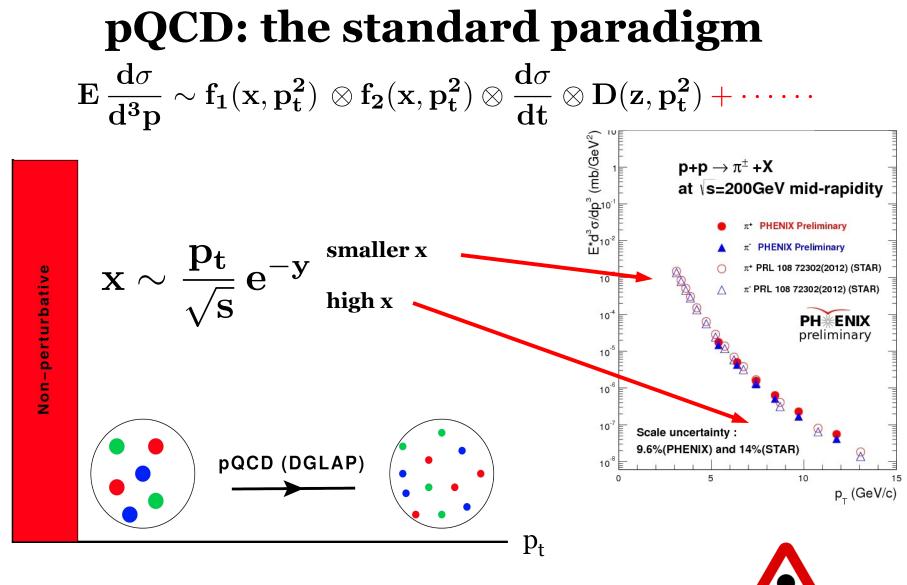
Photon-hadron/jet production and angular correlations in high energy proton-nucleus collisions

(from small to large x: toward a unified formalism)

Jamal Jalilian-Marian

Baruch College and City University of New York Graduate Center

based on: PRD102 (2020) 1, 014008 PRD99 (2019) 1, 014043 PRD96 (2017) 7, 074020 and work in progress

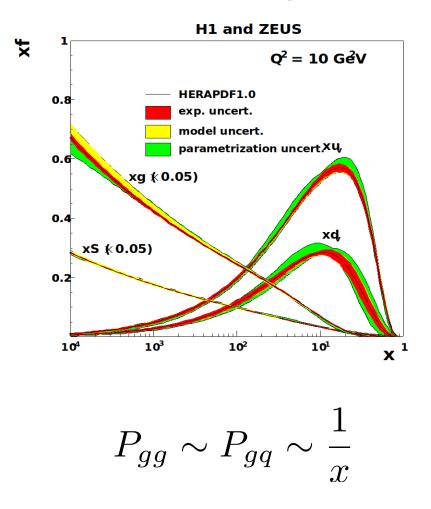


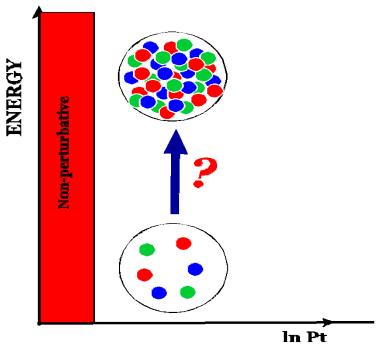
bulk of QCD phenomena happens at low pt (small x)

ENERGY



dynamics of universal gluonic matter: gluon saturation

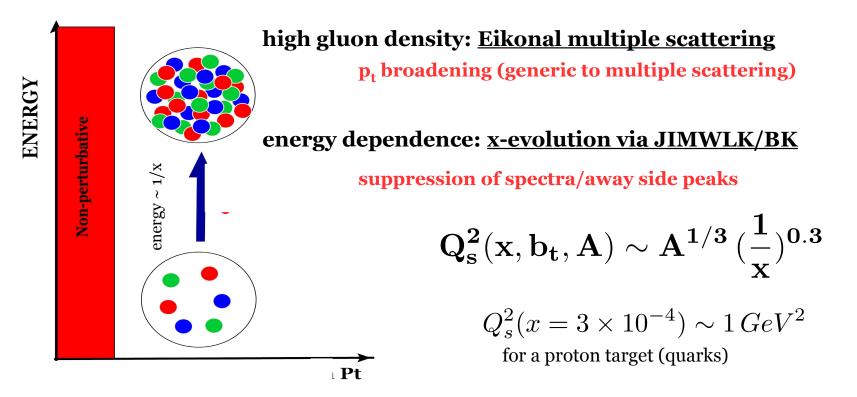




How does this happen?

