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ALICE

Recent results on 2-particle angular correlations from ALICE



[The 37th Winter Workshop on Nuclear Dynamics \(WWND2022\)](#)

Puerto Vallarta, Mexico

1. Overview of ALICE
2. Correlations & Techniques
3. Selected Recent Results
4. Summary

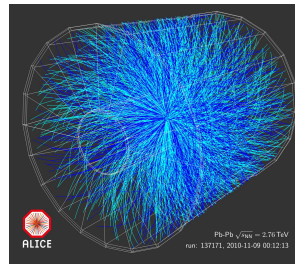
Sumit Basu (On behalf of the ALICE
Collaboration)

Lund University, Sweden

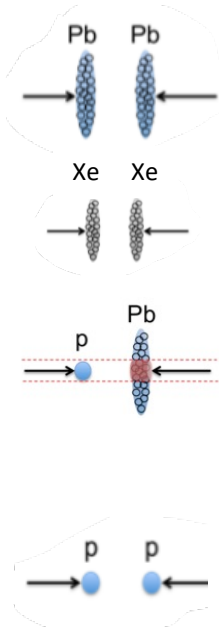


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Event Display **Pb-Pb**

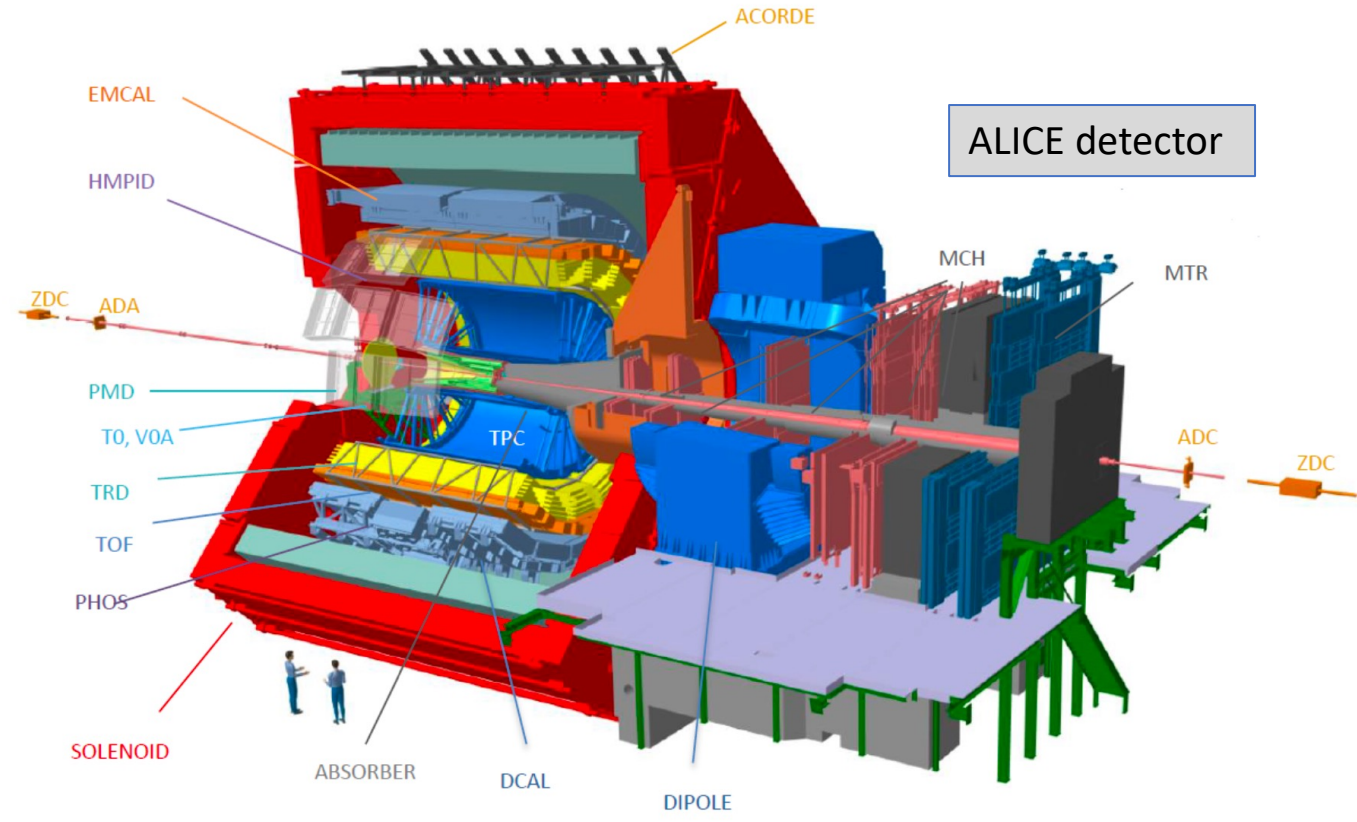
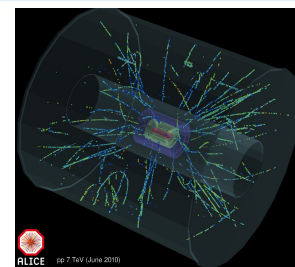


System	Year(s)	$\sqrt{s_{NN}}$ (TeV)
Pb-Pb	2010, 2011	2.76
	2015, 2018	5.02
Xe-Xe	2017	5.44
p-Pb	2013	5.02
	2016	5.02, 8.16
pp	2009-2013	0.9, 2.76, 7, 8
	2015, 2017	5.02
	2015-2018	13



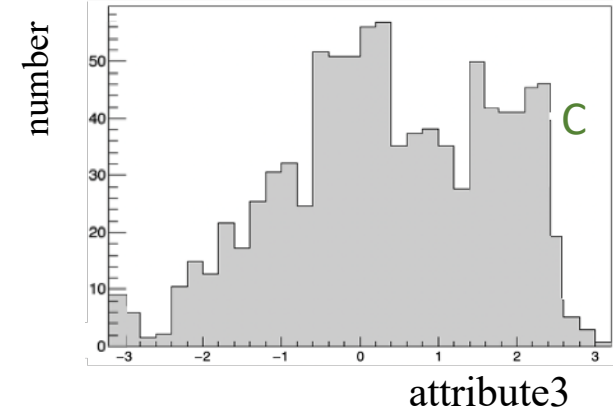
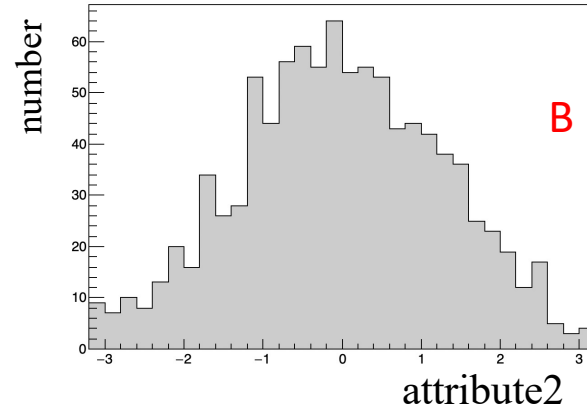
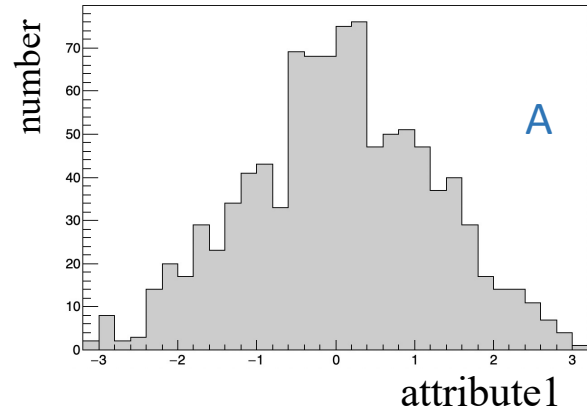
Run 1 **Run 2**

Event Display **pp**

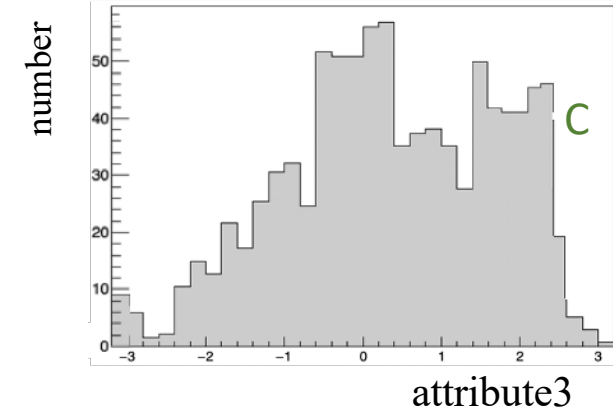
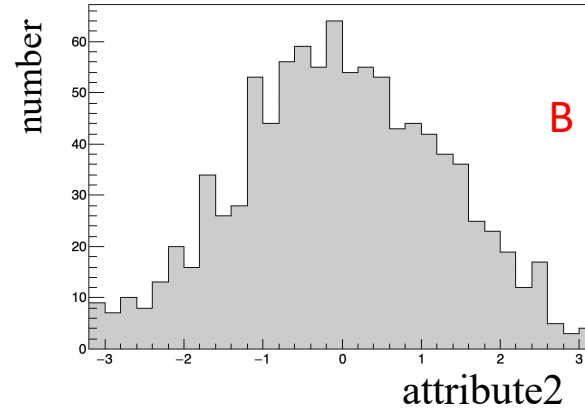
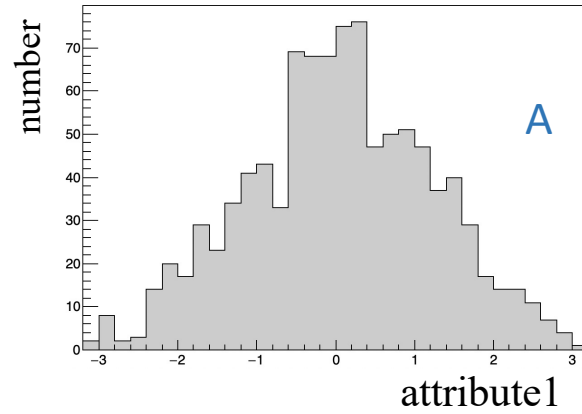


ALICE detector

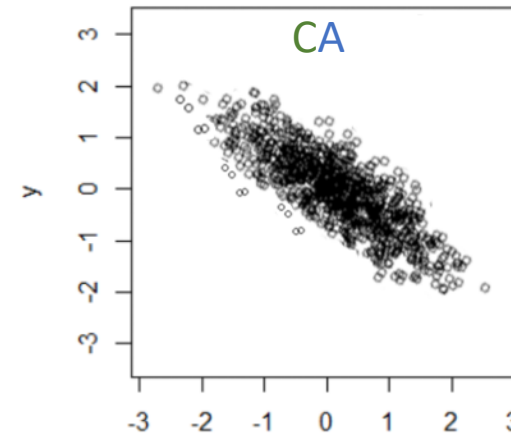
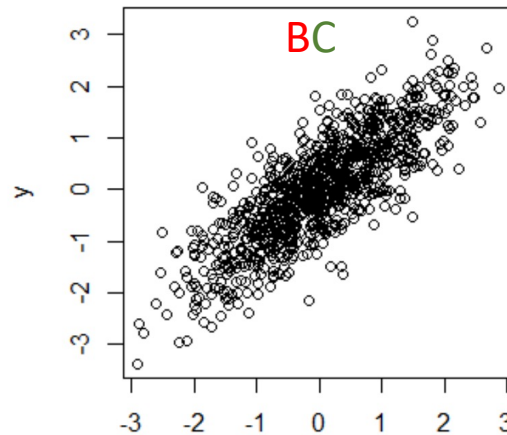
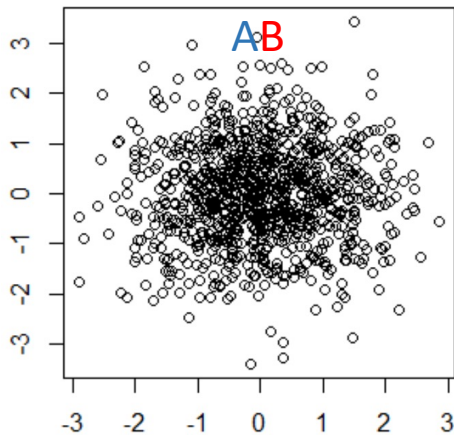
Imagine that we have a population of n things, each with multiple attributes



Imagine that we have a **population of n** things, each with multiple attributes

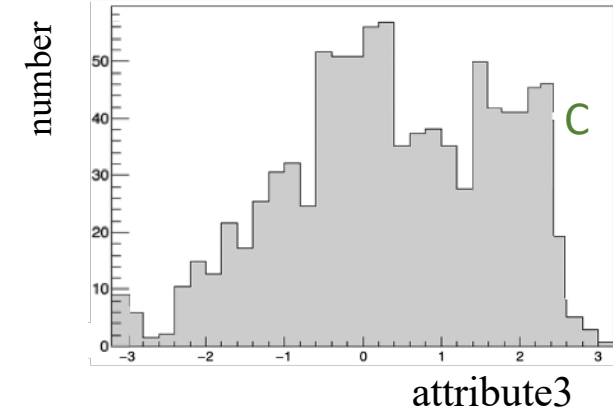
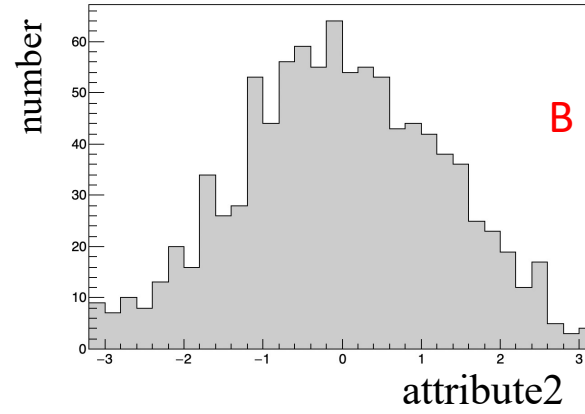
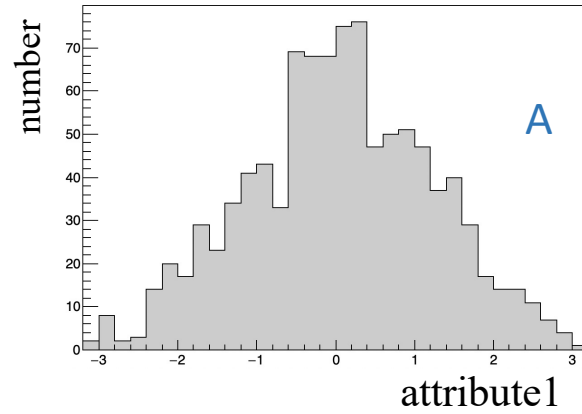


That is not all we can learn though – make the two-dimensional plot!

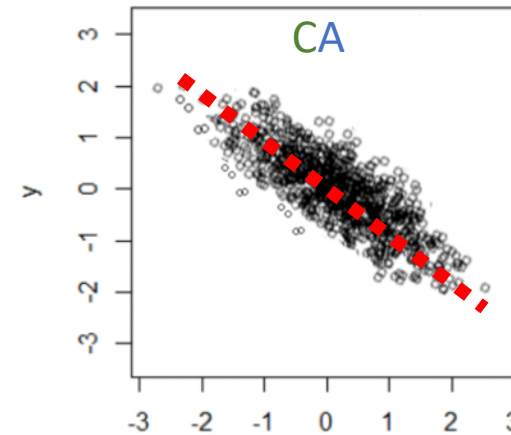
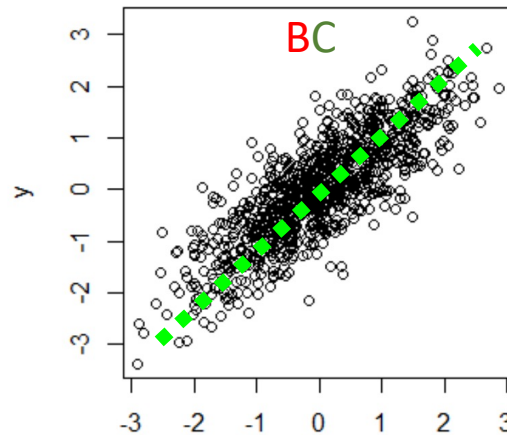
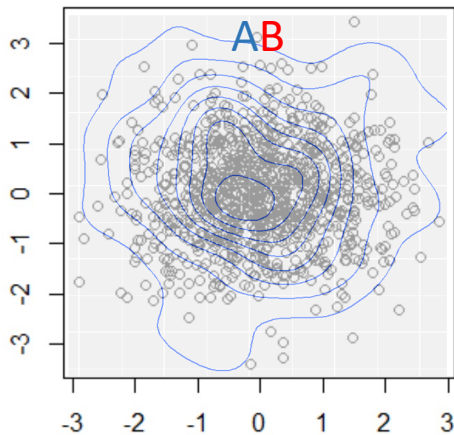


(“number axis” are out of the page)

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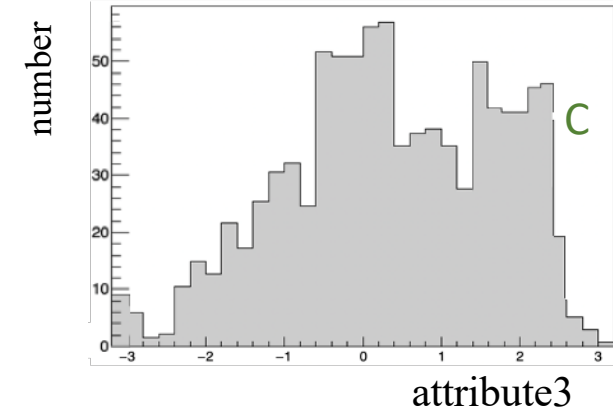
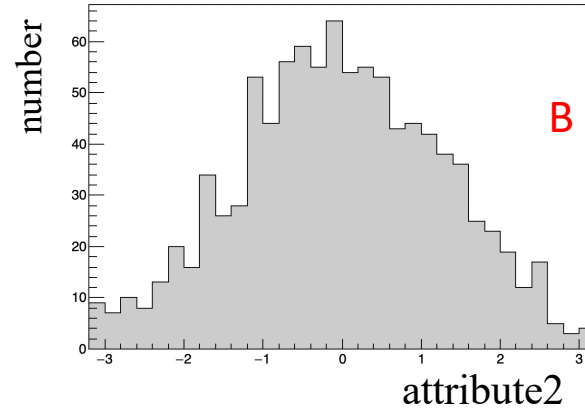
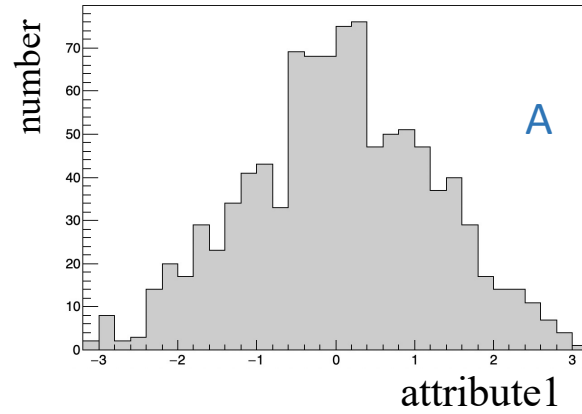


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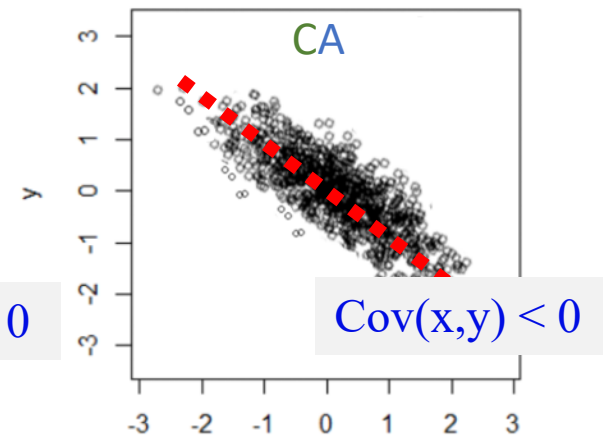
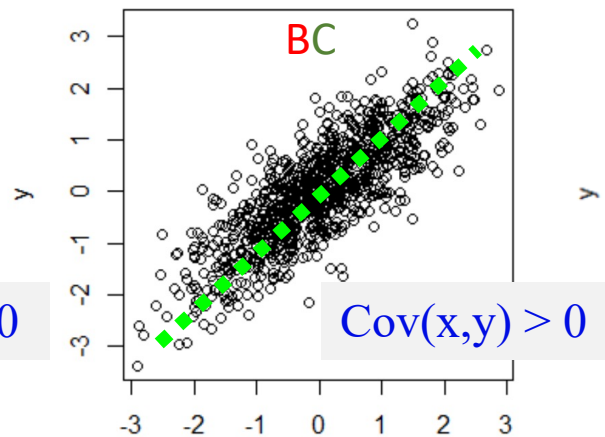
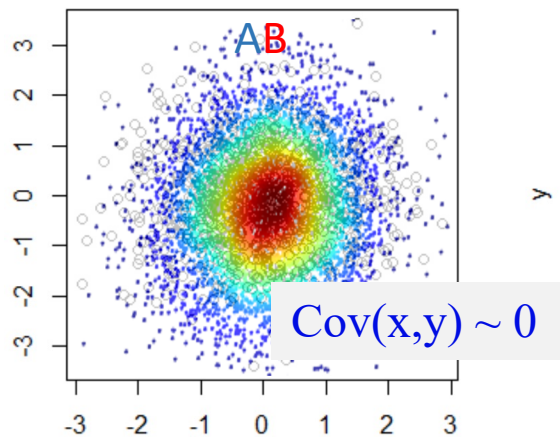


