dependence

$$\frac{v_n(p_T)}{\varepsilon_n} = \left(\frac{v'_n(p_T)}{\varepsilon'_n} \right) e^{-\frac{\beta}{\mathbb{R}} [n(n+\kappa p_T^2)] \left(1 - \frac{\mathbb{R}}{\mathbb{R}'} \right)},$$

$$\ln \left(\frac{v_n}{\varepsilon_n} \right) = -\frac{\beta}{\mathbb{R}} \left[n(n + \kappa p_T^2) \right], RT \equiv \mathbb{R} \propto \langle N_{\text{ch}} \rangle^{1/3}$$

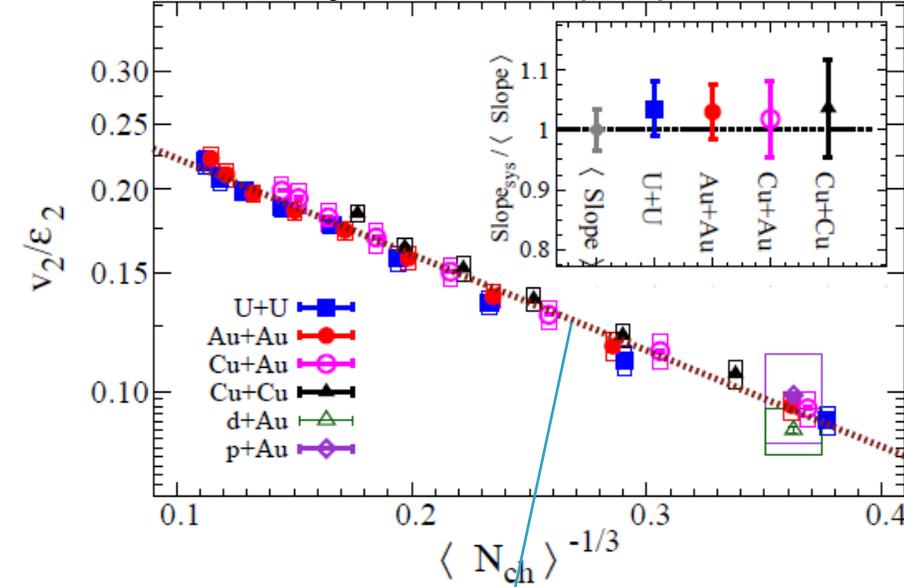


Centrality/ p_T -dependent viscous attenuation

Similar patterns for other beam energies

Collision-system & centrality dependence of p_T -integrated v_2/ε_2

STAR - Phys.Rev.Lett. 122 (2019) no.17, 172301



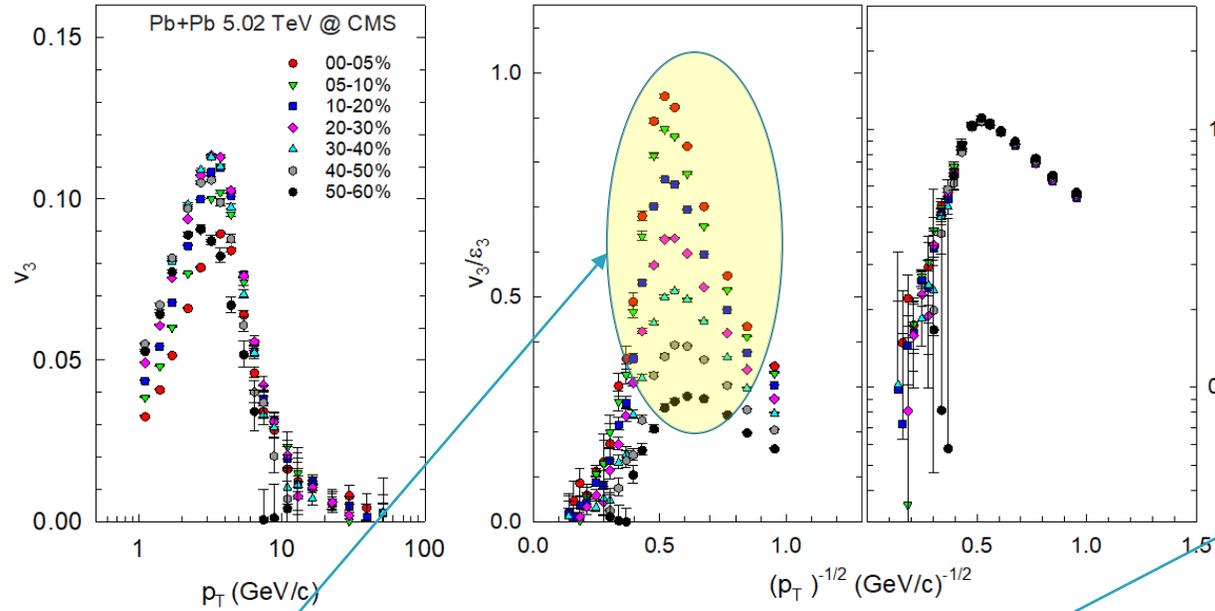
System-independent but RT -dependent viscous attenuation?

➤ Characteristic patterns of viscous attenuation validated

✓ $1/RT$ dependence of the viscous coefficient confirmed

❖ Scaling coefficient provides constraint for $\frac{\eta}{s}$ (T)

Anisotropy Scaling Functions

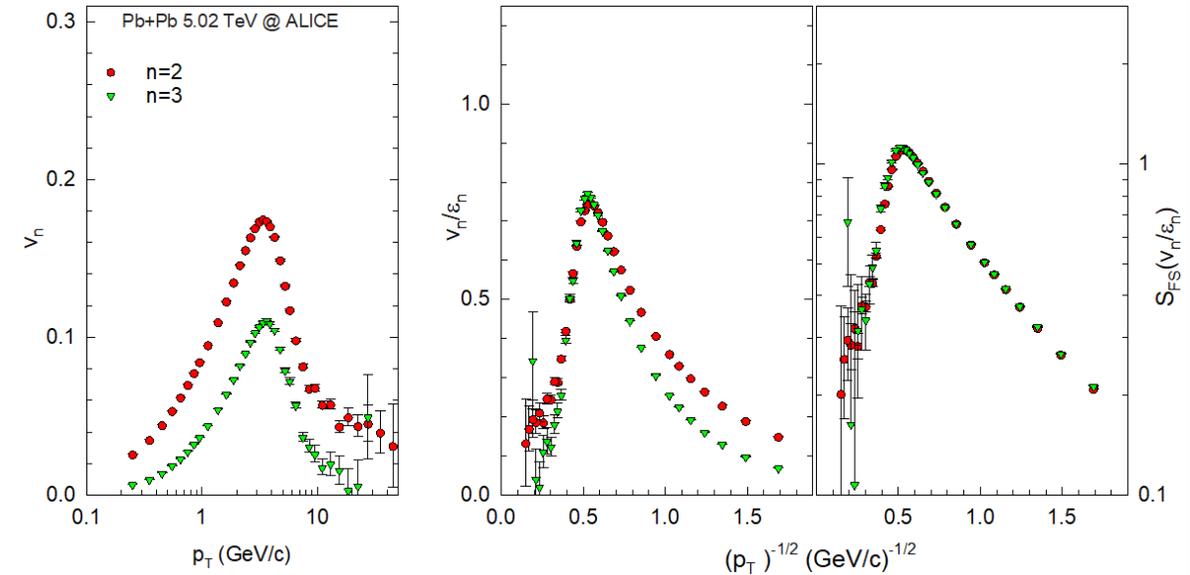
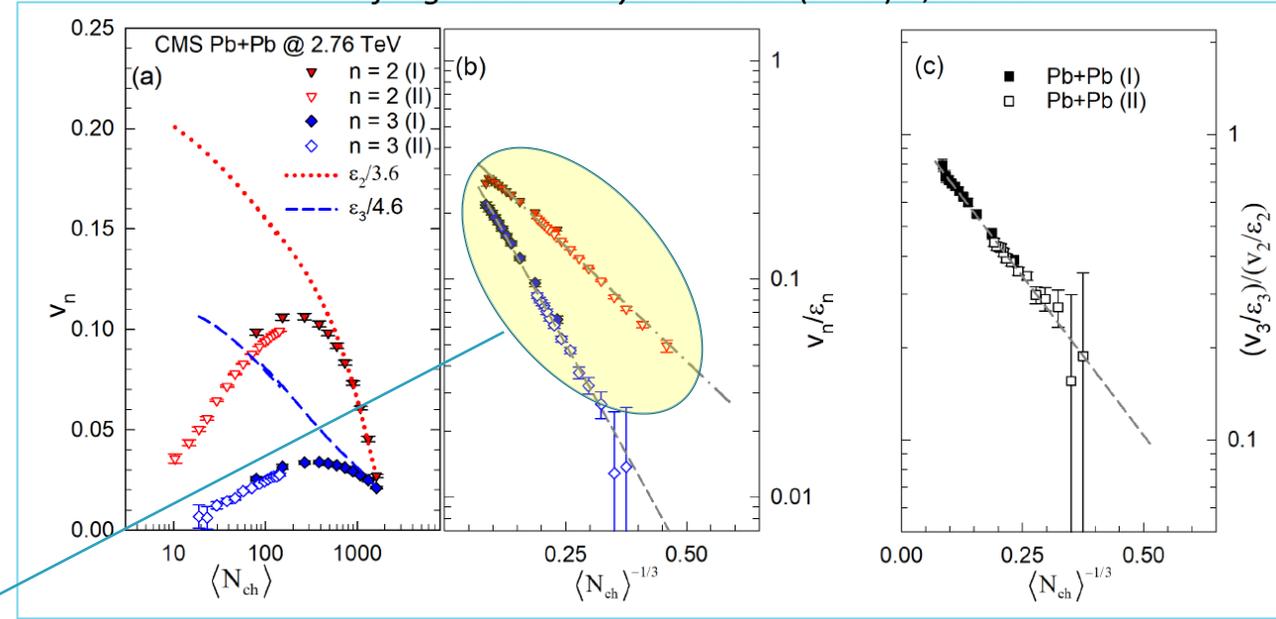


p_T - & RT -dependent viscous attenuation?