How do correlation functions evolve ? Is there a universal fixed point for the evolution ? Are there scaling laws ?

QCD at high energy/small x: gluon saturation



a framework for multi-particle production in QCD at small x/low p_t

Shadowing/Nuclear modification factor <u>Azimuthal angular correlations (photon-hadron,...)</u> Long range rapidity correlations (ridge,...) Initial conditions for hydro Thermalization (?)

 $\mathbf{x} \leq \mathbf{0.01}$ $\alpha_s \ln \left(x_v / x \right) \sim 1$

Eikonal approximation (dilute-dense scattering)

$$J_{a}^{\mu} \simeq \delta^{\mu-} \rho_{a}$$

$$D_{\mu} J^{\mu} = D_{-} J^{-} = 0$$

$$\partial_{-} J^{-} = 0 \quad (\text{in } A^{+} = \text{o gauge})$$

$$does \text{ not depend on } x^{-}$$
EOM: solution
$$A_{a}^{-} (x^{+}, x_{t}) \equiv n^{-} S_{a} (x^{+}, x_{t})$$

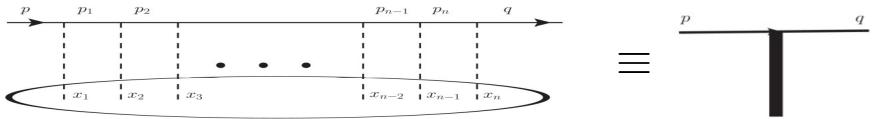
$$n^{\mu} = (n^{+} = 0, n^{-} = 1, n_{t} = 0)$$
recall (eikonal approx):
$$\bar{u}(q) \gamma^{\mu} u(p) \rightarrow \bar{u}(p) \gamma^{\mu} u(p) \sim p^{\mu}$$

$$\bar{u}(q) \mathcal{A} u(p) \rightarrow p \cdot A \sim p^{+} A^{-}$$

scattering of a quark from background color field

 $A_a^-(x^+, x_t)$

CGC: eikonal approximation (tree level)



$$i\mathcal{M}(p,q) = 2\pi\delta(p^+ - q^+)\bar{u}(q)\not h \int d^2x_t \, e^{-i(q_t - p_t)\cdot x_t} \left[V(x_t) - 1\right] u(p)$$

with $V(x_t) \equiv \hat{P} \exp\left\{ig \int_{-\infty}^{+\infty} dx^+ S_a^-(x^+, x_t)t_a\right\}$

scattering from <u>small x gluons</u> of the target can cause only a <u>small angle deflection</u>

DIS, proton-nucleus collisions: dipoles

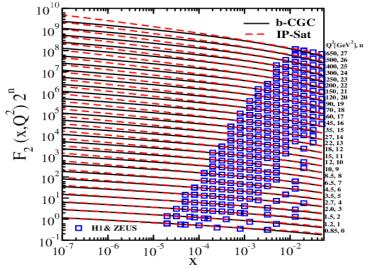
 $< Tr V(x_{\perp}) V^{\dagger}(y_{\perp}) >$

Pre-RHIC: all we know about saturation came from HERA structure function fits

HERA: large kinematic window for x structure functions: $x = x_{bi}$

Still debating whether saturation is seen at HERA!