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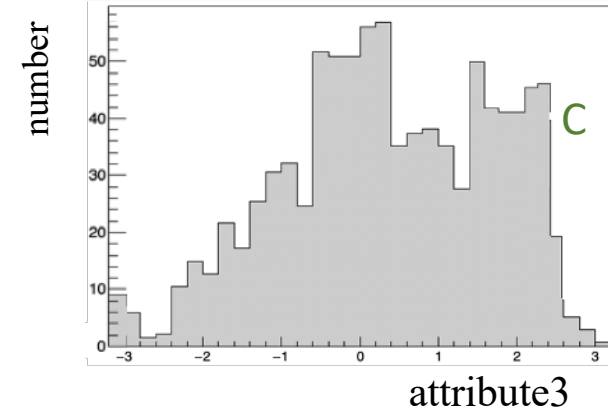
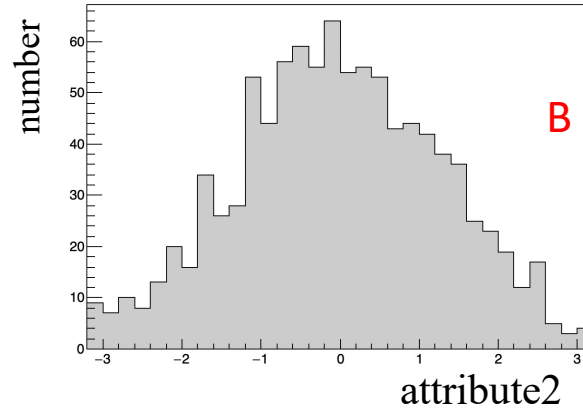
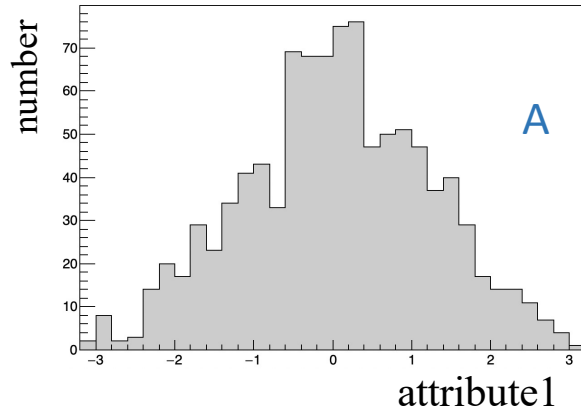


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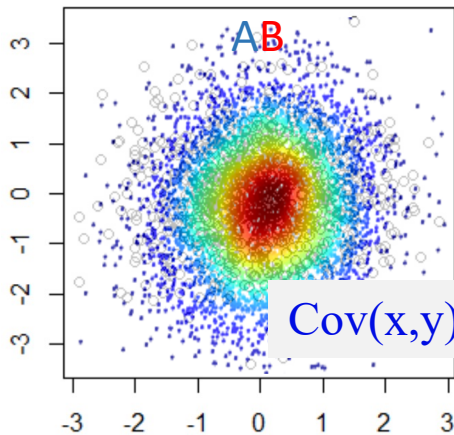
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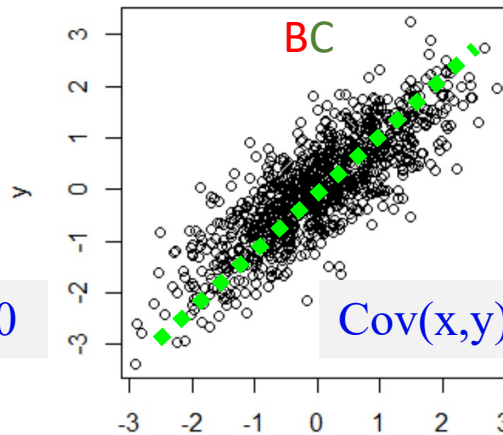
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Studying the correlation of the two variables has resulted in New knowledge!

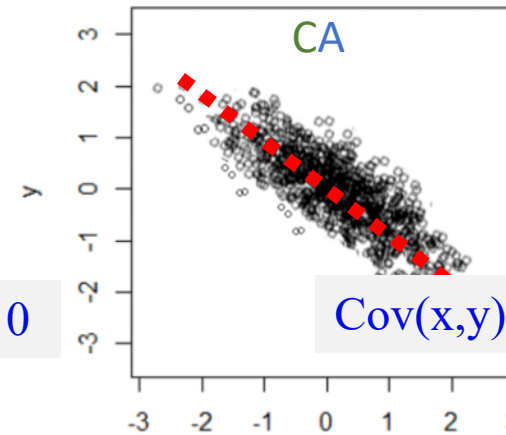
➔ ...knowing the value of one variable constrains the possible values of the other variable.



$Cov(x,y) \sim 0$



$Cov(x,y) > 0$

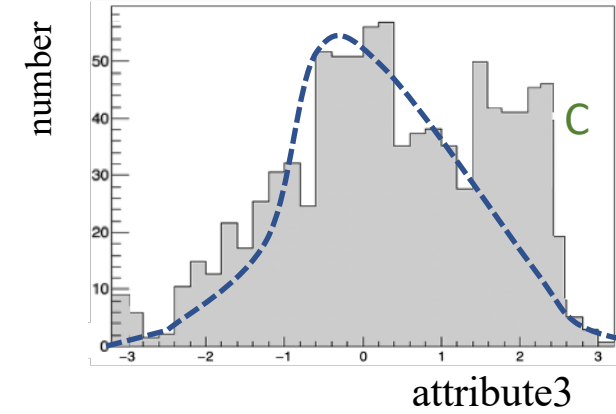
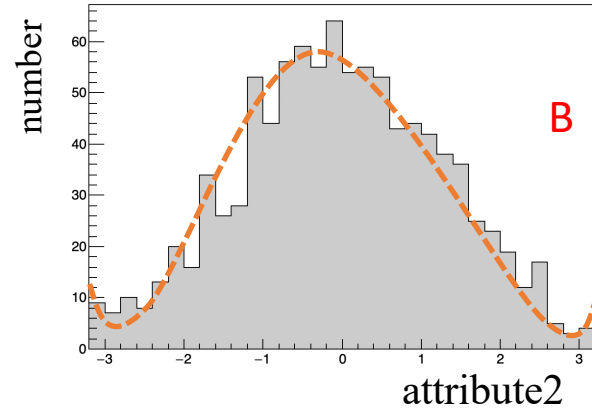
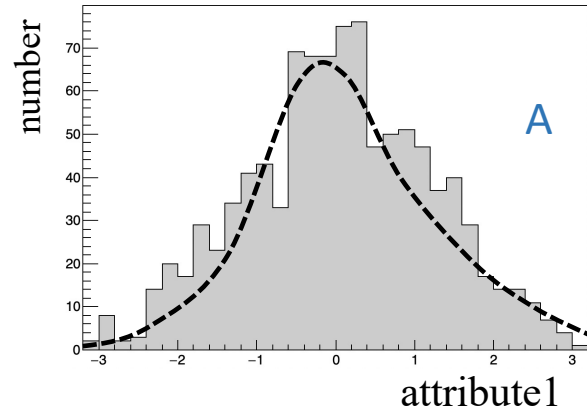


$Cov(x,y) < 0$

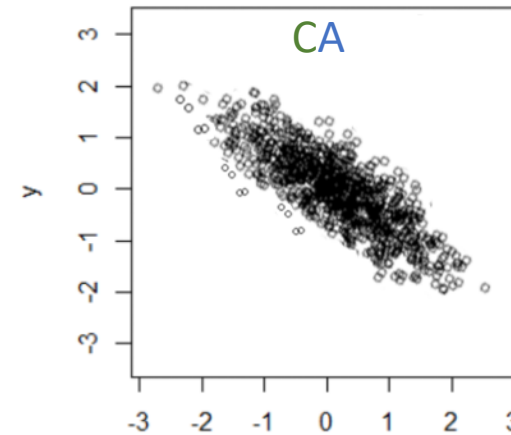
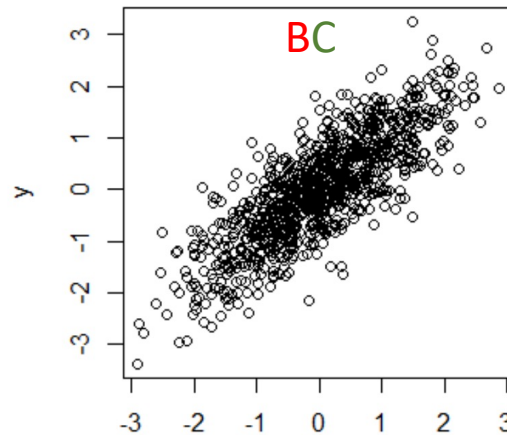
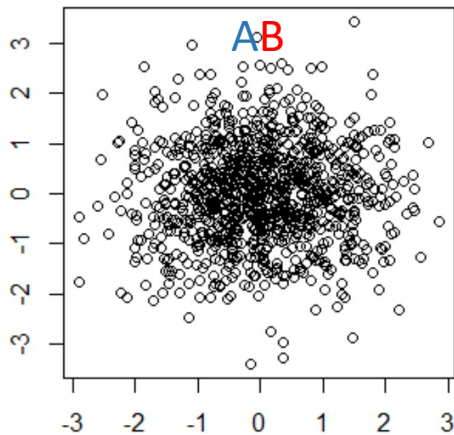
Cov. > 0	correlated
Cov. \sim 0	uncorrelated
Cov. < 0	anticorrelated

Quantify the amount of correlation with the **covariance**: $Cov(x, y) = \langle xy \rangle - \langle x \rangle \langle y \rangle$ ➔ for density $C_2 = \rho_2 - \rho_1 \rho_1$

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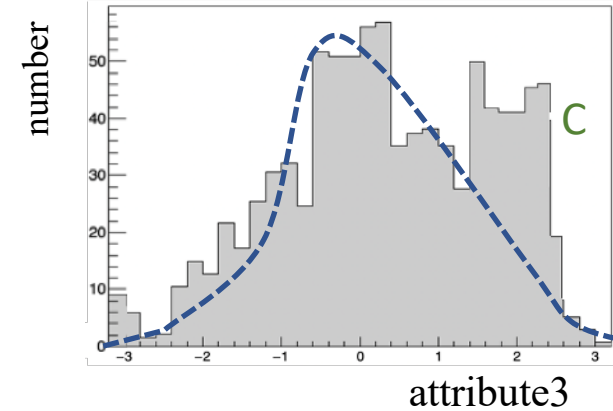
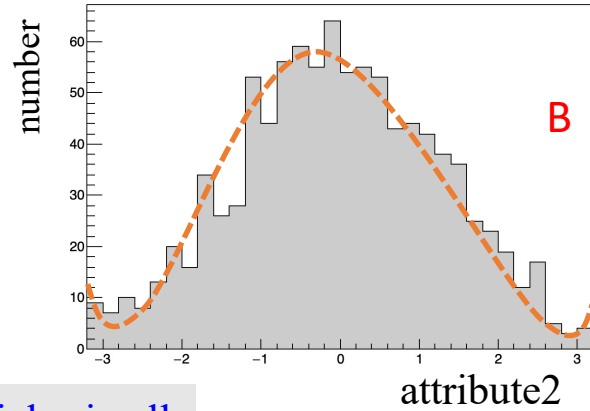
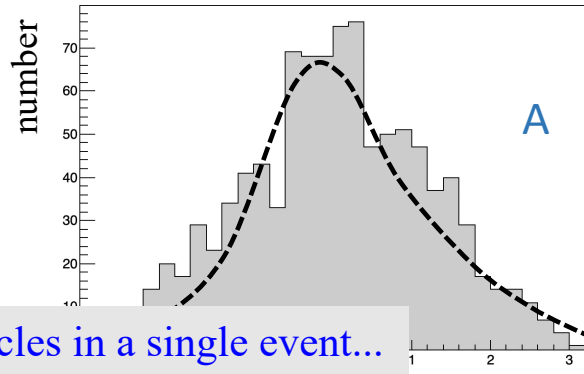


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Experimentally: The average values of specific powers of deviates give cumulants & cumulant ratios (or moments and moments products)....



No. of particles in a single event...

Average No. of particles in all "similar" events...

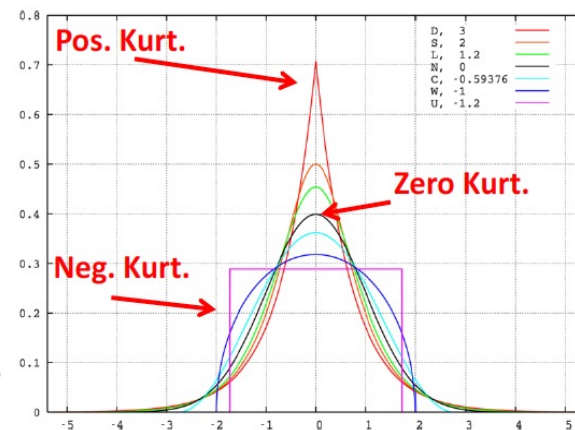
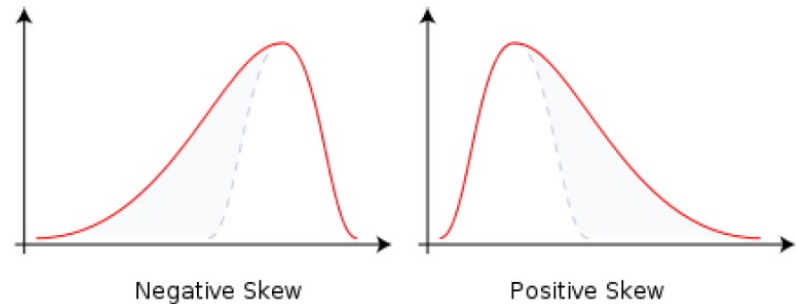
$$\delta x \equiv x - \langle x \rangle$$

$$\kappa_{2x} \equiv \langle \langle x^2 \rangle \rangle \equiv \langle (\delta x)^2 \rangle$$

$$\kappa_{3x} \equiv \langle \langle x^3 \rangle \rangle \equiv \langle (\delta x)^3 \rangle$$

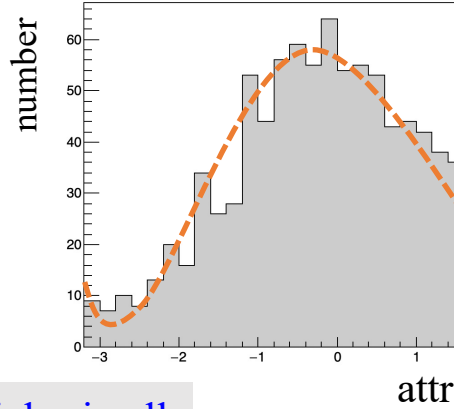
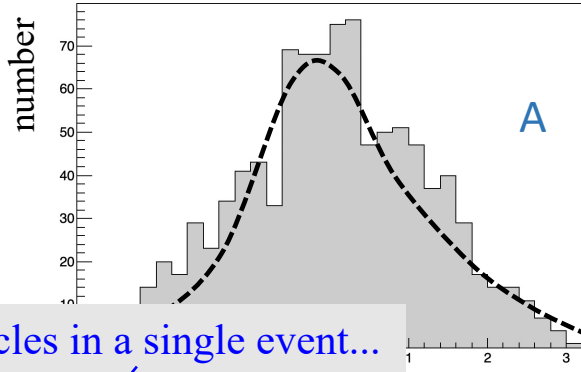
$$\kappa_{4x} \equiv \langle \langle x^4 \rangle \rangle \equiv \langle (\delta x)^4 \rangle - 3 \langle (\delta x)^2 \rangle^2$$

$$\text{skewness} = \frac{\kappa_3}{\kappa_2^{3/2}}, \text{ kurtosis} = \frac{\kappa_4}{\kappa_2^2}$$





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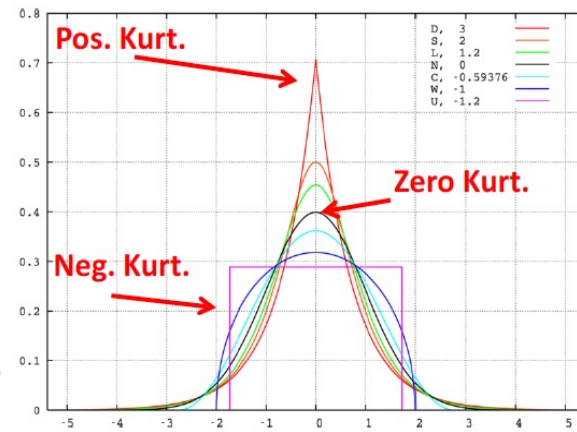
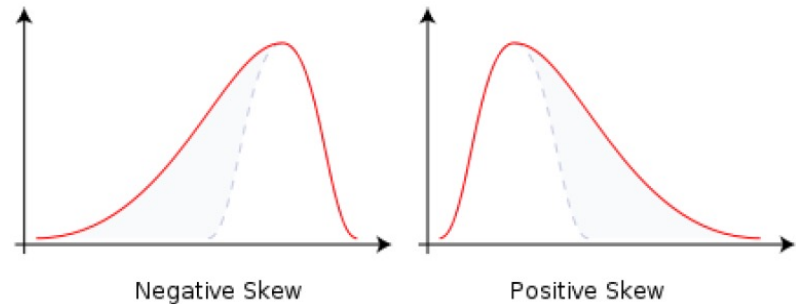
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Fluctuations are the integral of the Correlations

Multiplicity = Poisson + [deviations from Poisson]
Cumulants

$$K_1 = \langle N \rangle$$

$$K_2 = \langle N \rangle + c_2$$

$$K_3 = \langle N \rangle + 3c_2 + c_3$$

$$K_4 = \langle N \rangle + 7c_2 + 6c_3 + c_4$$

$$c_k = \int C_k(y_1, \dots, y_k) dy_1 \dots dy_k$$

"Correlation Functions"

$$C_2 = \rho_2 - \rho_1 \rho_1$$

$$C_3 = \rho_3 - 3\rho_2 \rho_1 + 2\rho_1 \rho_1 \rho_1$$

$$C_4 = \rho_4 - 4\rho_3 \rho_1 - 3\rho_2 \rho_2 + 12\rho_2 \rho_1 \rho_1 - 6\rho_1 \rho_1 \rho_1 \rho_1$$

Explicit subtraction of lower-order correlations...