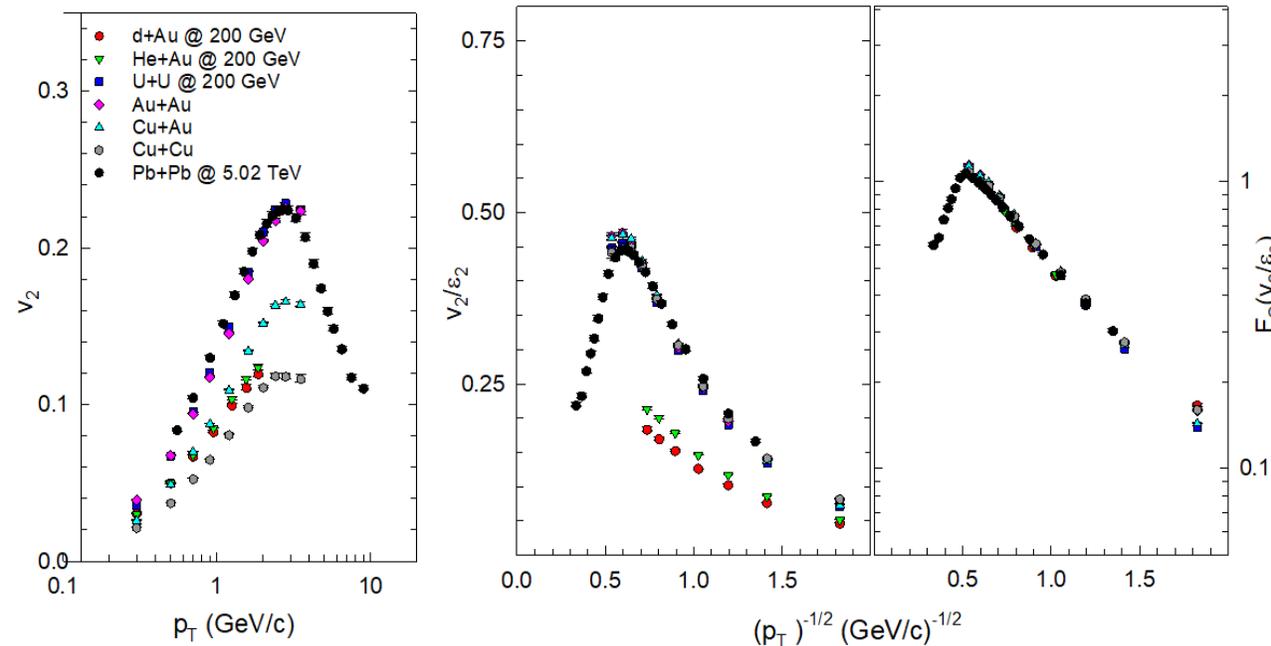
n - & RT -dependent viscous attenuation?

➤ Characteristic patterns of viscous attenuation validated

❖ Scaling coefficient provides constraint for $\frac{\eta}{s}$ (T)

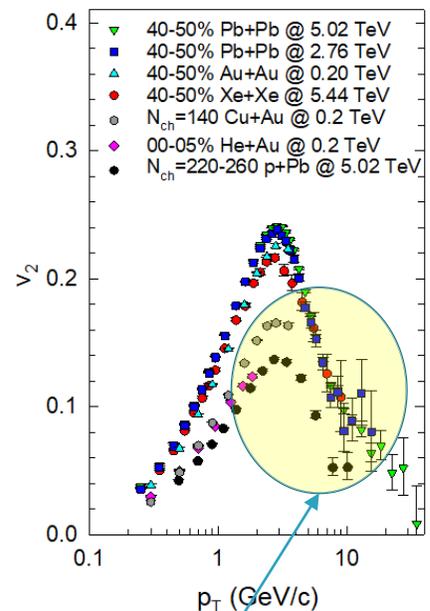


Anisotropy Scaling Functions – Systems & Energies

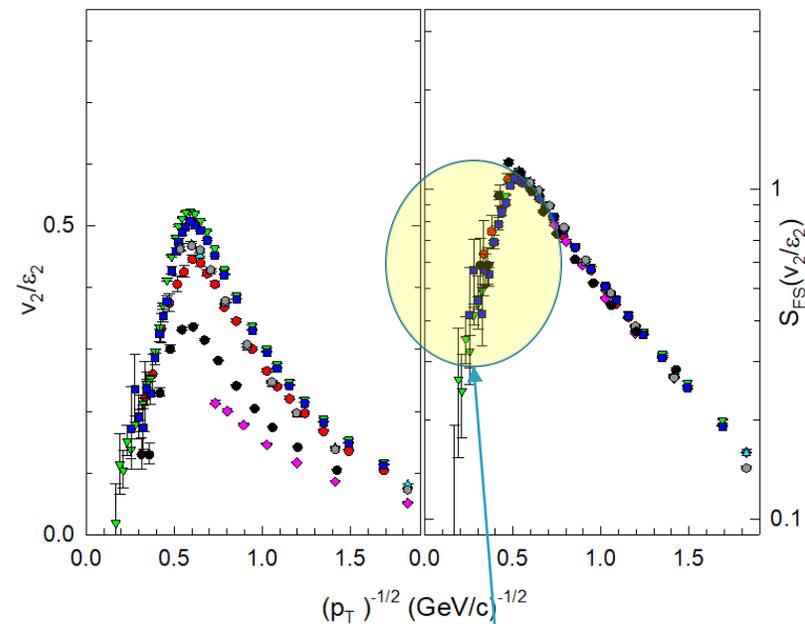


✓ Same $\langle N_{chg} \rangle$ for U+U, Pb+Pb, Au+Au, Cu+Au and Cu+Cu

✓ Different $\langle N_{chg} \rangle$ for d(³He)+Au



➤ High- p_T p+Pb



➤ Jet quenching in p+Pb?

➤ Indications for viscous attenuation and jet quenching across systems.

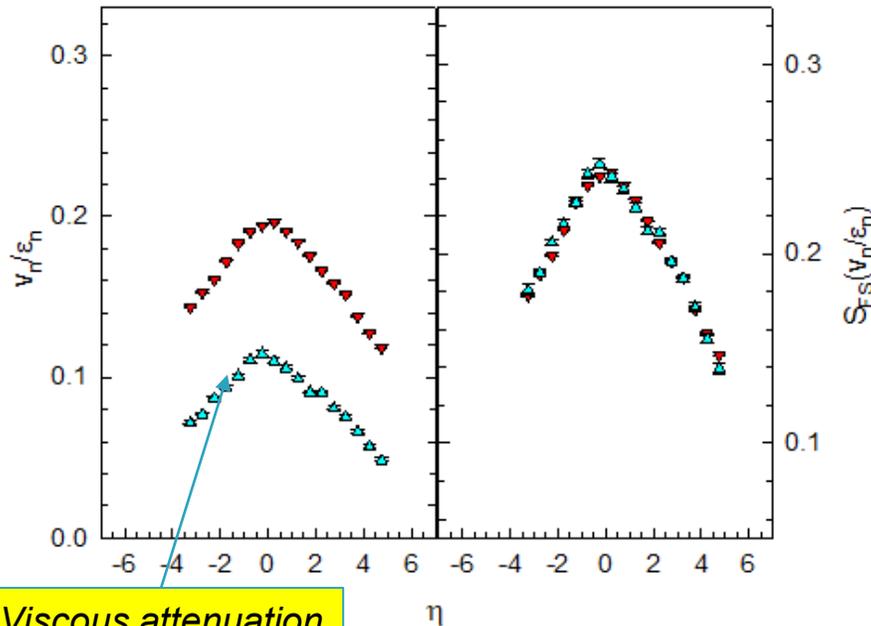
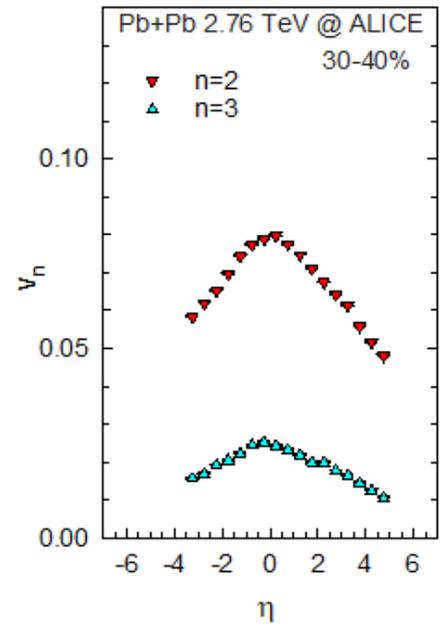
✓ Signal attenuation very important for small dimensionless sizes.