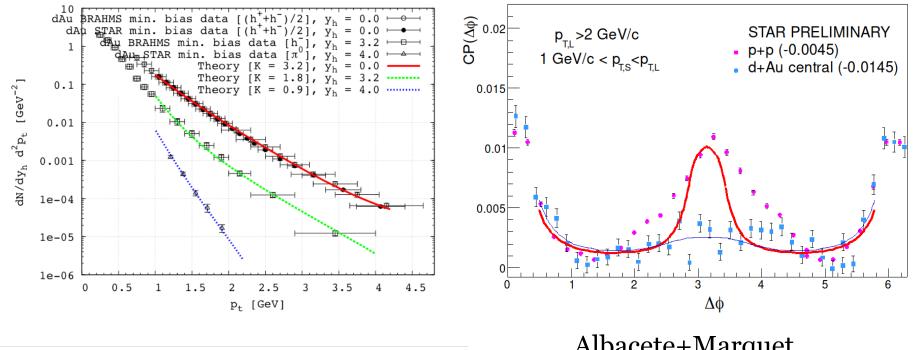
 $\frac{\text{Less-inclusive observables are more discriminatory}}{\mathbf{x} \neq \mathbf{x_{bj}}}$



A. Rezaeian and I. Schmidt, PRD88 (2013) 074016

CGC at RHIC

Single and double inclusive hadron production in dA collisions

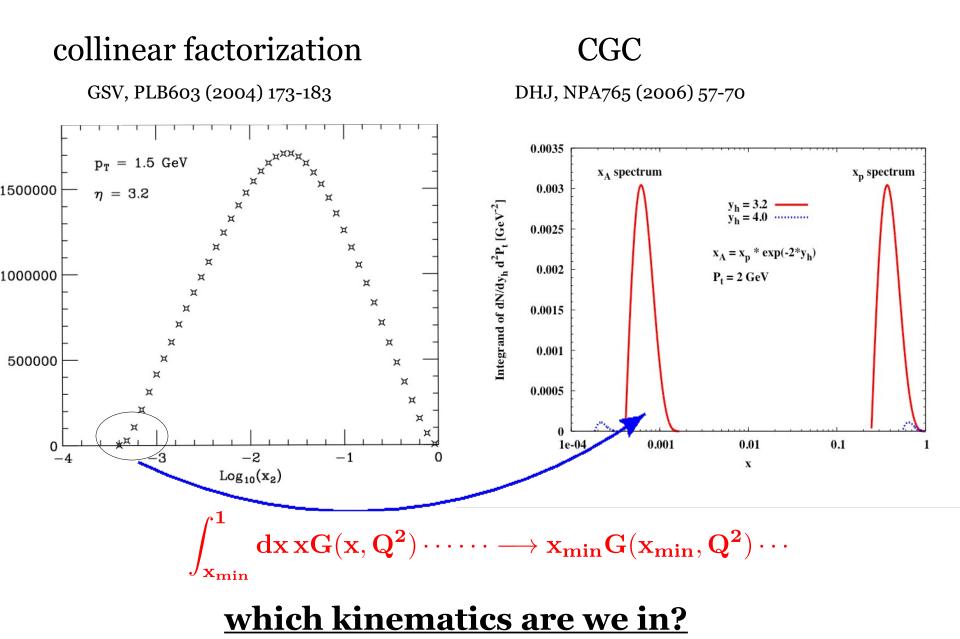


DHJ, NPA770 (2006) 57

Albacete+Marquet PRL105 (2010) 162301

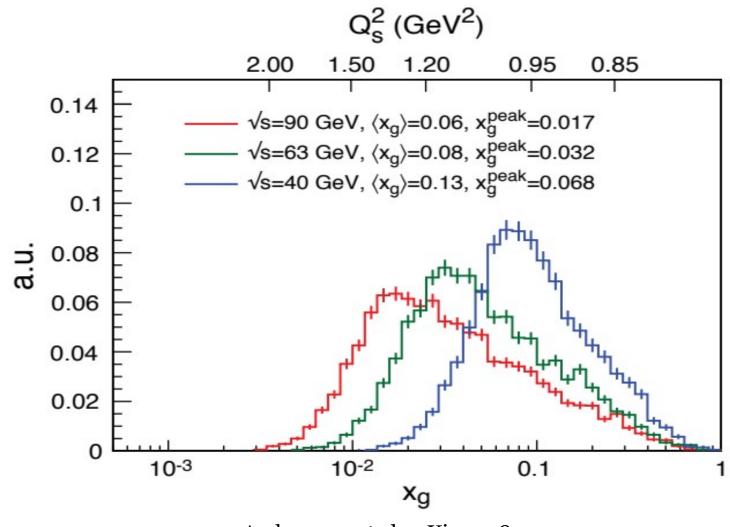
convolution over x: is x small enough?