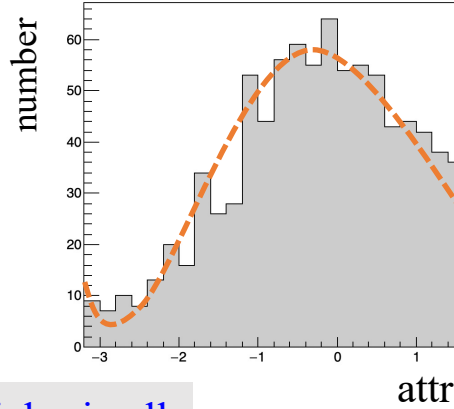
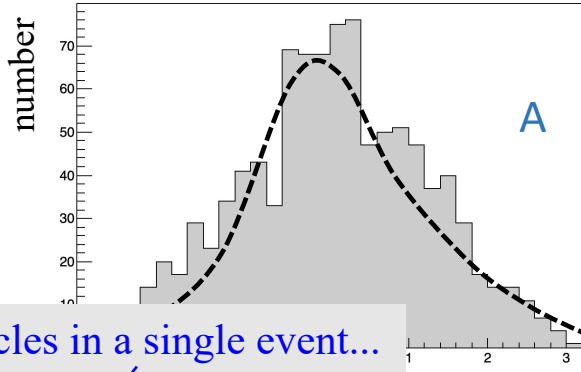
$$C_2 = \text{[Diagram: 2 yellow circles connected] - [Diagram: 2 blue circles unconnected]}$$

$$C_3 = \text{[Diagram: 3 green circles in a triangle] - 3 [Diagram: 2 yellow circles connected and 1 blue circle unconnected] + 2 [Diagram: 3 blue circles unconnected]}$$

$$C_4 = \text{[Diagram: 4 red circles in a square] - 4 [Diagram: 3 green circles in a triangle and 1 blue circle unconnected] - 3 [Diagram: 2 yellow circles connected and 2 blue circles unconnected] + 12 [Diagram: 2 yellow circles connected and 2 blue circles unconnected] - 6 [Diagram: 4 blue circles unconnected]}$$



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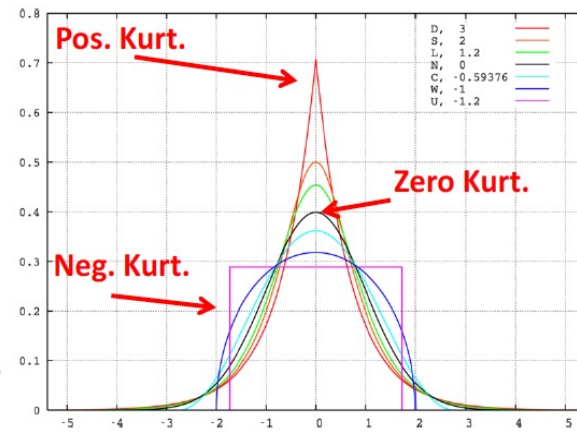
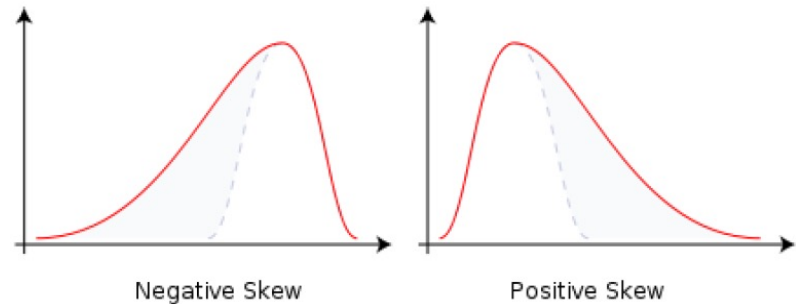
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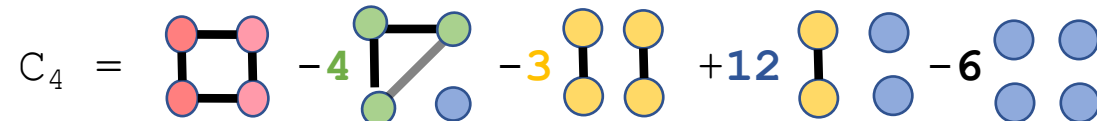
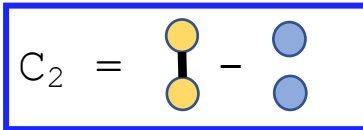
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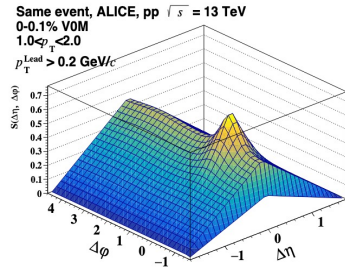
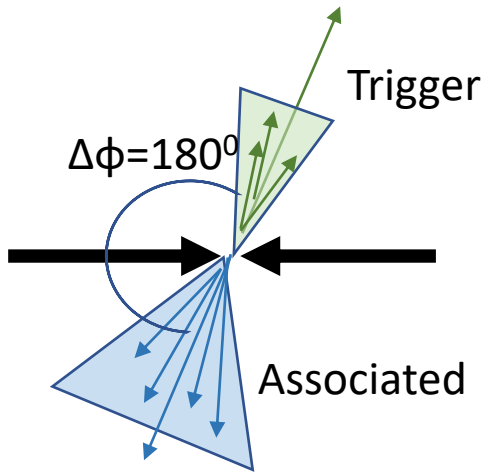
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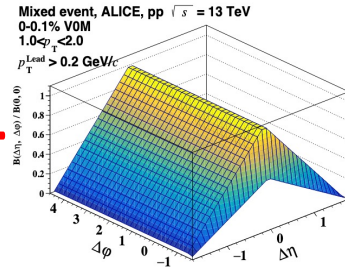
1. Angular correlation between trigger and associated particles is measured:

$$C(\Delta\phi, \Delta\eta) = \frac{1}{N_{\text{trig}}} \frac{d^2 N_{\text{assoc}}}{d\Delta\phi d\Delta\eta} = \frac{S(\Delta\phi, \Delta\eta)}{M(\Delta\phi, \Delta\eta)}$$

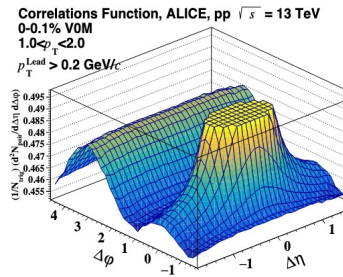
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−



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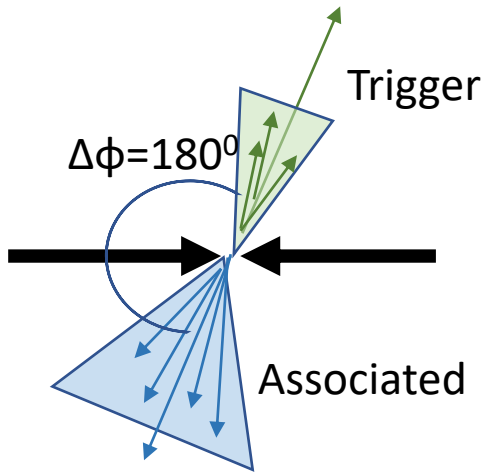


Technique

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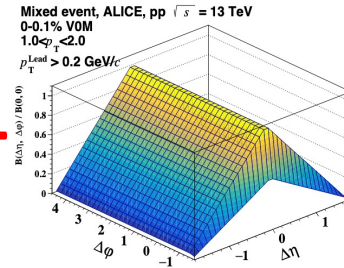
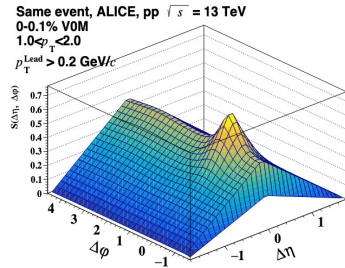
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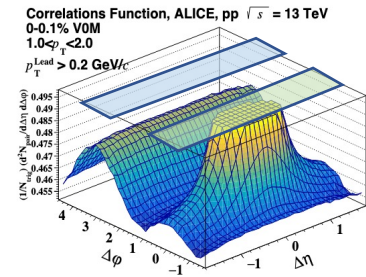


2. Background contribution is subtracted:

$$B(\Delta\phi) = B_0 \left(1 + 2 \sum_n V_n \cos(n\Delta\phi) \right)$$



Away Side Near Side



Result Extraction

I Width $\Delta\eta$ or width $\Delta\phi$ can be calculated

II Yield calculation: $Y_J^{\Delta\phi} = \int_{\Delta\phi_1}^{\Delta\phi_2} \frac{dN}{d\Delta\phi} d\Delta\phi$

III Yield and I_{AA} are calculated:

$$I_{AA} = \int I_{AA}(\Delta\phi) d\Delta\phi / \int I_{pp}(\Delta\phi) d\Delta\phi$$

Technique

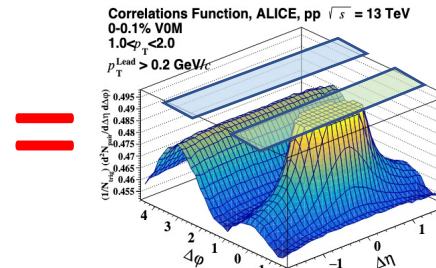
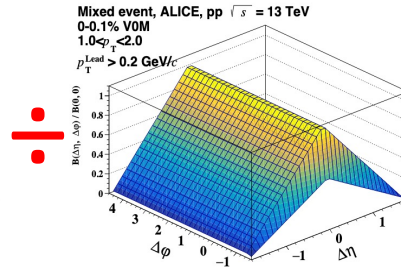
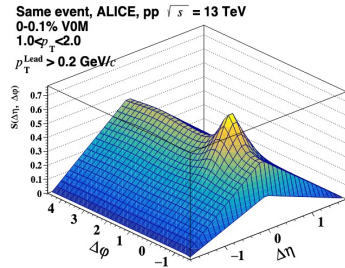
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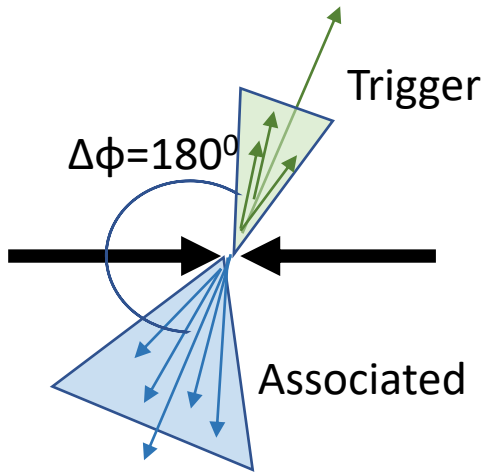
Result Extraction

I Width $\Delta\eta$ or width $\Delta\phi$ can be calculated

II Yield calculation: $Y_J^{\Delta\phi} = \int_{\Delta\phi_1}^{\Delta\phi_2} \frac{dN}{d\Delta\phi} d\Delta\phi$

III Yield and I_{AA} are calculated:

$$I_{AA} = \int I_{AA}(\Delta\phi) d\Delta\phi / \int I_{pp}(\Delta\phi) d\Delta\phi$$



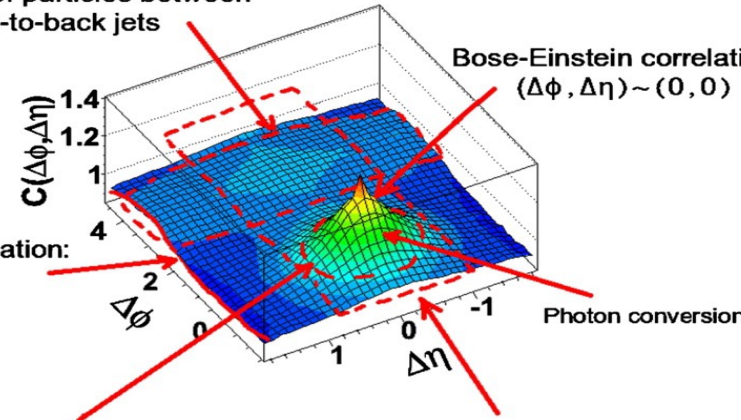
„Away-side” ($\Delta\phi \sim \pi$) jet correlations:
Correlation of particles between
back-to-back jets

Bose-Einstein correlations:
($\Delta\phi, \Delta\eta \sim (0, 0)$)

Momentum conservation:
 $\sim -\cos(\Delta\phi)$

„Near-side” ($\Delta\phi \sim 0$) jet peak:
Correlation of particles within
a single jet

Resonances, string fragmentation

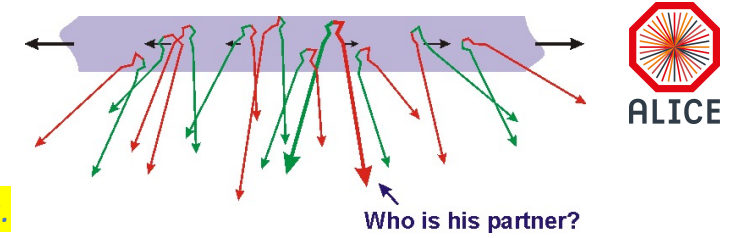


Physics Motivation

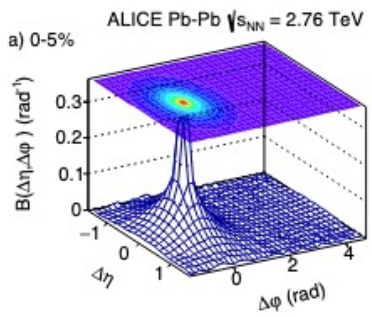
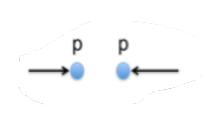
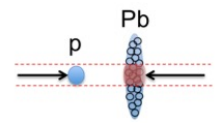
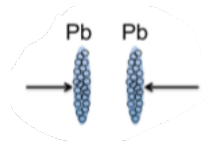
May want to scale C_2 by the number of uncorrelated pairs:

$$R_2 = \frac{C_2}{\rho_1 \times \rho_1} \quad \Rightarrow \quad R_2 = \frac{\rho_2}{\rho_1 \times \rho_1} - 1$$

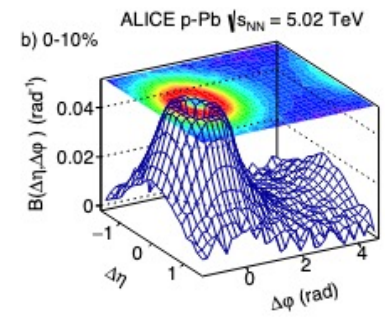
Balance Function (2-Particle correlations)



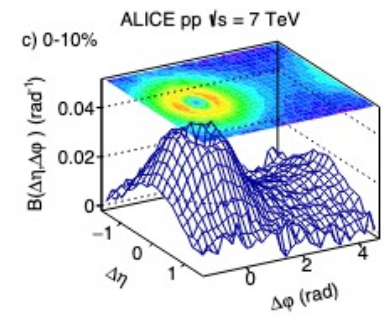
Conservation of quantum numbers.
 -> for each general charge, an opposite balancing charge produced at approx. the same space-time.



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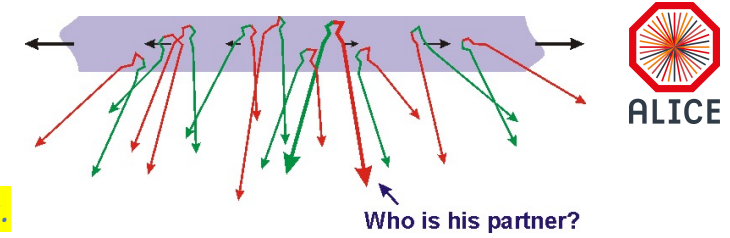


Phys. Rev. C 100, 044903 (2019)

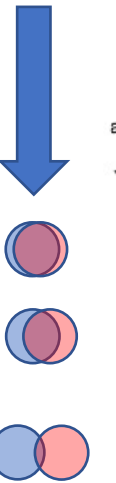


Old Result

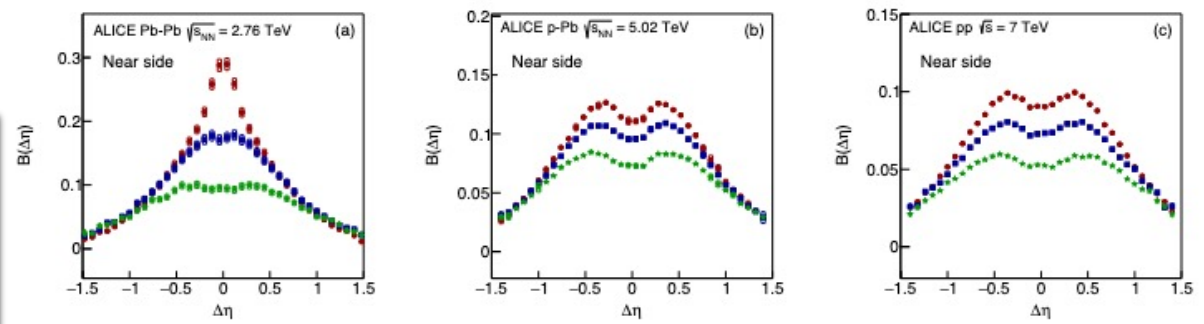
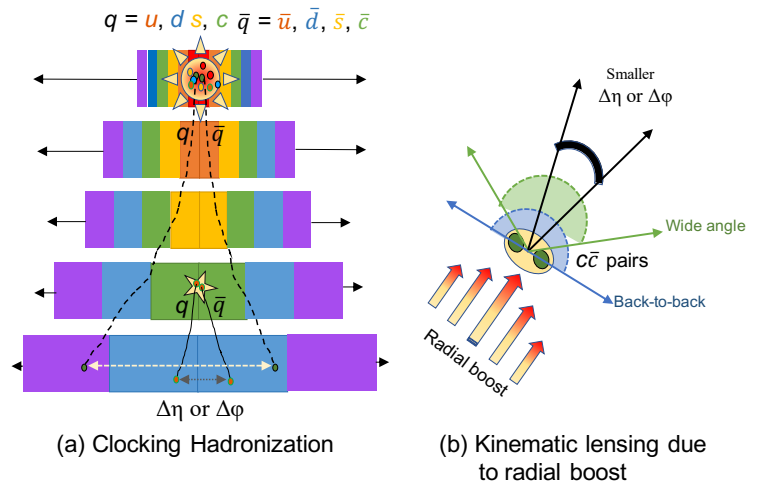
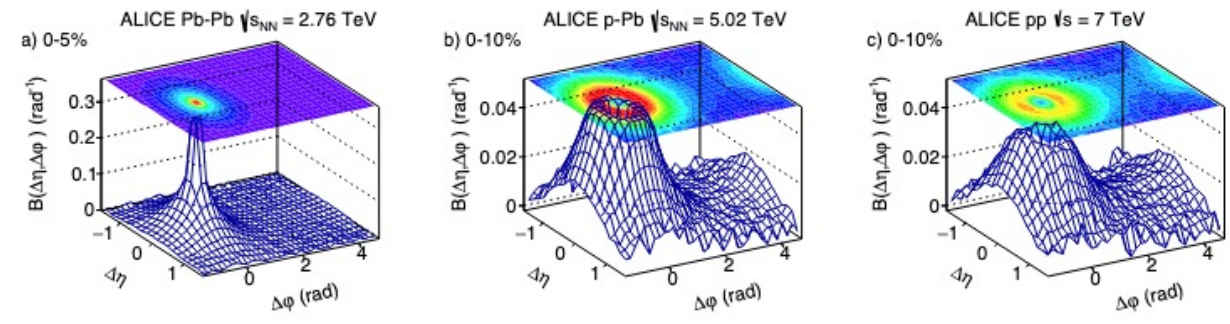
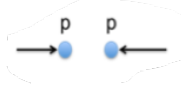
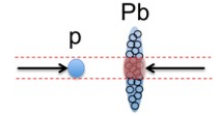
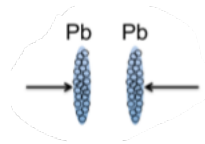
Balance Function (2-Particle correlations)



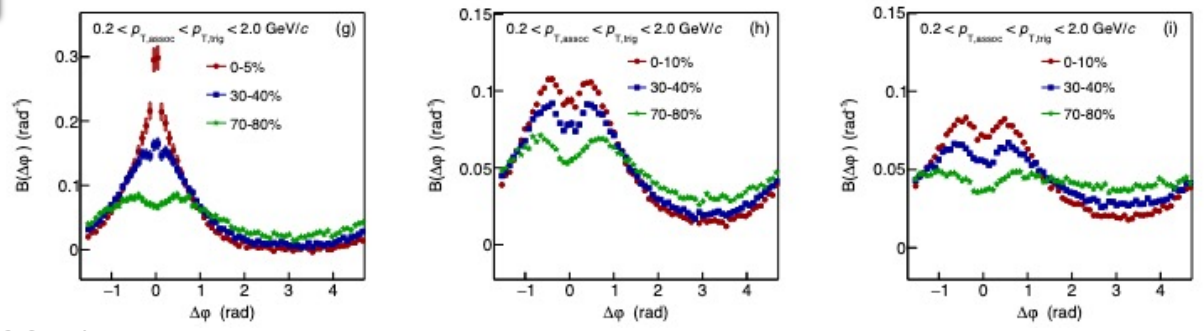
Conservation of quantum numbers.
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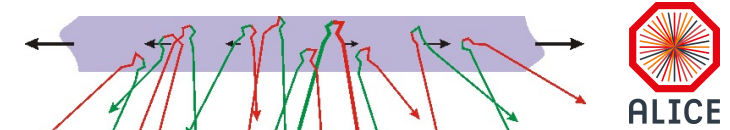
Old Result



Eur. Phys. J. C 76 (2016) 86 Phys. Rev. C 100, 044903 (2019)

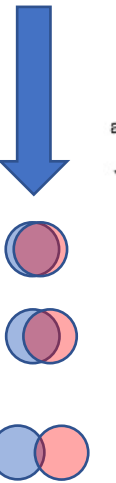


Balance Function (2-Particle correlations)