❖ *Scaling coefficients indicate;*

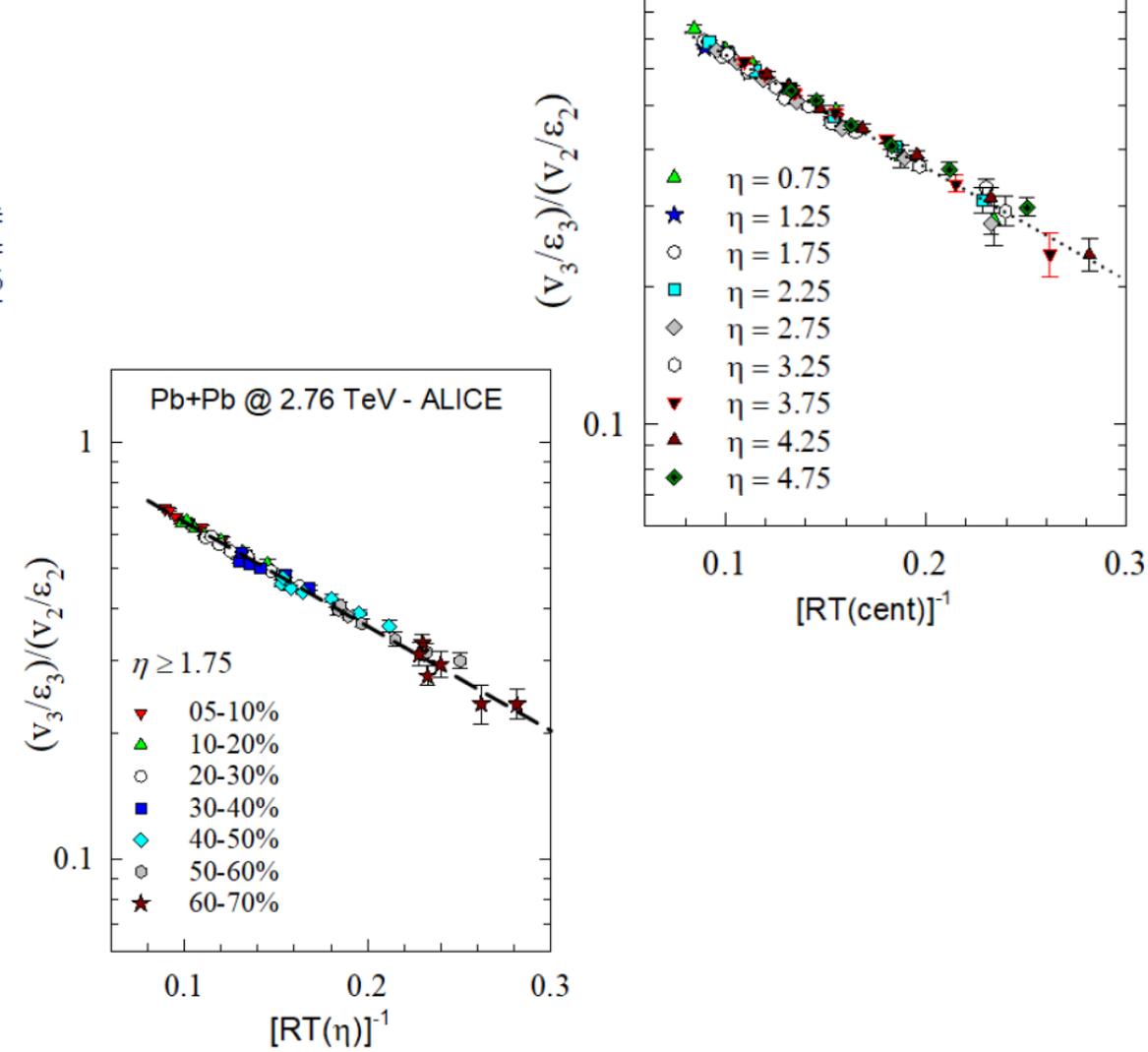
✓ *an increase in η/s from RHIC to LHC.*

✓ *A modest increase in η/s from large to small systems.*

η - dependence



Viscous attenuation difference



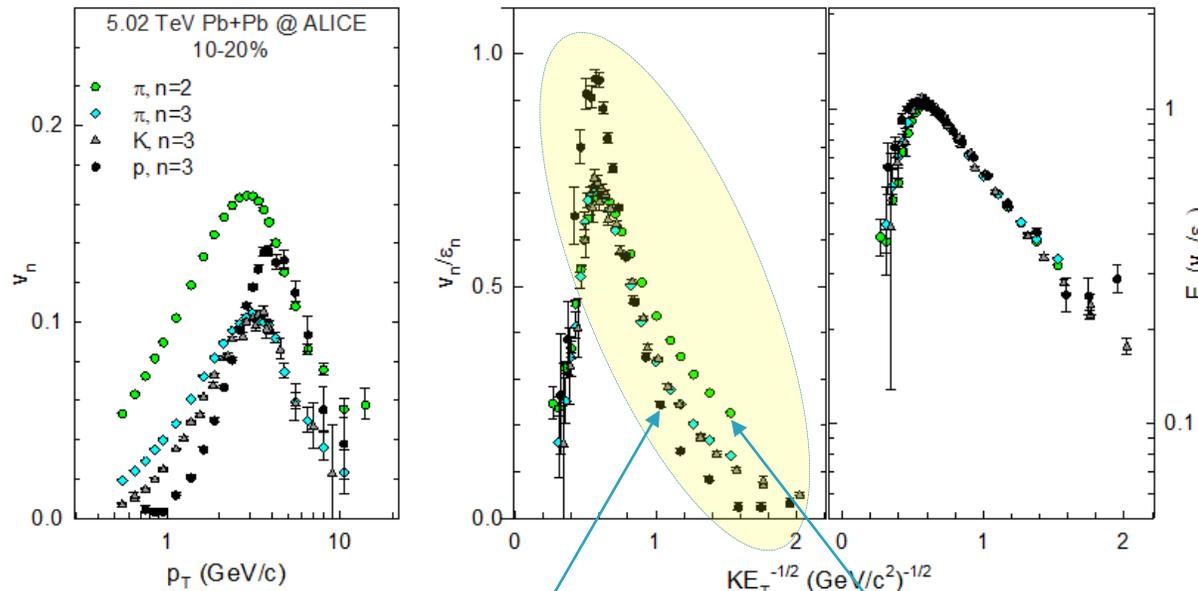
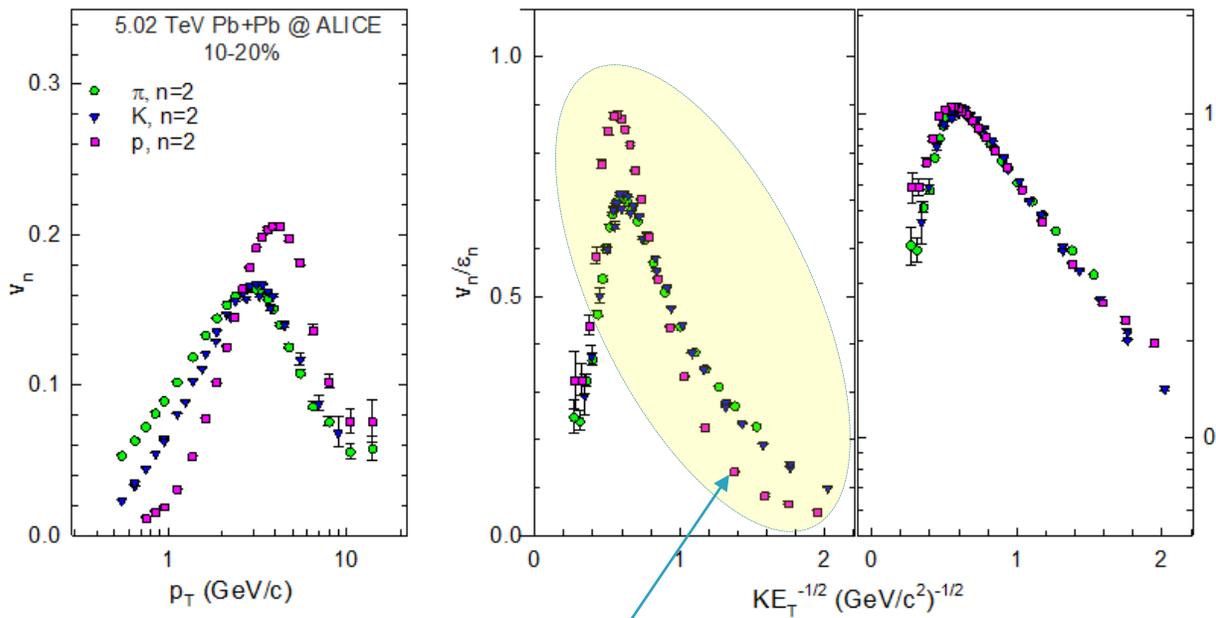
➤ η – dependent patterns of viscous attenuation validated

- ✓ Scaling coefficient provide a further constraint for $\frac{\eta}{s}$ (T)
- ✓ Factorization across η does not hold

Viscous_attenuation(cent, p_T , η) understood!

Anisotropy Scaling Function for PID species

Compare v_2 and v_3



Effects of expansion dynamics
 ✓ Scaling for this effect required

✓ Same $\frac{\eta}{s}$

Start @ 5 TeV

- ✓ $\mu_{B,S,I} \sim 0$
- ✓ PID-independent control variables
- ✓ PID-dependent expansion dynamics