Single inclusive pion production in pp at RHIC

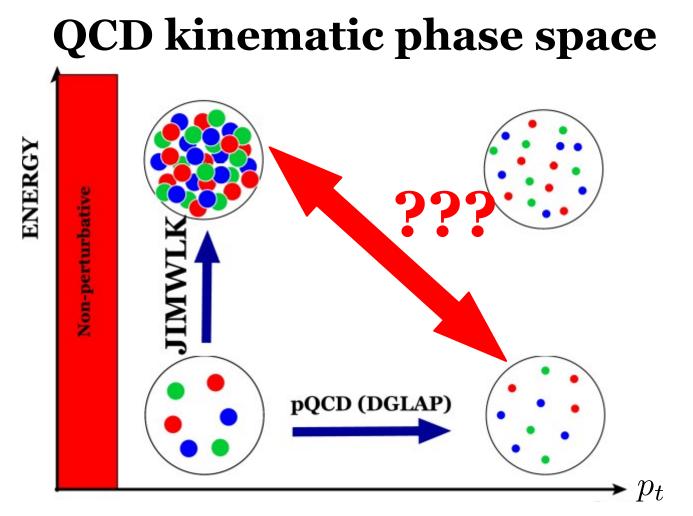


EIC

Kinematics of double inclusive hadron production



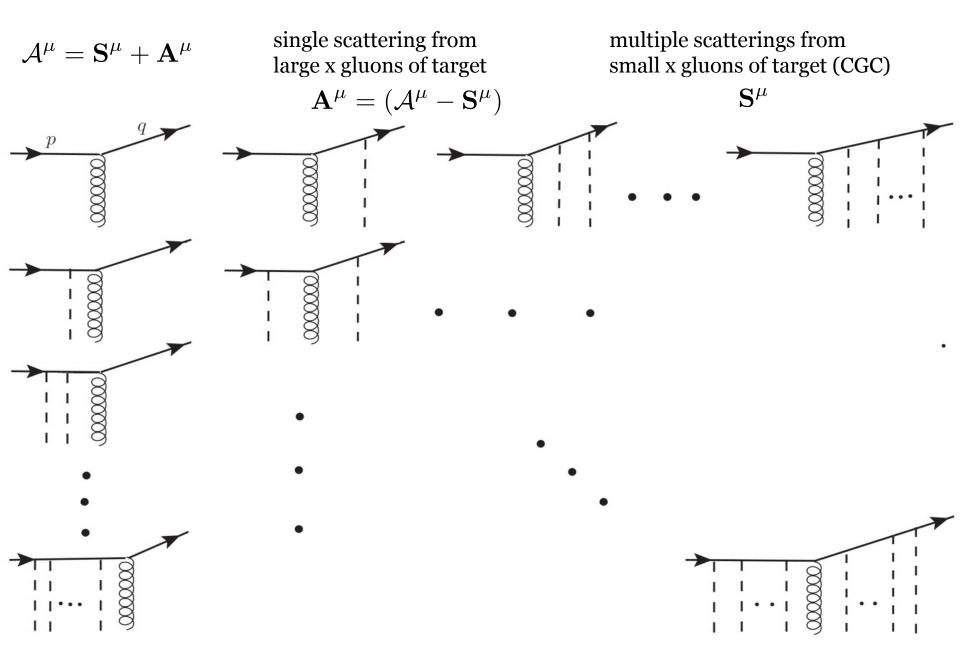
Aschenauer et al. arXiv:1708.01527



unifying saturation with high p_t (large x) pnysics?