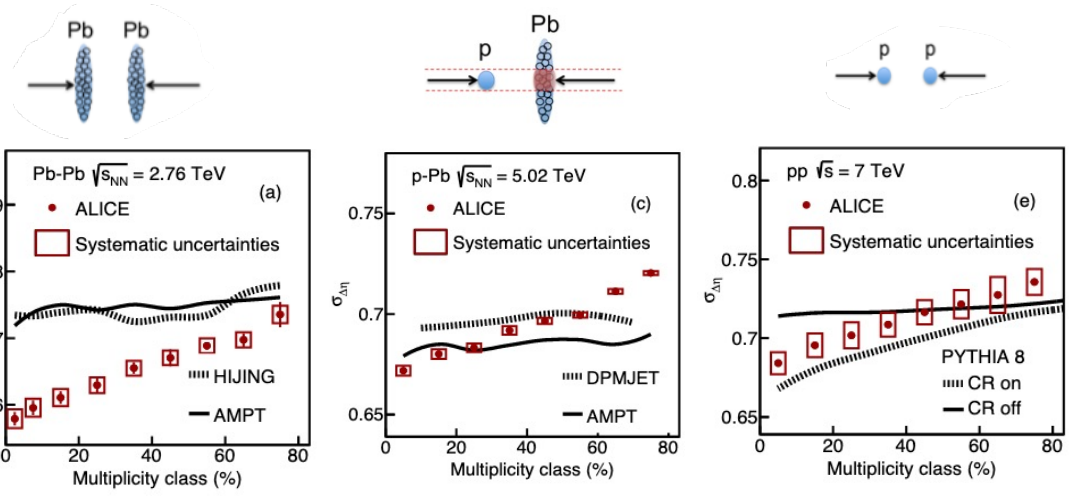
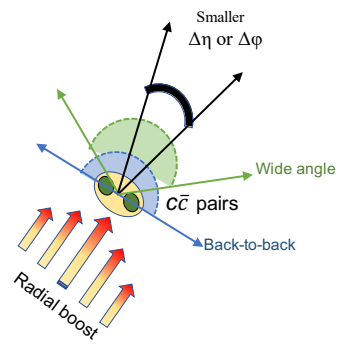
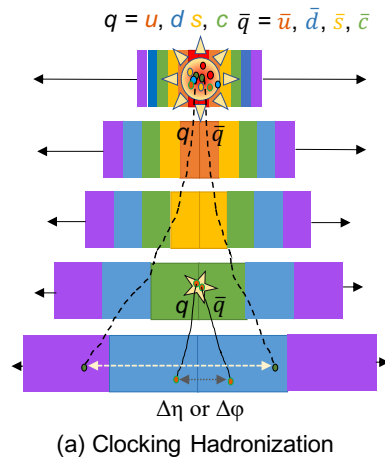
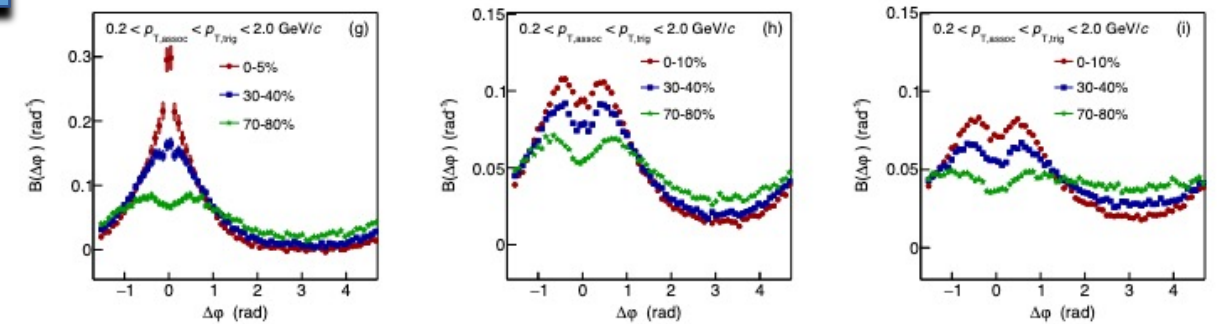
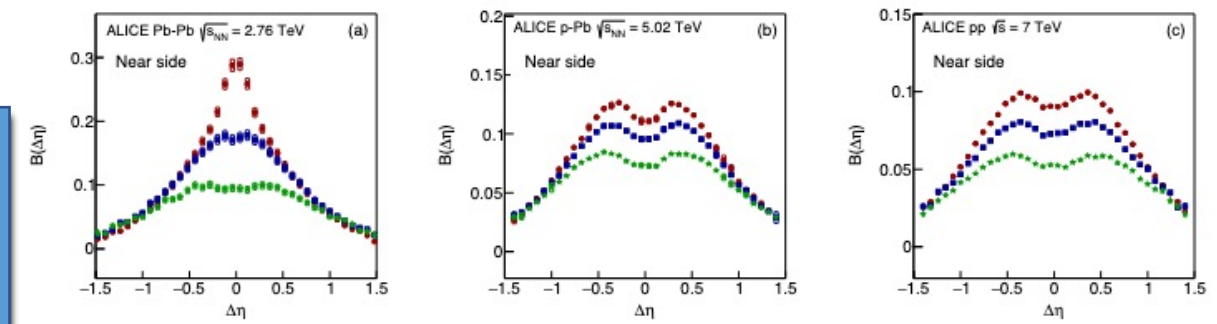
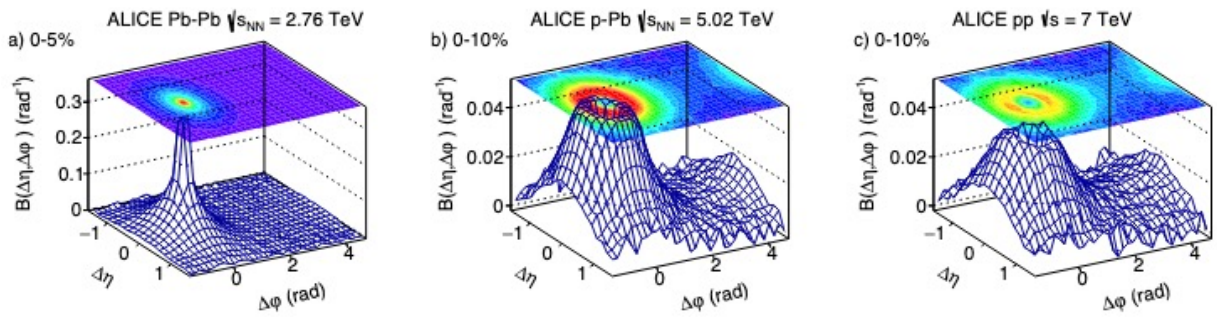
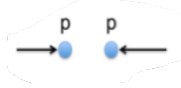
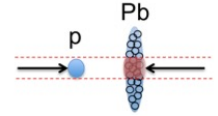
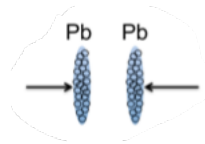


Conservation of quantum numbers.
 -> for each general charge, an opposite balancing charge produced at approx. the same space-time.

Who is his partner?



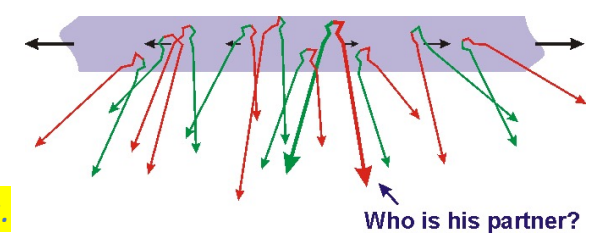
Old Result



Balance Function (2-Particle correlations)

Species Dependence

		h	π	k	p
Q	h	✓			
Q	π		?	?	?
Q	s		?	?	?
Q	B		?	?	?



Conservation of quantum numbers.
 -> for each general charge, an opposite balancing charge produced at approx. the same space-time.

Cumulant $C_2(x_1, x_2) = \rho_2(x_1, x_2) - \rho_1(x_1)\rho_1(x_2)$

$x \equiv \{y, \varphi, p_T\}$ $\rho(x) = \frac{1}{\sigma} \frac{d\sigma}{dx}$

Normalized Cumulant $R_2(x_1, x_2) = \frac{C_2(x_1, x_2)}{\rho_1(x_1)\rho_1(x_2)}$

R₂ is a robust observable!
 Single track efficiencies cancel out of the ratio

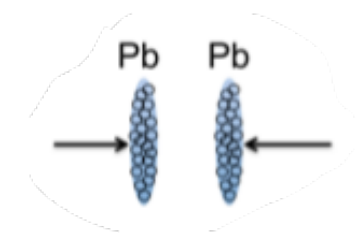
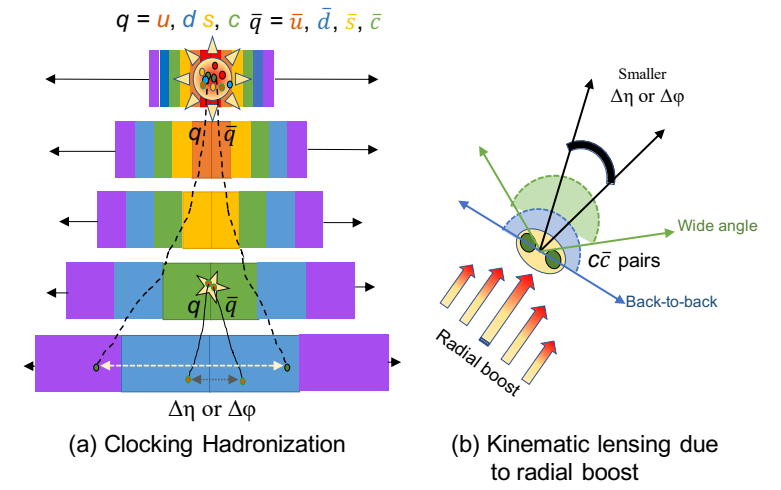
4 different charge combinations for R₂:
 (+ -), (- +), (+ +), and (- -)

Charge Independent (CI) combinations $CI = \frac{1}{2}\{LS + US\}$ $LS = \frac{1}{2}\{(++ + (--))\}$

Charge Dependent (CD) combinations $CD = \frac{1}{2}\{US - LS\}$ $US = \frac{1}{2}\{(++ + (--))\}$

R₂^{CD} is proportional to the Balance Function

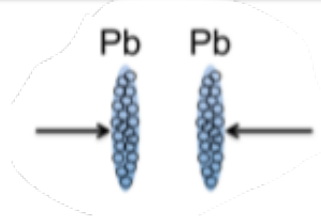
$$B(\Delta x) \approx \frac{dN_{ch}}{dx} R_2^{CD} = \frac{dN_{ch}}{dx} \frac{1}{2} [R_2^{+-} - R_2^{++} + R_2^{-+} - R_2^{--}]$$



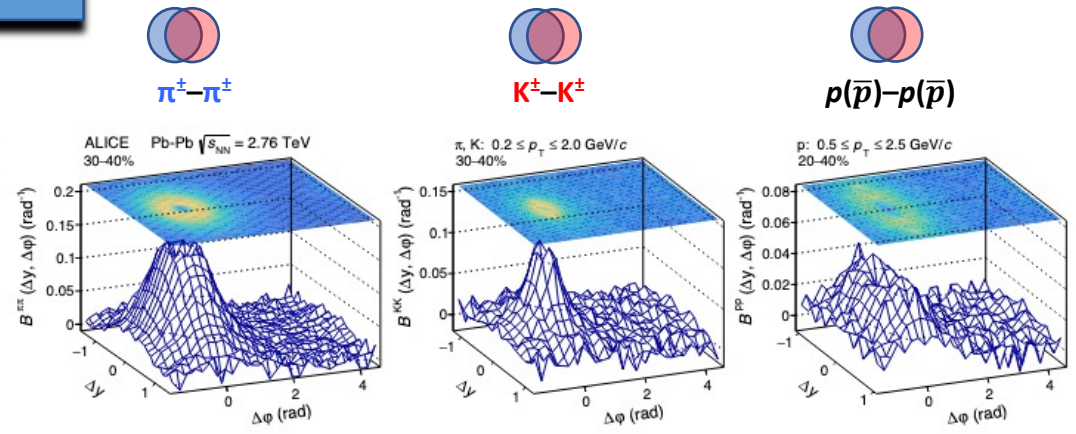
Balance Function (2-Particle correlations)

Species
Dependence

		h	π	k	p
Q	h	✓			
Q	π		?	?	?
Q	S		?	?	?
Q	B		?	?	?



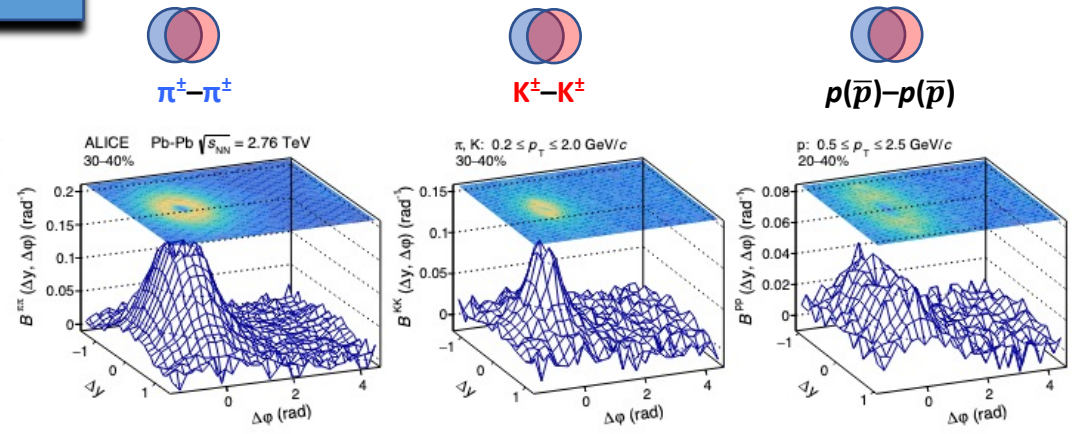
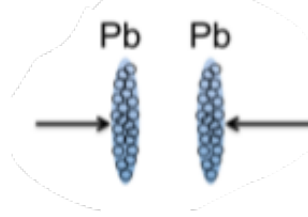
New
Result



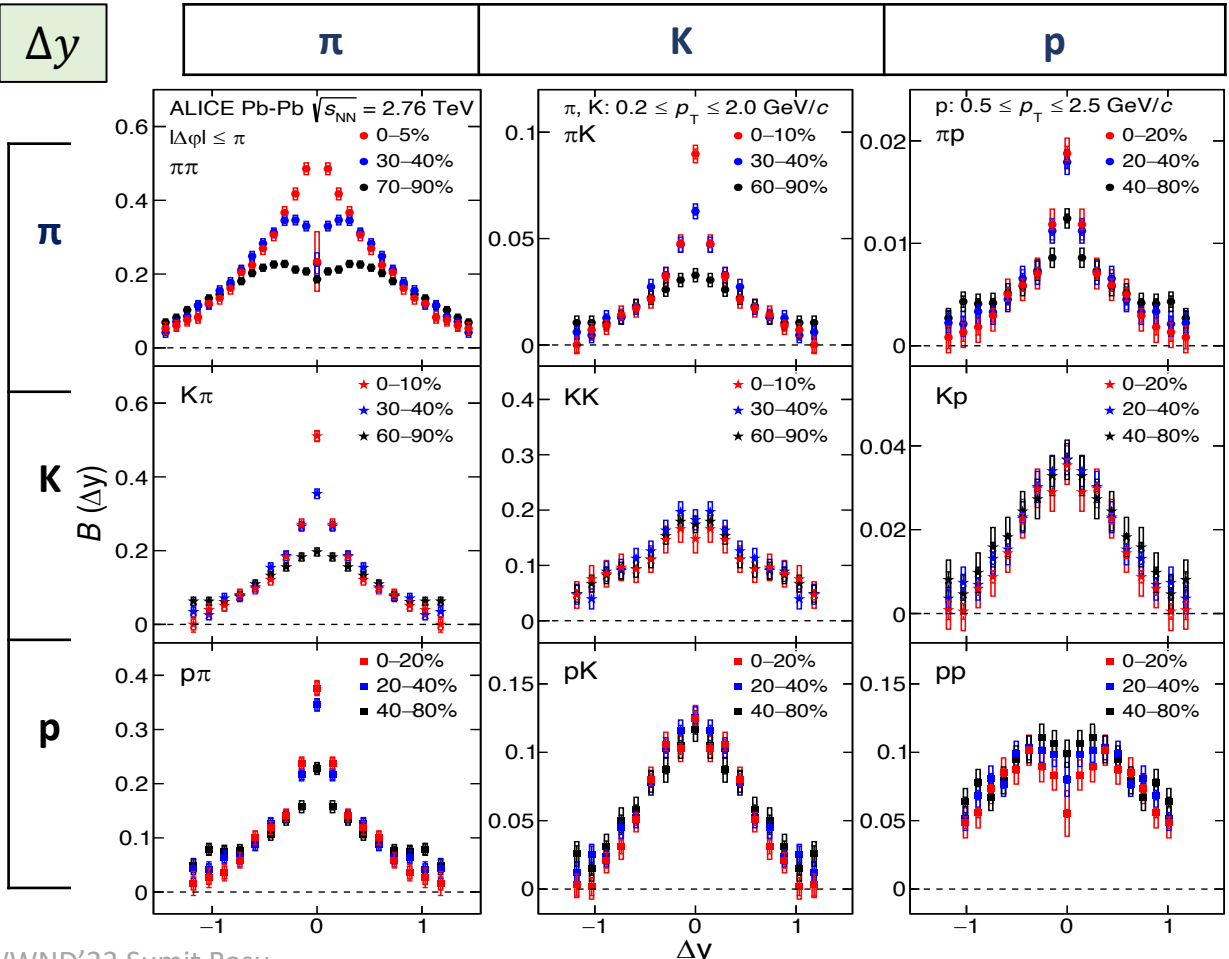
[arXiv:2110.06566](https://arxiv.org/abs/2110.06566) [nucl-ex]

Species
Dependence

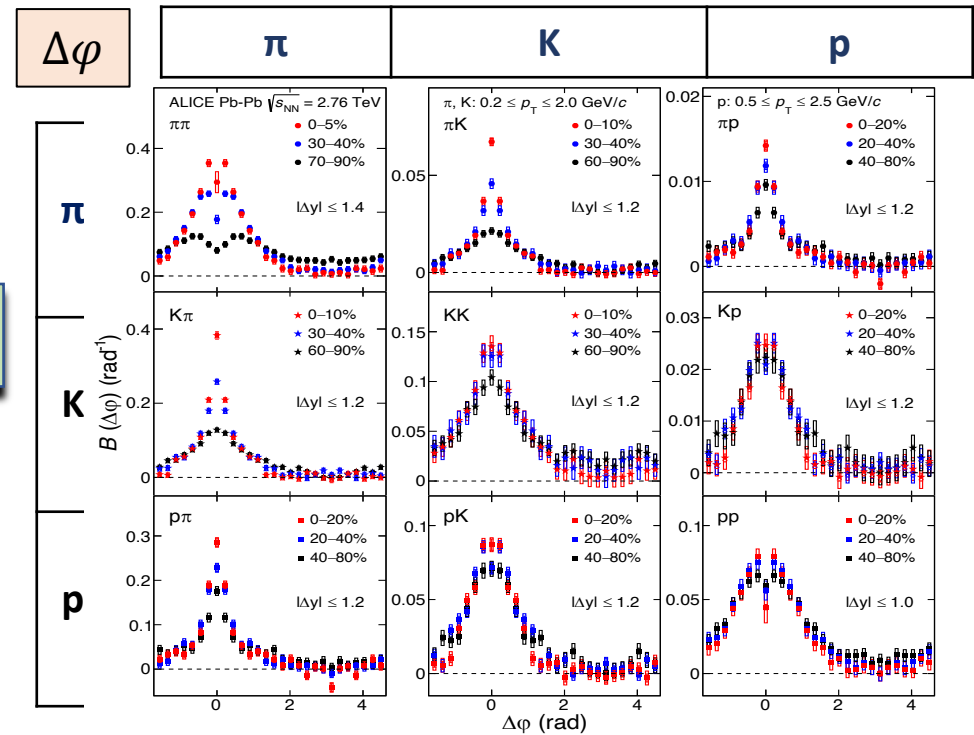
		h	π	k	p
Q	h	✓			
Q	π		?	?	?
Q	S		?	?	?
Q	B		?	?	?



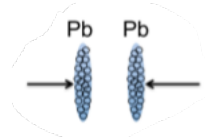
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New
Result



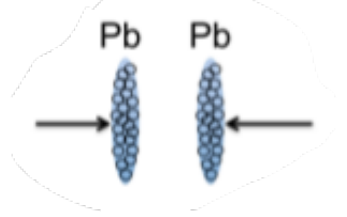
Balance Function (2-Particle correlations)



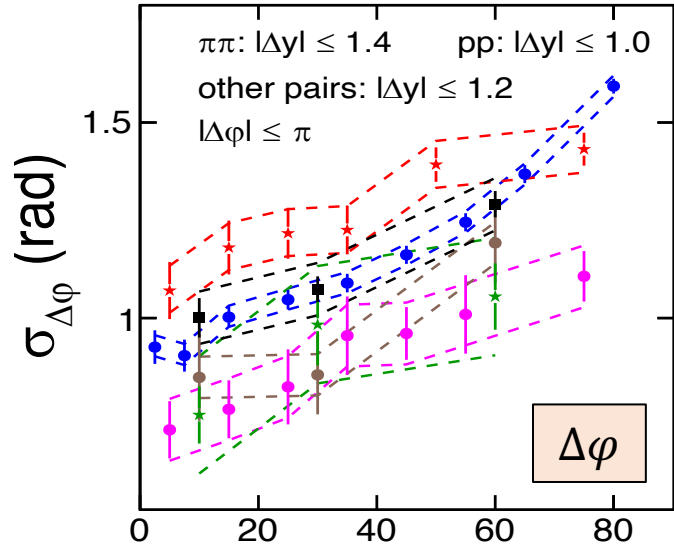
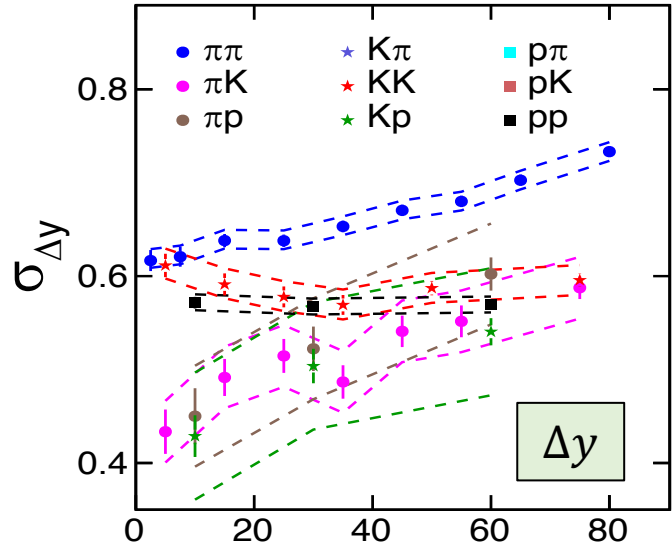
New Result

Species Dependence

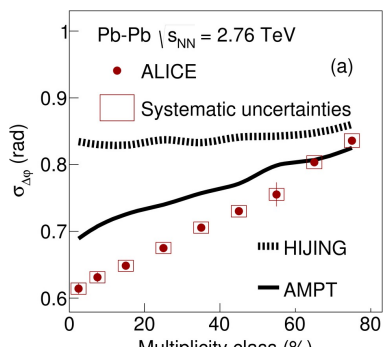
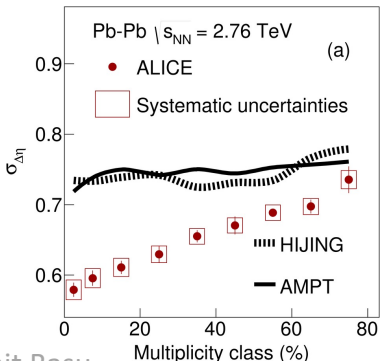
		h	π	k	p
Q	h	✓			
	π		?	?	?
	k		?	?	?
	p		?	?	?