Effects of expansion dynamics

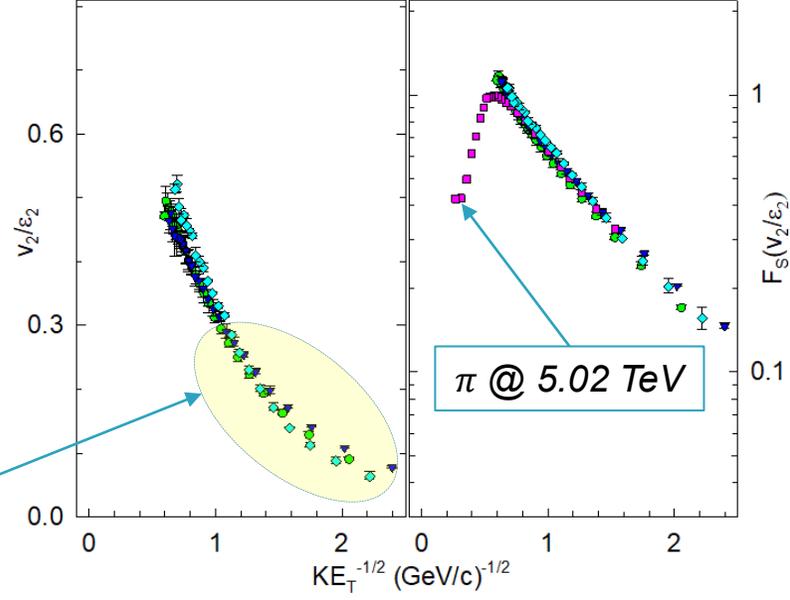
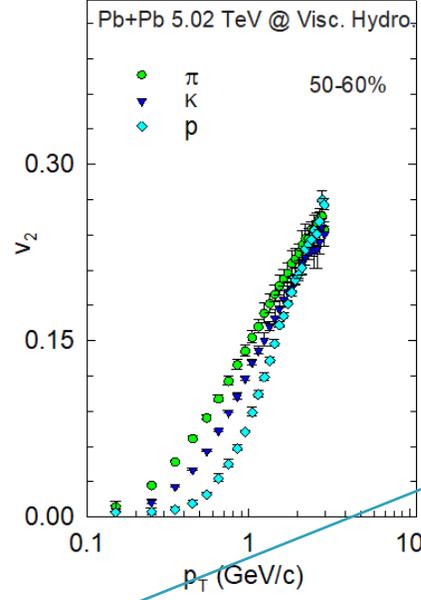
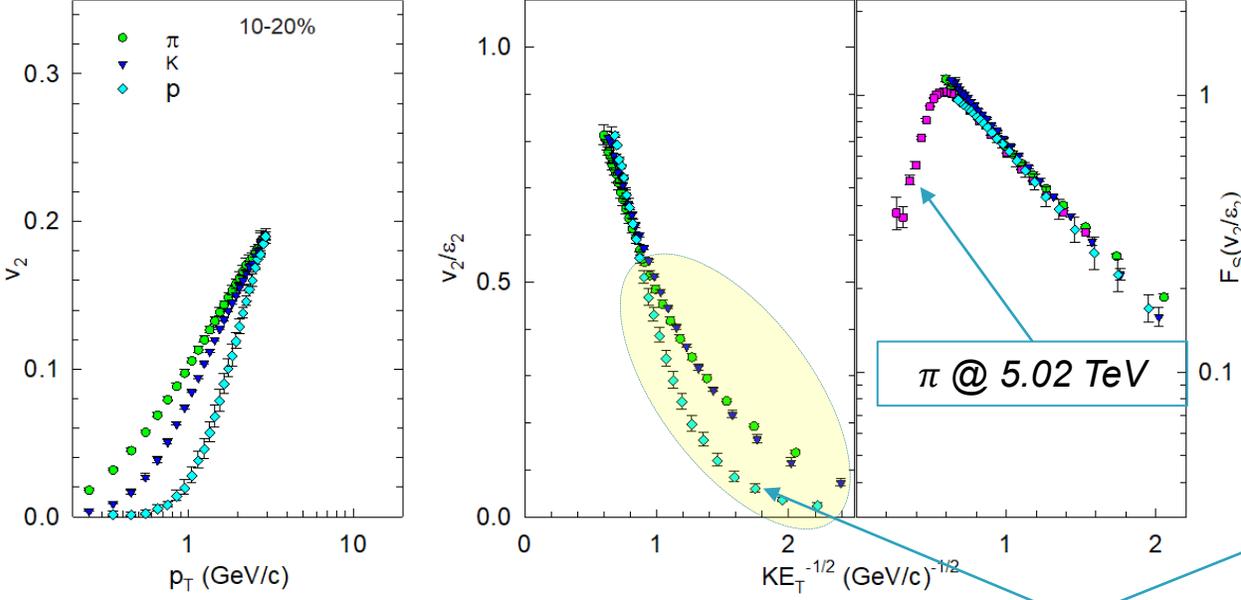
Effects of Viscous attenuation

➤ Scaling for viscous attenuation and expansion dynamics validated.

PID Scaling Functions – Further proof of principle

Simulated data [for identified particles] from Huichao Song et al.

Pb+Pb 5.02 TeV @ Hydrodynamics

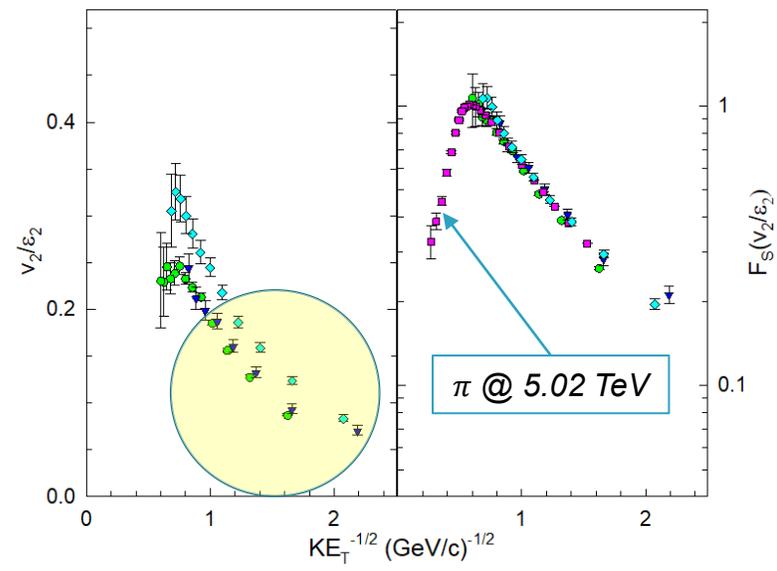
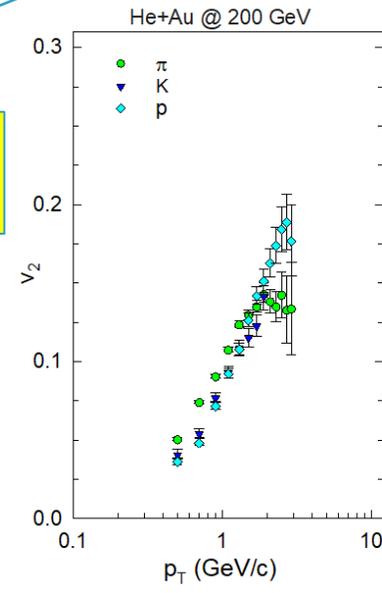


- ✓ $\mu_{B,S,I} \sim 0$
- ✓ PID-independent control variables
- ✓ PID-dependent expansion dynamics

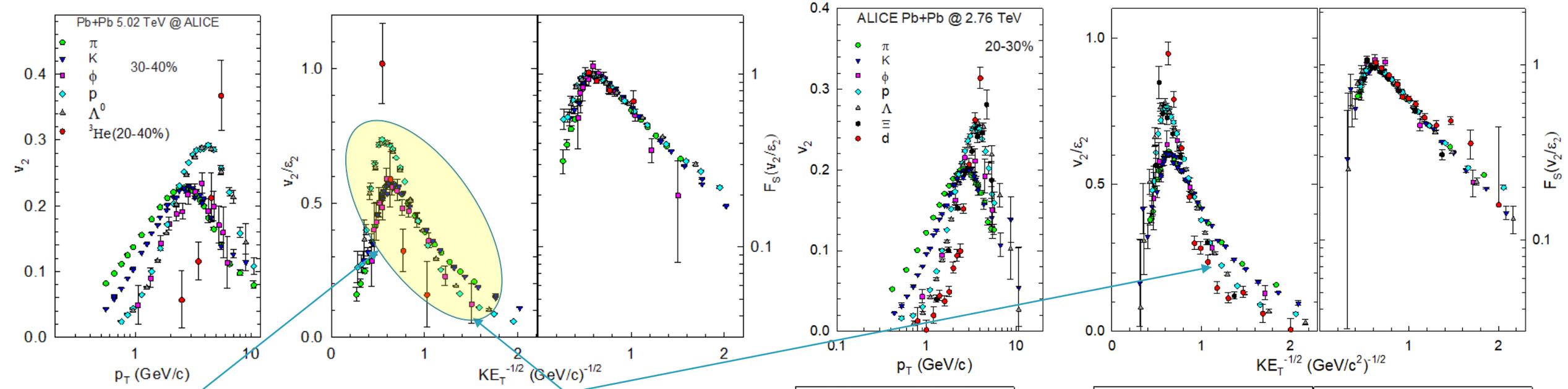
Diff. expansion dynamics
✓ Same η/s

➤ Scaling for viscous attenuation and different expansion dynamics validated.

✓ Expansion dynamics is centrality-, system- and beam-energy dependent



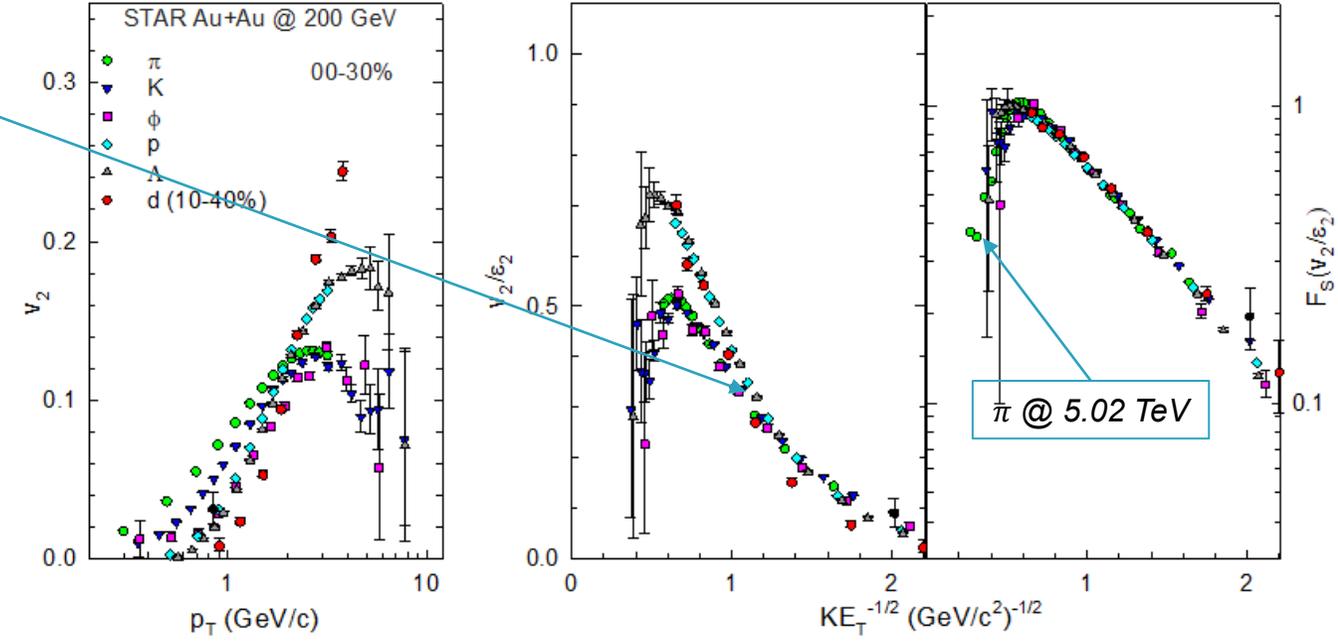
Anisotropy Scaling Functions – species across systems & energies



Diff. expansion dynamics
 ✓ Same η/s ?