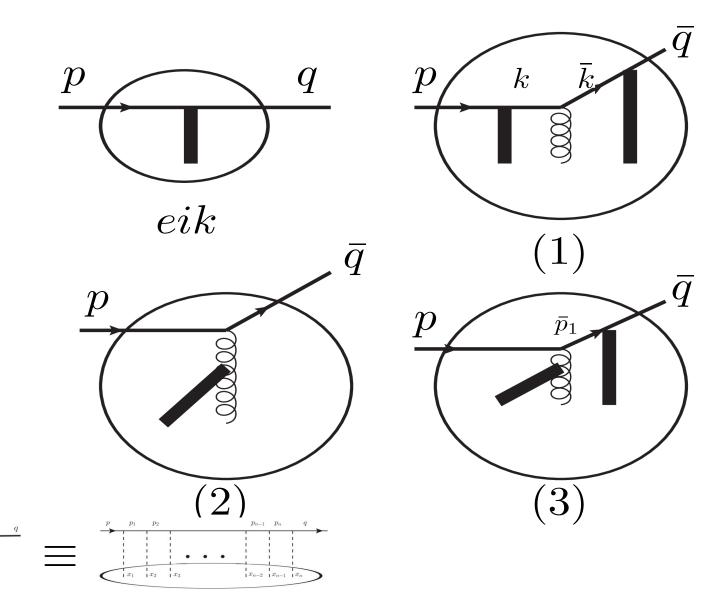
<u>kinematics of saturation: where is saturation applicable?</u> structure functions at all Q^2 high p_t and forward-backward correlations, spin physics, early time e-loss in heavy ion collisions,

Beyond eikonal approximation: longitudinal momentum exchange



Quark scattering: beyond small x approximation

large x partons of target can cause a <u>large-angle deflection</u> of the projectile



soft (eikonal) limit: $i\mathcal{M} \longrightarrow i\mathcal{M}_{eik}$

use spinor helicity formalism: helicity amplitudes

Including large x partons of the target leads to:

<u>longitudinal double spin asymmetries</u> (A_{LL})

baryon transport (beam rapidity loss),

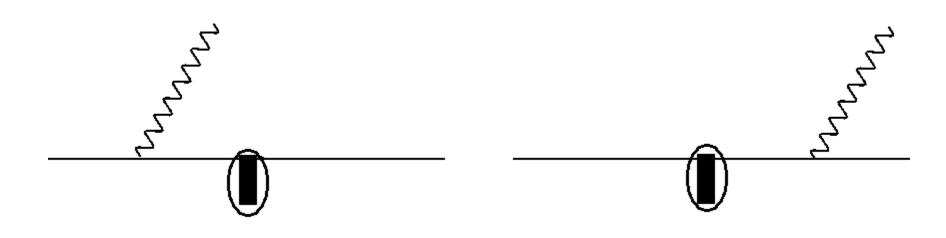
one-loop corrections: factorized cross section at all $x(p_t)$

gluon radiation

related problem: photon production

photon-hadron correlations: azimuthal angular correlations from low to high p_t forward-backward rapidity correlations

photon radiation: small x (eikonal approximation)