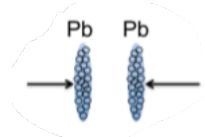
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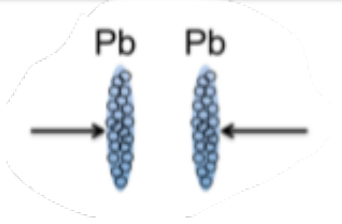
Balance Function (2-Particle correlations)



New Result

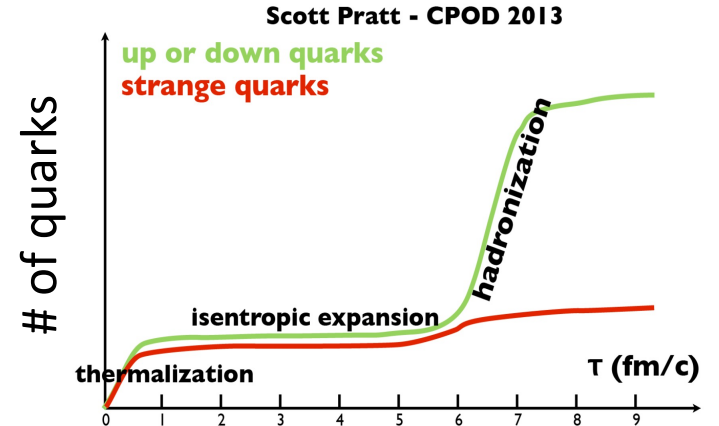
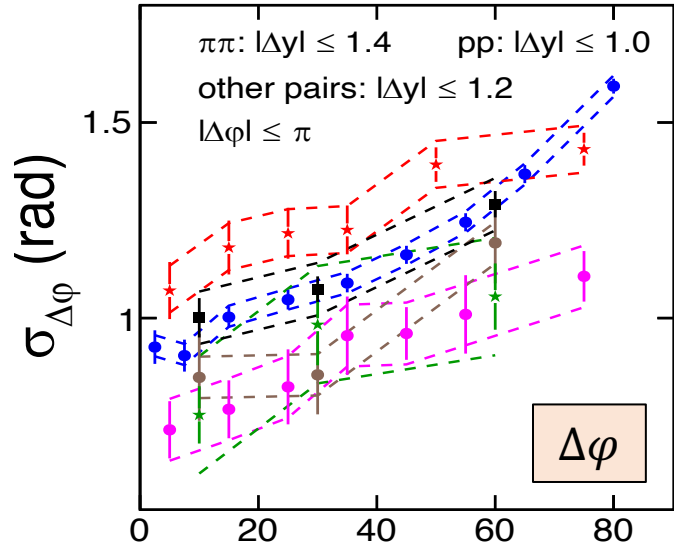
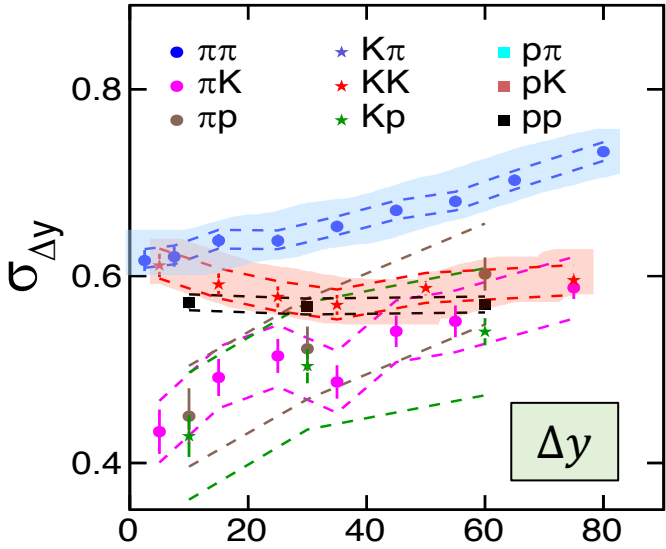
Species Dependence

		h	π	k	p
Q	h	✓			
Q	π		?	?	?
Q S	k		?	?	?
Q B	P		?	?	?



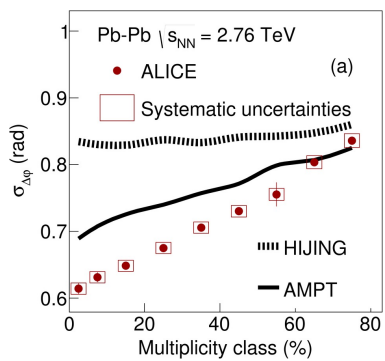
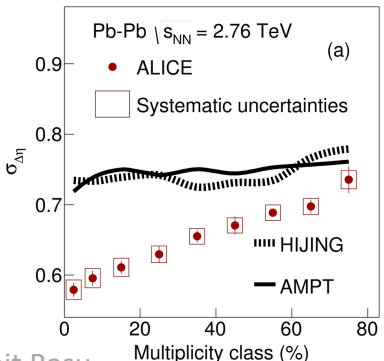
[arXiv:2110.06566](https://arxiv.org/abs/2110.06566) [nucl-ex]

- KK and pp widths no centrality dependence
- $\pi\pi$ and cross-species pairs narrow towards central collisions.

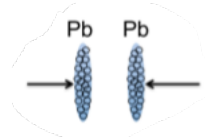


- Azimuthal narrowing for all species \rightarrow radial flow focusing
- Qualitatively consistent with radial flow and two-wave quark production

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Balance Function (2-Particle correlations)



New Result

Species Dependence

		h	π	k	p
Q	h	✓			
Q	π		?	?	?
Q	S		?	?	?
Q	B		?	?	?

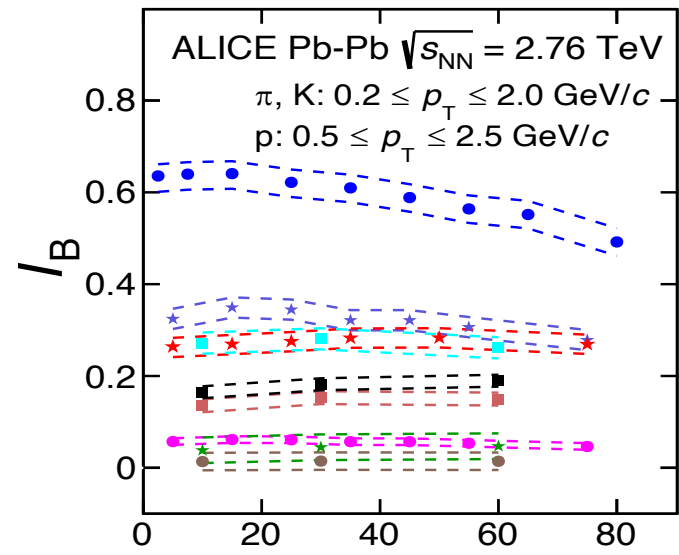
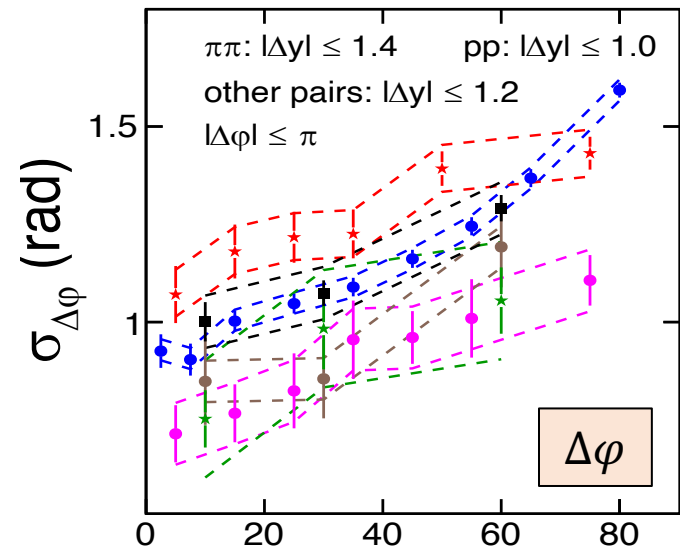
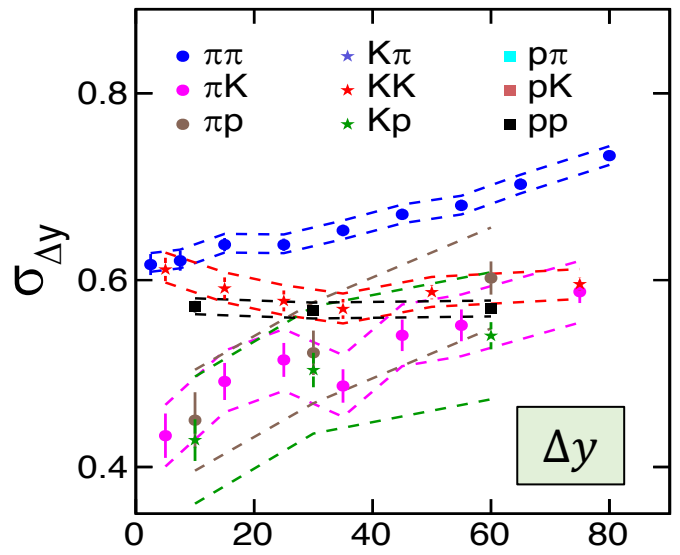
$$I_B = \int_{-\infty}^{\infty} B_{\pi^{\pm}h^{\pm}}(d\eta, d\phi) d\eta d\phi dp_T = 1$$

$$\equiv \int_{-\infty}^{\infty} B_{\pi^{\pm}\pi^{\pm}}(d\eta, d\phi) d\eta d\phi dp_T + \int_{-\infty}^{\infty} B_{\pi^{\pm}K^{\pm}}(d\eta, d\phi) d\eta d\phi dp_T + \int_{-\infty}^{\infty} B_{\pi^{\pm}p(\bar{p})}(d\eta, d\phi) d\eta d\phi dp_T$$

$$I_B = \int_{d\eta=-1.6}^{d\eta=1.6} \int_{d\phi=-\frac{\pi}{2}}^{\frac{3\pi}{2}} B_{\pi^{\pm}h^{\pm}}(d\eta, d\phi) d\eta d\phi dp_T \lesssim 1$$

arXiv:2110.06566 [nucl-ex]

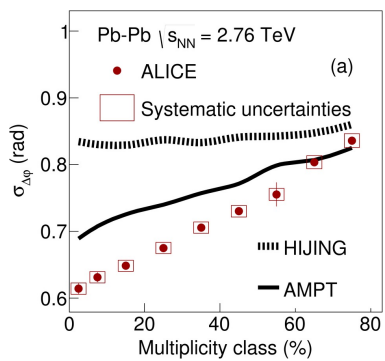
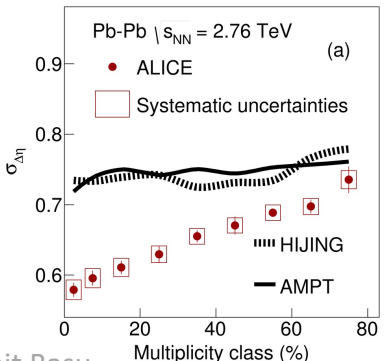
BF Integral



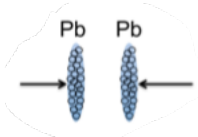
= Pairing Probabilities

Pairing probabilities are very different from single hadron ratios.
 → Kπ not larger than KK by K/π ratio;
 pp larger than pK.

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Balance Function (2-Particle correlations)



New Result

Species Dependence

		h	π	k	p
Q	h	✓			
Q	π		?	?	?
Q	S		?	?	?
Q	B		?	?	?

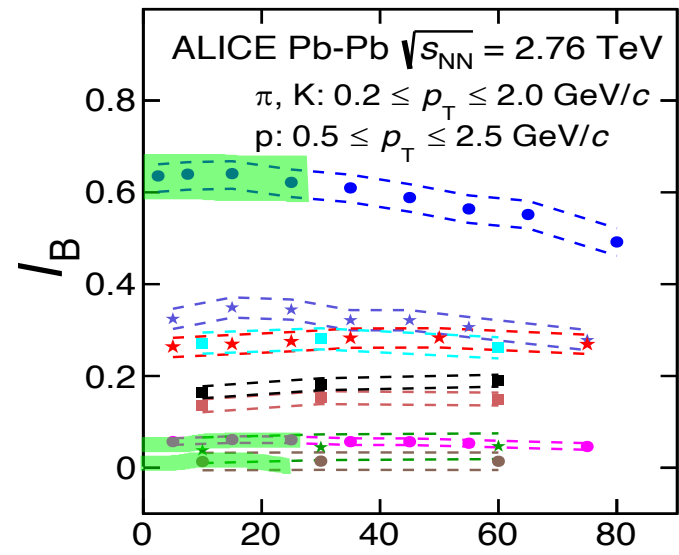
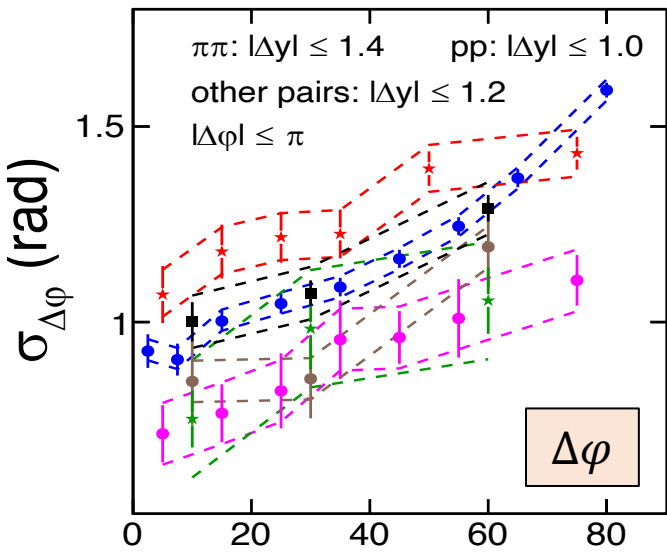
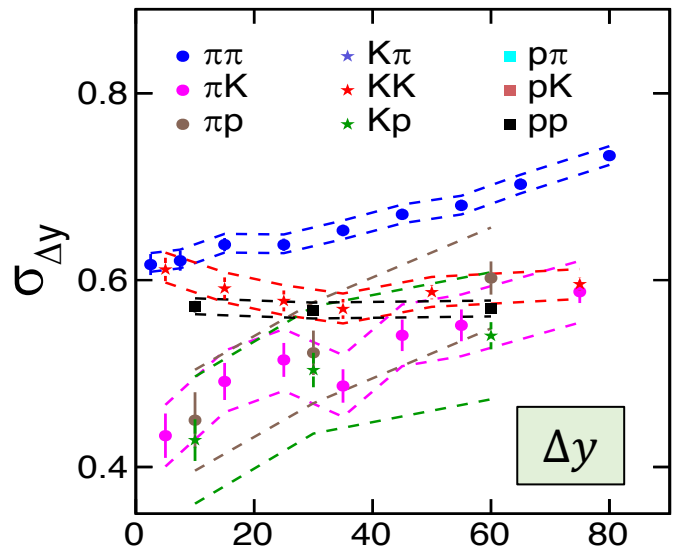
$$I_B = \int_{-\infty}^{\infty} B_{\pi^{\pm}h^{\pm}}(d\eta, d\phi) d\eta d\phi dp_T = 1$$

$$\equiv \int_{-\infty}^{\infty} B_{\pi^{\pm}\pi^{\pm}}(d\eta, d\phi) d\eta d\phi dp_T + \int_{-\infty}^{\infty} B_{\pi^{\pm}K^{\pm}}(d\eta, d\phi) d\eta d\phi dp_T + \int_{-\infty}^{\infty} B_{\pi^{\pm}p(\bar{p})}(d\eta, d\phi) d\eta d\phi dp_T$$

$$I_B = \int_{d\eta=-1.6}^{d\eta=1.6} \int_{d\phi=-\frac{\pi}{2}}^{d\phi=\frac{3\pi}{2}} B_{\pi^{\pm}h^{\pm}}(d\eta, d\phi) d\eta d\phi dp_T \lesssim 1$$

arXiv:2110.06566 [nucl-ex]

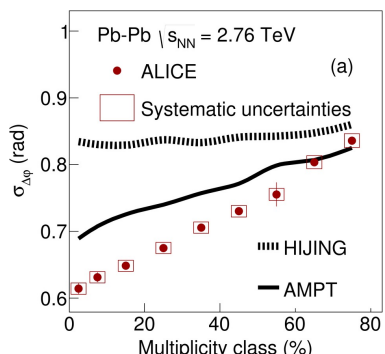
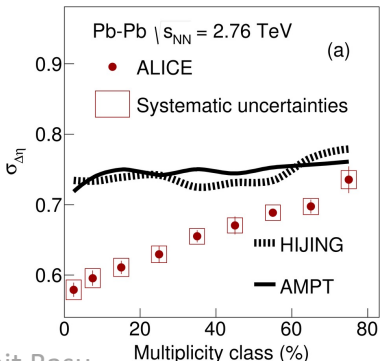
BF Integral



= Pairing Probabilities

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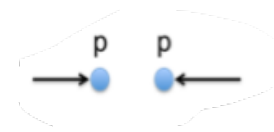
Eur. Phys. J. C 76 (2016) 86



$$I_B = \int_{d\eta=-1.6}^{d\eta=1.6} \int_{d\phi=-\frac{\pi}{2}}^{d\phi=\frac{3\pi}{2}} B_{\pi^{\pm}h^{\pm}}(d\eta, d\phi) d\eta d\phi dp_T \sim 0.8$$

→ Low p_T particles mostly balance by low p_T particles

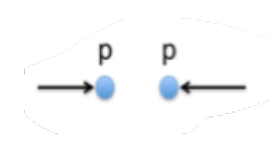
Ξ -hadron correlations in pp collisions at $\sqrt{s} = 13$ TeV



New Result

Trigger	$\Xi(\text{ssd})$	$\Xi(\text{ssd})$	$\Xi(\text{ssd})$	$\Xi(\text{ssd})$	$\Xi(\text{ssd})$
Associated (-)	$\pi^+(u\bar{d})$	$K^+(u\bar{s})$	$\bar{p}(\bar{u}\bar{u}\bar{d})$	$\bar{\Lambda}(\bar{u}\bar{d}\bar{s})$	$\bar{\Xi}(\bar{s}\bar{s}\bar{d})$
Background	$\pi^-(\bar{u}d)$	$K^-(u\bar{s})$	$p(uud)$	$\Lambda(uds)$	$\Xi(\text{ssd})$
(Subtracting same quantum number correlation)					

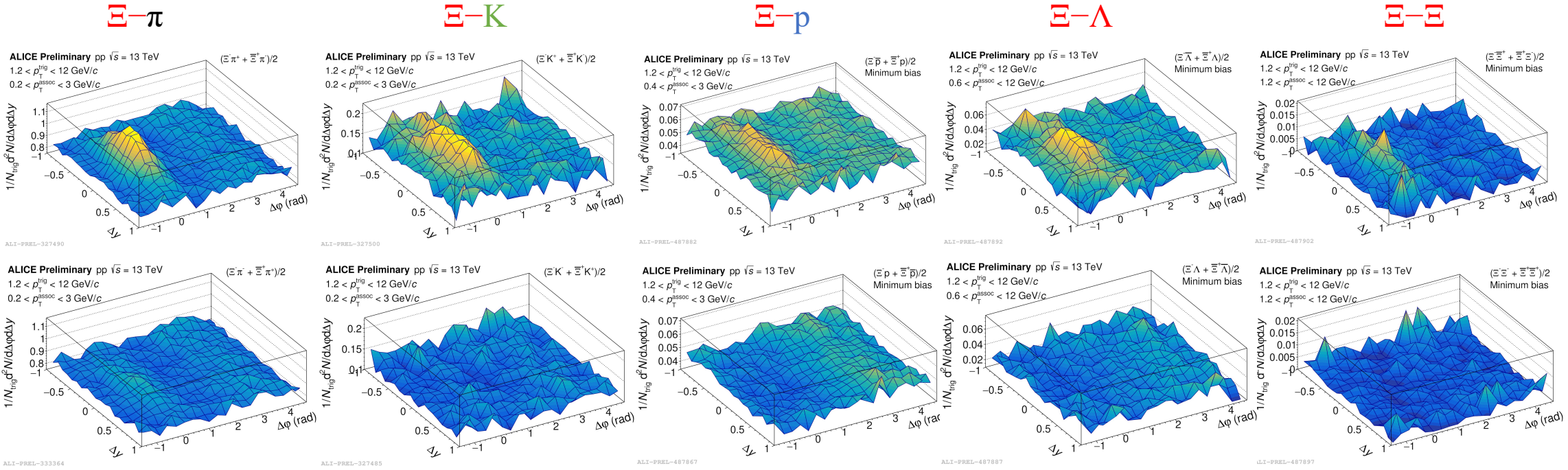
Ξ -hadron correlations in pp collisions at $\sqrt{s} = 13$ TeV