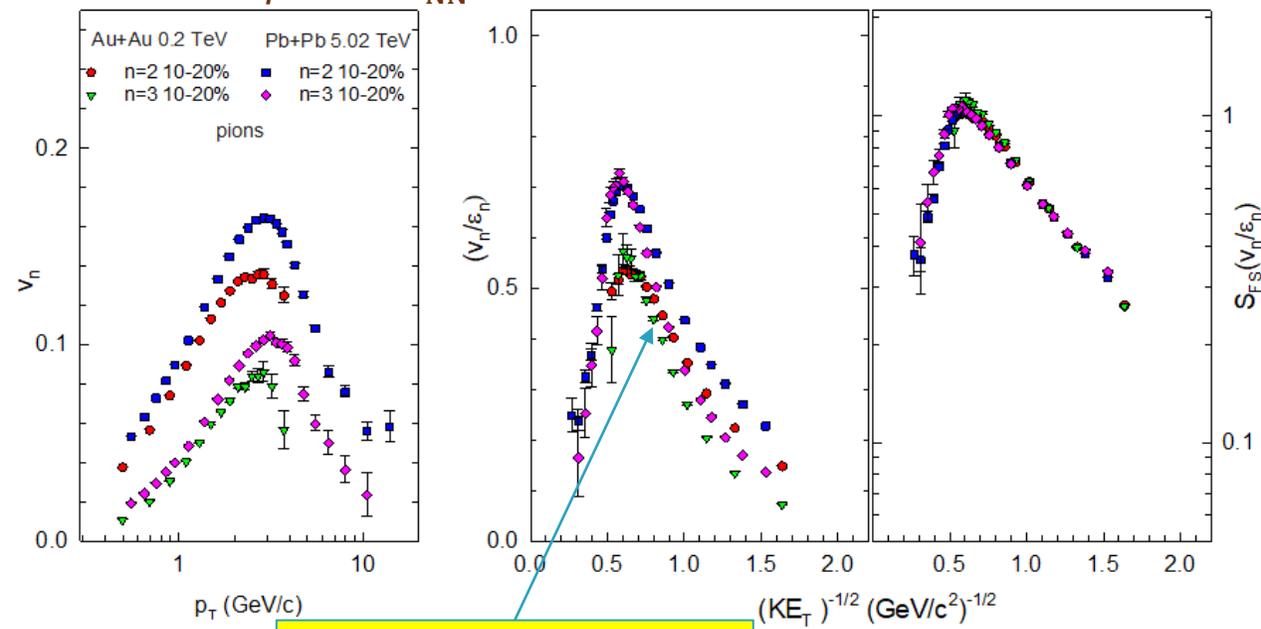
✓ 1 meson branch
 ✓ 2 baryon branches

- **Scaling for viscous attenuation and expansion dynamics validated for all beam energies studied**
 - ✓ **Scaling observed for all particle species (multiple mechanisms superfluous)**
 - ✓ **Baryon-number-dependent expansion dynamics - not mass**



PID Scaling Functions – Systems, species & Energies

Compare $\sqrt{s_{NN}} = 0.2$ and 5.02 TeV



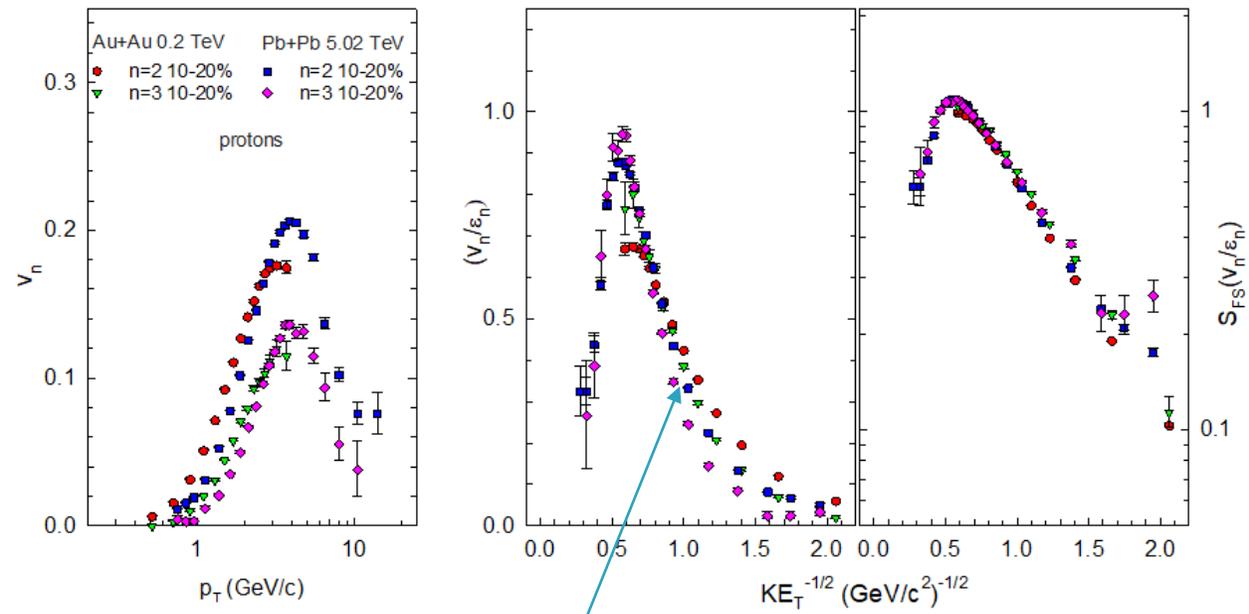
Indications for diff. η/s

➤ Scaling for viscous attenuation and expansion dynamics validated across beam energies and systems

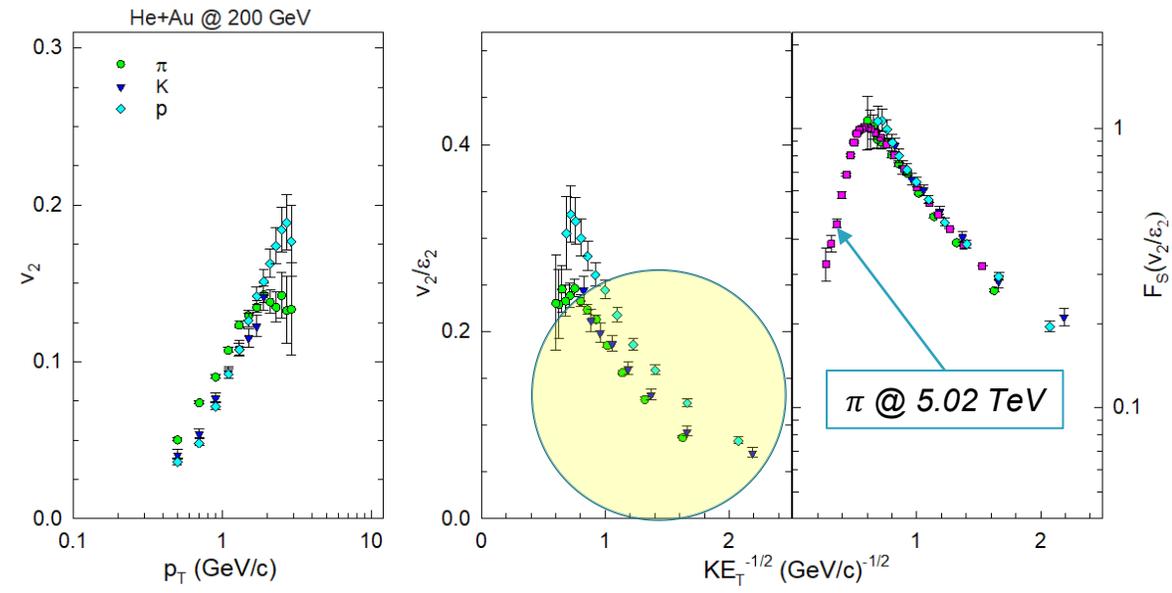
✓ Scaling observed for all particle species (multiple mechanisms superfluous)

❖ Scaling coefficients indicate;

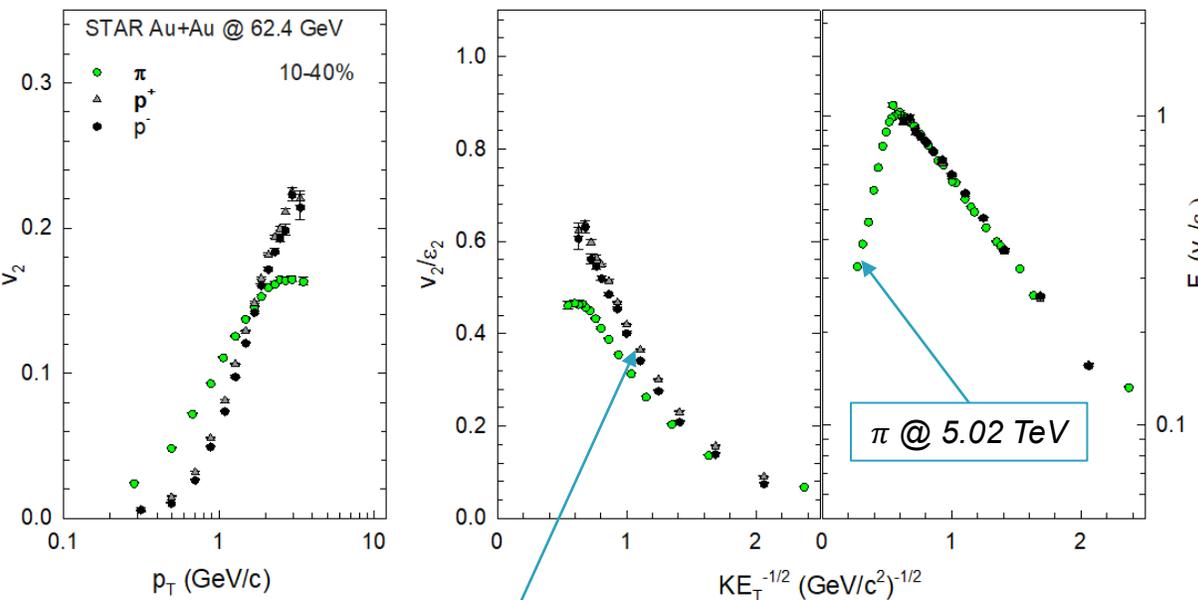
- ✓ an increase in η/s from RHIC to LHC.
- ✓ A modest increase in η/s from large to small systems.