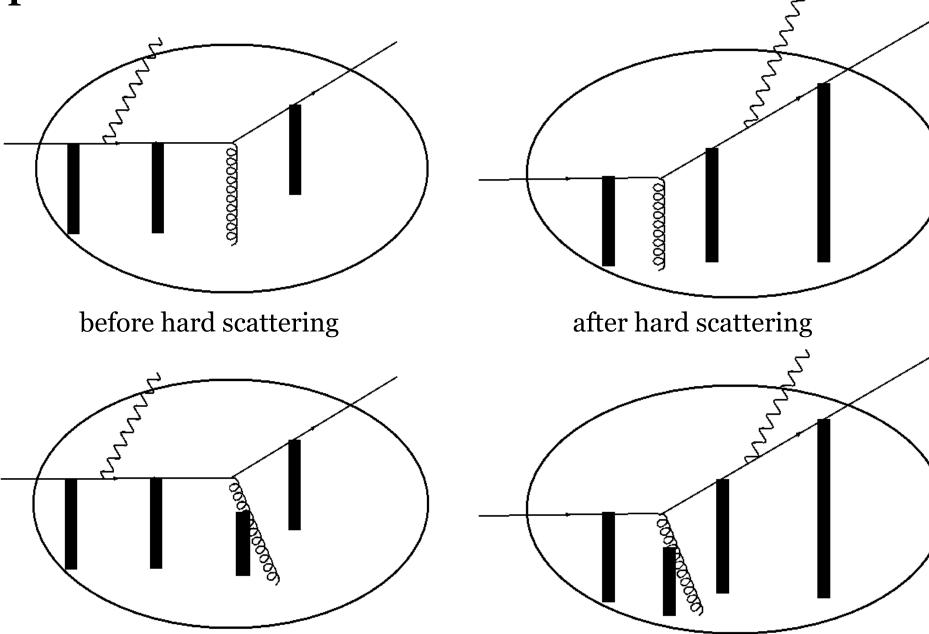


before quark scatters on the target

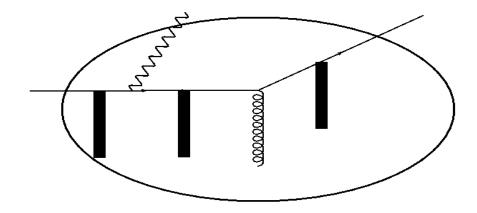
after quark scatters on the target

No radiation inside the target

photon radiation: all x

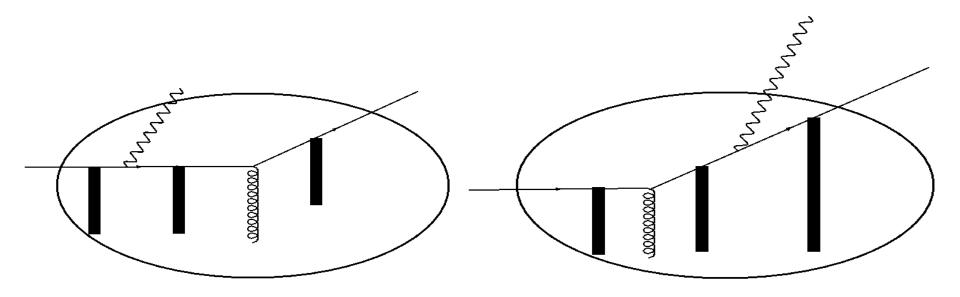


photon radiation: all x

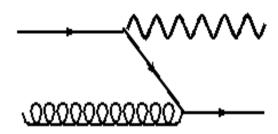


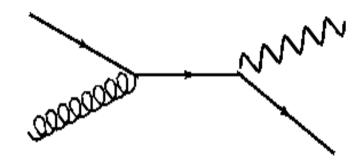
$$\begin{split} &i\mathcal{M}_{1}(p,q,l) = \\ ⪚ \int \frac{d^{2}k_{2t}}{(2\pi)^{2}} \frac{d^{2}k_{3t}}{(2\pi)^{2}} \frac{d^{2}\bar{k}_{1t}}{(2\pi)^{2}} \int d^{4}x \, d^{2}y_{1t} \, d^{2}y_{2t} \, d^{2}\bar{y}_{1t} \, dz^{+} \, \theta(x^{+} - z^{+}) \, e^{i(l^{+} + \bar{q}^{+} - p^{+})x^{-}} \\ &e^{-i(\bar{q}_{t} - \bar{k}_{1t}) \cdot \bar{y}_{1t}} \, e^{-i(\bar{k}_{1t} - k_{3t}) \cdot x_{t}} \, e^{-i(k_{3t} - k_{2t}) \cdot y_{2t}} \, e^{-i(l_{t} + k_{2t} - p_{t}) \cdot y_{1t}} \, \bar{u}(\bar{q}) \, \overline{V}(\bar{y}_{1t}; x^{+}, \infty) \, \frac{\not{n} \, \bar{k}_{1}}{2\bar{n} \cdot \bar{q}} \\ &\mathcal{A}(x) \, \left[\frac{\not{k}_{3}}{2n \cdot (p - l)} \, V(y_{2t}; z^{+}, x^{+}) \, \frac{\not{n} \, \not{k}_{2}}{2n \cdot (p - l)} + i \, \frac{\delta(x^{+} - z^{+})}{2n \cdot (p - l)} \not{n} \right] \\ &\not{\ell}(l) \, \frac{\not{k}_{1}}{2n \cdot p} \, V(y_{1t}; -\infty, z^{+}) \, \not{n} \, u(p) \end{split}$$

pQCD limit (single gluon exchange):