New Result

OS or US

SS or LS



→ Results challenge hadronization models:

Lund string breaking (PYTHIA)

standard:

with junctions:

with ropes:

Adolfsson et al. Eur. Phys. J. A 56, 288 (2020), Bierlich et al. J. High Energy. Phys. 2015, 148

EPOS:

central AA

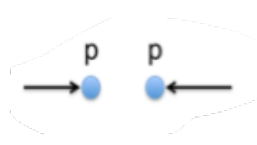
peripheral AA high mult pp,pA

low mult pp

core+corona model:
 core => hydro => flow + statistical decay
 corona => string decay

K. Werner. hal-02434245 (2019)

Ξ -hadron correlations in pp collisions at $\sqrt{s} = 13$ TeV

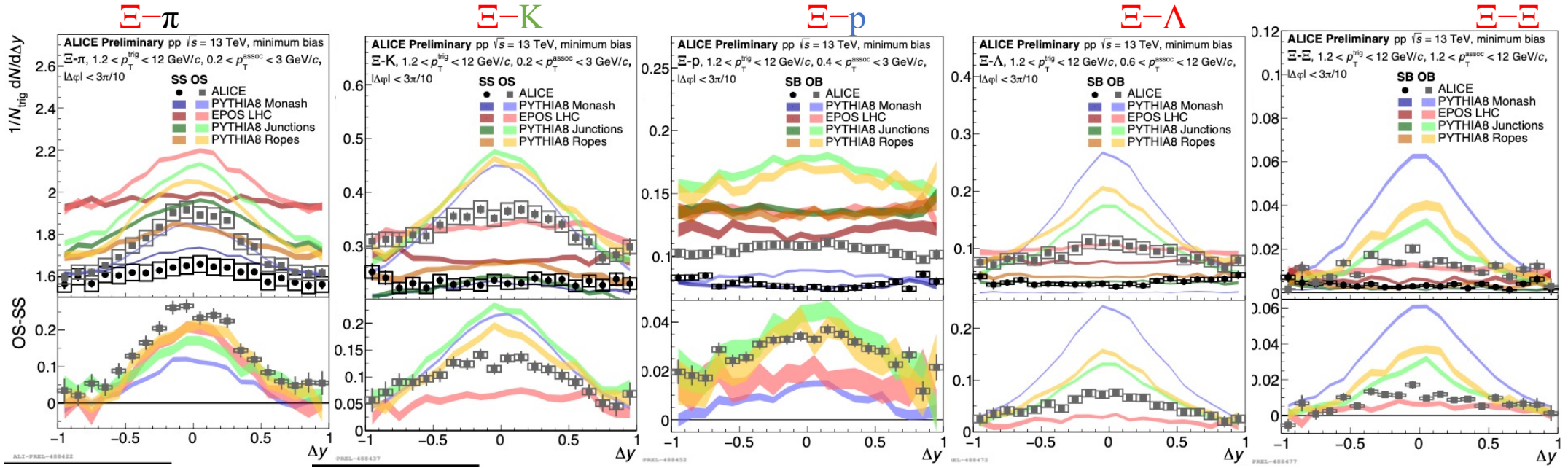


New
Result

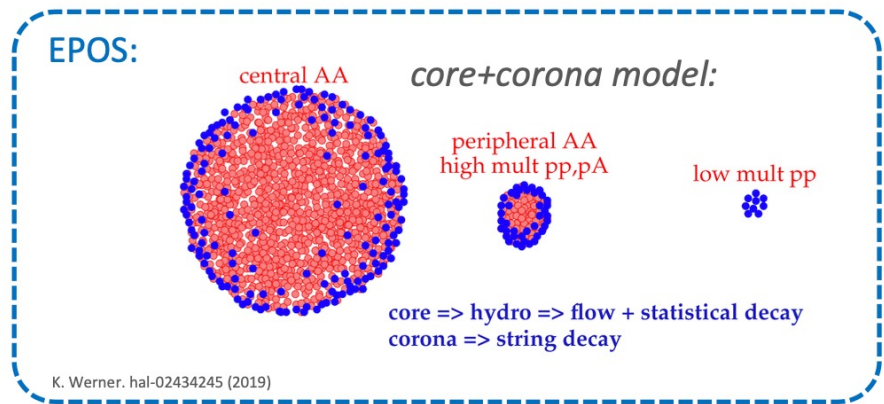
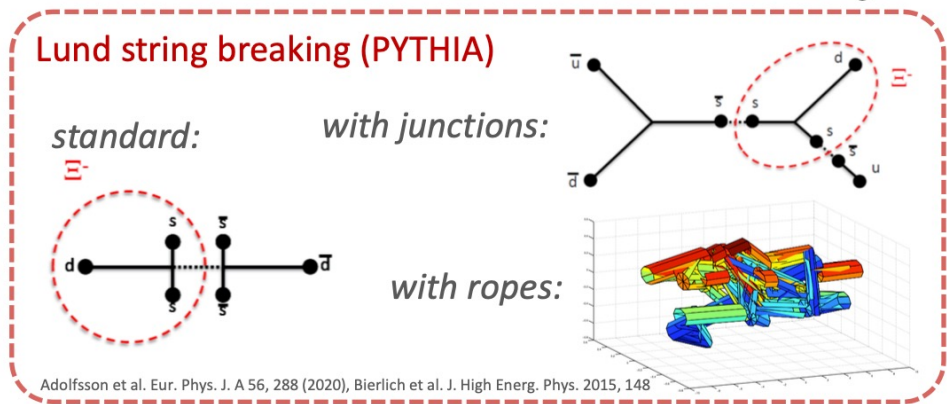
OS or US

OS-SS

Δy

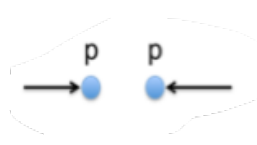


→ Results challenge hadronization models:



- Wider peak in data than in PYTHIA
- strange quarks produced at an earlier time
- Local conservation of quantum numbers → not implemented in EPOS

Ξ -hadron correlations in pp collisions at $\sqrt{s} = 13$ TeV

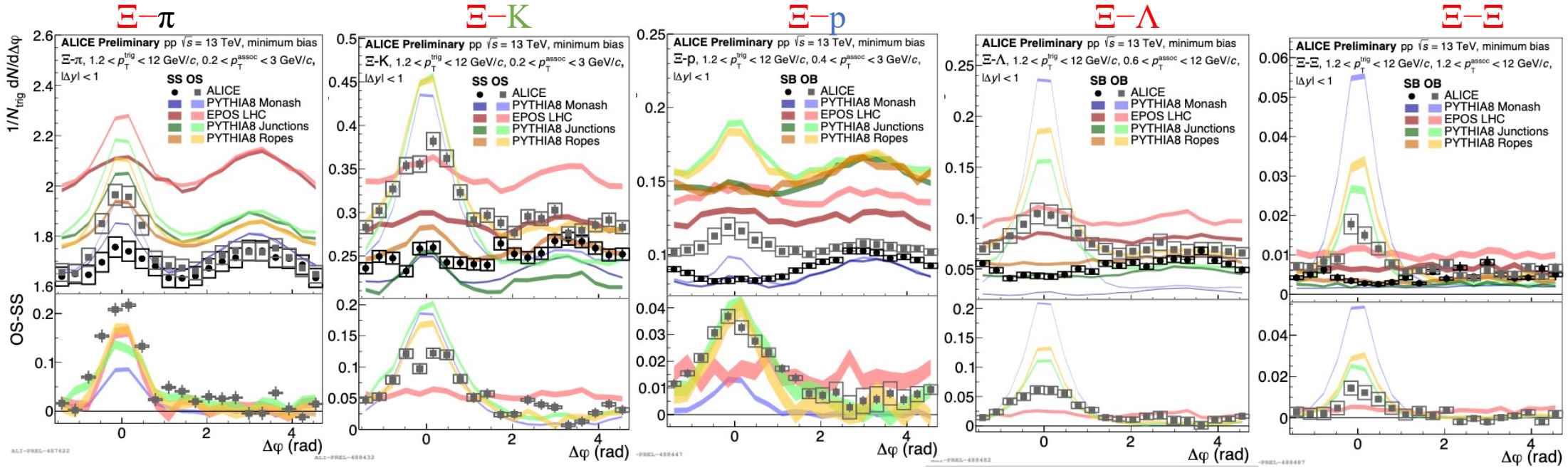


New
Result

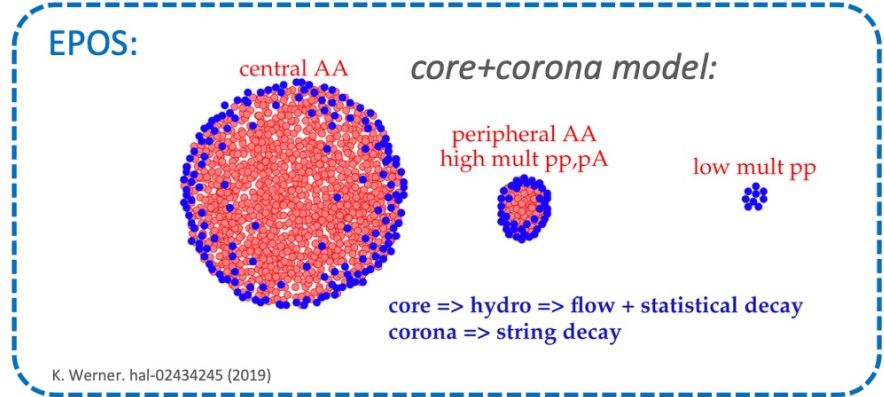
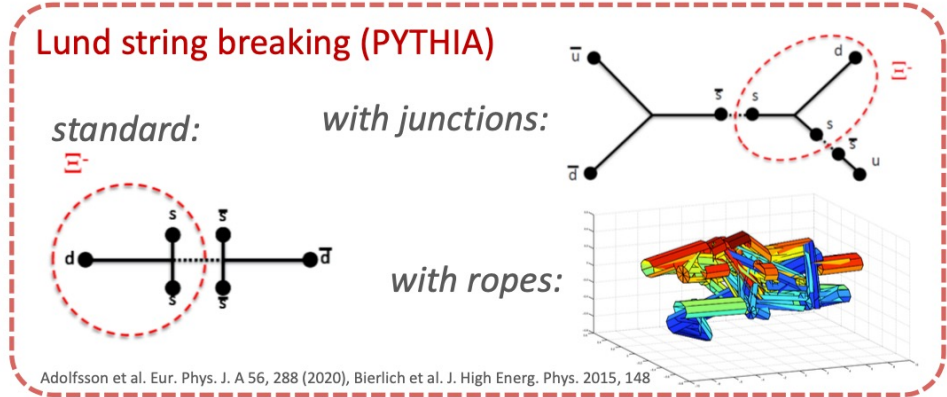
OS or US

OS-SS

$\Delta\phi$

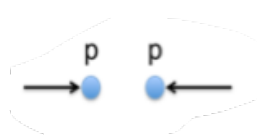


→ Results challenge hadronization models:



- Wider peak in data than in PYTHIA
- strange quarks produced at an earlier time
- Local conservation of quantum numbers → not implemented in EPOS
- Junction model reduces peak amplitude → favors this baryon production mechanism over diquark breaking

Ξ -hadron correlations in pp collisions at $\sqrt{s} = 13$ TeV

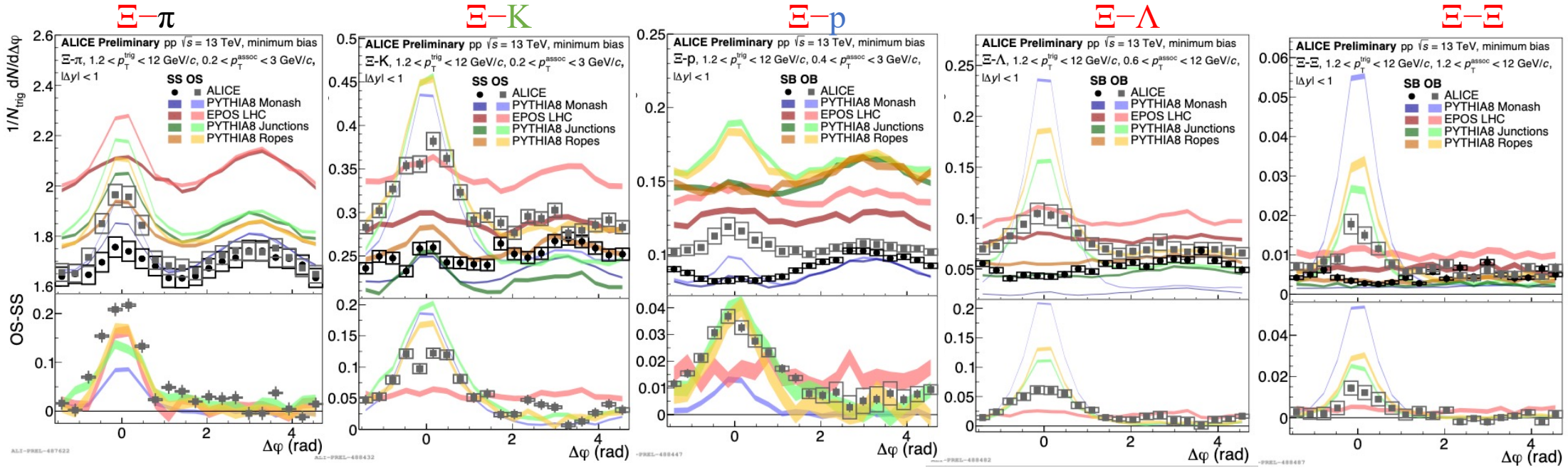


New Result

OS or US

OS-SS

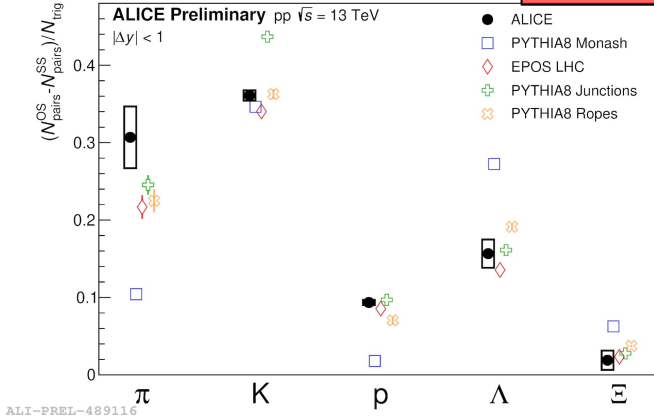
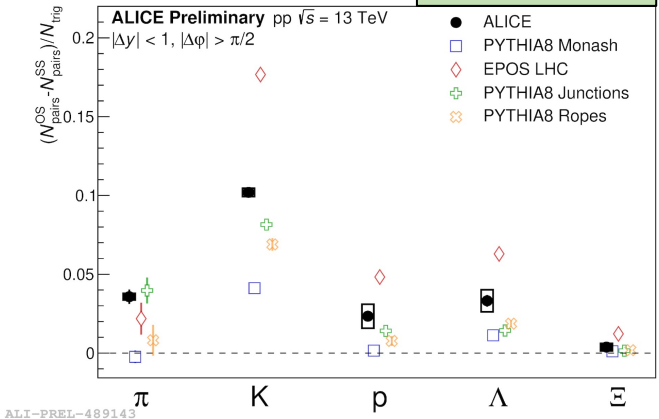
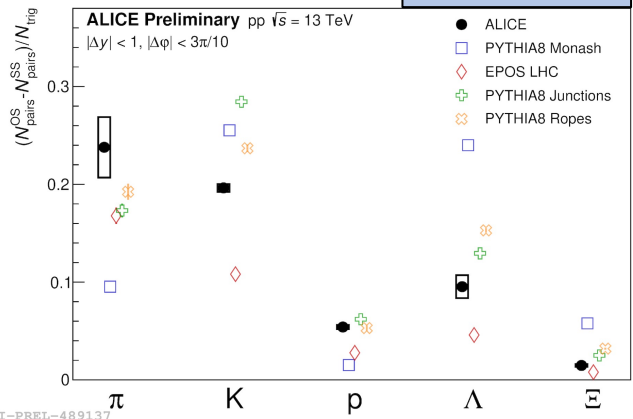
$\Delta\phi$



Near Side

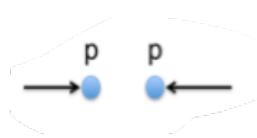
Away Side

Total

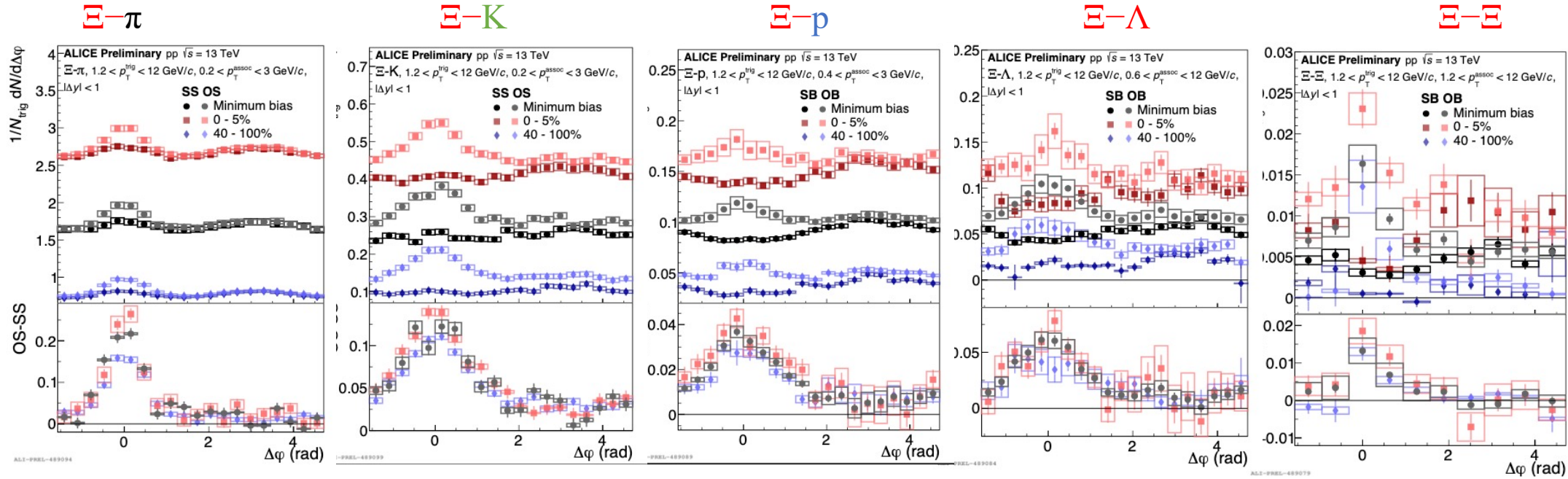


How electric charge, strangeness, and baryon number of the trigger are balanced in phase space.

Ξ -hadron correlations in pp collisions at $\sqrt{s} = 13$ TeV



New Result



Multiplicity dependence

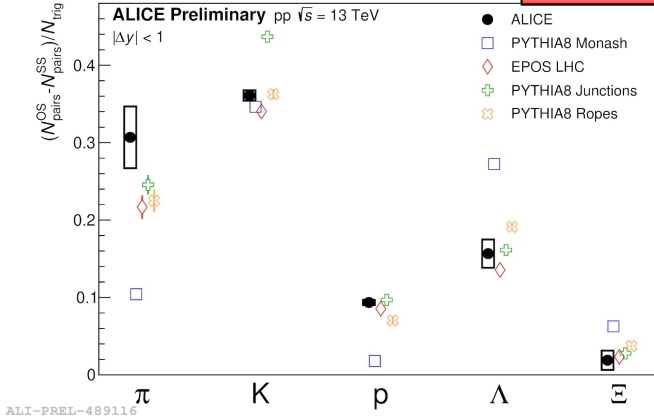
OS-SS

$\Delta\phi$

Red = high multiplicity
 Grey = minimum bias
 Blue = low multiplicity

- Similar relative difference seen for π , K, p
- No evidence of multiplicity dependence on Ξ production mechanism \rightarrow underlying quark distributions are the same
- quantitatively points towards common origin of baryon production

Total

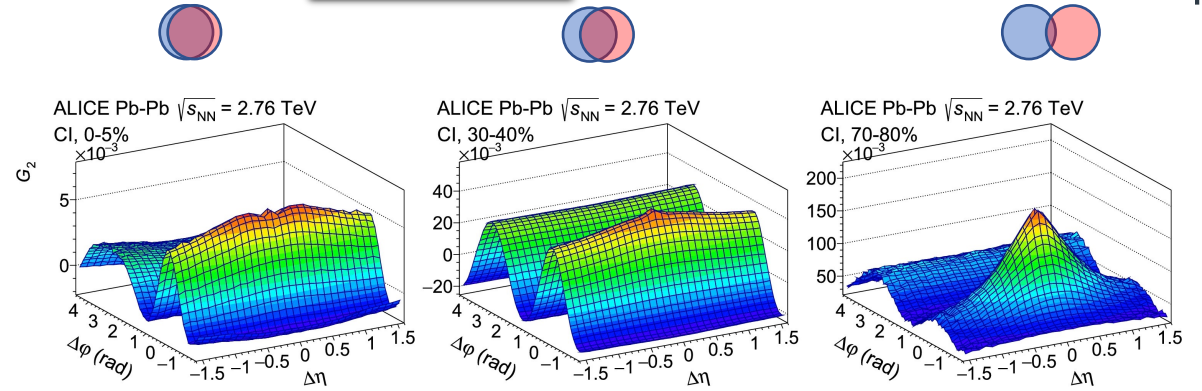


How electric charge, strangeness, and baryon number of the trigger are balanced in phase space.