Diff. η/s & expansion dynamics

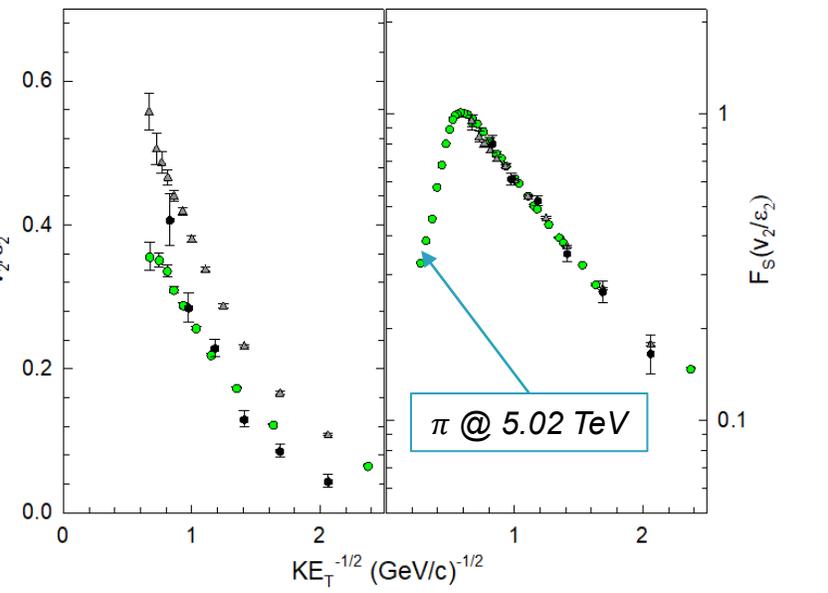
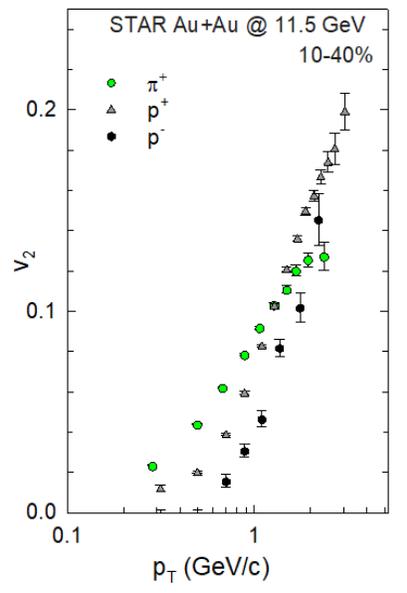
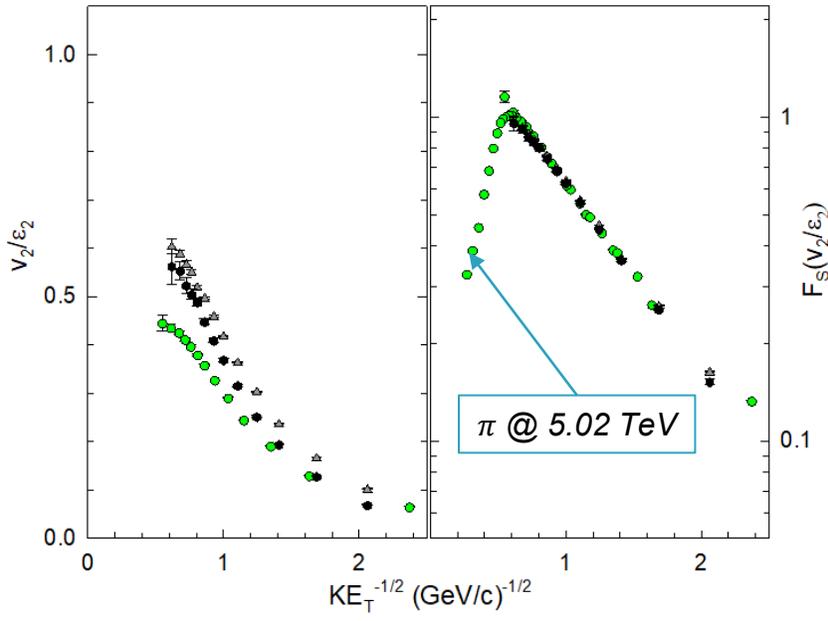
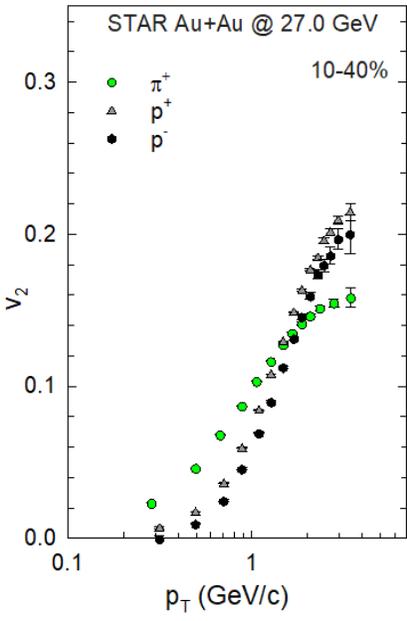


PID Scaling Functions – particles vs. anti-particles



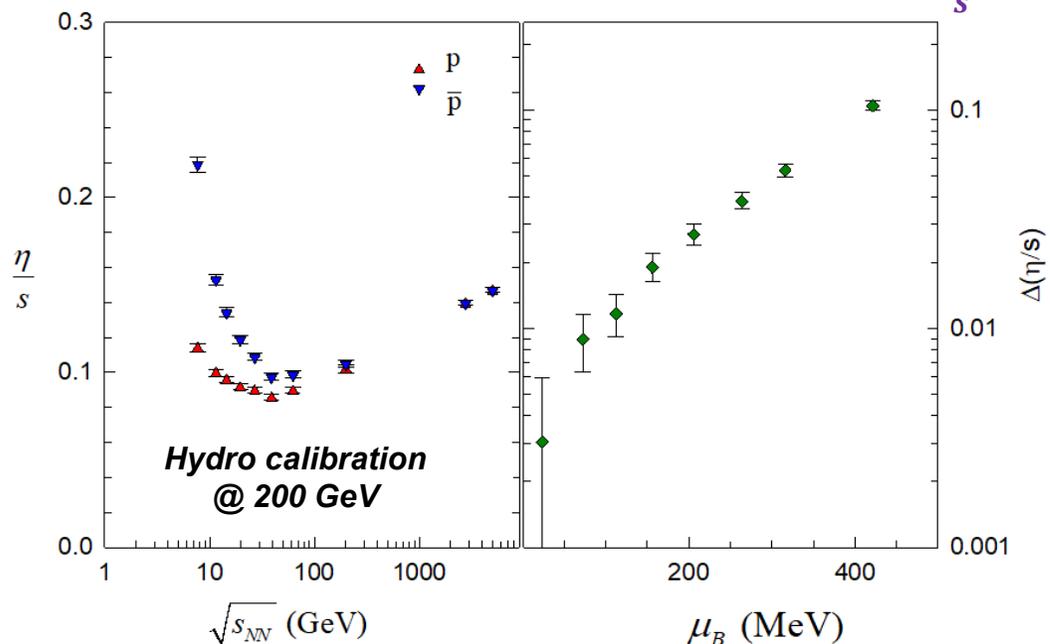
Particle/Anti-particle difference
 ✓ Same expansion dynamics

- **Scaling for viscous attenuation and expansion dynamics validated for particles and anti-particles across beam energies**
 - ✓ **Scaling coefficients indicate particle/anti-particle-dependent viscosity**

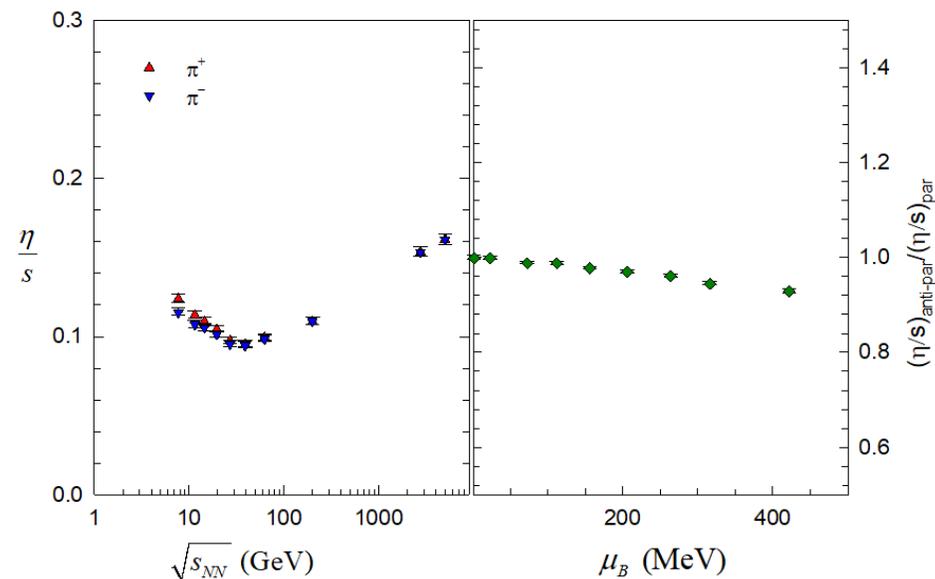
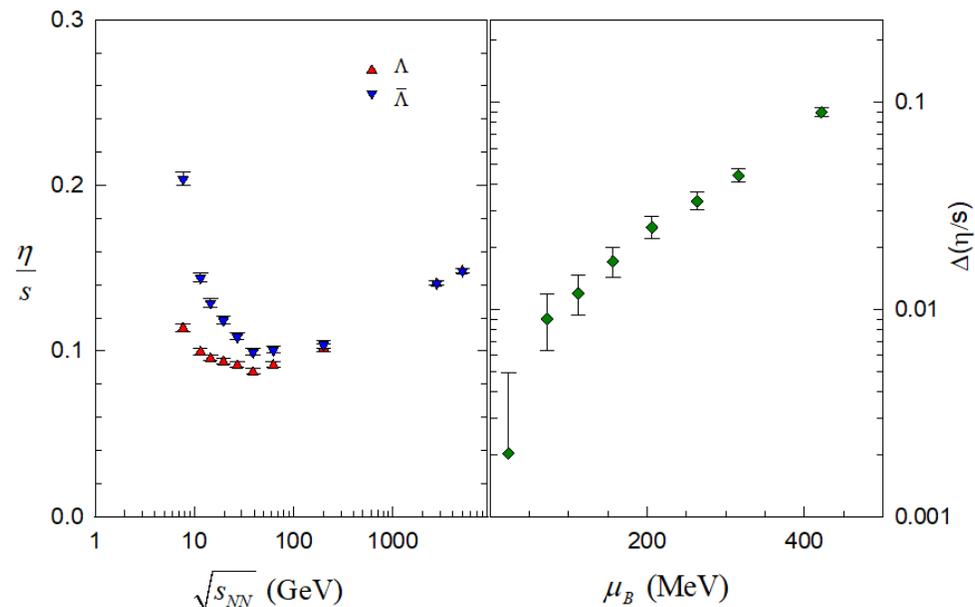