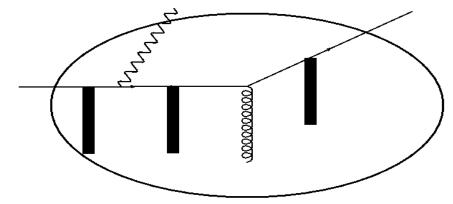


V = U = 1





photon radiation: helicity amplitudes



$$\mathcal{N}_{1-1} = \bar{u}(\bar{q}) \frac{\not{n} \, \bar{k}_1}{2\bar{n} \cdot \bar{q}} \mathcal{A}(x) \frac{k_3 \, \not{n} \, k_2 \not{\epsilon}(l) \, k_1 \, \not{n}}{2n \cdot p \, 2n \cdot (p-l) \, 2n \cdot (p-l)} \, u(p)$$

$$\mathcal{N}_{1-2} = \bar{u}(\bar{q}) \, \frac{\not{n} \, \bar{k}_1}{2\bar{n} \cdot \bar{q}} \, \mathcal{A}(x) \, \frac{\not{n} \, \epsilon(l) \, k_1 \, \not{n}}{2n \cdot p \, 2n \cdot (p-l)} \, u(p)$$

$$\mathcal{N}_{1-1}^{++} = \left(\mathcal{N}_{1-1}^{--}\right)^{\star} = -\sqrt{\frac{n \cdot p}{n \cdot (p-l)}} \frac{\left[n \cdot l \, k_{2\perp} \cdot \epsilon_{\perp}^{\star} - n \cdot (p-l) \, l_{\perp} \cdot \epsilon_{\perp}^{\star}\right]}{n \cdot l \, n \cdot (p-l)} \langle \bar{k}_{1}^{+} | \mathcal{A}(x) | k_{3}^{+} \rangle$$

$$\mathcal{N}_{1-2}^{++} = \left(\mathcal{N}_{1-2}^{--}\right)^{\star} = -\sqrt{\frac{n \cdot p}{n \cdot (p-l)}} \langle \bar{k}_{1}^{+} | \mathcal{A}(x) | n^{+} \rangle$$

$$\mathcal{N}_{1-1}^{+-} = \left(\mathcal{N}_{1-1}^{-+}\right)^{\star} = -\sqrt{\frac{n \cdot p}{n \cdot (p-l)}} \frac{\left[n \cdot p \, l_{\perp} \cdot \epsilon_{\perp} - n \cdot l \, k_{1\perp} \cdot \epsilon_{\perp}\right]}{n \cdot p \, n \cdot l} \langle \bar{k}_{1}^{+} | \mathcal{A}(x) | k_{3}^{+} \rangle$$

$$\mathcal{N}_{1-2}^{+-} = \mathcal{N}_{1-2}^{-+} = 0$$

So far

Classical CGC is generalized by including longitudinal momentum exchange (large angle scattering) with the target

beam rapidity loss

Helicity amplitudes for quark and photon production are evaluated

spin asymmetries

Relevant operators are identified

products of Wilson lines and large x gluon field computing expectation values?

Toward a factorized cross section at all $x(p_t)$

gluon radiation

Need to classify/regulate the divergences

Combining with small x

sharp boundary (x = 0.01)?
matching field strengths?

SUMMARY

CGC is a systematic approach to high energy collisions

strong hints from RHIC, LHC,...

toward precision: NLO, sub-eikonal corrections, ...

CGC breaks down at large x (high p_t)

a significant part of EIC/RHIC/LHC phase space is at large x transition from large x physics to CGC (kinematics?)

Toward inclusion of large x physics:

spin asymmetries
beam rapidity loss
particle production in both small and large p_t kinematics
two-particle correlations: from forward-forward to forward-backward
one-loop correction: both collinear and CGC factorization limits
need to clarify/understand: gauge invariance, initial conditions,