2-Particle Transverse momentum correlations

- Look for the diffusion of transverse momentum currents through their correlations

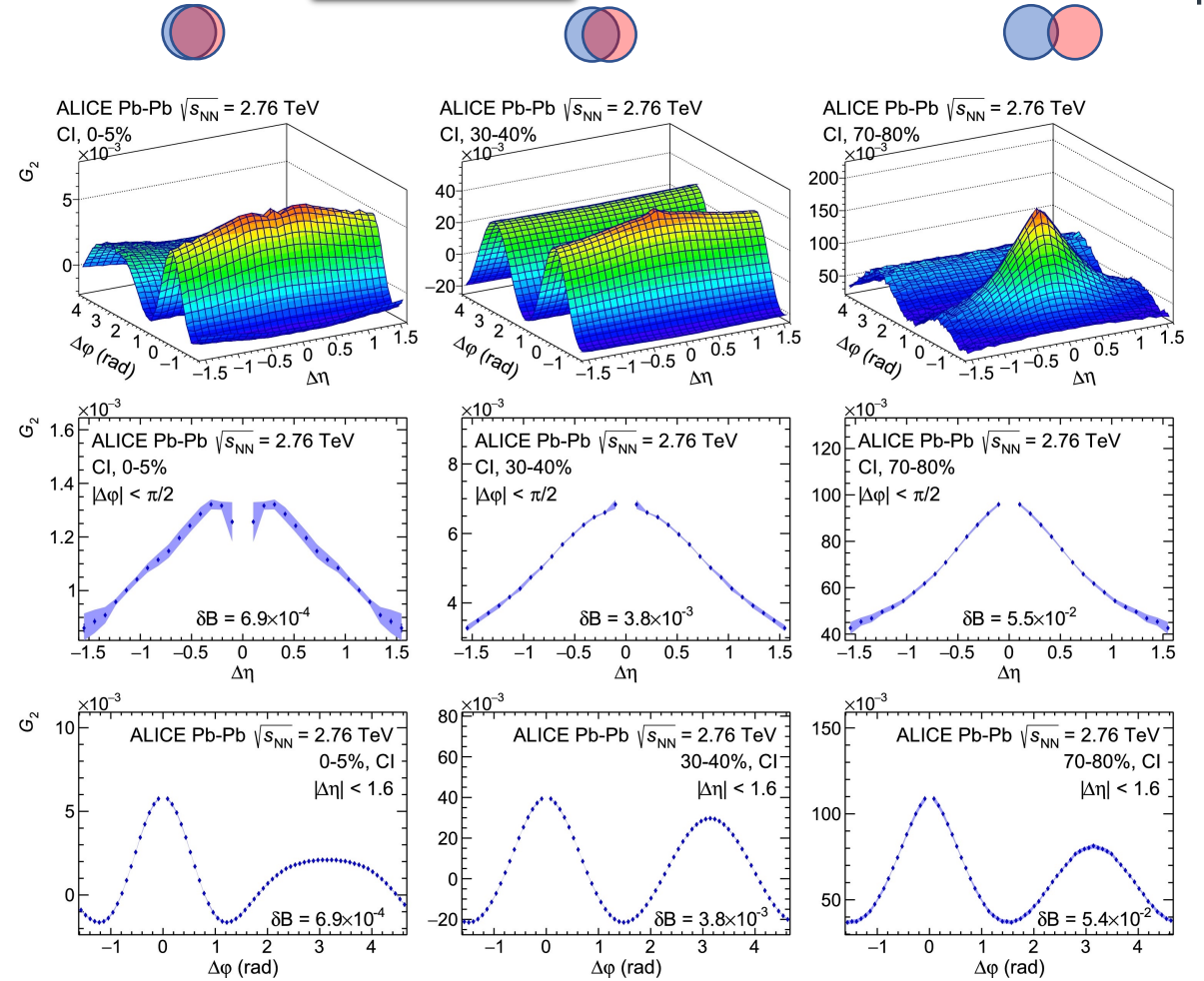
$$G_2(\Delta\eta, \Delta\varphi) = \frac{\langle \sum_i^{n_{1,1}} \sum_{j \neq i}^{n_{1,2}} p_{T,i} p_{T,j} \rangle}{\langle n_{1,1} \rangle \langle n_{1,2} \rangle} - \langle p_{T,1} \rangle \langle p_{T,2} \rangle$$



2-Particle Transverse momentum correlations

- Look for the diffusion of transverse momentum currents through their correlations

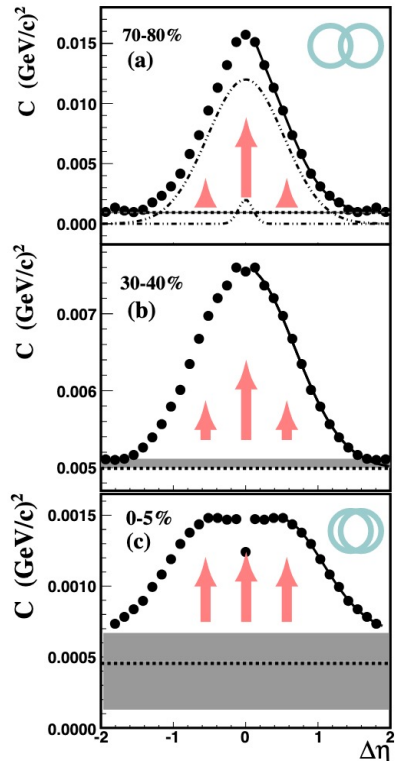
$$G_2(\Delta\eta, \Delta\varphi) = \frac{\langle \sum_i^{n_{1,1}} \sum_{j \neq i}^{n_{1,2}} p_{T,i} p_{T,j} \rangle}{\langle n_{1,1} \rangle \langle n_{1,2} \rangle} - \langle p_{T,1} \rangle \langle p_{T,2} \rangle$$



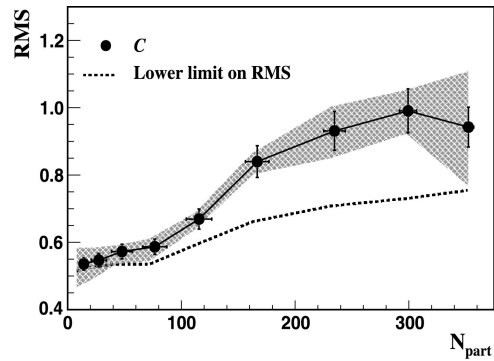
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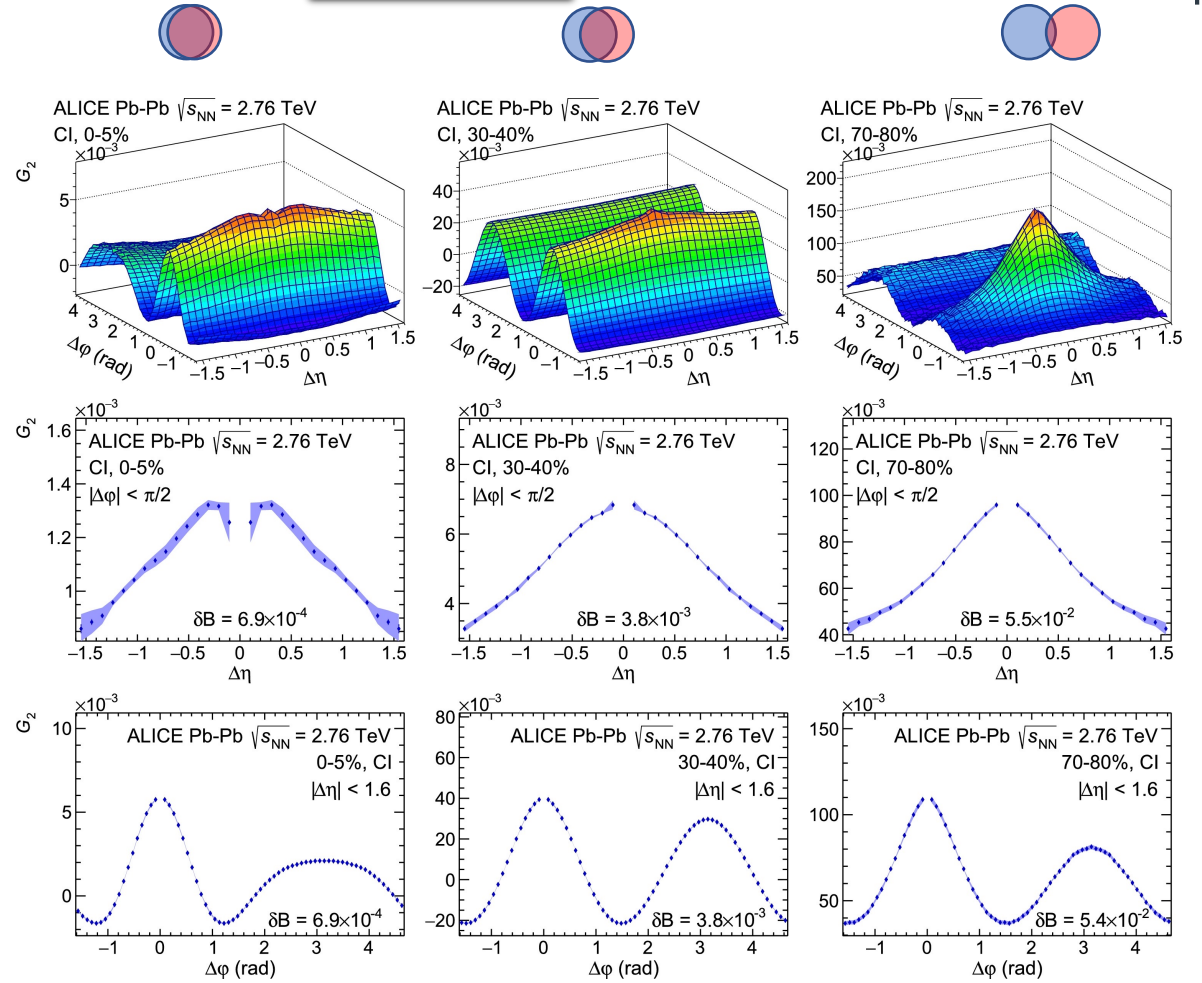
$$G_2(\Delta\eta, \Delta\varphi) = \frac{\langle \sum_i^{n_{1,1}} \sum_{j \neq i}^{n_{1,2}} p_{T,i} p_{T,j} \rangle}{\langle n_{1,1} \rangle \langle n_{1,2} \rangle} - \langle p_{T,1} \rangle \langle p_{T,2} \rangle$$



Au-Au $\sqrt{s_{NN}} = 200$ GeV



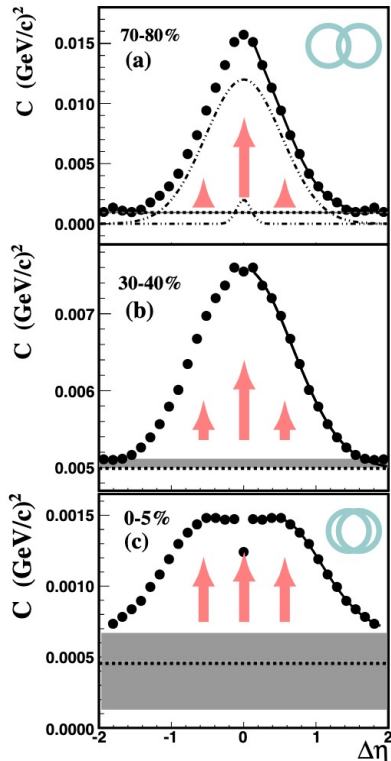
STAR, PLB 704, 467-473 (2011)



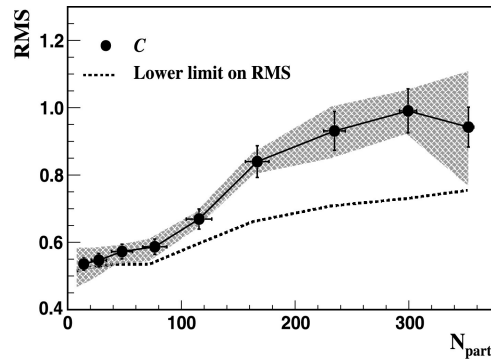
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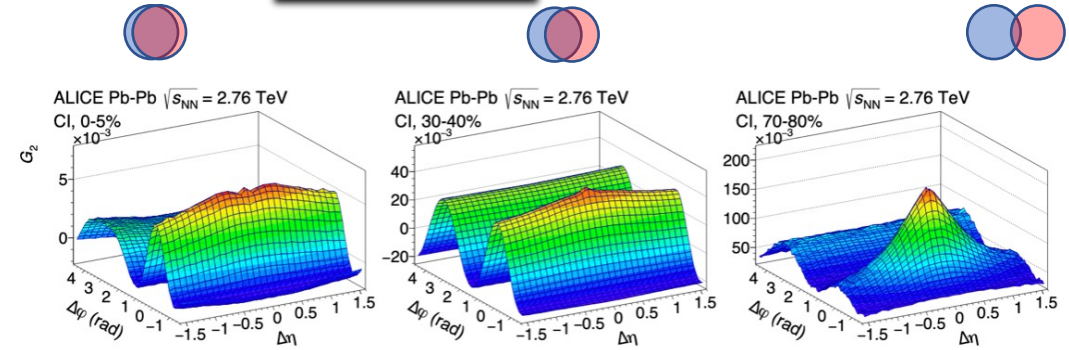
$$G_2(\Delta\eta, \Delta\varphi) = \frac{\langle \sum_i^{n_{1,1}} \sum_{j \neq i}^{n_{1,2}} p_{T,i} p_{T,j} \rangle}{\langle n_{1,1} \rangle \langle n_{1,2} \rangle} - \langle p_{T,1} \rangle \langle p_{T,2} \rangle$$



Au-Au $\sqrt{s_{NN}} = 200$ GeV

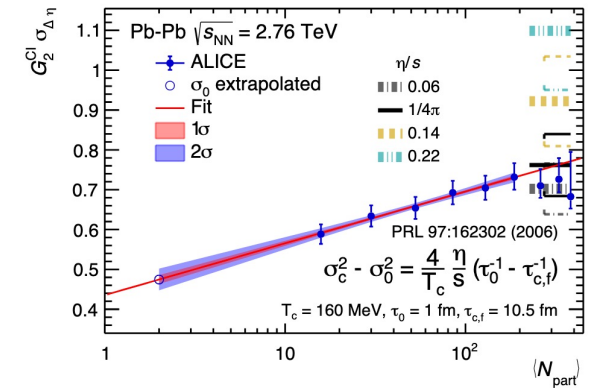
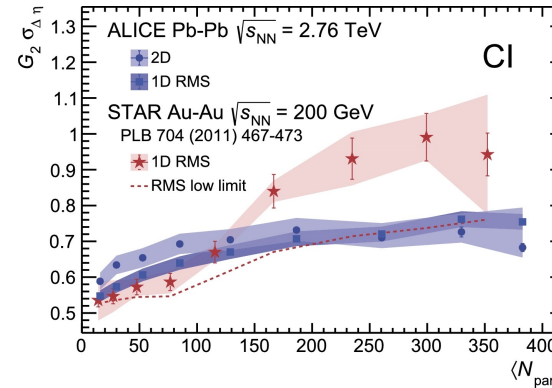


STAR, PLB 704, 467-473 (2011)

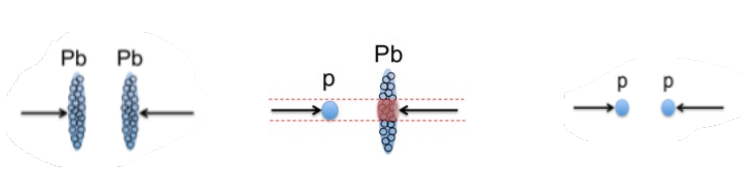


Fit with a generalized Gaussian in two dimensions

$$F(\Delta\eta, \Delta\varphi) = B + \sum_{n=2}^6 a_n \times \cos(n\Delta\varphi) + A \times \frac{\gamma_{\Delta\eta}}{2\omega_{\Delta\eta} \Gamma\left(\frac{1}{\gamma_{\Delta\eta}}\right)} e^{-\left|\frac{\Delta\eta}{\omega_{\Delta\eta}}\right|^{\gamma_{\Delta\eta}}} \times \frac{\gamma_{\Delta\varphi}}{2\omega_{\Delta\varphi} \Gamma\left(\frac{1}{\gamma_{\Delta\varphi}}\right)} e^{-\left|\frac{\Delta\varphi}{\omega_{\Delta\varphi}}\right|^{\gamma_{\Delta\varphi}}}$$



How about small systems? p+p & p+Pb collisions?

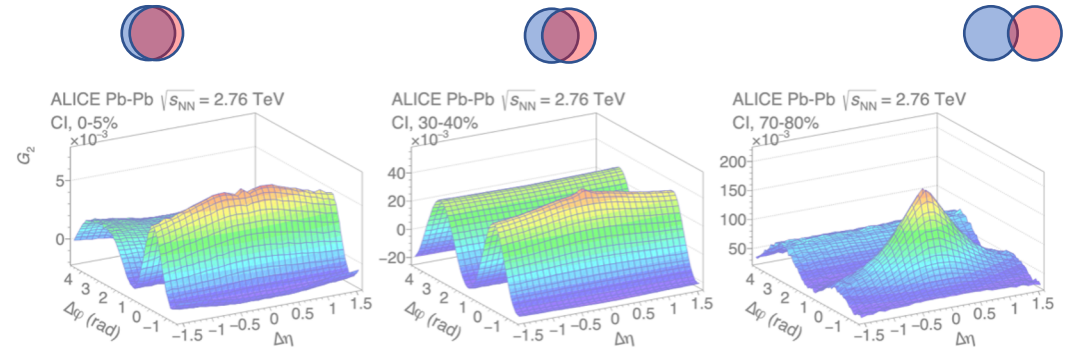


Pb-Pb 2.76 TeV

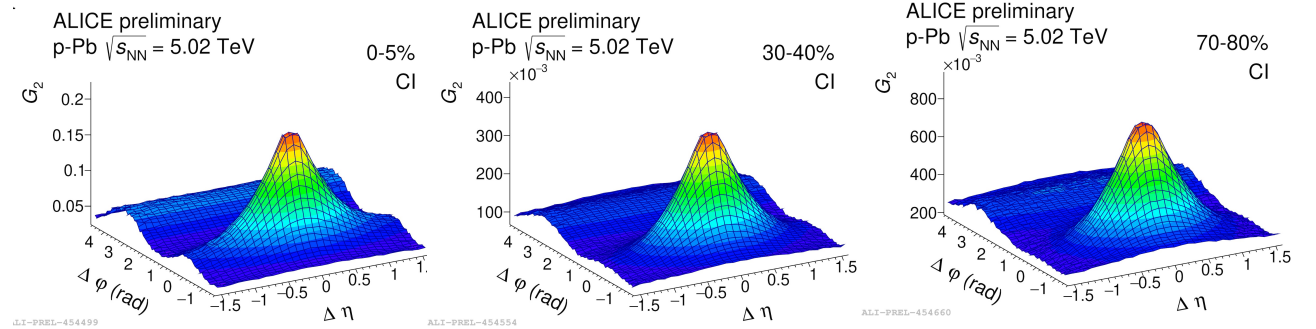
2-Particle Transverse momentum correlations

- Look for the diffusion of transverse momentum currents through their correlations