Beam-energy-dependent specific viscosity

$\sqrt{s_{NN}}$ dependence of the extracted values of $\frac{\eta}{s}$



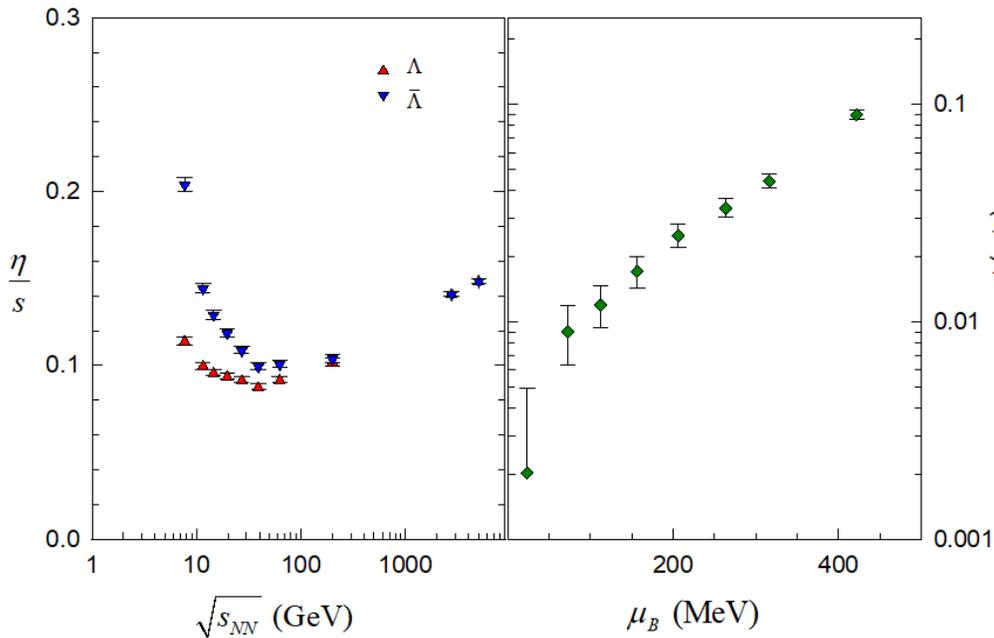
$$\frac{\eta}{s}(T, \mu_B, \mu_S, \mu_I)$$



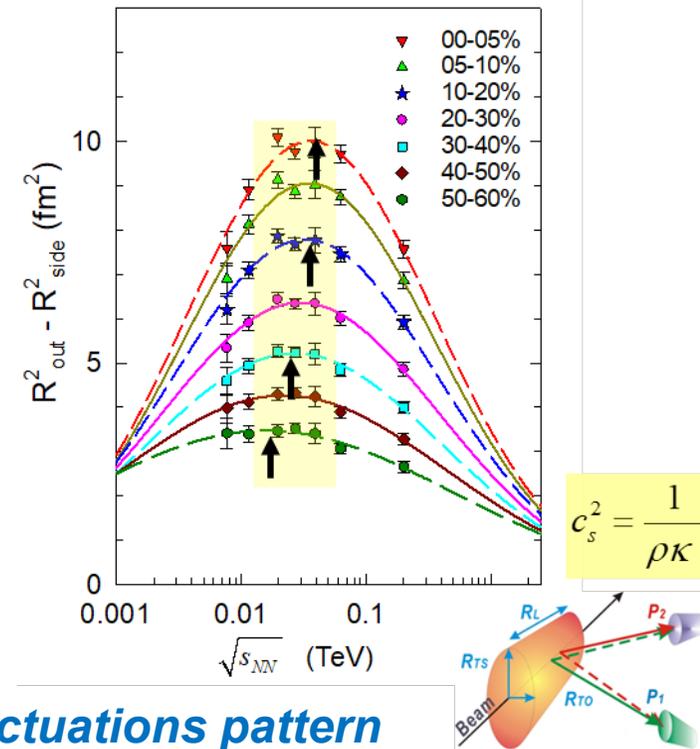
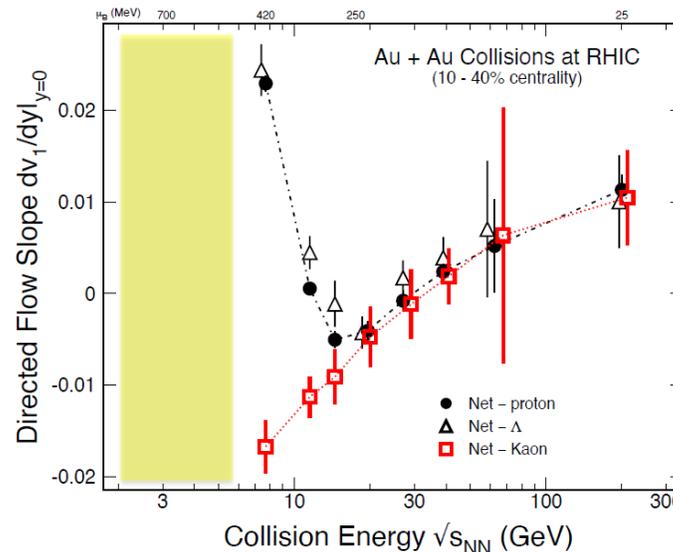
- Specific viscosity extracted across beam energies
 - ✓ Nonmonotonic patterns suggestive of critical behavior?
 - ✓ μ_B –dependent particle/anti-particle dependence

Charged currents drive particle/anti-particle viscosity difference

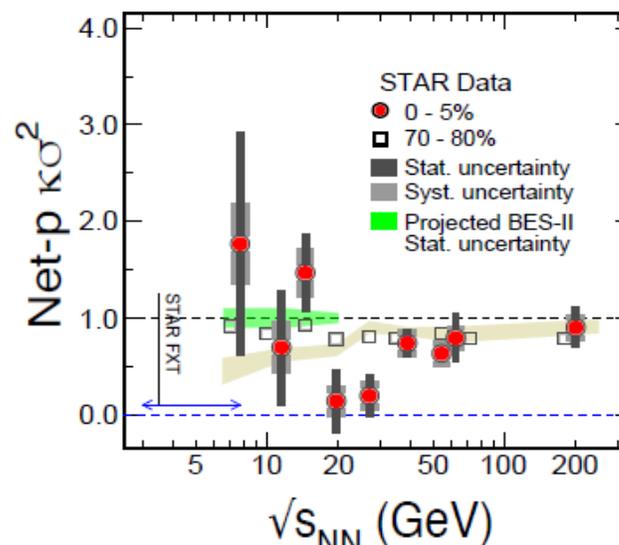
Non-monotonic patterns



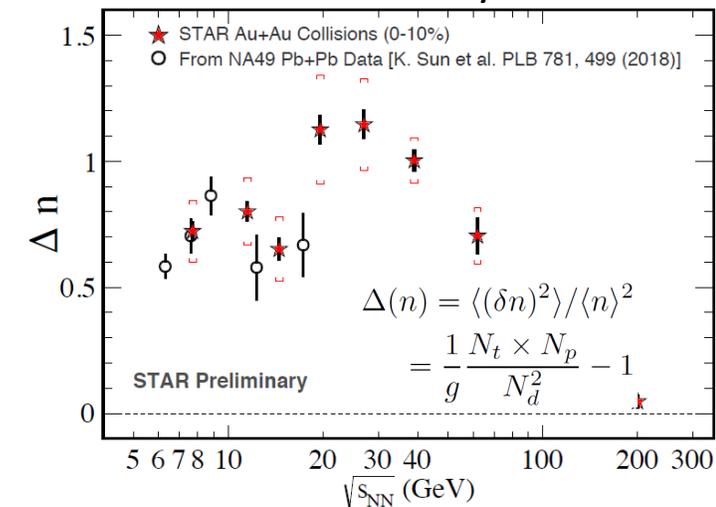
Other dynamics-driven non-monotonic patterns



Non-monotonic fluctuations pattern



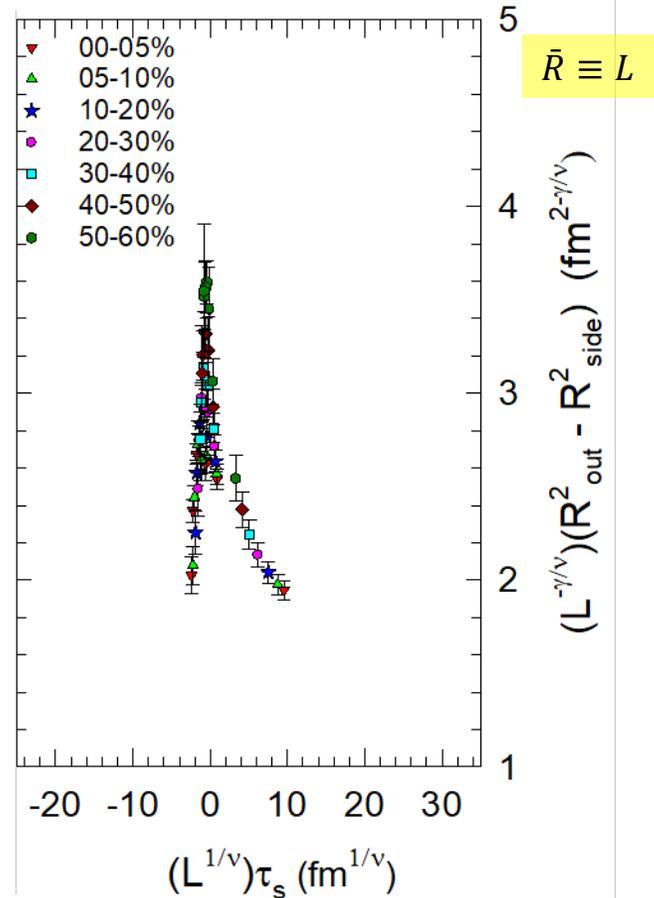
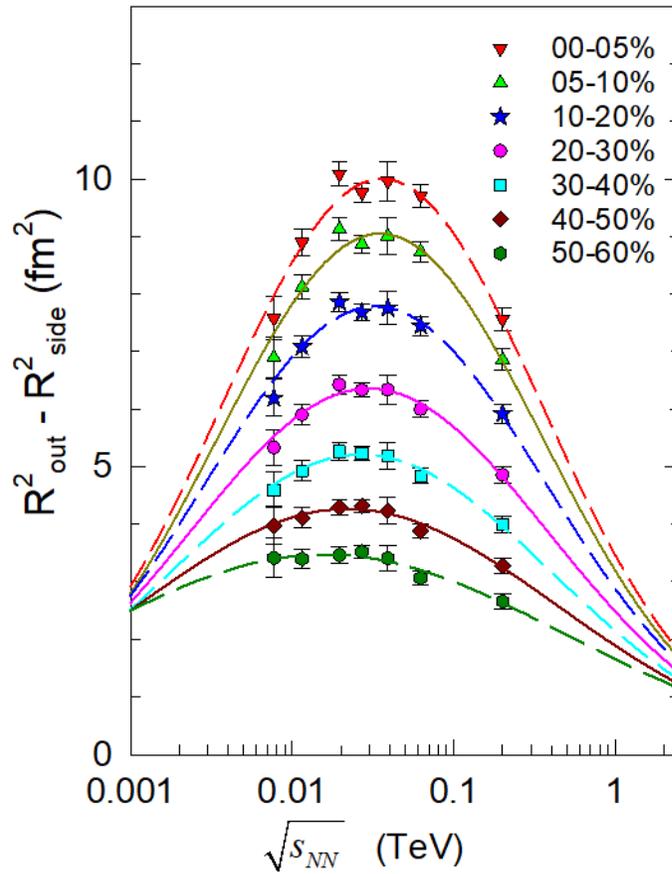
Neutron relative density fluctuations



- Anisotropy scaling functions indispensable for extraction of transport coefficients
 - ✓ Non-monotonic pattern observed → CEP?
 - ✓ Charged current dependence @ lower $\sqrt{s_{NN}}$
- Persistent non-monotonic pattern observed for several observables in similar $\sqrt{s_{NN}}$ range → CEP?

Succeptibility Scaling Function

$$L^{-\gamma/\nu} \chi(s, L) = f_2^s(sL^{1/\nu})$$



$$\sqrt{s_{CEP}} = 45 \text{ GeV}$$

$$\nu \sim 0.66$$

$$\gamma \sim 1.2$$

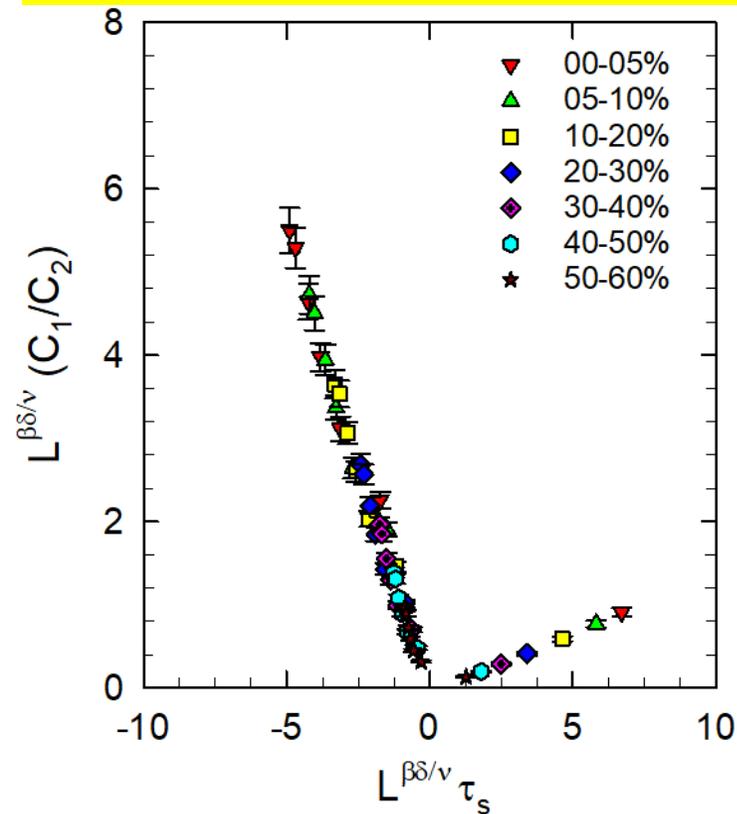
$$s = (\sqrt{s} - \sqrt{s_{CEP}}) / \sqrt{s_{CEP}}$$

Finite-size scaling leads to data collapse for 3D Ising critical exponents.

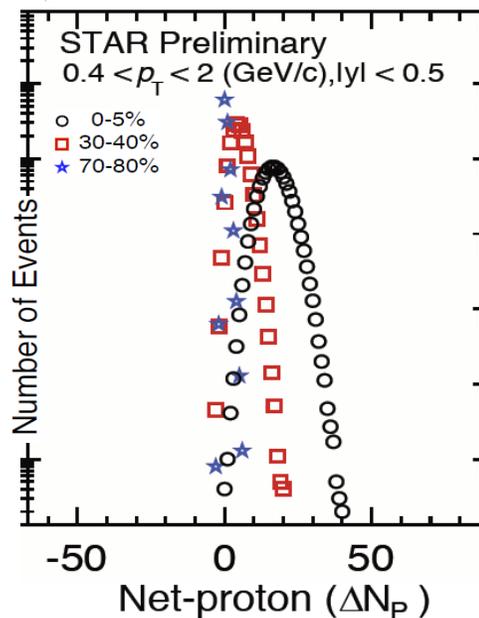
✓ **Slow quench in KZ language**

Scaling Function for Net baryon Susceptibility ratio

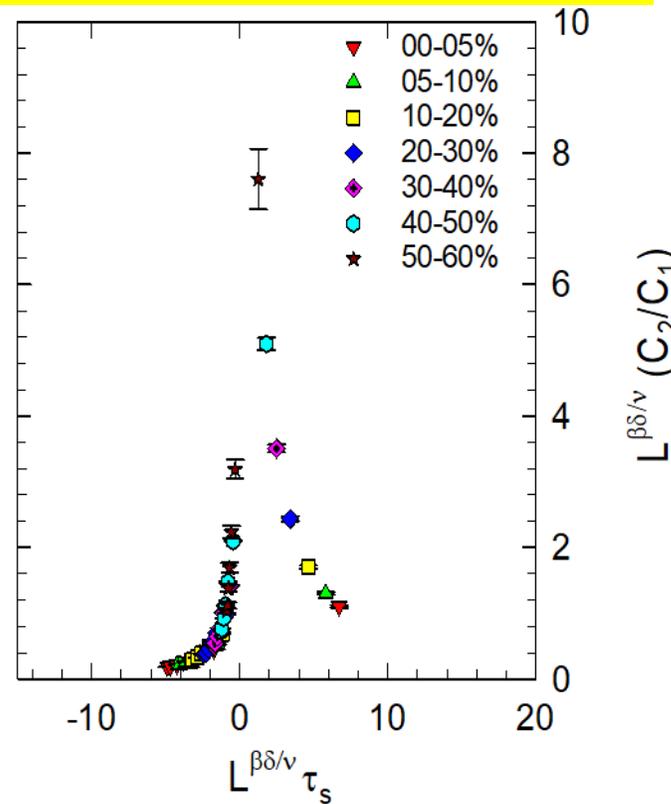
$$\chi(\mu_s, L) = L^{-\beta\delta/\nu} f_1^\mu(\mu_s L^{\beta\delta/\nu})$$



$C_n^{\Delta N_p}$ extracted
from distributions



$$\chi(\mu_s, L) = L^{\beta\delta/\nu} f_1^\mu(\mu_s L^{\beta\delta/\nu})$$



$$\nu \sim 0.66$$

$$\beta \sim 0.33$$

$$\delta \sim 4.8$$

$$\kappa_T \propto \frac{\langle C_2^{\Delta N_p} \rangle - \langle C_2^{\Delta N_p} \rangle^2}{\langle \Delta N_p \rangle} = \frac{C_2^{\Delta N_p}}{C_1^{\Delta N_p}}$$

Finite-size scaling leads to data collapse for 3D Ising critical exponents.

✓ Slow quench in KZ language

Summary

The $\sqrt{s_{\text{NN}}}$ -dependence of the anisotropy scaling functions for PID species can be used to:

- Delineate the respective influence of expansion dynamics and viscous attenuation
- Constrain $\frac{\eta}{s}(T, \mu_B, \mu_I, \mu_S)$?

The scaling functions extracted from the wealth of the flow data indicate:

- ❖ μ –dependent particle/anti-particle $\frac{\eta}{s}$ dependence
- ❖ non-monotonic patterns for;
 - ✓ $\frac{\eta}{s}(T, \mu_B)$,
 - ✓ $\frac{\eta}{s}(T, \mu_S)$,
 - ✓ $\frac{\eta}{s}(T, \mu_I)$,

Consistent with earlier indications for the CEP

End