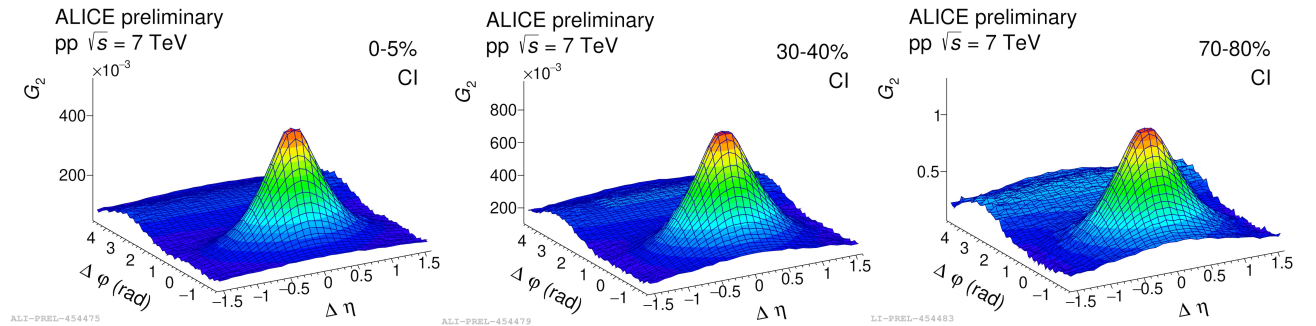
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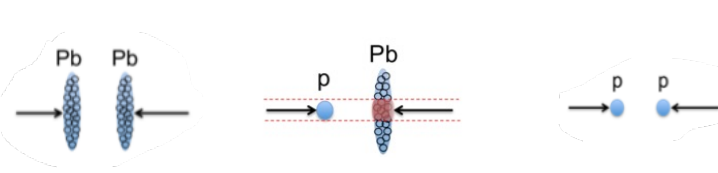


p-Pb 5.02 TeV



pp 7 TeV





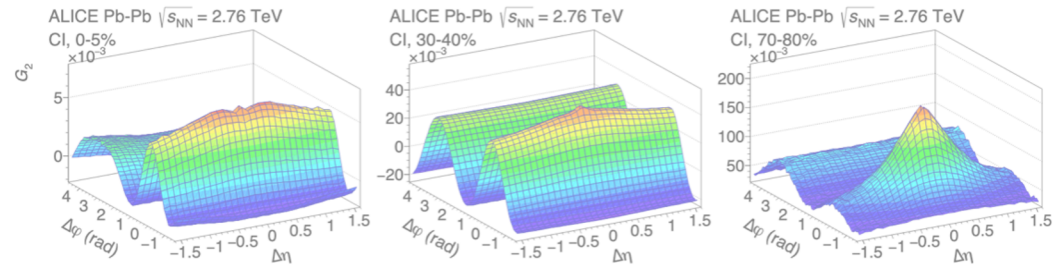
Pb-Pb 2.76 TeV

New Result

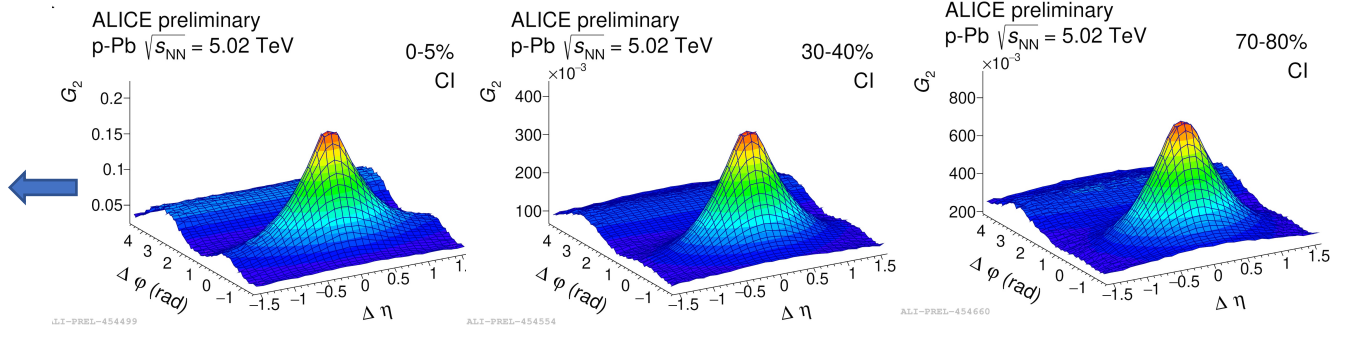
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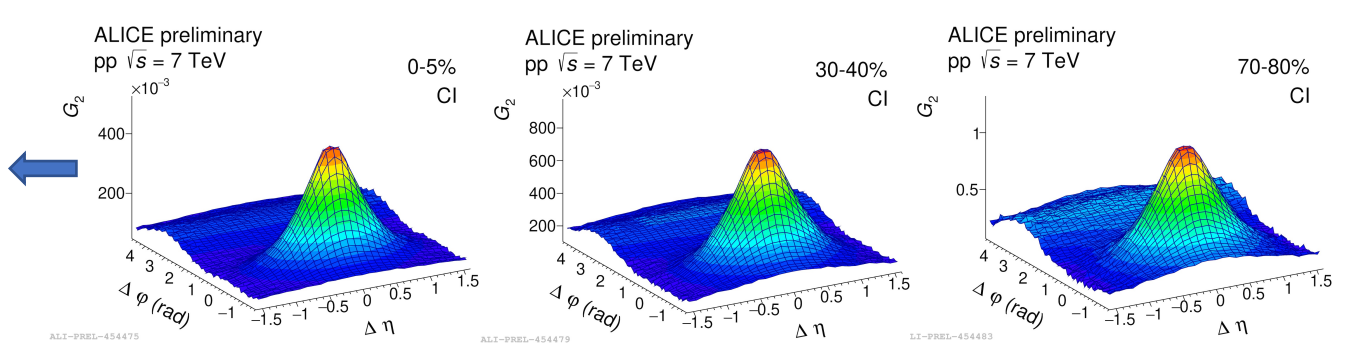
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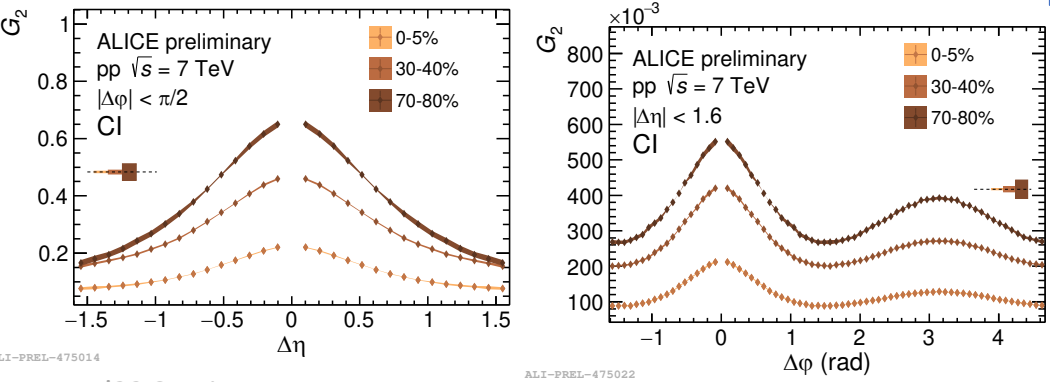
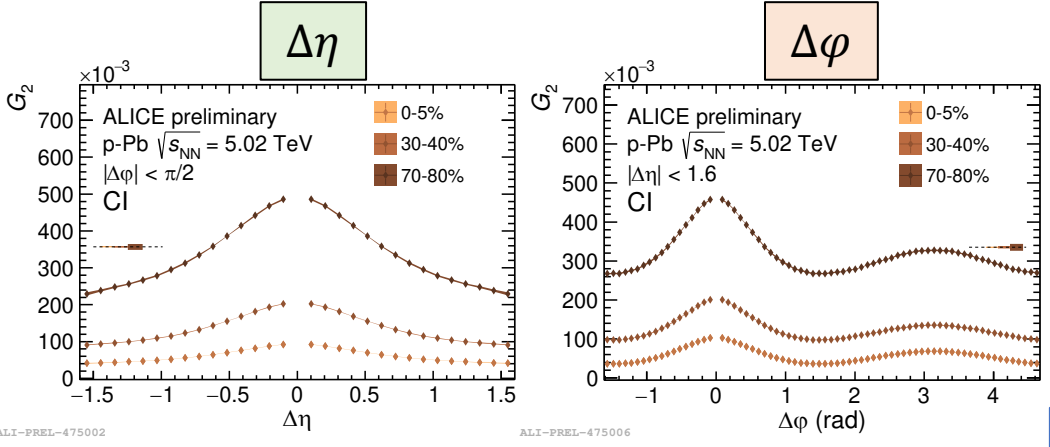
p-Pb 5.02 TeV

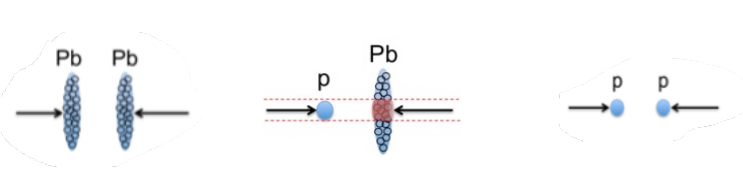


pp 7 TeV



Projection





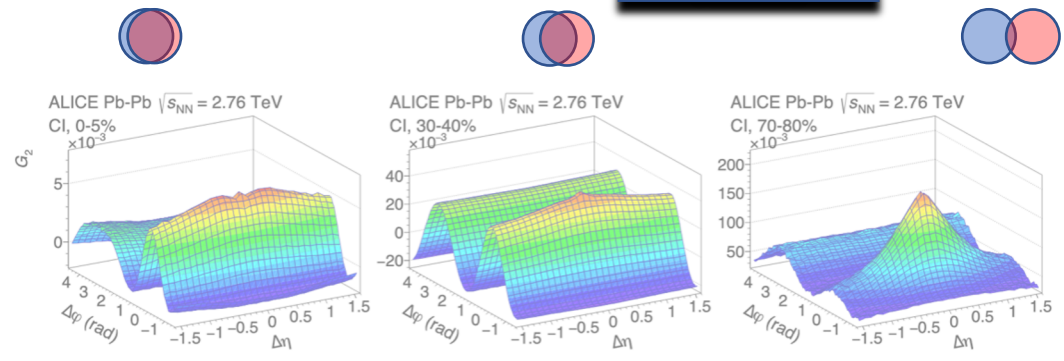
Pb-Pb 2.76 TeV

New Result

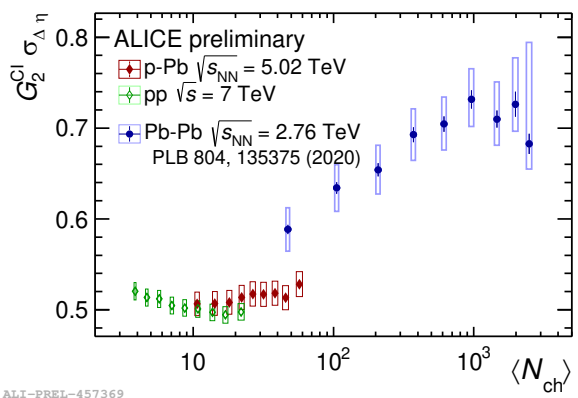
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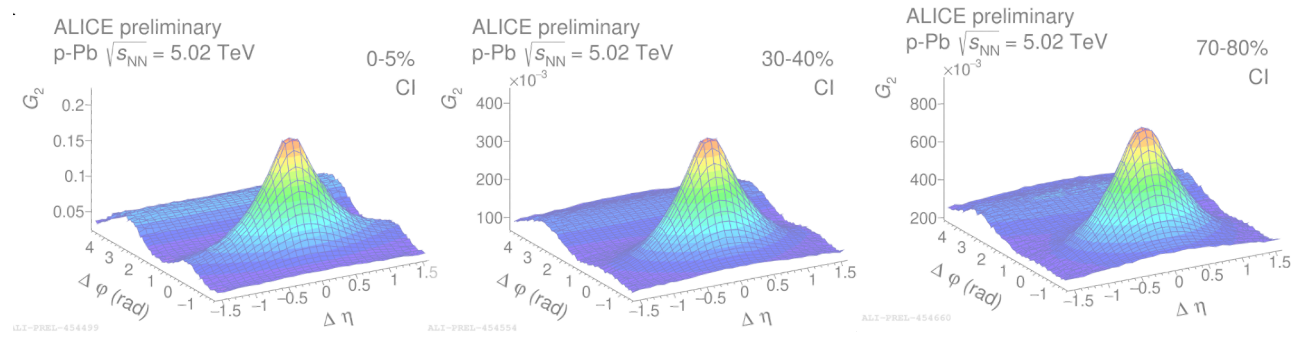
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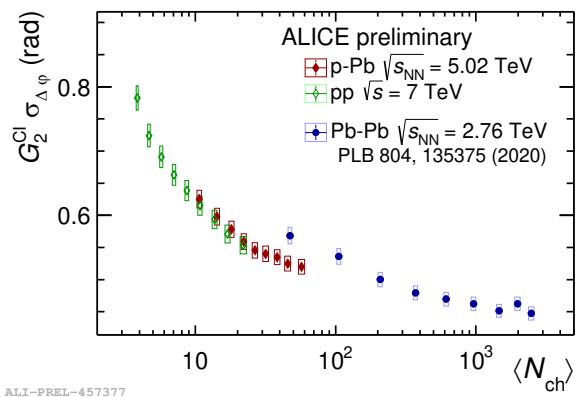
p-Pb 5.02 TeV



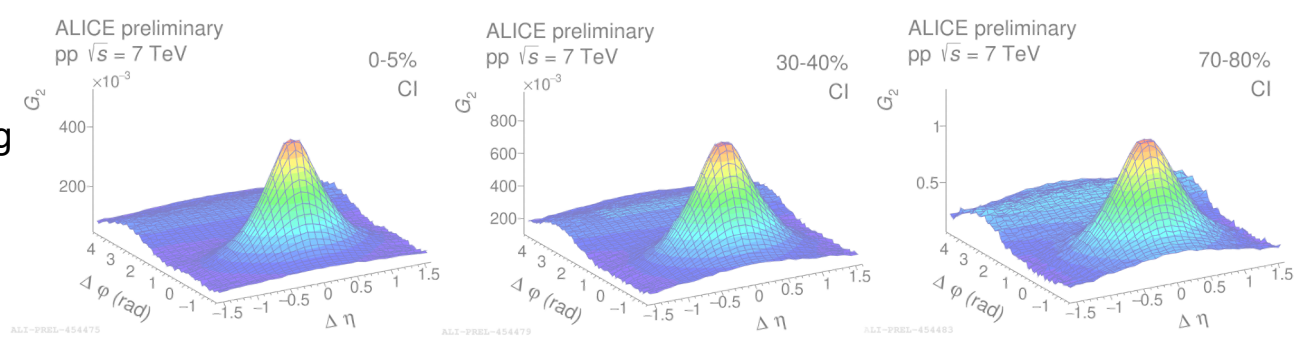
- Completely different longitudinal evolution

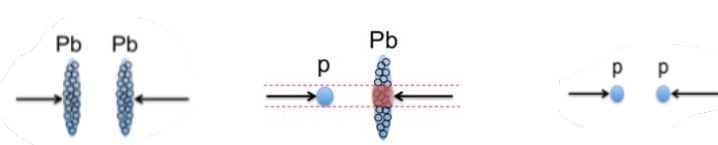


pp 7 TeV



- Azimuthal narrowing consistent along the three systems





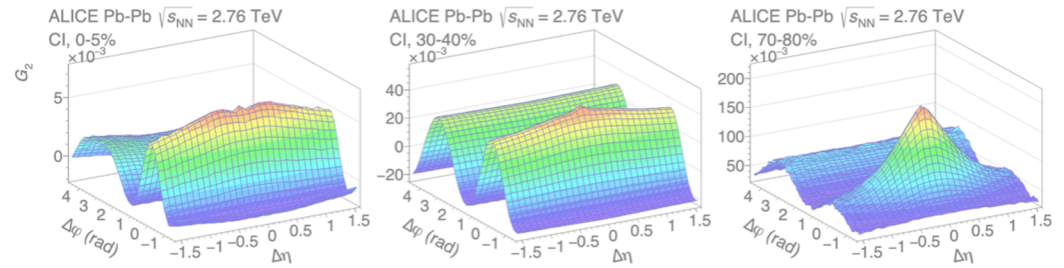
Pb-Pb 2.76 TeV

New Result

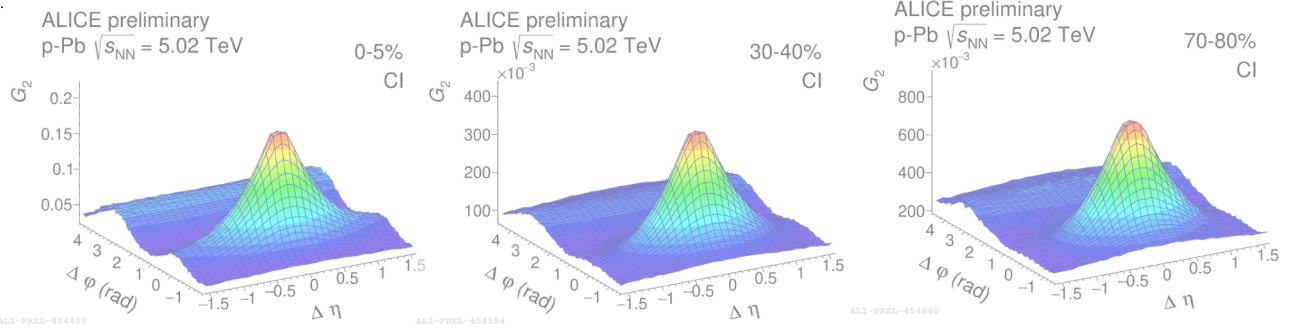
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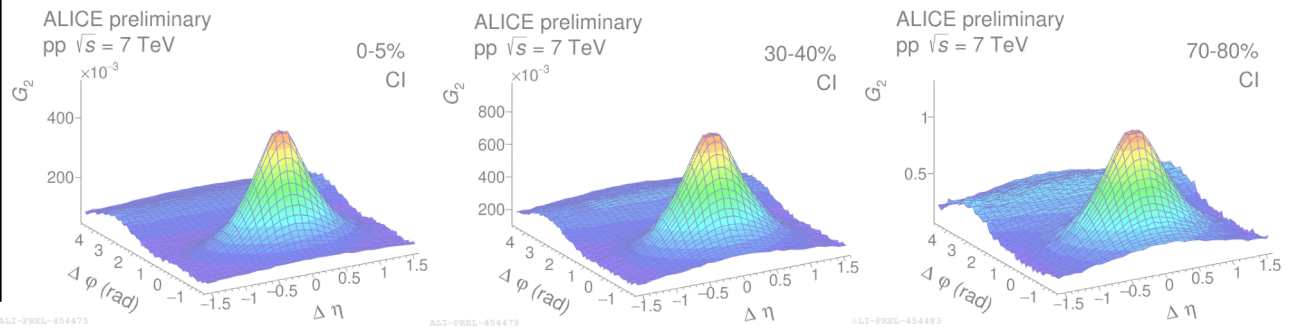
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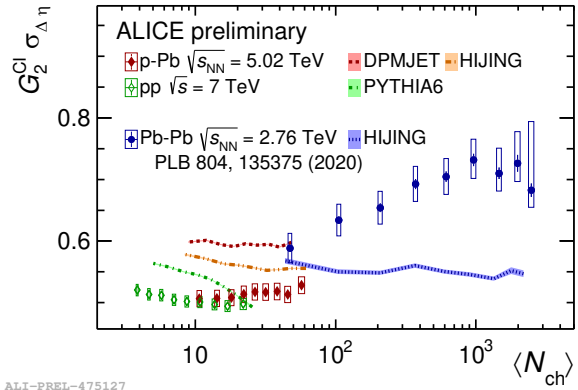
p-Pb 5.02 TeV



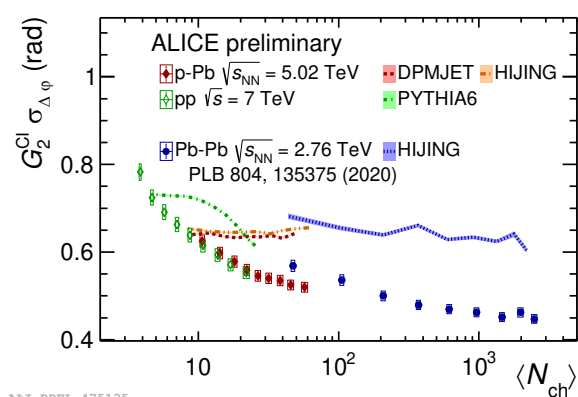
pp 7 TeV



MC Comparison



- Pythia8 Monash captures the azimuthal narrowing on both correlators



In small systems, \rightarrow fluid-like system produced is too small or too short lived (Or not produced quasi-equilibrated fluid)

- Balance function for Identified primary hadrons (π , K, p) pairs for Pb-Pb collision at $\sqrt{s_{NN}} = 2.76 \text{ TeV}$ are presented. Narrowing of azimuthal widths for all specie pairs \rightarrow Radial flow focusing (kinematic lensing), different width evolution behavior in $\Delta\eta$ \rightarrow Qualitatively consistent with radial flow and two-wave quark production mechanism.
- 2-particle correlation function for double strange baryon (Ξ -h) pairs for pp collision at $\sqrt{s} = 13 \text{ TeV}$ are presented. Multiplicity dependence very similar for all correlation measurements \rightarrow common origin of Ξ /strangeness production across multiplicity. Ξ -strangeness correlation peak is much wider in data than in PYTHIA \rightarrow Strange quarks are produced earlier in the event than from Lund string model alone. Local conservation of quantum numbers needs to be implemented in EPOS.
- Transverse momentum correlation for all collision system (pp, p-Pb, Pb-Pb) are presented. Azimuthal narrowing consistent in all three systems, changes behavior in the longitudinal dimension from narrowing to broadening when going from pp to p-Pb and to Pb-Pb \rightarrow Fluid-like system produced is too small or too short lived, in small system.
- ALICE continues to provide many interesting results on correlations and fluctuations, So Stay tuned!!! Currently, Run 3 preparations are ongoing, where many observables will benefit from more statistics and larger ALICE acceptance.

Thank you !!!